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Selecting background galaxies in weak-lensing analysis of galaxy clusters

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ABSTRACT

In this paper, we present a new method to select the faint, background galaxies used to derive the mass of galaxy clusters by weak lensing. The method is based on the simultaneous analysis of the shear signal, that should be consistent with zero for the foreground, unlensed galaxies, and of the colours of the galaxies: photometric data from the COSM_{ic} eV_{OL}ution Survey are used to train the colour selection. In order to validate this methodology, we test it against a set of state-of-the-art image simulations of mock galaxy clusters in different redshift [0.23–0.45] and mass [$0.5\text{--}1.55 \times 10^{15} M_{\odot}$] ranges, mimicking medium-deep multicolour imaging observations [e.g. *Subaru*, Large Binocular Telescope]. The performance of our method in terms of contamination by unlensed sources is comparable to a selection based on photometric redshifts, which however requires a good spectral coverage and is thus much more observationally demanding. The application of our method to simulations gives an average ratio between estimated and true masses of $\sim 0.98 \pm 0.09$. As a further test, we finally apply our method to real data, and compare our results with other weak-lensing mass estimates in the literature: for this purpose, we choose the cluster Abell 2219 ($z = 0.228$), for which multiband (*BVRi*) data are publicly available.

Key words: gravitational lensing: weak – galaxies: clusters: individual: A2219.

1 INTRODUCTION

Being the largest gravitationally bound structures in the Universe, galaxy clusters are the most powerful gravitational lenses on the sky. Their lensing signal is often detectable up to a few Mpc from the cluster centre. While strong lensing events such as multiple images of distant galaxies and gravitational arcs occur in the cluster cores, at larger radii lensing by clusters appears in the so-called weak-lensing regime. In this case, the lensing induced distortions of the galaxy shapes are tiny and the lensing signal is detectable only by averaging over ensembles of a sufficient number of lensed galaxies. The measurement of their coherent distortions can be used to constrain the gravitational potential of the cluster and therefore to map its mass distribution (Bartelmann & Schneider 2001).

Gravitational lensing offers several advantages compared to other methods to measure cluster masses. Indeed observations of the X-ray emission by the cluster intra-cluster medium or the kinematics of the cluster galaxies can be used to measure masses only under

the assumption of some sort of equilibrium between the baryons and the cluster gravitational potential. On the other end, there are several possible sources of biases in the weak-lensing mass estimates due to residual instrumental distortions, dilution, wrong assumptions on the source redshift distribution, large-scale structure and halo triaxiality which should be properly taken into account in the analysis (Hoekstra 2003; Corless & King 2007; Hoekstra et al. 2011; Giocoli et al. 2014; von der Linden et al. 2014; Hoekstra et al. 2015; Sereno et al. 2015).

Image simulations proved to be very helpful for understanding some of the limitations of the lensing analysis and for quantifying and possibly reduce these biases. The weak-lensing community has been involved in a series of large simulation-based challenges to identify the best performing algorithms for shear measurements (Erben et al. 2001; Heymans et al. 2006; Massey et al. 2007; Bridle et al. 2010; Kitching et al. 2012). Despite the fact that none of these experiments addressed the issue of shear measurability in regimes of relatively large shear signal (such as around clusters), these simulations resulted to be extremely useful to quantify measurement biases and to calibrate the methods. Meneghetti et al. (2010) used the image simulation software SKYLENS (Meneghetti et al. 2008) to mimic *Subaru* and *Hubble Space Telescope* observations of a set

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of mock galaxy clusters. These simulated observations were then processed with a Kaiser–Squires–Broadhurst (KSB) weak-lensing pipeline to measure the ellipticities of background galaxies (Kaiser, Squires & Broadhurst 1995; Luppino & Kaiser 1997; Hoekstra et al. 1998). Comparing different techniques to convert the shear measurements into mass estimates, it was possible to quantify the scatter induced by triaxiality and sub-structure. In a subsequent work, Rasia et al. (2012) used analogous image simulations to investigate the dependence of accurate mass measurements on the cluster environment.

In this paper, we specifically address the issue of the correct identification of background galaxies for the weak-lensing analysis. Unlensed sources in the sample used for the shape measurement cause the dilution of the lensing signal which is stronger at small radii due to the higher number density of cluster galaxies, leading to underestimate the cluster mass.

Several approaches have been proposed to achieve this goal, which generally employ multiband photometry. When photometric data are available in only two bands, the selection of background sources is generally done by identifying the galaxies redder than the cluster red-sequence (Broadhurst et al. 2005). A suitable colour cut has to be applied in order to avoid the contamination by faint red cluster galaxies which can result from dust reddening, intrinsic scatter in galaxy colours and measurement errors (Okabe et al. 2013). However, the number density of background galaxies obtained by such selection could not allow us to reach the accuracy in the statistical analysis required by the weak-lensing analysis of galaxy clusters. Galaxies bluer than the cluster red-sequence, chosen by properly defining a colour offset, can also be included as background galaxies but are highly contaminated by foreground cluster members (Medezinski et al. 2010; Huang et al. 2011; Okabe et al. 2013). If more than two bands are available, a more careful selection can be attempted by identifying the regions of the colour–colour (CC hereafter) diagrams populated by background galaxies (Medezinski et al. 2010; Umetsu et al. 2010, 2012; Medezinski et al. 2013).

Here, we propose a novel approach which optimises the selection of lensed background sources by employing the photometry of galaxies in the COSMic eVolution Survey (COSMOS). The colours of the COSMOS galaxies are used as training set for the separation of background/foreground galaxies in the images of the galaxy clusters under study. The method is based on the assumption that the galaxies in the COSMOS data set (Ilbert et al. 2009) are representative of the colours and of the redshift distribution of the global galaxy population in the Universe and that they allow us to characterize the distribution in the CC space of galaxies in different redshift ranges. Simultaneously, we tune the colour cuts on the basis of the amplitude of the shear signal of the background-selected galaxies. We use simulations made with the SKYLENS software to validate this method assuming different colour information and different lens masses and redshifts. We also made comparison with others approach, such as the selections based on the identification of the cluster red-sequence and on photometric redshifts. Finally, for each selection method we derive an estimate of the cluster mass, which we compare to the true mass of the simulated cluster. While testing the performance of the selection method, we also process the simulated data as done for real observations, including the process of measuring the shear from the galaxy images. This serves as a further validation test for the weak-lensing analysis pipeline that we recently used in Radovich et al. (2015).

To test our approach on real observations, we apply our method to the galaxy cluster Abell 2219, for which medium-deep *Subaru* and

LBT archival observations in the *BVRi* bands are available. Abell 2219 is a very well-known cluster for which several independent weak-lensing analyses have been published (Bardeau et al. 2007; Hoekstra 2007; Okabe et al. 2010).

The paper is structured as follows: in Section 2, we describe the image simulations used to validate the method. The code used to derive the galaxy ellipticities is described in Section 3. The selection method developed is outlined in Section 4 and compared with the method proposed by Medezinski et al. (2010). In Section 5, we describe the application of the selection method to simulated clusters using different combinations of photometric bands, cluster redshift and mass. In Section 5.2, we quantify the contamination of the selections, and in Section 5.4, we discuss how well the weak-lensing masses reproduce the true masses of the simulated galaxy clusters. In the second part of the paper, we show the analysis of Abell 2219 (Section 6). More precisely, we give an overview of the properties of Abell 2219 and of the previous lensing analyses of this cluster. Section 6.3 describes our data reduction. Section 6.4 outlines the selection of the background galaxies and the derivation of the cluster mass, which is subsequently compared to previous results in the literature. Finally, we summarize our conclusions in Section 7. Throughout this paper, we assume a Λ cold dark matter cosmology with $\Omega_{\Lambda} = 0.7$, $\Omega_{\text{M}} = 0.3$, $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$.

