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AN UNBIASED 1.3 mm EMISSION LINE SURVEY OF THE PROTOPLANETARY DISK ORBITING LkCa 15

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ABSTRACT

The outer (>30 AU) regions of the dusty circumstellar disk orbiting the ~2–5 Myr old, actively accreting solar analog LkCa 15 are known to be chemically rich, and the inner disk may host a young protoplanet within its central cavity. To obtain a complete census of the brightest molecular line emission emanating from the LkCa 15 disk over the 210–270 GHz (1.4–1.1 mm) range, we have conducted an unbiased radio spectroscopic survey with the Institut de Radioastronomie Millimétrique (IRAM) 30 m telescope. The survey demonstrates that in this spectral region, the most readily detectable lines are those of CO and its isotopologues ¹³CO and C¹⁸O, as well as HCO⁺, HCN, CN, C₂H, CS, and H₂CO. All of these species had been previously detected in the LkCa 15 disk; however, the present survey includes the first complete coverage of the CN (2–1) and C₂H (3–2) hyperfine complexes. Modeling of these emission complexes indicates that the CN and C₂H either reside in the coldest regions of the disk or are subthermally excited, and that their abundances are enhanced relative to molecular clouds and young stellar object environments. These results highlight the value of unbiased single-dish line surveys in guiding future high-resolution interferometric imaging of disks.

Key words: circumstellar matter – ISM: molecules – protoplanetary disks – stars: individual (LkCa 15) – stars: pre-main sequence

1. INTRODUCTION

To understand the origins of our own solar system, it is necessary to investigate the protoplanetary disks of young, solar-mass stars (T Tauri stars) that represent early solar analogs. The LkCa 15 star/disk system, which has been extensively studied at X-ray, optical, infrared, and (sub) millimeter wavelengths, represents a prime example of such an early solar analog (Thi et al. 2004; Öberg et al. 2010; Andrews et al. 2011; Isella et al. 2012; Skinner & Güdel 2013; Thalmann et al. 2014). The LkCa 15 system is associated with the Taurus-Auriga star-forming cloud complex, at a distance of about 140 pc (van den Ancker et al. 1998), and is an archetypal “transition disk” system (see, e.g., Williams & Cieza 2011). The K5 ($L_{\text{bol}} = 0.74 L_{\odot}$) central star is estimated to have an age of 2–5 Myr and a mass of $M = 1.0 M_{\odot}$ (Kenyon & Hartmann 1995; Simon et al. 2000). The LkCa 15 star/disk system is characterized by a mass accretion rate of about $3 \times 10^{-9} M_{\odot} \text{ yr}^{-1}$ (Hartmann et al. 1998), a partially dust-depleted inner region of ~50 AU in radius (Piétu et al. 2006; Espaillat et al. 2007; Andrews et al. 2011), and an infrared excess over the stellar photosphere that can be explained by the presence of hot dust within a few AUs from the central star (Espaillat et al. 2007; Mulders et al. 2010). The disk, which displays detectable molecular gas in Keplerian rotation out to ~900 AU (Piétu et al. 2007), is viewed at intermediate inclination ($i \sim 52^{\circ}$; Dutrey et al. 2011).

The ~50 AU central cavity constitutes one of the largest inner submillimeter “holes” observed among transition disks (Andrews et al. 2011). The discrepancy in the outer disk radius as inferred from interferometry of emission from dust and gas (150 AU versus 900 AU, respectively; Piétu et al. 2007; Isella et al. 2012) also suggests the density of mm-size dust grains drops precipitously in the outer disk, beyond ~150 AU. This radially stratified dust disk structure may reflect the effects of

recent or ongoing planet formation. Indeed, Kraus & Ireland (2012) reported the potential discovery of a $6 M_J$, ~1 Myr old protoplanet well within the submillimeter continuum cavity (i.e., at a distance of 16 AU from the central T Tauri star). It is unclear how a cavity as large as that observed within the LkCa 15 disk may be formed, but it is believed that it must be due to more than just dynamical interactions with a single planet, suggesting the presence of multiple planets and/or the need for additional clearing mechanisms (e.g., grain growth, photo-evaporation; Dodson-Robinson & Salyk 2011; Zhu et al. 2011).

Given its status as the first accreting T Tauri star that has been identified as potentially harboring a young protoplanet, the LkCa 15 system serves as a unique laboratory for constraining physical conditions within a planet-forming disk—one in which planet formation is likely at an advanced stage. Indeed, among protoplanetary disks, the LkCa 15 disk is especially chemically rich; it has previously been detected in several molecular transitions (Thi et al. 2004; Piétu et al. 2007; Öberg et al. 2010). Furthermore, Skinner & Güdel (2013) report the detection of LkCa 15 as a bright X-ray source with *Chandra* (intrinsic X-ray luminosity of $\log L_x \approx 30.4 \text{ ergs s}^{-1}$). As discussed in Kastner et al. (2014) and references therein, such strong X-ray emission (as well as extreme ultraviolet emission) from the central star likely has important consequences for disk heating, chemistry, accretion, and mass-loss—and, hence, for planet formation.

We have demonstrated that unbiased line surveys of evolved, chemically rich, protoplanetary disks, such as that orbiting LkCa 15, are necessary to establish the full inventory of readily detectable molecular species in disks (Kastner et al. 2014). Such observations are important both in providing constraints on the latest generations of disk models that incorporate stellar irradiation and gas-grain reactions, and in serving as guidance for subsequent interferometric imaging of molecular emission

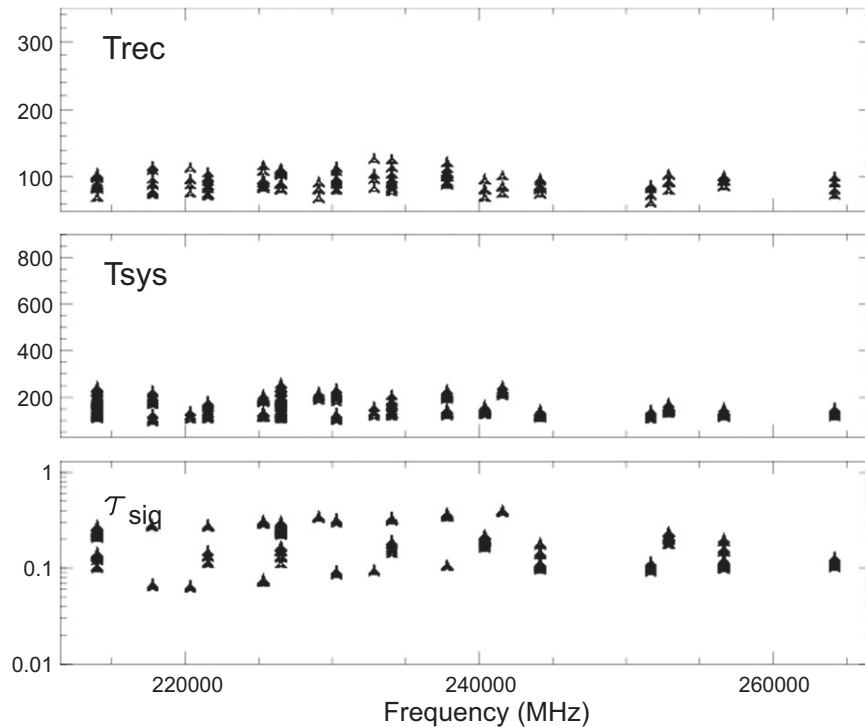


Figure 1. Receiver temperatures (top), system temperatures (middle), and signal-band opacities (bottom) as functions of frequency during the LkCa 15 spectral survey.

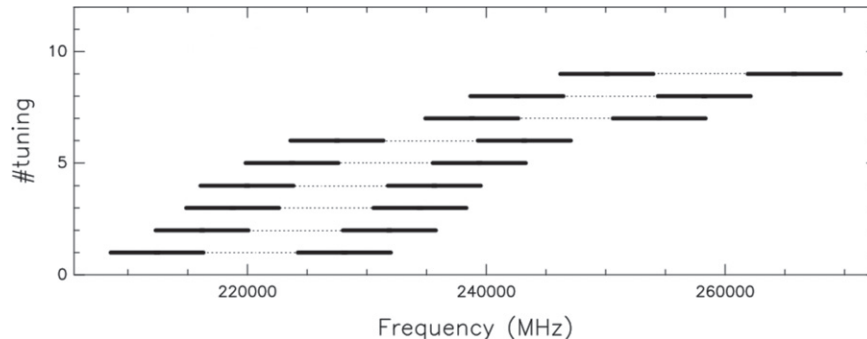


Figure 2. IRAM EMIR spectral coverage for the LkCa 15 spectral survey.

from protoplanetary disks (Kastner et al. 2015). With this as motivation, we have carried out such an unbiased radio molecular line survey of the LkCa 15 disk with the 30 m Institute de Radioastronomie Millimétrique (IRAM) radio telescope.

2. OBSERVATIONS AND DATA REDUCTION

The line survey of LkCa 15 reported here was carried out over the frequency range ~ 210 – 270 GHz with the IRAM 30 m telescope and its EMIR receivers in 2012 March during good weather conditions (~ 2.0 mm of precipitable water vapor). The EMIR receiver temperatures were fairly homogeneous over the frequency range surveyed (see Figure 1), with values in the range 80–100 K. The standard three-phase calibration procedure at the IRAM 30 m was used to obtain the zenith opacity every 10–15 minutes, with measured values between 0.1 and 0.3. The EMIR receivers were used along with the Fourier Transform Spectrometers in their 200 kHz (0.25 km s^{-1}) spectral resolution mode, providing ~ 16 GHz instantaneous bandwidth that was split into two chunks of ~ 8 GHz each (see

Figure 2). To ensure we obtained flat baselines and optimal atmospheric subtraction, the observations were performed in the wobbler switch mode, with a $70''$ throw. The pointing was checked approximately every 1–1.5 hr on a nearby continuum source. The focus was checked on a strong source every 2–2.5 hr. The frequency setups were chosen to ensure a factor 2 redundancy, except at the edges. However, bad weather conditions reduced the observing time, and the final frequency coverage was revised accordingly (see Figure 2). Consequently, nine setups have been observed so as to obtain complete coverage from 208.546 to 269.666 GHz. We note that having such a large quasiinstantaneous bandwidth not only minimizes inter-calibration issues between lines, but also minimizes pointing issues.