2 LENSING SIMULATIONS

To test and validate the selection method, we use realistic image simulations produced with SKYLENS (Meneghetti et al. 2008, 2010; Rasia et al. 2012). The code uses real galaxies taken from the *Hubble* ultra-deep field (HUDF) to create mock observations of patches of the sky with virtually any telescope. Being coupled with a ray-tracing code, it allows us to include the lensing effects from any mass distribution along the line of sight. As shown in Meneghetti et al. (2010) and in Rasia et al. (2012), it is particularly suitable to simulate weak and strong lensing effects by galaxy clusters.

The mock galaxy clusters used here are generated with the code MOKA (Giocoli et al. 2012). The cluster mass distributions have multiple components, namely a triaxial halo, modelled with a Navarro–Frenk–White (NFW) profile (Navarro, Frenk & White 1997), substructures described by truncated singular isothermal spheres (Metcalf & Madau 2001), and a brightest central galaxy (BCG) modelled with an Hernquist density profile (Hernquist 1990). More details about MOKA can be found in Giocoli et al. 2012.

To evaluate the performance of the selection algorithm, we need to populate the field of view around a mock galaxy cluster with cluster members, foreground and background galaxies. Most importantly, galaxy colours have to be as much realistic as possible. MOKA populates the dark matter sub-haloes using the halo occupation distribution technique. Stellar masses and *B*-band luminosities are assigned to each galaxy accordingly to the mass of the dark matter (sub-)halo within which it formed, following Wang et al. (2006). The morphological type and the spectral energy distribution (SED) are defined on the basis of the stellar mass and such to reproduce observed morphology–density and morphology–radius relations (van der Wel 2008). We verified that colours reproduce the expected red-sequences for clusters over a large range of redshifts.

The photometry of foreground and background galaxies is computed starting from the best-fitting spectral energy distribution (SEDs) determined by Coe et al. (2006) for the HUDF, and convolving them with filter transmission.

The redshift distribution of the input sources also matches that of the HUDF. Given the exposure times and the observational set-up

Table 1. Summary of the setups of the simulations produced.

Simulation	Exp. time (s)	z_l	Mass ($10^{15} M_{\odot}$)
Case 1	$B(720); V(1080); R(3330)$	0.23	1
Case 2	$B(720); V(1080); R(3330); i(3005.8)$	0.23	1
Case 3	$g(1800); r(3600); i(1800)$	0.23	1.55
Case 4	$g(1800); r(3600); i(1800)$	0.35	1.55
Case 5	$g(1800); r(3600); i(1800)$	0.45	1.55
Case 6	$g(1800); r(3600); i(1800)$	0.23	0.5
Case 7	$g(1800); r(3600); i(1800)$	0.23	0.75

used in our simulations, the median redshift is $z_m \sim 1$. All simulations in this work have an FOV of $30 \text{ arcmin} \times 30 \text{ arcmin}$. For simplicity, we adopt a spatially constant Gaussian point spread function (PSF) with full width at half-maximum (FWHM) 0.6 arcsec .

We produced several sets of images with different cluster mass and redshift, and combination of filters which we analyse in Section 5.

One set of simulations mimics the archival *Subaru* and LBT observations available for Abell 2219 which is analysed in the second part of the paper. More precisely, we simulate images in the *BVRi* bands with exposure times as listed in Table 4. We reproduce the telescope throughputs, according to information retrieved from the *Subaru* and LBT telescope web sites. In this case, the cluster simulated with MOKA has the same redshift of Abell 2219 and its virial mass is chosen to be $1 \times 10^{15} M_{\odot}$, consistent with previous mass measurements found in the literature for this cluster. We assume a concentration of 3.8, consistent with values expected for clusters of this mass on the basis of recent concentration–mass relations (Meneghetti et al. 2014). Using this set of simulations, we explore the effect of several selection techniques on the mass reconstruction, based on different filter combinations. Hereafter, the analysis of this data set with our selection method is coded with case 1 and case 2 if the photometry in *BVR* and *BVRi* bands is used, respectively, while analyses of this data set with alternative selection methods are coded with case 1–2.

Then we investigate how the method works under different cluster redshifts. In this case, we mimic *Subaru* Suprime-Cam observations in the *gri* bands. Here we consider the case $M_{\text{vir}} = 1.55 \times 10^{15} M_{\odot}$, $c_{\text{vir}} = 3.8$, and three cluster redshifts: $z_l = 0.23, 0.35, 0.45$ which are coded as cases 3, 4 and 5, respectively.

Finally, we investigate the mass dependence on the selection method. In this case, we simulate *Subaru* Suprime-Cam observations in the *gri* bands setting $z_l = 0.23$, $c_{\text{vir}} = 3.8$ while we vary the cluster mass: $M_{\text{vir}} = 0.5 \times 10^{15}$ and $0.75 \times 10^{15} M_{\odot}$; the cases 6 and 7, respectively. All simulations with photometry in *gri* bands have an exposure time in the *g* and *i* bands of 1800 s while we double it in the *r* band, which is used for the shear measurement. The setups of the simulations are listed in Table 1.

We derive the multiband photometry running SExtractor in dual-mode, with the *R*- or *r*-band images used as detection images. In the next section, we describe the shear measurement on these images that closely resembles that applied to the data analysis of A2219.

3 SHEAR MEASUREMENTS

The observed shear signal produced by the gravitational field of the clusters is measured using the KSB approach (Kaiser et al. 1995; Luppino & Kaiser 1997; Hoekstra et al. 1998). This method uses stars to remove from galaxy ellipticities the PSF contribute

due to the atmosphere and the telescope optics, giving a shear estimator uncontaminated by systematics. These distortions can be taken into account considering the PSF as made of two components: one isotropic, which smears images mimicking the effect introduced by seeing, and one anisotropic (due to guiding errors, co-addition, optics, bad focusing, etc) that introduces anisotropic distortions which mimic a lensing signal. The effect of the PSF anisotropy depends on the galaxy size and is stronger for the smallest sources. In the formalism described below, source ellipticity is a complex (two components) quantity. The correction of the observed ellipticity e_{obs} for the anisotropic component, as prescribed by the KSB method, is given by

$$e_{\text{aniso}} = e_{\text{obs}} - P^{\text{sm}} p, \quad (1)$$

$$p = e_{\text{obs}}^* P^{\text{sm}* -1}, \quad (2)$$

where p measures the PSF anisotropy and P^{sm} is the smear polarizability tensor given in Hoekstra et al. (1998). The tensor p can be estimated from images of stellar objects (the quantities with * symbol are computed on stars). The relation between the intrinsic ellipticity e of a galaxy and the reduced shear $g = \gamma / (1 - \kappa)$ will be

$$e_{\text{aniso}} = e + P^{\gamma} g, \quad (3)$$

where the tensor P^{γ} , called the pre-seeing shear polarizability tensor (Luppino & Kaiser 1997), describes the effect of seeing. P^{γ} has the following expression:

$$P^{\gamma} = P^{\text{sh}} - P^{\text{sm}} P^{\text{sh}*} P^{\text{sm}* -1} = P^{\text{sh}} - P^{\text{sm}} q, \quad (4)$$

with $q = P^{\text{sh}*} P^{\text{sm}* -1}$ and P^{sh} being the shear polarizability tensor (Hoekstra et al. 1998).

We use the KSB implementation described in Radovich et al. (2015). The main feature of our pipeline is that the PSF is modelled as a polynomial function of the position, evaluated with the PSFEX software (Bertin 2011), which extracts models of the PSF from the image. The p and q terms are derived at each galaxy's position using this PSF model. For each galaxy, the output quantity $e_{\text{iso}} = e_{\text{aniso}} P^{\gamma -1}$ is computed and used to derive the average reduced shear, being $\langle g \rangle = \langle e_{\text{iso}} \rangle$ under the assumption that the contribution of the intrinsic ellipticities vanishes when averaging over a large number of galaxies.