Data reduction was performed with the CLASS software within the GILDAS package.⁴ Residual bandpass effects were subtracted by applying a first order polynomial to each spectrum. We find the noise to be almost uniform (see

⁴ <http://www.iram.fr/IRAMFR/GILDAS>

Figure 1), with an average noise level of 7 mK (antenna temperature scale) in 0.27 km s^{-1} channels across the full bandwidth.

The output of the foregoing data acquisition and reduction sequences for LkCa 15 is a single spectrum that covers the full $\sim 210\text{--}270$ GHz frequency range. The results for line intensities and derived quantities described in Section 3 were obtained from these spectra, adopting a beam FWHM of $10''$. The line intensities are quoted on the main beam temperature scale, using $T_{\text{mb}} = \frac{F_{\text{eff}}}{B_{\text{eff}}} T_A^*$, where the appropriate values of F_{eff} (forward efficiency) and B_{eff} (main beam efficiency) were found by linear interpolation of the values that are provided in the IRAM 30 m documentation. These line intensities can be converted to Jy km s^{-1} using $3.906 \frac{F_{\text{eff}}}{A_{\text{eff}}} \text{ Jy K}^{-1}$ for S/T_A^* , the point source sensitivity, with $F_{\text{eff}} = 0.88\text{--}0.94$ and $A_{\text{eff}} = 0.41\text{--}0.49$ (aperture efficiency), where these represent the values for the IRAM 30 m telescope in the 1.3 mm window.⁵

3. RESULTS AND ANALYSIS

3.1. Molecular Line Inventory

We used the spectral line identification methods available in the WEEDS extension of the GILDAS⁶ software tools, along with visual inspection, to compile a list of molecular transitions that are readily detectable and measurable in the 210–270 GHz IRAM spectrum of LkCa 15. A summary of these molecular transitions is presented in Tables 1 and 2. The spectral regions covering all the lines we detected are displayed in Figure 3. A listing of all molecular transitions previously detected and searched for in LkCa 15 is presented in Table 3. An inspection of Table 1 shows that, in the spectral region surveyed, the most readily detectable lines are those of CO and its isotopologues ^{13}CO and C^{18}O , as well as HCO^+ , HCN, CN, C_2H , CS, and H_2CO . All of these species had been previously detected in the LkCa 15 disk. In most cases, however, the detections were widely spaced in time and/or made with different telescope/instrument combinations, potentially resulting in significant calibration uncertainties where, e.g., measurement of line ratios is concerned.

Absent such uncertainties we can, e.g., exploit our near-simultaneous coverage of CO and its isotopologues ^{13}CO and C^{18}O to estimate the optical depths of the observed transitions of CO and ^{13}CO (Section 3.4). In addition, the survey includes first measurements of the full suite of hyperfine transitions of CN $N = 2 \rightarrow 1$ and C_2H $N = 3 \rightarrow 2$, whose analysis yields estimates of optical depth and determinations of excitation temperature for these molecules (Section 3.5). We note that the $J = 5 \rightarrow 4$ transition of CS is detected here for the first time in the LkCa 15 disk. We also detect multiple transitions of H_2CO , one of which (the $J = 3_{13} \rightarrow 2_{12}$ line at 211.211 GHz) had not been previously detected in this disk.

3.2. Line Intensity Measurements

Because the LkCa 15 disk is viewed at intermediate inclination ($i \sim 52^\circ$; Dutrey et al. 2011), it displays molecular line profiles that are double-peaked. Thus, we fit the observed profile with a Keplerian disk model described in Section 2 of Beckwith & Sargent (1993) and used by Kastner et al.

(2008, 2010, 2014) to analyze molecular emission of protoplanetary disks observed with a single-dish radio telescope. The profile parameters are: the peak line intensity (T_{mb}), the Keplerian velocity at the outer edge of the disk (v_d , equivalent to half of the peak-to-peak velocity width), the disk radial temperature profile power law component (q , where $T(r) \propto r^{-q}$), and the outer disk density cutoff (p_d). When we have obtained a good fit to the observed line profile, as determined by the reduced χ^2 of the fit, we can then integrate under the model line profile and obtain the total line intensity (I). We stress that the purpose of applying such a (parametric) model is to obtain line intensity measurements, and not necessarily to infer disk structure; such inferences are, in any case, inherently difficult given only single-dish observations.

For all lines with high signal-to-noise, all four parameters were allowed to vary to find the best-fitting line profile, as determined by minimizing the reduced χ^2 of the fit. To fit the low signal-to-noise lines, we adopt the model parameters for v_d , q , and p_d from the ^{12}CO line and only allow T_{mb} to remain free. To estimate upper limits and tentative detections, we fixed all parameters, except T_{mb} , to the values that best fit the ^{12}CO line, and then determined the maximum line intensity that could be present given the observed noise level. The upper limit (at the 1σ level) was then determined by integrating over the resulting model line profile. The results of these line profile fits are listed in Table 1; the upper limits listed are at the 3σ level. In Figures 4–6, we display the observed and best-fit model line profiles for the transitions of CO isotopologues, HCO^+ , HCN, CS, and H_2CO ; in all of these cases, all four model profile parameters were allowed to vary. We also obtain tentative detections for DCN and DCO^+ , for which three of the four parameters in the fit were fixed. The spectra of these tentative detections were rebinned and are presented in Figure 7. For comparison, a summary of previous molecular line intensity measurements for LkCa 15 is provided in Table 3.

3.3. Column Density Estimates

We estimate source-averaged column densities for the (pure rotational) molecular transitions detected using the methods described in Goldsmith & Langer (1999) and Kastner et al. (2014), where we use the results for the integrated line intensities (Table 1), along with molecular line reference data from The Cologne Database for Molecular Spectroscopy (CDMS).⁷ These methods are briefly summarized here.

In the case of optically thin emission ($\tau \ll 1$), the total column density for linear molecules—for which the degeneracy of the upper rotational level is given by $g_u = 2J_u + 1$ and the energy of the lower level is $E_l = h B_{\text{rot}} J_l (J_l + 1)$ —can be approximated by

$$N_{\text{tot}} = \frac{8\pi\nu^3}{c^3} \frac{Q_{\text{rot}}(T_{\text{ex}})}{A_{\text{ul}} g_u} \frac{e^{E_l/kT_{\text{ex}}}}{1 - e^{T_{\text{ul}}/T_{\text{ex}}}} \frac{W}{\Delta J_\nu} f_{\text{dilution}}, \quad (1)$$

where ν is the frequency of the transition, W is the integrated line intensity (in K cm s^{-1}), A_{ul} is the Einstein coefficient for the transition, $T_{\text{ul}} = \frac{h\nu}{k}$ is the temperature equivalent of the transition energy, $f_{\text{dilution}} = \frac{\Omega_{\text{mb}} + \Omega_s}{\Omega_s}$ is the main-beam dilution correction factor (with Ω_{mb} and Ω_s the main-beam and source solid angles, respectively), and $\Delta J_\nu = J_\nu(T_{\text{ex}}) - J_\nu(T_{\text{bg}})$ where

⁵ See <http://www.iram.es/IRAMES/mainWiki/Iram30mEfficiencies>.

⁶ See Maret et al. (2011).

⁷ See <http://www.astro.uni-koeln.de/cdms>.

Table 1
LkCa 15: Molecular Species Detected in the IRAM 30 m Line Survey

Species	Transition	ν (GHz)	T_{mb} (K)	v_d^a (km s ⁻¹)	I (K km s ⁻¹) ^f
¹² CO	$J = 2 \rightarrow 1$	230.5380000	0.542	1.128	1.441
¹³ CO	$J = 2 \rightarrow 1$	220.3986765	0.201	1.184	0.550
C ¹⁸ O	$J = 2 \rightarrow 1$	219.5603568	0.038	1.124	0.092
CN	$N = 2 \rightarrow 1$	226.6321901	0.036	1.413	0.108
		226.6595584	0.079	1.084	0.226
		226.6636928	0.040	1.228	0.118
		226.6793114	0.042	1.251	0.112
		226.8741908	0.183	1.431	0.681 ^b
		226.8747813
		226.8758960
		226.8874202	0.020	1.762	0.081
		226.8921280	0.015	1.188	0.062
		244.9356435	0.064	1.245	0.182
CS	$J = 5 \rightarrow 4$	262.0042266	0.049	1.128	0.269 ^c
		262.0064034
		262.0648433	0.033	1.128	0.152 ^d
		262.0673312
HCN	$J = 3 \rightarrow 2$	265.8861800	0.238	1.411	0.732
		267.5575260	0.334	1.298	1.024
DCN	$J = 3 \rightarrow 2$	217.2386307	0.019	1.128	0.152 ^c
DCO ⁺	$J = 3 \rightarrow 2$	216.1125766	0.015	1.128	0.118 ^c
H ₂ CO	$J = 3_{13} \rightarrow 1_{12}$	211.2114680	0.056	1.300	0.147
		218.2221920	0.032	1.134	0.085
		218.4756320	<0.008	1.128	<0.060
		225.6977750	0.058	1.093	0.136

^a As defined by Kastner et al. (2008).

^b Sum of integrated intensities of hyperfine structure lines in range 226.87419–226.87590 GHz.

^c Sum of integrated intensities of hyperfine structure lines in range 262.00426–262.00649 GHz.

^d Sum of integrated intensities of hyperfine structure lines in range 262.06498–262.06747 GHz.

^e Tentative detection.

^f These line intensities can be converted to Jy km s⁻¹ using $3.906 \frac{F_{eff}}{A_{eff}} \text{ Jy K}^{-1}$ for S/T_A^* , the point source sensitivity, with $F_{eff} = 0.92$ and $A_{eff} = 0.46$ (aperture efficiency), where these represent mean values for the IRAM 30 m telescope in the 1.3 mm window.