The main quantities described by the KSB approach are weighted through a Gaussian window function of scalelength θ to beat down the noise. The size of the window function has a crucial role since, if it is not properly chosen, noise will bias the ellipticity measurements. We find the size of the window function that maximizes the ellipticity signal-to-noise ratio (equation 5), defined by the equation 16 in Erben et al. (2001):

$$SNe(\theta) = \frac{\int I(\theta) W(|\theta|) d^2\theta}{\sigma_{\text{sky}} \sqrt{\int W^2(|\theta|) d^2\theta}}; \quad (5)$$

as described by Huang et al. (2011).

Hoekstra et al. (1998) first pointed out that quantities from stellar objects should be calculated with the same scale as the object to be corrected. This is done in our implementation of the KSB analysis. The selection of stars used to correct the distortions introduced by the PSF components is made in the magnitude MAG AUTO¹ versus $\delta = \text{MU MAX} - \text{MAG AUTO}$ space, where MU MAX is the

¹ SExtractor parameter indicating Kron-like elliptical aperture magnitude.

SETRACTOR parameter indicating the peak surface brightness above the background. The quantity δ is used as the estimator of the object size to separate stars from galaxies. When we analyse simulations, we use stars with r or R magnitudes in the range [15, 19] to correct for the effects introduced by PSF anisotropy and seeing.

The tangential and the cross components of the reduced shear, g_T and g_\times , are obtained from the quantities $e_{1\text{iso}}$ and $e_{2\text{iso}}$ used, respectively, as g_1 and g_2 in equations (6) and (7):

$$g_{T(i)} = -g_{1(i)} \cos 2\varphi - g_{2(i)} \sin 2\varphi; \quad (6)$$

$$g_{\times(i)} = -g_{1(i)} \sin 2\varphi + g_{2(i)} \cos 2\varphi. \quad (7)$$

Here (i) indicates the i th galaxy and φ the position angle determined assuming the position of the BCG as the cluster centre. The radial profiles are obtained by averaging both shear components in annuli centred on the BCG.

We exclude from the catalogue galaxies with $\text{SNe} < 5$ (our choice of $\text{SNe} < 5$ corresponds to roughly signal to noise ratio (SNR) < 10 , if we use the definition $\text{SNR} = \text{FLUX AUTO}/\text{FLUXERR AUTO}$, where FLUX AUTO and FLUXERR AUTO are SETRACTOR quantities) and galaxies with corrected ellipticities $e_{1\text{iso}} > 1$ or $e_{2\text{iso}} > 1$, for which the ellipticity measurement is not meaningful. The performances of our KSB code are extensively discussed in Radovich et al. (2015). In addition, in the appendix, we give an estimate of the bias in the shear measurement using the set of simulations tailored on observations of Abell 2219 as described in the Section 2. We find that our KSB pipeline on average underestimates the measured shear by about 5 per cent with respect to the input values. Therefore we apply a calibration factor of 1.05 to the measured ellipticities and use the corrected ellipticities to derive cluster masses.

4 SELECTION OF THE BACKGROUND SAMPLE

In this section, we describe the technique that we developed for the selection in CC space assuming that observations in at least three bands can be used.

Medezinski et al. (2010) did an extensive work on the usage of colours to identify the background galaxies for the weak-lensing analysis. Using Subaru Suprime-Cam observations of the three clusters A1703 ($z = 0.258$), A370 ($z = 0.375$), RXJ 1347-11 ($z = 0.451$), they derive the CC diagrams of the galaxies in these fields. More precisely, they used $r-i$ versus $g-r$ (A1703); $R-z$ versus $B-R$ (A370) and $R-z$ versus $V-R$ (RXJ 1347-11).

In Medezinski et al. (2010) (fig. 1), the cluster member population is identified in CC space adding the information on the mean distance of all objects from the cluster centre in a given cell of the CC diagram. The region in the CC space, where the mean distance is lower, corresponds to an overdensity of galaxies comprising the cluster red-sequence and blue later type cluster members. The identification of foreground and background populations is done looking at the galaxy overdensities in CC space and at the corresponding weak-lensing signal. In particular, CC diagrams (fig. 2 in Medezinski et al. 2010) appear to be characterized by a clump of red objects with a rising weak-lensing signal at small radii corresponding to background sources in the redshift range $0.5 < z < 2$. This clump is located in the lower right side of the CC diagrams. Two others clumps lying blueward of the cluster in CC space are populated by blue objects. The redder one is located near the centre of the CC diagrams and is populated by foreground unlensed sources since they show a little weak-lensing signal and no clustering near the cluster centre in their spatial distribution. Galaxies belonging to the

bluer clump are also unclustered near to the cluster centre and have a weak-lensing signal similar to that of the red background galaxies indicating that these blue objects are background galaxies in the redshift range $1 < z < 2.5$. This clump is located in the left-hand side of the CC diagrams.

Here we propose a variation of the method adopted by Medezinski et al. (2010). More precisely, we introduce two novelties in the selection process, namely:

(i) we use the galaxies in the COSMOS, including their colours and photometric redshifts, to identify the regions of the CC diagrams populated by the galaxies at redshift higher than the cluster. COSMOS colours are used as a training set for the separation of background/foreground galaxies in the images of the galaxy clusters under study. We use the COSMOS photometric redshift catalog of Ilbert et al. (2009). The redshift measurements are based on the photometry in 30 broad, intermediate and narrow bands spanning a range of wavelengths from the UV to the mid-IR. Based on a sub-sample of spectroscopically confirmed sources, the accuracy of the photo- z s at magnitudes $i_{AB}^+ < 24$ and redshifts $z < 1.25$ is $\sigma_{\Delta z/(1+z_s)} = 0.012$, where $\Delta z = z_s - z_p$ is the difference between the spectroscopic and the photometric redshifts. Photo- z s up to $z \sim 2$ are less accurate: $\sigma_{\Delta z/(1+z_s)} = 0.06$ at $i_{AB}^+ \sim 24$. The source redshift distribution is characterized by median redshifts $z_m = 0.66$ and 1.06 for sources in the magnitude ranges $22 < i_{AB}^+ < 22.5$ and $24.5 < i_{AB}^+ < 25$, respectively.

(ii) A fine-tuning in the selection of background galaxies is done by maximizing the amplitude of the shear profile, in particular in the inner radial bins, where the dilution of the shear by cluster members is higher.

The implementation of our method is based on the following steps.

(i) We determine on the COSMOS ($col1$, $col2$) CC plane the line ($a \times col1 + col2 + c = 0$) separating background and foreground galaxies at the cluster's redshift. This line is defined as the tangent to the ellipse fitting the 20 per cent density contour in the foreground galaxies' CC distribution, parallel to the ellipse major axis.

(ii) For each i -galaxy observed in the cluster field with colours $col1_i$, $col2_i$, we compute its distance from this line:

$$d_i = \frac{a \times col1_i + col2_i + c}{\sqrt{a^2 + 1}}; \quad (8)$$

and make an initial separation between foreground ($d_i \geq 0$) and background ($d_i < 0$) galaxies.

(iii) In a similar way, we can define in the COSMOS CC plane additional colour cuts that allow us to improve the separation of foreground and background galaxies at the cluster's redshift as detailed later in Section 5.

(iv) We define a threshold in magnitude below which uncertainties in the measured fluxes are too high to derive reliable colours, but which most likely come from background galaxies: galaxies fainter than this threshold are included in the background sample.

(v) We measure the shear signal (equation 6) from the background galaxies and fine-tune the a , c coefficients to maximize its amplitude. To this end, we randomly vary a around ± 50 per cent of its initial value, while we randomly vary c of ± 0.2 around its initial value. Then we make a new selection of background galaxies $d_i(a_n, c_n) < 0$, with a_n and c_n being the coefficient values in the n th realization: as a result of this iterative procedure, we select the values of a , c that maximize the amplitude of the shear signal in the three inner radial bins. Simultaneously, we also check that

the tangential shear profile of the foreground and cluster member galaxies is consistent with zero at all radii.