Table 2
Transitions of CN and C₂H

Species (Rotational Transition)	Transition	ν^a (GHz)	A_{ul}^a (s ⁻¹)	g_u^a	Relative Intensities ^a	I (K km s ⁻¹)
CN ($N = 2 \rightarrow 1$)	$J = 3/2 \rightarrow 1/2, F = 3/2 \rightarrow 3/2$	226.6321901	4.26×10^{-5}	4	0.0533	0.106
	$J = 3/2 \rightarrow 1/2, F = 5/2 \rightarrow 3/2$	226.6595584	9.47×10^{-5}	6	0.1776	0.221
	$J = 3/2 \rightarrow 1/2, F = 1/2 \rightarrow 1/2$	226.6636928	8.47×10^{-5}	2	0.0529	0.116
	$J = 3/2 \rightarrow 1/2, F = 3/2 \rightarrow 1/2$	226.6793114	5.27×10^{-5}	4	0.0659	0.110
	$J = 5/2 \rightarrow 3/2, F = 5/2 \rightarrow 3/2$	226.8741908	9.62×10^{-5}	6	0.1805	0.667 ^b
	$J = 5/2 \rightarrow 3/2, F = 7/2 \rightarrow 5/2$	226.8747813	1.14×10^{-4}	8	0.2860	...
	$J = 5/2 \rightarrow 3/2, F = 3/2 \rightarrow 1/2$	226.8758960	8.59×10^{-5}	4	0.1074	...
	$J = 5/2 \rightarrow 3/2, F = 3/2 \rightarrow 3/2$	226.8874202	2.73×10^{-5}	4	0.0342	0.080
	$J = 5/2 \rightarrow 3/2, F = 5/2 \rightarrow 5/2$	226.8921280	1.81×10^{-5}	6	0.0340	0.060
	C ₂ H ($N = 3 \rightarrow 2$)	$J = 7/2 \rightarrow 5/2, F = 4 \rightarrow 3$	262.0042600	5.32×10^{-5}	9	0.3213
$J = 7/2 \rightarrow 5/2, F = 3 \rightarrow 2$		262.0064820	5.12×10^{-5}	7	0.2404	...
$J = 5/2 \rightarrow 3/2, F = 3 \rightarrow 2$		262.0649860	4.89×10^{-5}	7	0.2296	0.160 ^d
$J = 5/2 \rightarrow 3/2, F = 2 \rightarrow 1$		262.0674690	4.47×10^{-5}	5	0.1501	...

^a Values of frequencies, A_{ul} , g_u , and theoretical relative intensities of hyperfine transitions of CN and C₂H obtained from the CDMS database. CN line frequencies were measured in the laboratory by Skatrud et al. (1983) and the fitted values from the CDMS database are from Müller et al. (2000).

^b Sum of integrated intensities of hyperfine structure lines in range 226.87419–226.87590 GHz.

^c Sum of integrated intensities of hyperfine structure lines in range 262.00426–262.00649 GHz.

^d Sum of integrated intensities of hyperfine structure lines in range 262.06498–262.06747 GHz.

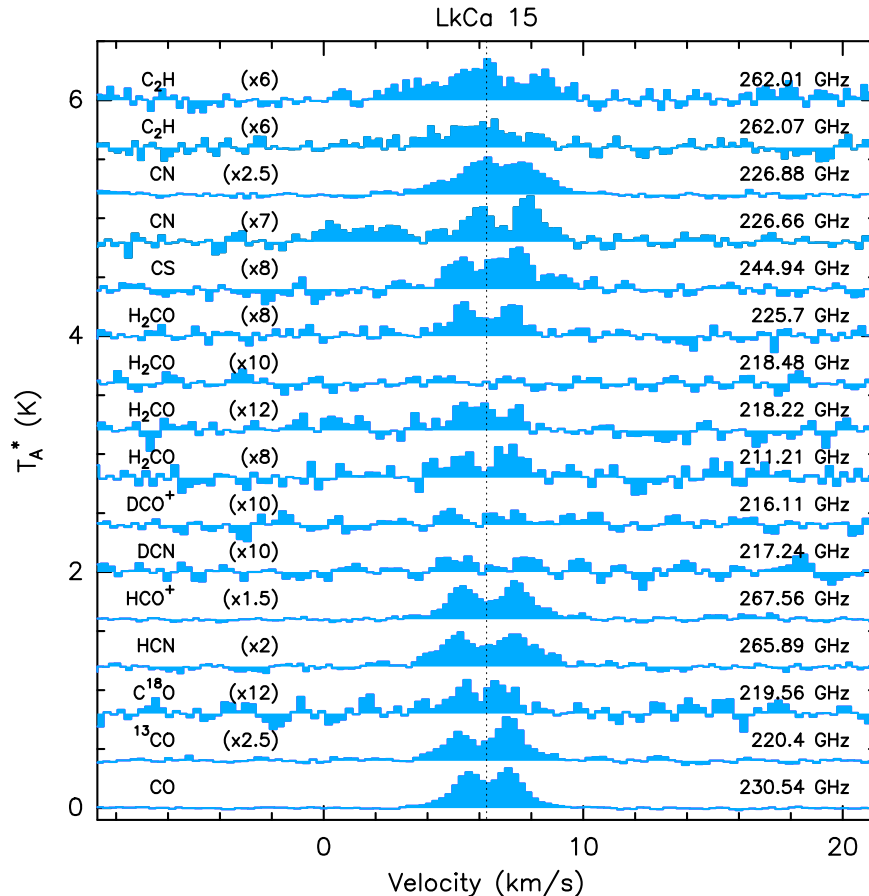


Figure 3. Spectra of all molecular transitions that have been studied in our IRAM 30 m line survey of LkCa 15. Ordinate is velocity with respect to the local standard of rest (in km s^{-1}) and abscissa is line flux (in K); for clarity, spectral regions (other than ^{12}CO) have been shifted upwards, and lines have been rescaled.

$J_\nu = \frac{h\nu}{k \exp(h\nu/kT) - 1}$ and T_{bg} is the temperature of the cosmic microwave background radiation. For optically thick emission ($\tau \gtrsim 1$), a correction factor of $\frac{\tau}{1 - e^{-\tau}}$ must be applied to Equation (1).

For the simple (diatomic) case where the rotational level is not split by higher-order effects (e.g., fine, hyperfine structures) the calculation of the total column density is straightforward. The temperature dependent total partition function then reduces to the rotational partition function, which can be approximated by

$$Z(T_{\text{ex}}) \approx Q_{\text{tot}}(T_{\text{ex}}) \approx Q_{\text{rot}}(T_{\text{ex}}) \approx \frac{1}{3} + \frac{kT_{\text{ex}}}{hB_{\text{rot}}}, \quad (2)$$

where T_{ex} is the excitation temperature of the gas and B_{rot} is the molecular rotational constant (see, e.g., Mangum & Shirley 2015). Equation (2) only applies to the case of linear (and approximately linear) molecules, such that the molecules have only one (or only one dominant) moment of inertia, or cases in which the hyperfine structures are not resolved. In the (more complex) case of the H_2CO molecule, we adopt the values of the partition function provided by the CDMS database and interpolate to determine the value for the partition function at arbitrary temperatures.

Using the foregoing methods, we obtain source-averaged column densities N_X for a range of excitation temperatures ($T_{\text{ex}} = 4, 6, 10, 20, 30 \text{ K}$) for the species that are assumed to be

optically thin. Molecular data (e.g., ν , B_{rot} , A_{ul}) were obtained from the CDMS database. In Table 4, we list these column densities for the foregoing excitation temperatures, along with column density estimates for ^{13}CO (2–1). Although the ^{13}CO (2–1) emission is evidently optically thin (Section 3.4), the column density estimates for ^{13}CO (2–1) were derived from the column density obtained from the C^{18}O integrated intensity, adopting isotopic abundance ratios $^{12}\text{C}:^{13}\text{C} \sim 68$ (Milam et al. 2005) and $^{16}\text{O}:^{18}\text{O} \sim 480$ (Scott et al. 2006).

For all molecules except CN and C_2H (see Section 3.5), the values of column density (Table 4) were obtained assuming the characteristic emitting region is a uniform disk of angular radius $R = 4''$, based on results from interferometric studies (e.g., Piétu et al. 2007; Öberg et al. 2010; Chapillon et al. 2012a). Hence, we have adopted a beam dilution factor⁸ of $f_{\text{dilution}} \approx 1.8$ for all molecules in the disk, corresponding to the half power beam width of the IRAM 30 m telescope at 230 GHz (10/7). This simplification introduces modest systematic errors in the values of f_{dilution} and, hence, column density at the extremes of the frequency range surveyed (i.e., a $\sim 6\%$ underestimate at 210 GHz and a $\sim 30\%$ overestimate at 260 GHz).

⁸ In Kastner et al. (2014), we mistakenly quote the inverse of the beam dilution factor.