Despite the limited size of the COSMOS field, several tests show that using the COSMOS colours as starting point for the identification of distant sources is safe. For example, if we apply the colour cuts based on the COSMOS photometry to select high-redshift galaxies in the (HUDF), we find a low contamination by low-redshift sources. The method can be also applied to data observed with a different filter set. In the examples discussed later in this paper, we use the software ZEBRA (Feldmann et al. 2006) to derive the SEDs of the COSMOS galaxies and we use them to calculate magnitudes in other filters by convolving with the filter transmission curves.

In Section 5, we show the application of this approach to simulated images of clusters in different redshift and mass ranges exploring how it works with some combinations of filters. The fraction of background/foreground misidentified galaxies is derived and discussed in the Section 5.2.

5 ANALYSIS OF SIMULATED DATA

In this section, we show the results of the application of our selection method to simulations. First, we consider different combinations of cluster redshifts (z_1) and filters: $z_1 = 0.23$ (*BVR*, *BVRi*, and *gri*; cases 1, 2, 3), $z_1 = 0.35$ (*gri*; case 4) and $z_1 = 0.45$ (*gri*; case 5). In these cases, the mass of the lens is $M_{\text{vir}} = 10^{15} M_{\odot}$ (cases 1, 2) and $M_{\text{vir}} = 1.55 \times 10^{15} M_{\odot}$ (case 3, 4, 5). Cases 1, 2 are tested also against alternative methods to identify the background galaxies. Finally, we complement these simulations with additional two test cases, where we keep the lens redshift fixed at $z_1 = 0.23$ and assume the lens mass to be $M_{\text{vir}} = 0.5 \times 10^{15}$ and $0.75 \times 10^{15} M_{\odot}$ (cases 6 and 7). The filter combination used in these last two cases is the same used in the test case 3.

5.1 Selection process

As outlined above, in all these scenarios, we start by deriving the CC diagrams of the COSMOS galaxies for the corresponding filter combinations. As an example, Fig. 1 shows the $r-i$ versus $g-r$

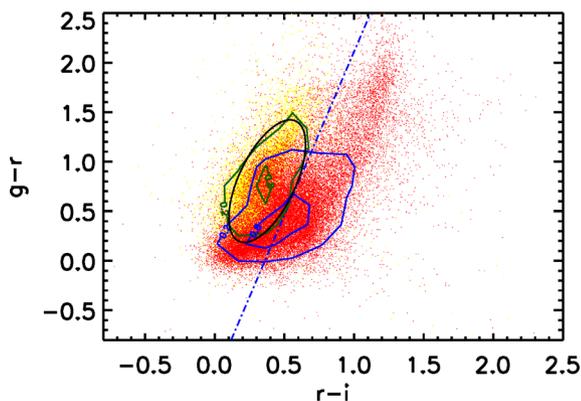


Figure 1. First step of the selection process to be applied to simulated data: for the test case 5, outlined in Section 5.1, we show the CC diagram of COSMOS galaxies and the identification of background galaxies based on the photometric redshift estimates of Ilbert et al. (2009). The galaxies at redshift higher and lower than the cluster redshift ($z_1 = 0.45$) are indicated by red and yellow dots, respectively (the corresponding number density contours are given by the blue and green solid lines).

diagram of the COSMOS galaxies and the foreground/background galaxy separation in the case of a cluster at redshift $z_1 = 0.45$ (case 5). Galaxies with photo- $z_s > 0.47$ are plotted in red, while those with photo- $z_s \leq 0.47$ are plotted in yellow. A margin of 0.02 in redshift is used to account for photo- z accuracy when we select galaxies at a redshift higher than the cluster redshift. Density contours of background and foreground galaxies are shown in blue and in green, respectively. Also displayed is the ellipse fitting the 20 per cent foreground contour; the dot-dashed line is the tangent to this ellipse.

Fig. 2 summarizes the full procedure. The dot-dashed lines in the left-hand panels correspond to the initial values of the a , c coefficients, derived for all the test cases here considered. Obviously, the results depend on the filter combination and on the lens redshift. The coefficients a and c obtained for test cases 1–5 are listed in the Table 2. Test cases 6 and 7 are omitted because the results are identical to test case 3 (the selection based on COSMOS data does not depend on the lens mass). The second step of our procedure consists in the fine-tuning of the selection based on the amplitude of the shear profile derived from the galaxies which are classified as background sources. Therefore, we consider now the CC diagrams built using all galaxies in the simulated cluster fields. We start by identifying the background sources based on the lines derived in the first step. For each galaxy, we compute the distance d_i from these lines and looking at the amplitude of the shear profile of sources with $d_i < 0$ we derive new coefficients (a , c) as explained in the Section 4. The adjusted lines separating the background from the foreground and cluster galaxies are given for the test cases 1–5 by the dashed lines in the left-hand panels of Fig. 2. The lines derived from the COSMOS data are also shown for comparison (dot-dashed lines). The optimized coefficients after the fine tuning with the shear signal are listed in the Table 2. Given that fine tuning is done on the CC diagrams, which also include the cluster galaxies, being the lens different, the lines optimized for the test cases 6 and 7 differ slightly from that derived for the test case 3.

The final results of the selection process are shown in the left-hand panels of Fig. 2 for the test cases 1–5, where background, foreground, and cluster galaxies are shown in magenta, green, and orange, respectively. The detailed classification of the sources in the CC diagrams corresponding to each of the test cases is given below.

cases 1, 2: When using the colours $\text{col1} = V-R$ and $\text{col2} = B-R$ to separate background, cluster and foreground galaxies, we find that the selection which maximizes the amplitude of the tangential reduced shear profile is obtained using the following criteria for the background galaxies:

- (i) $d_i \geq 0$, $24.2 \leq R < 26$ mag;
- (ii) $d_i < 0$, $21 < R \leq 26$ mag;
- (iii) $d_i \geq 0$, $V-R < 0.1$ and $21 < R < 24.2$ mag.

The foreground and the cluster galaxies are separated as follows:

- (i) foreground galaxies: $d_i \geq 0$, $0.7 \leq B-R < 1.3$, $0.12 \leq V-R < 0.35$ and $R < 24.2$ mag;
- (ii) cluster galaxies: $d_i \geq 0$, $0.2 < V-R < 0.6$, $1.3 \leq B-R < 2.2$ and $R < 24.2$ mag.

If the selection is performed in the colour space $R-i$ versus $B-V$, the criteria (i) and (ii) above obviously remain valid. In this case, the background sample is composed by galaxies satisfying the criteria (i) and (ii) and by those satisfying the following conditions: $d_i \geq 0$, $R-i < 0.03$ and $21 < R < 24.2$ mag.