Table 3
LkCa 15: Molecular Transitions Observed to Date

Species	Transition	ν (GHz)	I^b (K km s ⁻¹) ^a	I (Jy km s ⁻¹) ^c	References
¹² CO	$J = 2 \rightarrow 1$	230.5380000	1.42, 1.82	13.94, ... ^e , 12.5, ... ^e , ... ^e	1, 2, 3, 9, 10, 11, 12
	$J = 3 \rightarrow 2$	345.7959899	1.17, ... ^l		2, 13
	$J = 6 \rightarrow 5$	691.4730763	1.9		13
¹³ CO	$J = 1 \rightarrow 0$	110.2013541		... ^e , 6.39	9, 10
	$J = 2 \rightarrow 1$	220.3986765	0.53	... ^e	1, 9
	$J = 3 \rightarrow 2$	330.5879601	0.39, ... ^l		2, 13
C ¹⁸ O	$J = 1 \rightarrow 0$	109.7821734		1.90	10
	$J = 2 \rightarrow 1$	219.5603568	0.09, <0.20		1, 2
CN	$J = 3 \rightarrow 2$	329.3305453	<0.14		2
	$N = 2 \rightarrow 1$	226.6321901	0.11		1
		226.6595584	0.22		1
		226.6636928	0.12		1
		226.6793114	0.11	6.80	1, 3
		226.8741908	0.67 ^g		1
		226.8747813	... ^g	1.16, ... ^e	1, 3, 6
		226.8758960	... ^g		1
		226.8874202	0.08		1
		226.8921280	0.06		1
CS	$J = 3\frac{7}{2} \rightarrow 2\frac{5}{2}$	340.2485764	0.67		2
	$J = 3 \rightarrow 2$	146.9690287	... ^e		15
	$J = 5 \rightarrow 4$	244.9356435	0.18		1
	$J = 7 \rightarrow 6$	342.8830000	<0.08 ^f		2
CCS	$J = 7, 7 \rightarrow 6, 8$	90.6863810	<0.010 ^f		7
	$J = 7, 8 \rightarrow 6, 7$	93.8701070	<0.008 ^f		7
C ₂ H	$J = 11, 12 \rightarrow 10, 11$	144.2448364	0.061		7
	$J = 1 \rightarrow 0$	87.3286240		... ^e	4
	$J = 2 \rightarrow 1$	174.6632220		... ^e	4
	$J = 2 \rightarrow 1$	174.6676850		... ^e	4
	$N = 3 \rightarrow 2$	262.0042266	0.28 ^h		1
		262.0064034	... ^h		1
		262.0648433	0.16 ⁱ		1
HCN		262.0673312	... ⁱ		1
	$J = 1 \rightarrow 0$	88.6318470		... ^e	6
	$J = 3 \rightarrow 2$	265.8861800	0.78	5.52	1, 3
	$J = 4 \rightarrow 3$	354.5054759	0.25		2
HCO ⁺	$J = 1 \rightarrow 0$	89.1885230		... ^e , 3.30, ... ^e , ... ^e	9, 10, 11, 12
	$J = 3 \rightarrow 2$	267.5575260	1.10	5.19	1, 3
	$J = 4 \rightarrow 3$	356.7341340	0.26		2
H ¹³ CO ⁺	$J = 1 \rightarrow 0$	86.7542884		<0.88	10
	$J = 4 \rightarrow 3$	346.9983360	<0.13		2
DCN	$J = 3 \rightarrow 2$	217.2386307	<0.15	0.41	1, 3
DCO ⁺	$J = 3 \rightarrow 2$	216.1125766	<0.11	0.51	1, 3
	$J = 5 \rightarrow 4$	360.1697771	<0.10		2
N ₂ H ⁺	$J = 1_1 \rightarrow 0_1$	93.1718800		0.15 ^j , 3.83 ^j	8, 10
	$J = 1_2 \rightarrow 0_1$	93.1737000		... ^j	8
	$J = 1_0 \rightarrow 0_1$	93.1761300		... ^j	8
	$J = 3 \rightarrow 2$	279.5117010		0.71	3
	$J = 4 \rightarrow 3$	372.6725090	<0.10		2
H ₂ D ⁺	$J = 1_{10} \rightarrow 1_{11}$	372.4213850	<0.10		2
H ₂ CO	$J = 2_{12} \rightarrow 1_{11}$	140.8395020	0.17		2
	$J = 3_{13} \rightarrow 2_{12}$	211.2114680	0.14		1
	$J = 3_{03} \rightarrow 2_{02}$	218.2221920	0.08, 0.14	0.66	1, 2, 3
	$J = 3_{22} \rightarrow 2_{21}$	218.4756320	<0.06, <0.10		1, 2
	$J = 3_{12} \rightarrow 2_{11}$	225.6977750	0.13, 0.10		1, 2
	$J = 4_{31} \rightarrow 3_{30}$	291.3804880		1.12	3
	$J = 5_{15} \rightarrow 4_{14}$	351.7686450	0.29		2
HC ₃ N	$J = 10 \rightarrow 9$	90.9790230	<0.010 ^f		7
	$J = 12 \rightarrow 11$	109.1736340	<0.015 ^f		7
	$J = 16 \rightarrow 15$	145.5609460	0.017 ^f		7
CH ₃ OH	$J = 2_K \rightarrow 1_K$	96.741371	<0.05		2
	$J = 4_2 \rightarrow 3_1$ E ⁺	218.440063	<0.10		2

Table 3
(Continued)

Species	Transition	ν (GHz)	I^b (K km s ⁻¹) ^a	I (Jy km s ⁻¹) ^c	References
c-C ₃ H ₂	$J = 5_K \rightarrow 4_K$	241.791352	<0.10		2
	$J = 2_{1,2} \rightarrow 1_{0,1}$	85.3388930	... ^d		5
	$J = 6_{1,6} \rightarrow 5_{0,5}$	217.8221480	... ^d		5
	$J = 6_{0,6} \rightarrow 5_{1,5}$	217.8221480	... ^d		5
SO	$J = 2_{2,3} \rightarrow 1_{1,2}$	99.2998700	... ^d		15
	$J = 3_4 \rightarrow 2_3$	138.1786000	... ^d		5
	$J = 5_6 \rightarrow 4_5$	219.9494420	... ^d		5
H ₂ S	$J = 1_{10} \rightarrow 1_{01}$	168.7627624	... ^d		15

References: (1) Thi et al. (2004), (2) Öberg et al. (2010), (3) Henning et al. (2010), (4) Fuente et al. (2010), (5) Chapillon et al. (2012a), (6) Chapillon et al. (2012b), (7) Dutrey et al. (2007), (8) Piétu et al. (2007), (9) Qi et al. (2003), (10) Duvert et al. (2000), (11) Simon et al. (2000), (12) Thi et al. (2001), (13) van Zadelhoff et al. (2001), (14) Dutrey et al. (2011).

^a Integrated line intensities are quoted in main beam temperature scale using $T_{\text{mb}} = \frac{F_{\text{eff}}}{B_{\text{eff}}} T_A^*$. See <http://www.iram.es/IRAMES/mainWiki/Iram30mEfficiencies>.

^b These values can be converted to Jy km s⁻¹ using $3.906 \frac{F_{\text{eff}}}{A_{\text{eff}}} \text{Jy K}^{-1}$ for S/T_A^* , the point source sensitivity, with $F_{\text{eff}} = 0.92$ and $A_{\text{eff}} = 0.46$. See <http://www.iram.es/IRAMES/mainWiki/Iram30mEfficiencies>.

^c Integrated line intensities quoted in Jy km s⁻¹ are for measurements by interferometers.

^d These transitions were searched for, but not detected.

^e These transitions were detected, but line intensities were not tabulated in the given references.

^f These values were converted from Jy km s⁻¹ using $3.906 \frac{F_{\text{eff}}}{A_{\text{eff}}} \text{Jy K}^{-1}$ for S/T_A^* , the point source sensitivity, with $F_{\text{eff}} = 0.95$ and $A_{\text{eff}} = 0.63$. See <http://www.iram.es/IRAMES/mainWiki/Iram30mEfficiencies>.

^g Sum of integrated intensities of hyperfine structure lines in range 226.87419–226.87590 GHz.

^h Sum of integrated intensities of hyperfine structure lines in range 262.00426–262.00649 GHz.

ⁱ Sum of integrated intensities of hyperfine structure lines in range 262.06498–262.06747 GHz.

^j Sum of integrated intensities of hyperfine structure lines in range 93.17188–93.17613 GHz.^k Integrated intensities are corrected for main-beam efficiency.^l Profiles fitted with two Gaussians, but the total integrated intensity is not tabulated.

3.4. Opacities in CO Lines

We can take advantage of the detection of C¹⁸O to estimate the optical depths of the ¹³CO (2–1) and ¹²CO (2–1) transitions. The optical depth of ¹³CO may be estimated by comparing the expected and measured line ratios for the CO isotopologues. We expect

$$R = \frac{1 - \exp(-\tau_{^{13}\text{CO}})}{1 - \exp(-\tau_{^{13}\text{CO}}/X)}, \quad (3)$$

where R is the measured ¹³CO:C¹⁸O line ratio and X is the ¹³CO:C¹⁸O abundance ratio. Equation (3) follows directly from the equation of radiative transfer, if it is assumed that the contribution from the background radiation is negligible and the excitation temperatures of ¹³CO and C¹⁸O are identical. We adopt a value of $X = 7$ for the abundance ratio, ¹³C:¹²C/¹⁸O:¹⁶O, based on the values ¹²C:¹³C = 68 and ¹⁶O:¹⁸O = 480 (as determined from measurements of CN and CO isotopologues for the local interstellar medium and the Solar System, respectively; Milam et al. 2005; Scott et al. 2006). So, if the emission from ¹³CO is optically thin ($\tau_{^{13}\text{CO}} \ll 1$), we would expect to obtain $R \approx 7$. The measured ratio for the $J = 2 \rightarrow 1$ transitions for LkCa 15, $R \approx 6.0$, therefore confirms that the emission in the ¹³CO (2–1) line is optically thin, with an implied optical depth of $\tau_{^{13}\text{CO}} \sim 0.38$. Similar results were obtained by Qi et al. (2003) from observations of LkCa 15 in the $J = 1 \rightarrow 0$ transitions of ¹³CO and C¹⁸O. Then, given the assumptions for the isotopic

abundances of carbon, we expect $\tau_{^{12}\text{CO}} \approx 68\tau_{^{13}\text{CO}}$. Hence, we estimate an optical depth of ~ 26 for ¹²CO (2–1), i.e., the emission in this transition of ¹²CO from the LkCa 15 disk is highly optically thick.

We caution that the foregoing estimates of ¹³CO optical depths are potentially subject to large uncertainties. These CO opacity estimates assume that there is no selective photo-dissociation or chemical fractionation across the line profile. Finally (and perhaps least significantly), the ¹³CO line profile is slightly asymmetric (see Figure 4), which suggests that the telescope was mispointed by a few arcseconds for that particular observation, and hence, the intensity of the ¹³CO line may be a slight underestimate.

3.5. CN and C₂H Hyperfine Component Analysis: Optical Depths, Excitation Temperatures, and Column Densities

The CN and C₂H rotational emission lines display hyperfine splitting (HFS) due to interactions between electron and nuclear spins (e.g., Ziurys et al. 1982, see Table 2). We performed measurement and analysis of the LkCa 15 CN $N = 2 \rightarrow 1$ and C₂H $N = 3 \rightarrow 2$ spectra—which are displayed in Figures 8 and 9, respectively—using a new adaptation of the classical HFS fitting method not yet implemented in the CLASS software. This adaptation uses two Gaussians of the same width to fit the double-peaked emission features in our spectrum that result from the inclined Keplerian disk. The two-Gaussian HFS code makes the simplifying assumption that all the hyperfine transitions within a given rotational transition share the same excitation temperature. Fitting all the hyperfine components simultaneously results in the determination of the

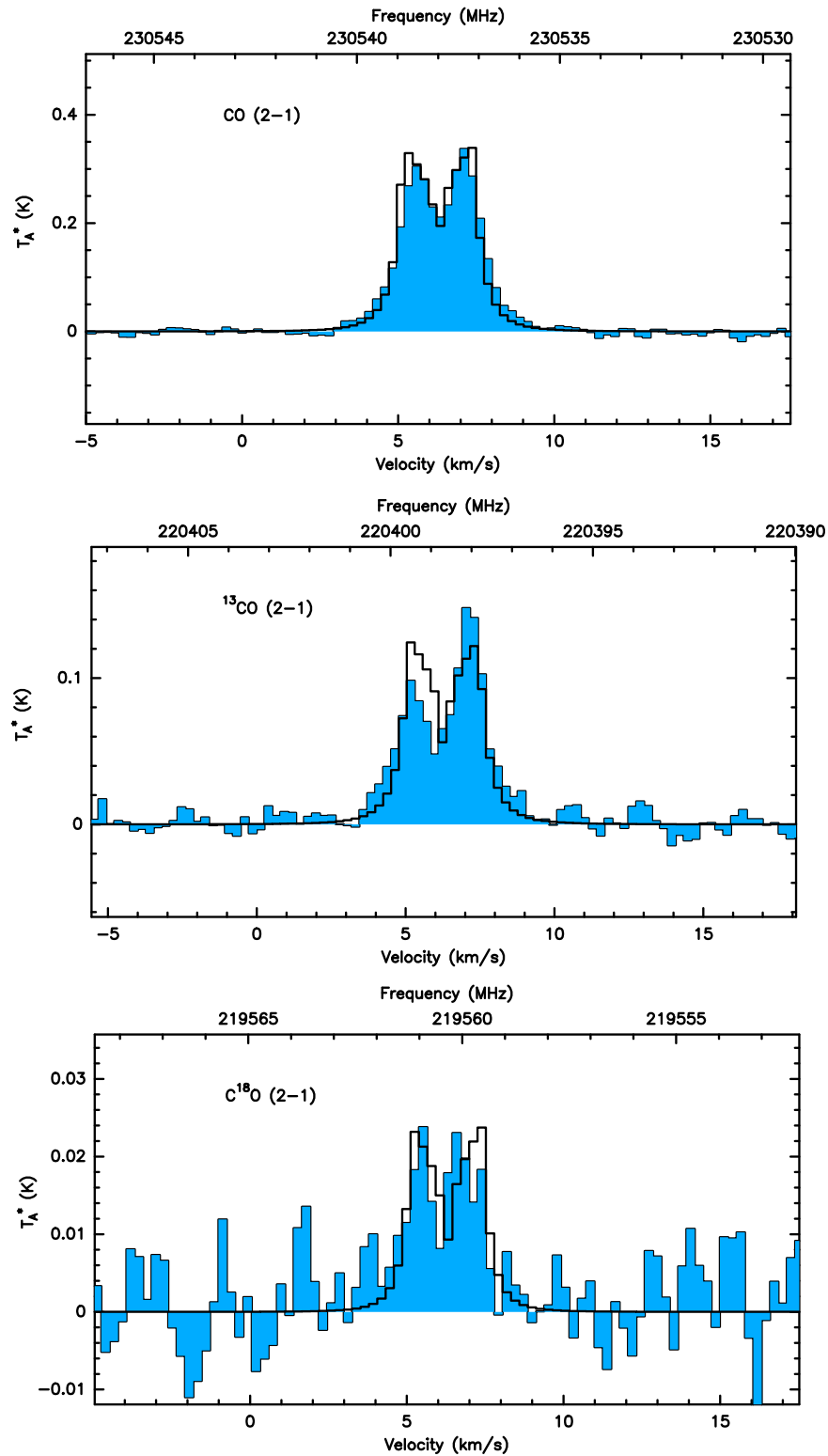


Figure 4. IRAM 30 m spectra of ^{12}CO ($J = 2 \rightarrow 1$) (top), ^{13}CO ($J = 2 \rightarrow 1$) (middle), and C^{18}O ($J = 2 \rightarrow 1$) (bottom) detected in LkCa 15 (blue histogram) overlaid with parametric Keplerian model line profiles (see Section 3.2). Ordinate is velocity with respect to the local standard of rest (in km s^{-1} , with frequencies in MHz at the top of the panel) and abscissa is line flux (in K).

opacity, τ , and excitation temperature, T_{ex} , of the rotational transition (see Kastner et al. 2014). Figures 8, 9, and Table 5 summarize the results of this analysis for the IRAM 30 m measurements of the CN and C_2H emission from LkCa 15.

The determination of the total partition functions Q_{tot} for the CN and C_2H radicals must account for the nuclear-spin

statistics, as discussed in Kastner et al. (2014). The column density is then computed as

$$N_{\text{tot}} = \frac{8\pi\nu^3}{c^3} \frac{e^{E_l/kT_{\text{ex}}}}{1 - e^{-T_{\text{ul}}/T_{\text{ex}}}} \int \tau_{\text{ul}} d\nu \frac{Q_{\text{tot}}(T_{\text{ex}})}{A_{\text{ul}} g_u}, \quad (4)$$

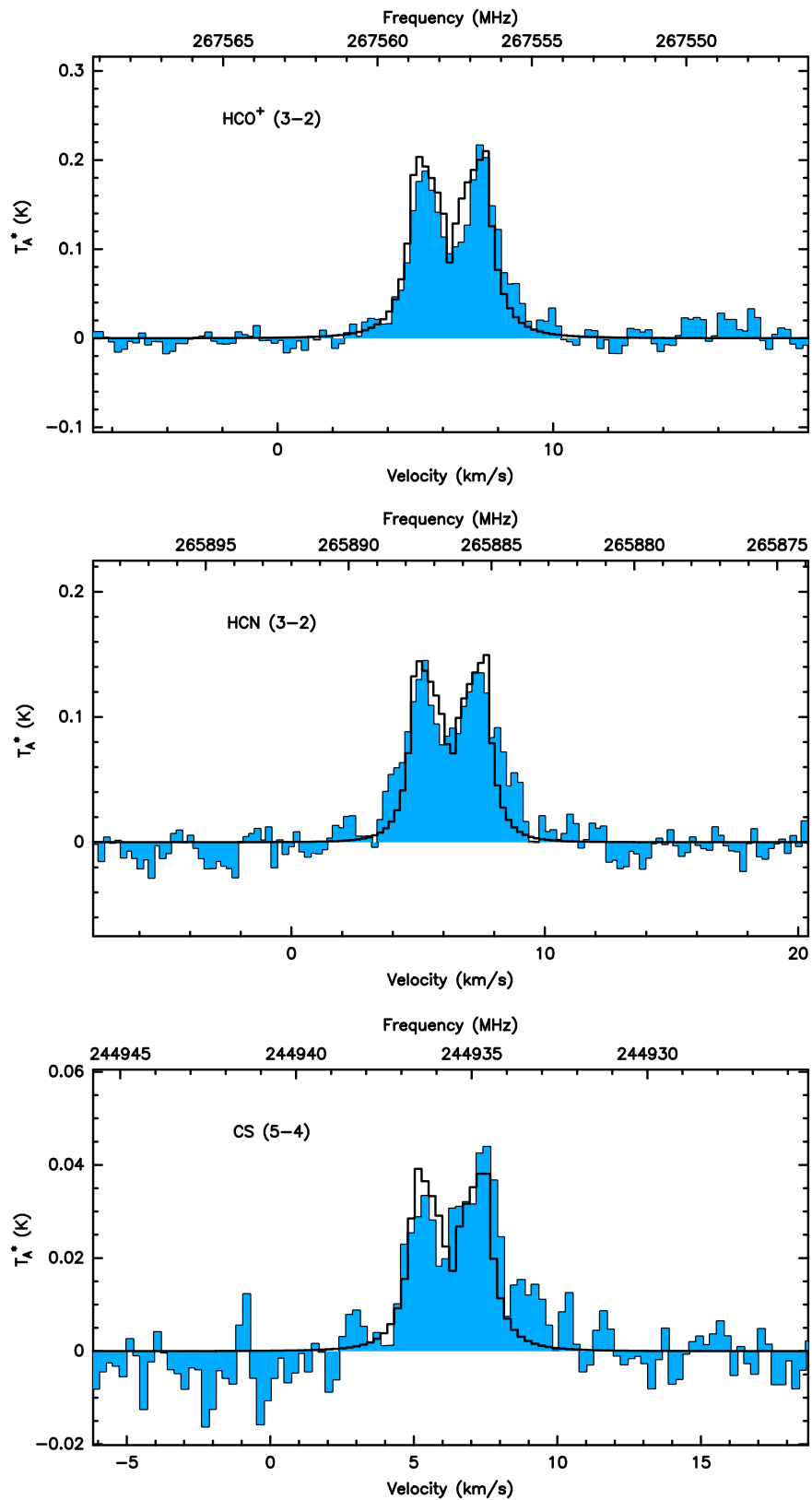


Figure 5. IRAM 30 m spectra of HCO^+ ($J = 3 \rightarrow 2$) (top), HCN ($J = 3 \rightarrow 2$) (middle), and CS ($J = 5 \rightarrow 4$) (bottom) detected in LkCa 15 (blue histogram) overlaid with parametric Keplerian model line profiles (see Section 3.2). Ordinate is velocity with respect to the local standard of rest (in km s^{-1} , with frequencies in MHz at the top of the panel) and abscissa is line flux (in K).