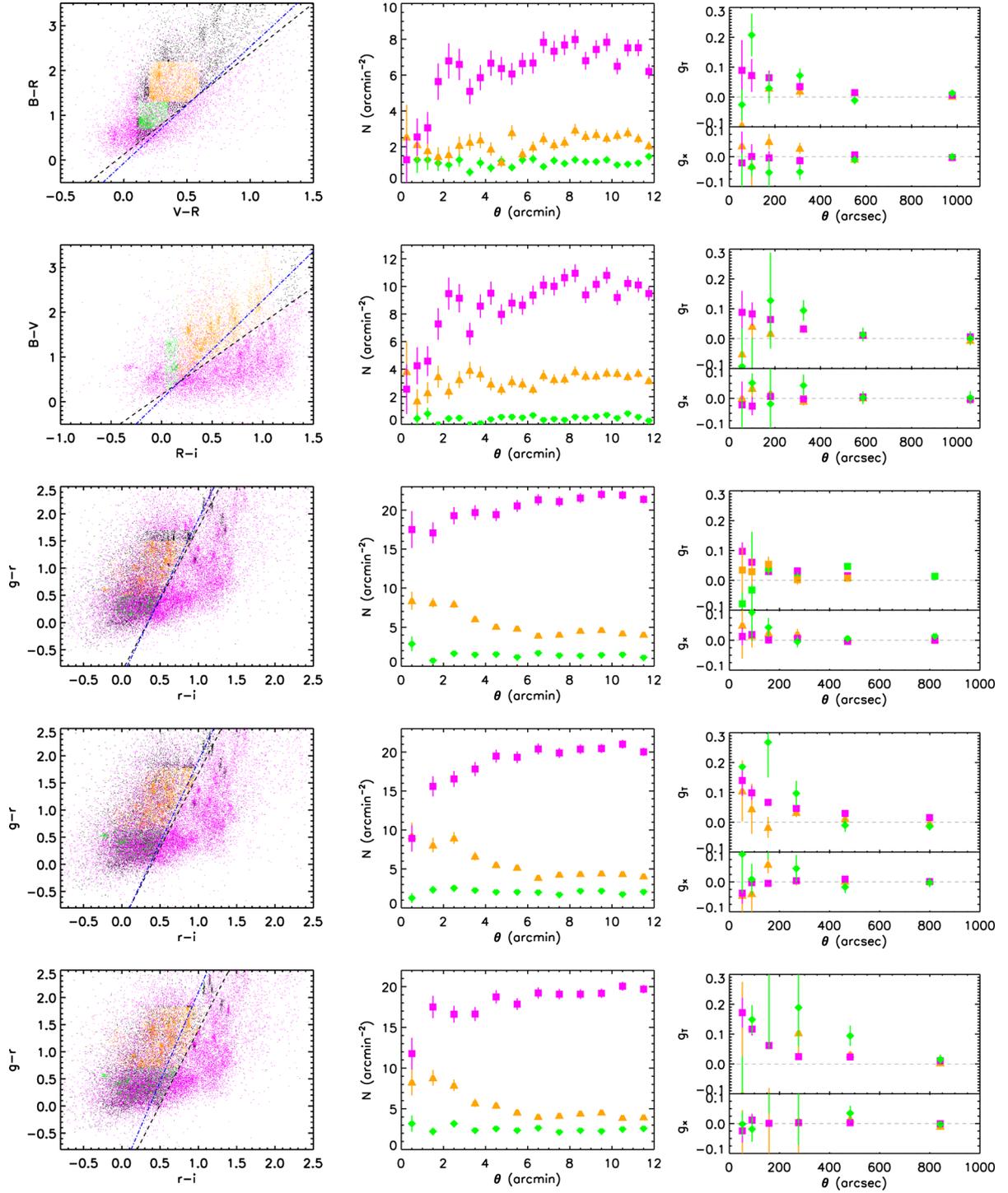


Figure 2. Second step of the selection process applied to the simulations. The five rows of panels refer to the test cases 1–5 for which we give a lengthy description in Section 5.1. The left-hand panels show the identification of background (magenta), foreground (green), and cluster galaxies (orange) in colour space. Unclassified galaxies are indicated by black dots. The dot–dashed lines in the left-hand panels correspond to the initial separation between foreground and background galaxies based on COSMOS data. Dashed lines correspond to this separation after the fine-tuning with the shear signal. In the middle panels, we show the radial number density profiles of background, foreground, and cluster galaxies. The results for the three galaxy sub-samples are shown using the same colour-coding as in the left-hand panels. The right-hand panels show the corresponding radial profiles of the tangential (g_T) and cross components (g_x) of the reduced shear.

In addition, we identify the foreground and the cluster galaxies as follows:

(i) foreground galaxies: $d_i \geq 0$, $0.03 \leq R-i < 0.16$ and $R < 24.2$ mag.

(ii) cluster galaxies: $d_i \geq 0$, $0.16 \leq R-i < 1.1$ and $R < 24.2$.

cases 3, 6, 7: in these cases, the cluster redshift is $z_1 = 0.23$ and the filter combination is gri . However, the cluster mass is different. The selection is performed in colour space $col1 = r-i$ and $col2 = g-r$.

Table 2. Coefficients of the line separating the background from the foreground and cluster member galaxies obtained following our procedure outlined above for each test case analysed. Column 1: simulated case; Column 2: line coefficients found using Cosmos photometry; Column 3: line coefficients found after the optimization with the amplitude of the shear profile of background sources.

Simulation	Initial coefficients	Optimized coefficients
Case 1	$a = -2.6, c = 0.07$	$a = -2.25, c = -0.12$
Case 2	$a = -2.2, c = -0.06$	$a = -1.60, c = -0.16$
Case 3	$a = -2.9, c = 1.00$	$a = -2.65, c = 0.90$
Case 4	$a = -3.0, c = 1.10$	$a = -2.75, c = 1.07$
Case 5	$a = -3.3, c = 1.20$	$a = -2.75, c = 1.35$
Case 6	-	$a = -2.69, c = 0.96$
Case 7	-	$a = -2.67, c = 0.93$

In the test case 3, the background sources are selected following the criteria:

- (i) $d_i \geq 0, 25 \leq r \leq 26$ mag;
- (ii) $d_i < 0, 22 < r \leq 26$ mag;
- (iii) $d_i \geq 0, 22 < r < 25$ mag, $g-r > 1.7$.

The foreground galaxies and the cluster members are instead identified as follows:

- (i) foreground galaxies: $d_i \geq 0, 0.1 < g-r < 0.5, r < 24.5$ mag;
- (ii) cluster galaxies: $d_i \geq 0, 0.5 \leq g-r \leq 1.5, r < 24.5$ mag.

We find that these criteria change very little when considering the test cases 6 and 7. case 4: in the case of this cluster at $z_1 = 0.35$, the criteria to identify the background galaxies are very similar to those found for the test case 3. Indeed, the filter combination is identical and the selection must be modified only accounting for the shift of the cluster galaxies in the CC diagram. The optimal criteria for identifying the background galaxies are:

- (i) $d_i \geq 0, 25 \leq r \leq 26$ mag;
- (ii) $d_i < 0, 22 < r \leq 26$ mag;
- (iii) $d_i \geq 0, 22 < r < 25$ mag, $g-r > 1.8$.

The foreground and cluster galaxies are instead selected based on the conditions:

- (i) foreground galaxies: $d_i \geq 0, 0.1 < g-r < 0.65, r < 24.5$ mag;
- (ii) cluster galaxies: $d_i \geq 0, 0.65 \leq g-r \leq 1.78, r < 24.5$ mag.

case 5: finally, in the case of the cluster at the highest redshift among those we have considered ($z_1 = 0.45$), operating in the same colour space of cases 3 and 4, we optimize the identification of the galaxies in the cluster background by applying the following colour and magnitude cuts:

- (i) $d_i \geq 0, 25 \leq r \leq 26$ mag;
- (ii) $d_i < 0, 22 < r \leq 26$ mag;
- (iii) $d_i \geq 0, 22 < r < 25$ mag, $g-r > 1.85$.

Again, we distinguish between cluster members and foreground galaxies:

- (i) foreground galaxies: $d_i \geq 0, 0.1 < g-r < 0.7, r < 24.5$ mag;
- (ii) cluster member galaxies: $d_i \geq 0, 0.7 \leq g-r \leq 1.82, r < 24.5$ mag.

In all cases, for a fraction of galaxies, the identification is uncertain. In Fig. 2, background galaxies that are located too close to foreground galaxies in the CC diagrams, are marked as black dots. The radial profiles of the galaxy number densities and of the

tangential and cross components of the reduced shear, measured in each of the three sub-samples of background, foreground, and cluster galaxies, are shown in the middle and right-hand panels of Fig. 2. Again, we use the magenta, green, and orange colours to distinguish between the three classes of sources.

As expected, in all cases, we notice that the number density of background galaxies drops at small radii. The decrement of number counts is in large part due to magnification, which becomes stronger when approaching the cluster critical lines (Broadhurst et al. 2008). In addition, a fraction of sources at the smallest distances from the cluster centre remains hidden behind the brightest cluster members.