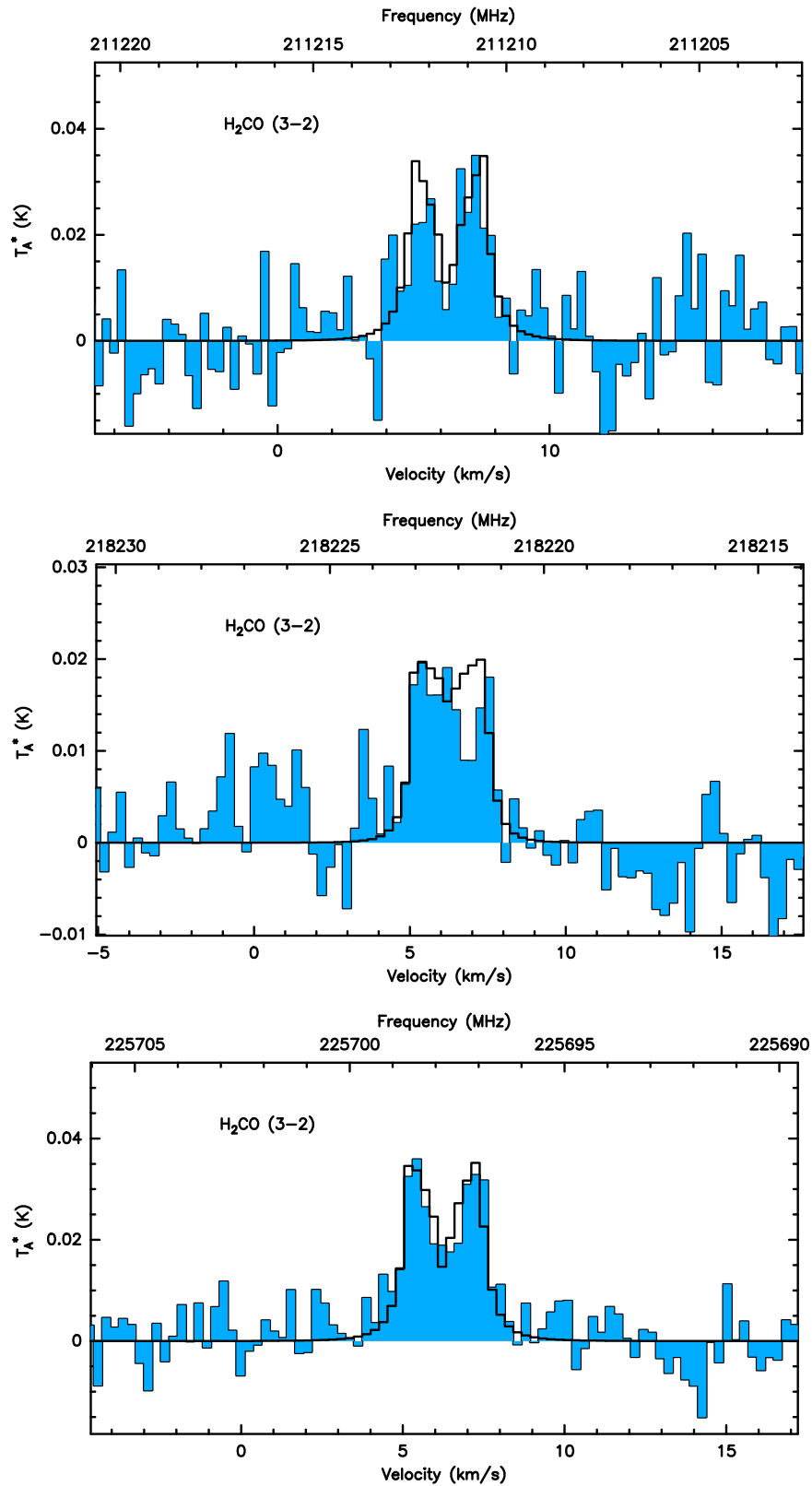


Figure 6. IRAM 30 m spectra of H_2CO ($J = 3_{13} \rightarrow 2_{12}$) (top), H_2CO ($J = 3_{03} \rightarrow 2_{02}$) (middle), and H_2CO ($J = 3_{12} \rightarrow 2_{11}$) (bottom) detected in LkCa 15 (blue histogram) overlaid with parametric Keplerian model line profiles (see Section 3.2). Ordinate is velocity with respect to the local standard of rest (in km s^{-1} , with frequencies in MHz at the top of the panel) and abscissa is line flux (in K).

where all parameters refer to the particular hyperfine transition of interest. We assume the characteristic emitting region is a uniform disk of angular radius $R = 5''$, based on results from

interferometric studies (e.g., Öberg et al. 2010; Chapillon et al. 2012a). We adopt a beam dilution correction factor $f_{\text{dilution}} \approx 1.1$, but the results we obtain are fairly insensitive to the

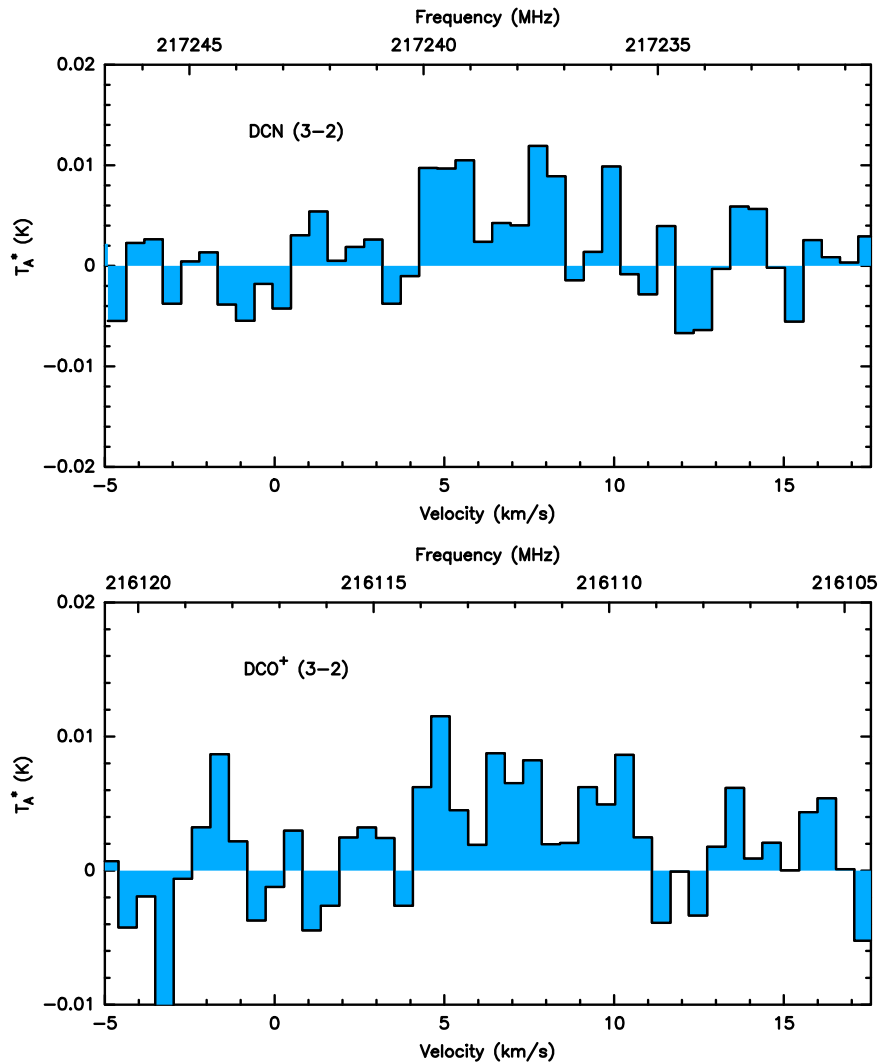


Figure 7. IRAM 30 m smoothed spectra of DCN ($J = 3 \rightarrow 2$) (top) and DCO⁺ ($J = 3 \rightarrow 2$) (bottom) detected in LkCa 15 (blue histogram; see Section 3.2). Ordinate is velocity with respect to the local standard of rest (in km s⁻¹, with frequencies in MHz at the top of the panel) and abscissa is line flux (in K).

assumed source solid angle. Following this procedure, we obtain values of $T_{\text{ex}} \sim 3.9$ K and $\tau \sim 2.2$ for CN, and $T_{\text{ex}} \sim 2.9$ K and $\tau \sim 18$ for C₂H, leading to estimates of $N_{\text{tot}}(\text{CN}) \sim 0.5 \times 10^{14}$ cm⁻² and $N_{\text{tot}}(\text{C}_2\text{H}) \sim 2 \times 10^{16}$ cm⁻² for the LkCa 15 disk (Table 5). The results for τ indicate the emission in both lines is optically thick, where the reader is reminded that the resulting values of τ represent the total opacity over the HFS. Hence, the derived values of T_{ex} imply either that the gas is subthermally excited or that the CN and C₂H emission originates from disk gas at very low temperature. We caution, however, that the low T_{ex} for CN and C₂H rest on the assumed beam filling factor. A less extended disk in these molecules would imply a higher T_{ex} .

4. DISCUSSION

In Table 4, we list fractional molecular abundances relative to ¹³CO, $N(X)/N(^{13}\text{CO})$, as obtained from the calculated source-averaged molecular column densities for LkCa 15 (Section 3.3). We present values of $N(X)/N(^{13}\text{CO})$ that we obtained by adopting $T_{\text{ex}} = 4, 6, 10, 20, 30$ K. We caution that

these column density estimates may have large systematic uncertainties, given the uncertainties in the assumed source parameters and, in particular, the assumed source emitting region solid angles in the various molecules (Aikawa et al. 2003; Piétu et al. 2007; Öberg et al. 2010; Chapillon et al. 2012a). Nevertheless, our results yield column densities that are overall consistent with previous studies of the LkCa 15 disk. This consistency is apparent from Table 6, in which we present a comparison of our column densities and column densities determined by previous studies. In most cases, our column density estimates agree with previous estimates to within a factor ~ 3 . Apart from CN and C₂H (see below), the largest discrepancies appear in the comparison with results for ¹³CO, C¹⁸O, and HCO⁺ published by Qi et al. (2003); the large differences between column density results in these cases can be attributed to the contrasting beam sizes of OVRO ($\sim 3''$) and IRAM ($\sim 12''$).