On the contrary, at least in the test cases 3, 4 and 5, the number density profiles of the cluster members have the opposite trend, rising towards the centre of the cluster. This is less evident in the test cases 1 and 2, corresponding to the lowest cluster redshift among those we have investigated. In these cases, the profiles are quite flat and a small increment of galaxy counts is seen only in the innermost radial bin.

The tangential shear profiles of the background galaxies nicely rise towards small radii, as expected in the case of centrally concentrated lenses. The profiles of the cross component of the shear are all consistent with zero, indicating that the measurements are not affected by systematics.

The tangential shear profiles of foreground and cluster galaxies should be consistent with zero. We find this not to be the case in several of the test cases. In the next section, we will see that the origin of this behaviour is due to the contaminations of these two samples by background galaxies. We want to stress that the goal of our selection method is to obtain a sample of background galaxies with a low contamination by unlensed sources, this implies to lose a fraction of background galaxies lying in the same regions of the CC space populated by foreground and cluster members.

5.2 Contamination

Working with simulated data, we can easily quantify how precisely we are able to distinguish between background, foreground, and cluster galaxies. We quantify the performance in terms of contamination of the sub-sample of background galaxies by other kind of galaxies. Although identifying the cluster and the foreground galaxies is not the main goal of this analysis, we can verify also what is the level of contamination by background galaxies in the sub-samples of foreground and cluster galaxies.

We compute the contamination of the background sample by foreground galaxies as

$$f_{\text{for}} = \frac{n_{\text{fb}}}{n_{\text{b}}}, \quad (9)$$

where n_{fb} is the number of galaxies with redshift $z_t \leq z_1$ (z_t input redshift) identified as background galaxies and n_{b} is the total number of background-selected galaxies.

Similarly, the contamination of the background sample by cluster galaxies is given by

$$f_{\text{clus}} = \frac{n_{\text{cb}}}{n_{\text{b}}}, \quad (10)$$

where n_{cb} is the number of cluster galaxies mis-identified as background sources.

We also define the fraction of background galaxies which are incorrectly assigned to the foreground sample,

$$f_{\text{backf}} = \frac{n_{\text{bf}}}{n_{\text{f}}}, \quad (11)$$

Table 3. Summary of the simulation analysis. Column 1: simulated case; Column 2: selection method. Selection in colour space is used in the first seven cases. The bottom three rows refer to the analysis of the cluster field used in test cases 1–2 when the selection of the background galaxies is done with alternative methods: based on photometric redshifts, magnitude cut, and identification of the cluster red-sequence. Columns 3 and 4: fractions of foreground galaxies, f_{for} , and member galaxies, f_{clus} , contaminating the sample of background galaxies. Column 5: total fraction of non-background galaxies contaminating the sample of background galaxies; Columns 6 and 7: fractions of background galaxies contaminating the samples of foreground and cluster galaxies; Column 8: redshift of the lens; Column 9: mass of the lens; Column 10: number density of background galaxies used in the weak-lensing analysis; Column 11: mass recovered from the weak-lensing analysis; Column 12: mean value of β .

Simulation	Sel. method	f_{for}	f_{clus}	f_{tot}	f_{back_f}	f_{back_c}	z_1	Mass ($10^{15} M_{\odot}$)	Density (arcmin^{-2})	Lensing mass ($10^{15} M_{\odot}$)	(β)
Case 1	$V-R$ versus $B-R$	8%	1%	9%	63%	64%	0.23	1	7	$0.96^{+0.16}_{-0.12}$	0.74
Case 2	$R-i$ versus $B-V$	6%	1%	7%	51%	49%	0.23	1	10	$0.85^{+0.11}_{-0.10}$	0.74
Case 3	$r-i$ versus $g-r$	6%	1%	7%	64%	38%	0.23	1.55	22	$1.40^{+0.08}_{-0.08}$	0.74
Case 4	$r-i$ versus $g-r$	8%	1%	9%	58%	40%	0.35	1.55	20	$1.50^{+0.08}_{-0.09}$	0.60
Case 5	$r-i$ versus $g-r$	10%	1%	11%	35%	32%	0.45	1.55	19	$1.50^{+0.11}_{-0.10}$	0.46
Case 6	$r-i$ versus $g-r$	6%	1%	7%	65%	36%	0.23	0.5	22	$0.55^{+0.06}_{-0.05}$	0.74
Case 7	$r-i$ versus $g-r$	6%	1%	7%	64%	37%	0.23	0.75	22	$0.86^{+0.07}_{-0.06}$	0.74
Case 1–2	z_{phot}	5%	1%	6%	37%	38%	0.23	1	8	$1.11^{+0.14}_{-0.14}$	0.62
Case 1–2	$22 \leq R \leq 26$	12%	1%	13%	–	–	0.23	1	13	$0.78^{+0.07}_{-0.08}$	0.72
Case 1–2	VR	9%	1%	10%	–	–	0.23	1	8	$0.83^{+0.11}_{-0.10}$	0.73

where n_{bf} is the number of galaxies with $z_t > z_1$ mis-identified as foreground galaxies and n_f is the total number of foreground galaxies, and the fraction of background galaxies contaminating the sub-sample of cluster galaxies,

$$f_{\text{back}_c} = \frac{n_{\text{bc}}}{n_c}, \quad (12)$$

where n_{bc} is the number of galaxies with $z_t > z_1$ erroneously assigned to the sub-sample of cluster members, which contains n_c galaxies.

These quantities are listed in columns 3–7 of Table 3 for the seven test cases described above (column 5 shows the total contamination, resulting from the both foreground and cluster galaxies, $f_{\text{tot}} = f_{\text{for}} + f_{\text{clus}}$). Using the method that we have developed, we find that the contamination by cluster galaxies is always very small, of the order of ~ 1 per cent in all the test cases. Instead, the major contaminants of the sample of background galaxies are sources in the cluster foreground. In this case, the contamination ranges between ~ 6 per cent and ~ 10 per cent, being larger for lenses at larger redshift, as expected. Given that the total level of contamination is always $\lesssim 11$ per cent, we can conclude that the method works well.

The level of contamination of the sub-samples of foreground and cluster galaxies allows us to interpret some of the results shown in Fig. 2. As noted earlier, the tangential shear profile measured from galaxies in these two sub-samples often varies as a function of radius, while, in absence of lensing, the profiles should be flat and consistent with zero. From Table 3, we see that a significant number of background galaxies are indeed mis-identified as cluster members or foreground galaxies. The contamination of these two sub-samples varies from ~ 32 per cent to ~ 64 per cent. The level of contamination depends not only on the lens redshift, but also on the filter combination used to define the colour space. For example, considering the test case 1, we see that the number of background galaxies included in the samples of foreground or cluster galaxies amounts to ~ 63 – 64 per cent of the total. The addition of the i band allows us to reduce the contamination to ~ 50 per cent. Given that our goal is to optimize the selection of background sources for the subsequent lensing analysis, we consider the contamination of the sub-sample of foreground and cluster galaxies a fair price to pay.

It is also important to note that the method produces the same results when applied to lenses of different masses. This is evident comparing the results for the test cases 3, 6, and 7. While in Section 5.1, we showed that changing the mass of the lens only weakly impacted on the choice of the coefficients a and b , we see here that the contamination of the sample of background sources is also independent on the lens. We have to bear in mind however that the same prescriptions were used to populate the lenses with galaxies and to assign them an SED.

5.3 Comparison to other selection methods

An interesting question is how the performance of the method proposed here compares to that of other techniques to identify lensed galaxies behind clusters. To answer this question, we implement some of these other methods and apply them to our simulations.

We consider here three further approaches, which we apply to the lens at $z_1 = 0.23$ used in the test cases 1 and 2 above. The first is based on the identification of the cluster members through the location of the red-sequence in the colour–magnitude diagram (see e.g. Okabe et al. 2010). Even in this case, some boundaries on both sides of the red-sequence must be defined in order to separate the candidate background sources from the cluster members. This is done again by maximizing the amplitude of the tangential shear profile derived from the sources ending up in the background sample. Therefore we identify the cluster red-sequence in the $V-R$ versus R colour–magnitude diagram and we apply the Okabe’s method.