It should be noted that for LkCa 15, we derive column densities for CN and especially for C₂H that are relatively high; in the latter case, the inferred source-averaged column density is similar to that of ¹³CO. Specifically, we derive $N(\text{CN})/N(^{13}\text{CO})$ and $N(\text{C}_2\text{H})/N(^{13}\text{CO})$ values of approximately 0.01 and 3.4, respectively, adopting $T_{\text{ex}} = 4$ K for ¹³CO. The former

Table 4
Column Densities^a

Species (Transition)	ν (GHz)	A_{ul} (s ⁻¹)	T_{ex} (K)	$Z(T)$	N_X (cm ⁻²)	$N_X/N_{^{13}\text{CO}}$
¹³ CO (2–1) ^b	220.398686	...	30	...	6.98×10^{14}	...
			20	...	6.18×10^{14}	...
			10	...	7.28×10^{14}	...
			6	...	1.41×10^{15}	...
			4	...	4.65×10^{15}	...
C ¹⁸ O (2–1)	219.560357	6.012×10^{-7}	30	11.721	9.97×10^{13}	...
			20	7.925	8.82×10^{13}	...
			10	4.129	1.04×10^{14}	...
			6	2.610	2.02×10^{14}	...
			4	1.852	6.64×10^{14}	...
CS (5–4)	244.9356	2.981×10^{-4}	30	25.852	9.91×10^{11}	1.42×10^{-3}
			20	17.346	1.20×10^{12}	1.95×10^{-3}
			10	8.840	3.65×10^{12}	5.01×10^{-3}
			6	5.437	2.49×10^{13}	1.76×10^{-2}
			4	3.736	3.92×10^{14}	8.43×10^{-2}
HCN (3–2)	265.886431	8.356×10^{-4}	30	14.439	1.12×10^{12}	1.60×10^{-3}
			20	9.737	1.16×10^{12}	1.87×10^{-3}
			10	5.035	2.18×10^{12}	2.99×10^{-3}
			6	3.154	7.84×10^{12}	5.55×10^{-3}
			4	2.214	5.50×10^{13}	1.18×10^{-2}
HCO ⁺ (3–2)	267.557633	1.453×10^{-3}	30	14.351	9.12×10^{11}	1.31×10^{-3}
			20	9.678	9.47×10^{11}	1.53×10^{-3}
			10	5.006	1.80×10^{12}	2.47×10^{-3}
			6	3.137	6.54×10^{12}	4.63×10^{-3}
			4	2.202	4.65×10^{13}	9.99×10^{-3}
DCN (3–2)	217.2386	4.575×10^{-4}	30	17.598	$<2.67 \times 10^{11}$	$<3.82 \times 10^{-4}$
			20	11.843	$<2.55 \times 10^{11}$	$<4.14 \times 10^{-4}$
			10	6.088	$<3.83 \times 10^{11}$	$<5.25 \times 10^{-4}$
			6	3.786	$<1.02 \times 10^{12}$	$<7.24 \times 10^{-4}$
			4	2.635	$<5.04 \times 10^{12}$	$<1.09 \times 10^{-3}$
DCO ⁺ (3–2)	216.1126	7.658×10^{-4}	30	17.688	$<1.23 \times 10^{11}$	$<1.76 \times 10^{-4}$
			20	11.903	$<1.17 \times 10^{11}$	$<1.90 \times 10^{-4}$
			10	6.118	$<1.75 \times 10^{11}$	$<2.40 \times 10^{-4}$
			6	3.804	$<4.64 \times 10^{11}$	$<3.29 \times 10^{-4}$
			4	2.647	$<2.27 \times 10^{12}$	$<4.88 \times 10^{-4}$
H ₂ CO (3–2)	211.2115	2.271×10^{-4}	30	73.902	1.43×10^{12}	2.05×10^{-3}
			20	46.495	1.07×10^{12}	1.73×10^{-3}
			10	19.089	7.50×10^{11}	1.03×10^{-3}
			6	8.126	6.75×10^{11}	4.78×10^{-4}
			4	2.645	6.41×10^{11}	1.38×10^{-4}
H ₂ CO (3–2)	218.2222	2.818×10^{-4}	30	73.902	7.24×10^{11}	1.04×10^{-3}
			20	46.495	5.46×10^{11}	8.84×10^{-4}
			10	19.089	3.89×10^{11}	5.34×10^{-4}
			6	8.126	3.56×10^{11}	2.52×10^{-4}
			4	2.645	3.45×10^{11}	7.43×10^{-5}
H ₂ CO (3–2)	218.4756	1.571×10^{-4}	30	73.902	$<9.27 \times 10^{11}$	$<1.33 \times 10^{-3}$
			20	46.495	$<6.99 \times 10^{11}$	$<1.13 \times 10^{-3}$
			10	19.089	$<4.98 \times 10^{11}$	$<6.83 \times 10^{-4}$
			6	8.126	$<4.56 \times 10^{11}$	$<3.23 \times 10^{-4}$
			4	2.645	$<4.43 \times 10^{11}$	$<9.52 \times 10^{-5}$
H ₂ CO (3–2)	225.6978	2.772×10^{-4}	30	73.902	5.14×10^{11}	7.36×10^{-4}
			20	46.495	3.89×10^{11}	6.30×10^{-4}
			10	19.089	2.81×10^{11}	3.87×10^{-4}
			6	8.126	2.63×10^{11}	1.86×10^{-4}
			4	2.645	2.61×10^{11}	5.61×10^{-5}

^a Beam-averaged column densities assuming optically thin emission and adopting source radii (for all molecules) of 4". See Section 3.3.

^b ¹³CO column densities derived directly from the C¹⁸O intensity assuming $[^{16}\text{O}:^{18}\text{O}]/[^{12}\text{C}:^{13}\text{C}] = 7$ (Milam et al. 2005; Scott et al. 2006).

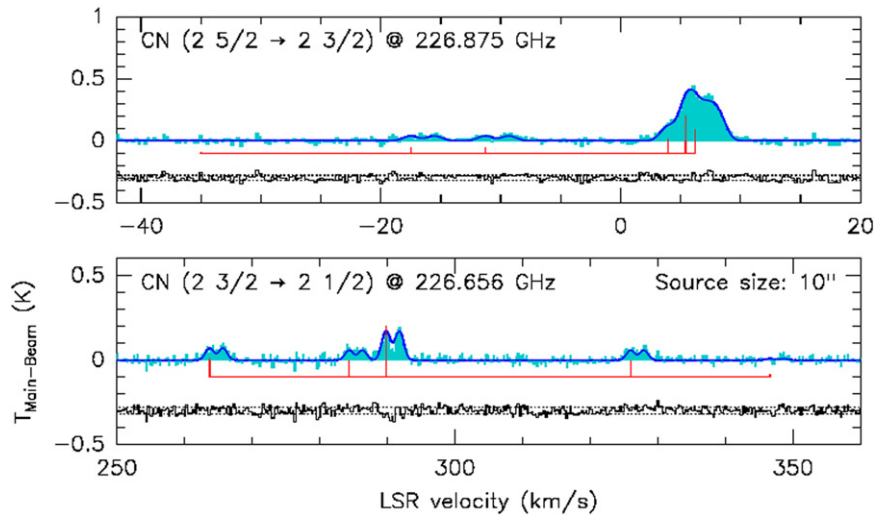


Figure 8. Full spectra of hyperfine components of CN observed toward LkCa 15. The red lines indicate the positions and relative intensities of the hyperfine components (see Table 2). The spectra are overlaid with the best-fit models obtained via the method described in Section 3.5. The residuals from the fit are below the spectra.

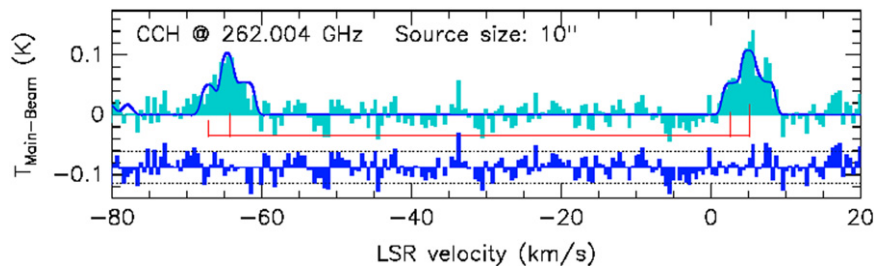


Figure 9. Full spectrum of hyperfine components of C_2H observed toward LkCa 15. The red lines indicate the positions and relative intensities of the hyperfine components (see Table 2). The spectrum is overlaid with the best-fit models obtained via the method described in Section 3.5. The residuals from the fit are below the spectrum.

result is somewhat smaller than that quoted for LkCa 15 in Kastner et al. (2014), where the difference reflects our refined treatment of CN optical depth and excitation temperature (Section 3.5), as well as different assumptions regarding beam filling factors. In deriving our column densities, we assumed that all the molecular emission is distributed within a smooth disk, but clearly this may not be the case. If the CN or C_2H emission were instead restricted to specific regions of the disk, we could be severely underestimating the column densities of these molecular species. Thus, it is imperative to obtain high signal-to-noise interferometric images to study the surface brightness distribution for these highly abundant molecules.

It is possible that high-energy (far-UV (FUV) and/or X-ray) radiation from the central star may be enhancing the abundances of CN and C_2H in T Tauri disks (Dutrey et al. 1997; Kastner et al. 1997, 2008, 2014; Thi et al. 2004; Henning et al. 2010). Suggestions about the influence of X-rays on disk chemistry were made by Kastner et al. (2008), Thi et al. (2004), and Salter et al. (2011) to explain enhanced HCN and HCO^+ line fluxes, and by Kastner et al. (2014) to explain the enhanced abundances that were found for CN and C_2H for TW Hya and V4046 Sgr.

We also find very low values of excitation temperature for CN and C_2H ($T_{ex} \sim 3-5$ K) in the LkCa 15 disk. These determinations of low T_{ex} may indicate that the emission from CN and C_2H arises from regions that are deep within the disk. Hence, as noted by Kastner et al. (2014), it may be that the

abundances of CN and C_2H are being enhanced by X-rays as opposed to extreme FUV photons, since X-rays can penetrate deeper into disks (e.g., Skinner & Güdel 2013; Cleeves et al. 2013). Indeed, Skinner & Güdel (2013) found that LkCa 15 is an X-ray luminous source, with a spectral energy distribution that contains contributions from both soft and hard components. This combination may increase the abundances of species (like C_2H) that are the dissociation products of complex organic molecules. Alternatively, the CN and C_2H molecules may reside in tenuous regions of the disk where collision probabilities are low, such that the molecules are subthermally excited.

Our estimates of $N_{tot}(CN)$ and $N_{tot}(C_2H)$ for the LkCa 15 disk (Table 5) are similar to those found for the TW Hya and V4046 Sgr disks (Kastner et al. 2014, their Table 5). Hence, it appears that large relative abundances of C_2H (as well as CN; see Chapillon et al. 2012a) may be a common feature of evolved protoplanetary disks. The excitation temperatures that we find for CN and C_2H ($T_{ex} \sim 3-5$ K) in the LkCa 15 disk are similar to, though even lower than, those determined by Kastner et al. (2014) for TW Hya and V4046 Sgr; furthermore, we have inferred relatively large values for the optical depths of CN and C_2H for all three of these disks. Followup Submillimeter Array imaging of TW Hya has revealed that the C_2H emission emanates from a ring-like structure, with an inner radius (~ 45 AU) that appears to lie between the CO “snow line” (as traced by the ~ 30 AU diameter hole in N_2H^+)

Table 5
Results of CN and C₂H Hyperfine Structure Analysis

Species	v_0^c (km s ⁻¹)	Δv^c (km s ⁻¹)	τ^a	T_{ex} (K)	$N_{\text{tot}}/Q_{\text{tot}}^b$ (cm ⁻²)	$Q_{\text{tot}}(T_{\text{ex}})^c$	N_{tot}^d (cm ⁻²)	
CN	5.42 ± 0.03	7.38 ± 0.03	1.40 ± 0.03	2.23 ± 0.45	3.9 ^{+0.2} _{-0.2}	4.9 ^{-0.3} _{+0.2} × 10 ¹²	10.97 ^{+0.42} _{-0.42}	0.54 ^{-0.02} _{+0.01} × 10 ¹⁴
C ₂ H	5.18 ± 0.11	7.47 ± 0.07	1.39 ± 0.06	18.1 ± 0.3	2.9 ^{-0.1} _{+0.1}	2.3 ^{-0.1} _{+0.0} × 10 ¹⁵	7.1 ^{+0.18} _{+0.18}	1.6 ^{-0.31} _{+0.41} × 10 ¹⁶

^a Total opacity obtained as the sum of individual hyperfine line opacities τ_{ul} , where $\tau_{\text{ul}} = R. I. \times \tau$.

^b From Equation (4), Section 3.5.

^c From a linear interpolation to tabulated partition function values obtained from the CDMS database.

^d Column densities assuming a source diameter of 10'' for both molecules. See Section 3.5.

^e To obtain the two component Keplerian fit, we first fit one of the Doppler-shifted components and then required that the line widths of the two components be identical.

Table 6
Column Density Comparisons

Species	This Study		Previous Studies		Reference
	N_X (cm ⁻²)	T_{ex} (K)	N_X (cm ⁻²)	T_{ex} (K)	
¹³ CO	9.97 × 10 ¹³	30	1.11 × 10 ¹⁶	30	2
	6.18 × 10 ¹⁴	20	3.6 × 10 ¹⁴	25	3
	7.28 × 10 ¹⁴	10	2.6 ± 0.2 × 10 ¹⁵	13.7	4
C ¹⁸ O	9.97 × 10 ¹³	30	3.31 × 10 ¹⁵	30	2
	1.20 × 10 ¹²	20	5.1 × 10 ¹²	25	3
CS	3.65 × 10 ¹²	10	<1.2 × 10 ¹³	10	6
	2.49 × 10 ¹³	6	8.7 ± 1.6 × 10 ¹²	8.7	7
	1.16 × 10 ¹²	20	1.8 × 10 ¹²	25	3
HCN	2.18 × 10 ¹²	10	4.3 × 10 ¹²	10	6
	7.84 × 10 ¹²	6	10.6 ± 1.5 × 10 ¹²	7.0 ± 0.6	8
	9.12 × 10 ¹¹	30	9.25 × 10 ¹²	30	2
HCO ⁺	9.47 × 10 ¹¹	20	3.3 × 10 ¹¹	25	3
	9.47 × 10 ¹¹	20	8.0 ± 0.5 × 10 ¹²	19	4
	1.80 × 10 ¹²	10	7.8 × 10 ¹¹	10	6
DCO ⁺	<1.17 × 10 ¹¹	20	<2.9 × 10 ¹¹	25	3
	<1.75 × 10 ¹¹	10	<1.1 × 10 ¹¹	10	6
H ₂ CO	(0.39–1.07) × 10 ¹²	20	(0.7–1.9) × 10 ¹³	25	1
	(0.39–1.07) × 10 ¹²	20	7.2 × 10 ¹²	20	3
CN	(0.52–0.55) × 10 ¹⁴	3.9 ± 0.2	1.5 × 10 ¹³	25	3
	(0.52–0.55) × 10 ¹⁴	3.9 ± 0.2	3.4 × 10 ¹³	10	6
	(0.52–0.55) × 10 ¹⁴	3.9 ± 0.2	58 ± 5 × 10 ¹²	8.8 ± 0.3	8
C ₂ H	(1.29–2.01) × 10 ¹⁶	2.9 ± 0.1	2.9 ± 1.1 × 10 ¹³	6.3 ± 1.4	5

References: (1) Aikawa et al. (2003), (2) Qi et al. (2003), (3) Thi et al. (2004), (4) Piétu et al. (2007), (5) Henning et al. (2010), (6) Fuente et al. (2010), (7) Dutrey et al. (2011), (8) Chapillon et al. (2012a).

and the outer boundary of large-grain emission at ~60 AU (Kastner et al. 2015). Given the similarity of the TW Hya and LkCa 15 disks in terms of the excitation and optical depth of C₂H, it is likely that the C₂H (and possibly CN) emission from LkCa 15 also originates from a ring-like structure.

We present fractional abundances with respect to ¹³CO, $N(X)/N(^{13}\text{CO})$ in Table 7, where we cite the results of Thi et al. (2004) as a cross-check of our results. This table illustrates that CN and C₂H are enhanced in evolved T Tauri star disks relative to other astronomical objects, especially with respect to young stellar objects (YSOs) and even the expanding, C-rich envelope of the classical carbon star IRC+10216, in which the abundances of carbon-bearing molecules are strongly enhanced. Thus, our line survey of LkCa 15 reinforces previous assertions (e.g., Kastner et al. 2014 and references therein) that there must be some mechanism that is enhancing the abundances of CN and C₂H in protoplanetary disks.

5. SUMMARY

We have used the IRAM 30 m telescope and EMIR fast Fourier transform spectrometers to conduct a comprehensive mm-wave emission line survey of the circumstellar disk orbiting the nearby, pre-main sequence (T Tauri) star LkCa 15 over the 210–267 GHz (1.4–1.1 mm) range. We find that lines of ¹²CO, HCO⁺, HCN, ¹³CO, CN, C₂H, CS, H₂CO, and C¹⁸O constitute the strongest molecular emission from the LkCa 15 disk in the spectral region surveyed. The $J = 5 \rightarrow 4$ transition of CS and the $J = 3_{13} \rightarrow 2_{12}$ transition of H₂CO are detected for the first time in the LkCa 15 disk. We use simultaneous measurements of CO and its isotopologues (i.e., ¹³CO and C¹⁸O) to estimate the optical depths of the transitions of CO and ¹³CO observed; we find that ¹³CO emission is optically thin, and emission from ¹²CO is highly optically thick. The survey also includes first measurements of the full suite of hyperfine transitions of CN $N = 2 \rightarrow 1$ and C₂H $N = 3 \rightarrow 2$. Modeling of these CN and C₂H hyperfine

Table 7
Comparison of Fractional Abundances Relative to ^{13}CO , $N(X)/N(^{13}\text{CO})$

Molecule	Protoplanetary Disk				YSO				Carbon Star IRC +10216
	LkCa 15	TW Hya	V4046 Sgr	IRAS 4A	IRAS 4B	IRS 7B	IRAS 16293		
	This Work ^a	1	2	2	3	3	4 ^b	5 ^b	
HCN	1.87(-3)	5.0(-3)	2.6(-3)	1.7(-3)	1.4(-3)	7.1(-4)	5.6(-4)	6.6(-4)	0.56
HCO ⁺	1.53(-3)	9.17(-4)	2.4(-4)	3.3(-4)	5.8(-4)	8.4(-4)	2.5(-3) ^b
H ₂ CO	0.63–2.43(-3)	0.39–4.72(-3)	2.6(-3)	3.1(-3)	8.4(-4)
CS	1.95(-3)	...	5.8(-3)	<7.0(-3)	1.1(-3)	7.9(-4)	9.6(-4)	1.8(-3)	0.01
CN	0.01	0.04	0.40	0.08	9.6(-4)	4.8(-5)	0.01
C ₂ H	3.4	...	21	8	3.2(-3)	1.3(-4)	0.09

References. (1) Thi et al. (2004), (2) Kastner et al. (2014), (3) Blake et al. (1995), (4) Watanabe et al. (2012), (5) Schöier et al. (2002), (6) Groesbeck et al. (1994), (7) Pulliam et al. (2011).

^a Abundances computed assuming $T_{\text{ex}} = 20$ K, except for CN ($T_{\text{ex}} \approx 3.9$ K) and C₂H ($T_{\text{ex}} \approx 2.9$ K), where we adopt $N(^{13}\text{CO})$ at $T_{\text{ex}} = 4$ K.

^b Values adapted from Table 5 of Schöier et al. (2002; IRAS16293), Table 4 of Watanabe et al. (2012; IRS 7B), and Table 5 of Pulliam et al. (2011; IRC +10216), assuming $N(\text{H}_2)/N(^{13}\text{CO}) = 6 \times 10^5$.

complexes indicates that the emission from both species is optically thick and suggests either that the emission from these molecules originates from very cold (<10 K) regions of the disk, or that the molecules reside in tenuous upper atmosphere regions of the disk and are subthermally excited. We also find that the column densities for CN and C₂H are comparable to those of isotopologues of CO, implying that these molecules are very abundant in the disk of LkCa 15.

As with our molecular line surveys of TW Hya and V4046 Sgr (Kastner et al. 2014), the results of our molecular line survey of LkCa 15 stress the need for additional molecular line surveys of protoplanetary disks that cover a range of pre-MS evolutionary states; interferometric imaging of protoplanetary disks in transitions of CN and C₂H, as well as other potential molecular tracers of UV and/or X-ray radiation; and detailed modeling of T Tauri star-disk systems that includes the effects of high-energy radiation on disk chemistry and is constrained by observations of pre-MS UV and X-ray fields. In general, we note that interferometric spectral surveys will become very important in the upcoming years. Unbiased line surveys carried out with interferometers would allow for the simultaneous determination of the chemical inventory of disks and the emitting region sizes and morphologies of these molecules.

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