We further consider a selection based on a simple magnitude cut, selecting the galaxies with $22 \leq R \leq 26$. The idea behind this simple method is that the faintest galaxies are likely to be in the cluster background, while the brightest ones are in large fraction foreground and cluster galaxies.

Finally, we consider the identification of background sources based on photometric redshifts. Of course, the precision of the measurements depends on the number of bands available and on the depth of the corresponding exposures. In order to make a fair comparison to our previous results, we use the simulated $BVRi$ photometric data produced for the test cases 1 and 2. We use the ZEBRA code in maximum-likelihood mode to fit the source SEDs

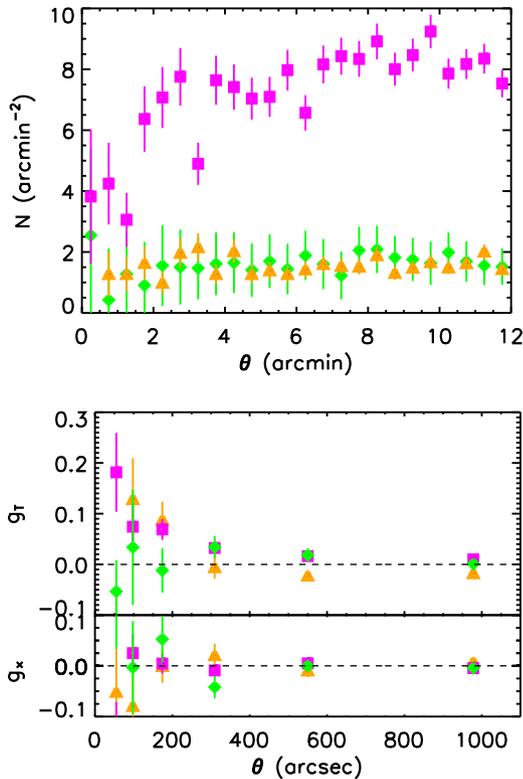


Figure 3. Selection with photometric redshifts. Background galaxies having $z_{\text{phot}} > 0.27$ are marked with magenta squares, cluster galaxies having $0.19 \leq z_{\text{phot}} \leq 0.27$ with orange triangles and foreground galaxies having $z_{\text{phot}} < 0.19$ with green diamonds. The rise towards the centre of the g_t profile of the cluster member galaxies is due to the limited accuracy of the photometric redshifts derived.

with a library of templates and derive the photometric redshifts (Feldmann et al. 2006). Given the limited number of bands used, we decide not to include galaxies with photometric redshifts higher than $z_{\text{phot}} = 2$ and with measurement error $\sigma_z/(1+z) > 0.1$. We select as background galaxies those having $z_{\text{phot}} > 0.27$.

Fig. 3 shows the radial profiles of the galaxy number counts (upper panel) and of the tangential and cross components of the reduced shear (bottom panel) for the sub-samples of background, foreground, and cluster galaxies identified using the photometric redshift technique. Comparing to the test case 2 in Fig. 2, we notice that our conservative cuts lead to a lower number density of background galaxies ($\sim 8 \text{ gal arcmin}^{-2}$) compared to the test case 2 ($\sim 10 \text{ gal arcmin}^{-2}$).

The tangential shear profile derived from the sub-sample of background sources is steeper than measured in test cases 1 and 2. We also notice that the corresponding profile obtained using those galaxies classified as cluster members grows towards the cluster centre, leading to the suspect that this sub-sample is contaminated by background galaxies.

In the last three rows of Table 3, we list the contamination estimates obtained by analysing the simulations with the three alternative methods outlined above. It appears that the selection based on the photometric redshifts leads to a contamination of the sample of background sources similar to that found in test case 2 (the total contamination is ~ 6 per cent). The contaminations of the sub-samples of foreground and cluster galaxies are smaller using the photometric redshift selection than found using our selection in colour space (~ 37 – 38 per cent versus 49 – 51 per cent).

The methods based on magnitude cuts and on the identification of the red-sequence lead to define sub-samples of background sources which have the highest contamination (at the level of ~ 10 – 13 per cent). In these two cases, we did not attempt to separate foreground and the cluster galaxies.

5.4 Mass estimates

As a last test based on simulations, we quantify now how the selection of background sources impacts on the estimate of the lens mass.

Mass estimates are obtained by fitting the reduced tangential shear profiles. Obviously, only the sources classified as background galaxies are used for this analysis. The fitting model is based on the NFW density profile (Navarro et al. 1997):

$$\rho(r) = \frac{\rho_s}{r/r_s(1+r/r_s)^2}. \quad (13)$$

The model depends on two parameters, namely the scale radius r_s and the characteristic density ρ_s . Most commonly, the NFW profile is parametrized in terms of the concentration, $c = r_{\text{vir}}/r_s$, and of the virial mass:

$$M_{\text{vir}} = 4\pi\rho_s r_s^3 \left[\ln(1+c) - \frac{c}{1+c} \right]. \quad (14)$$

The analytic formulas describing the radial profile of the shear for such a model can be found in (Bartelmann 1996). The fit is performed minimizing the log-likelihood function (equation 15; Schneider, King & Erben 2000):

$$l_\gamma = \sum_{i=1}^{N_\gamma} \left[\frac{|\epsilon_i - g(\vec{\theta}_i)|^2}{\sigma^2[g(\vec{\theta}_i)]} + 2 \ln \sigma[g(\vec{\theta}_i)] \right]; \quad (15)$$

where ϵ_i are the observed ellipticities of the background sources, $g(\vec{\theta}_i)$ is the reduced shear predicted by the model at the position $\vec{\theta}_i$, and $\sigma[g(\vec{\theta}_i)] = (1 - |g(\vec{\theta}_i)|^2)\sigma_e$. The variance of the intrinsic source ellipticity is assumed to be $\sigma_e = 0.3$. The computation of the cluster mass requires to know the critical surface density Σ_{cr} of the lens,

$$\Sigma_{\text{cr}} = \frac{c^2}{4\pi G} \frac{D_S}{D_L D_{LS}} = \frac{c^2}{4\pi G} \frac{1}{D_L \beta}; \quad (16)$$

where D_S , D_L and D_{LS} are the angular diameter distances between the observer and the sources, between the observer and the lens, and between the lens and the source, respectively. While D_L is fixed for a given lens, the sources are not all at the same redshift. They have a redshift distribution which should be accounted for. The dependence of Σ_{cr} on the source redshift z_s is contained in the function $\beta = D_{LS}/D_S = \beta(z_s)$. We adopt the approximation of a constant $\langle \beta \rangle$ for all galaxies. Its value, $\langle \beta \rangle = \frac{1}{N} \sum_{i=1}^N \left(\frac{D_{LS_i}}{D_{S_i}} \right)$, is computed using the COSMOS photometric redshifts, applying to the COSMOS catalogue the colour and magnitude cuts which led to the identification of the background galaxies. The values of $\langle \beta \rangle$ in all test cases are listed in column 12 of Table 3. We verify that the $\langle \beta \rangle$ value computed in this way is equivalent to that obtained assuming that all galaxies lie at the same mean redshift $\langle z_s \rangle$: $\beta = \frac{D_{LS}(\langle z_s \rangle)}{D_S(\langle z_s \rangle)}$. In column 11 of Table 3, we summarize the measurements of the virial masses for all the test cases considered. In all cases, the fit is performed keeping the concentration constant and equal to the input value used to generate the lens mass distributions with MOKA ($c = 3.8$). Thus the fit of the tangential reduced shear profiles is done assuming only one free parameter, namely the virial mass. The fit is performed in the radial range of $30 \text{ arcsec} < \theta < 1000 \text{ arcsec}$. The

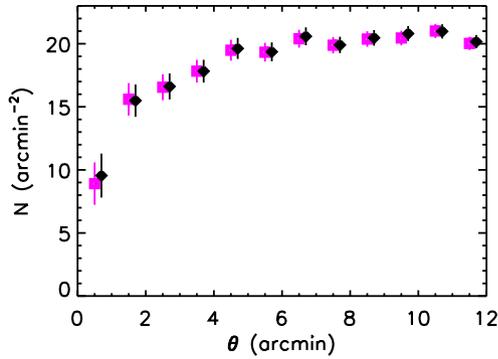


Figure 4. Radial number density profiles of background galaxies using the BCG as the shear centre (magenta squares) and adding an offset of 30 arcsec from the BCG position (black squares). For clarity of the plot, we apply an offset of 0.2 arcmin to the black squares.

results show that the masses obtained from the galaxies selected in colour space differ from the true masses by $\lesssim 10$ –15 per cent. On average, the ratio between estimated and true masses is $\sim 0.98 \pm 0.09$, thus consistent with no mass bias. The values of $\langle\beta\rangle$ based on the COSMOS photometric redshifts are similar in all the test cases where the lens has the same redshift ($\langle\beta\rangle \sim 0.74$). This value is consistent with the true value obtained from the input source redshift distribution, when truncated to mimic the depth of the simulated observations. As expected, $\langle\beta\rangle$ is smaller for lenses at higher redshift (test cases 4 and 5) when the depth of the observation is fixed.

In Table 3, we report also the mass estimates obtained in the three cases where we have used alternative methods to identify the background galaxies (bottom rows). The selections based on the identification of the cluster red-sequence and on the magnitude cuts give the highest contaminations by foreground and cluster galaxies, thus resulting in masses that under-estimate the true masses more significantly (−17 per cent and −22 per cent, respectively). When we select sources based on photometric redshifts derived with the ZEBRA code, we find that the value of $\langle\beta\rangle$ estimated from the photometric redshifts themselves is smaller than the true value ($\langle\beta\rangle \sim 0.62$). In part, this is due to the choice of using only sources with $z_{\text{phot}} < 2$. However, we notice that, for a significant fraction of galaxies, the photometric redshifts underestimate the true redshifts. Indeed, the median photometric redshift of the selected sample is $z_{\text{phot, med}} = 0.86$, while the median true redshift of the same galaxies is $z_{\text{true, med}} = 0.99$. This results in underestimating β . Since this appears in the denominator of equation (16), the mass of the lens is overestimated by ~ 11 per cent.

5.5 Effects of mis-centring

Finally we verify if the selection of background sources can be significantly affected by the choice of the cluster centre, which is used to measure the tangential shear profile of the background galaxies. For this purpose, we repeat the procedure by assuming a centre which is off with respect to the true centre (BCG position) by ~ 30 arcsec. The radial density profiles of background sources selected in both cases of correct and wrong identification of the cluster centre are shown in Fig. 4. They are barely distinguishable, indicating that the mis-centring has little effect on the identification of the background galaxies.

6 ABELL 2219

Having evaluated the performance of the method on simulated data, we apply our procedure to select the galaxies in the background of Abell 2219. The choice of this cluster is motivated by the fact that this is one of the best-studied gravitational lenses in the literature. Some of the previous analyses will be discussed in the following sections.

6.1 General properties

Abell 2219 is a galaxy cluster at $z_1 = 0.228$ with X-ray luminosity $L_x = 2 \times 10^{45}$ ergs s^{-1} in the 0.1–2.4 keV energy range (Ebeling et al. 1998). This makes it one of the most luminous X-ray clusters in the northern sky. Its X-ray surface brightness distribution has quite elongated structure.

The core is dominated by a massive cD galaxy.

Smail et al. (1995) reported the discovery of two systems of giant arcs which allowed them to constrain the (dark) matter distribution in the central region. They found that this is not well aligned with the gas, interpreting this result as the signature of an on-going merger event. On the basis of other multiple images identified by Smith et al. (2005), Richard et al. (2010) estimated a projected mass within a radius $R < 250$ kpc $M_{\text{SL}} = 2.33 \pm 0.23 \times 10^{14} M_{\odot}$. The *Einstein* radius was found to be (15.6 ± 0.6) arcsec.

The study of the internal dynamics of the cluster on the basis of Telescopio Nazionale Galileo spectroscopic data (Boschin et al. 2004) gives a value for the line-of-sight velocity dispersion of $\sigma_v = 1438^{+109}_{-86}$ km s^{-1} , while the mass within the virial region is $\sim 2.8 \times 10^{15} h^{-1} M_{\odot}$. On the basis of the multiwavelength analysis obtained using both optical and X-ray data, Boschin et al. (2004) also suggested that the cluster is not dynamically relaxed and likely to be a merging system.

At X-ray wavelengths, *Chandra* observations reveal the presence of a large-scale shock front (Million & Allen 2009) with $kT \lesssim 16$ keV, high pressure and entropy which can be associated with a merger-driven shock front. The X-ray mass of Abell 2219 was estimated by Mahdavi et al. (2008) on the basis of data available in the *Chandra* X-ray Observatory public archive, obtaining a value of $M_{500} = 17.32 \pm 6.50 \times 10^{14} M_{\odot}$, and by Mantz et al. (2010) who obtained a value of $M_{500} = 18.9 \pm 2.5 \times 10^{14} M_{\odot}$ from the analysis of *Chandra* and *ROSAT* images taken with the position sensitive proportional counters.

6.2 Previous weak-lensing analyses

We briefly summarize some results of previous weak-lensing analyses of Abell 2219. Bardeau et al. (2007) investigated this cluster as part of a sample of eleven X-ray luminous clusters selected from the XBACs catalogue (Ebeling et al. 1996) in the narrow redshift range at $z = 0.21 \pm 0.04$. The weak-lensing analysis was based on wide-field images obtained with the CFHT12k camera at the Canada–France–Hawaii–Telescope (CFHT) in the *B* (5400 s), *R* (6300 s), and *I* (3000 s) bands. The background galaxies were selected by applying the magnitude cut $22.2 < R < 25.3$, followed by the colour selection $R-I \gtrsim -0.7$ to isolate galaxies above the cluster red-sequence. The resulting source number density is 8–10 gal arcmin^{-2} . The mass and the concentration were estimated to be $M_{200} = 2094 \pm 435 \times h_{70}^{-1} 10^{12} M_{\odot}$ and $c_{200} = 3.84 \pm 0.99$.

Hoekstra (2007) performed an independent weak-lensing analysis using the same data of Bardeau et al. (2007). After selecting galaxies in the magnitude range $21 < R < 25$, he used the

photometry in B and R bands to remove galaxies that lie on the cluster red-sequence. The shear profile was corrected for the residual contamination by increasing the observed shear by the factor $1 + f_{\text{cg}}(\theta)$, where $f_{\text{cg}}(r)$ is the fraction of cluster galaxies at radius r . The value of f_{cg} as function of radius was determined by stacking several other clusters. Using a value of $\langle\beta\rangle = 0.54$ inferred from the redshift distribution of the *Hubble Deep Field* (Fernández-Soto, Lanzetta & Yahil 1999), and fitting the tangential shear profile with an NFW model, the resulting virial mass was $11.3^{+3.2}_{-2.7} \times h^{-1} 10^{14} M_{\odot}$. The fit was done assuming one free parameter (the mass) and using the relation between concentration and mass proposed by Bullock et al. (2001).

Okabe et al. (2010) also performed a weak-lensing analysis of Abell 2219 as part of their study of 30 clusters from the Local Cluster Sub-structure Survey sample. They used only *Subaru* data in R and V bands. The exposure times were 24 and 18 min, respectively. The selection of background sources was done combining the magnitude cut $22 \leq R \leq 26$ with colour cuts to identify galaxies redder and bluer than the cluster red-sequence. The optimal selection was found by maximizing the mean amplitude of the tangential shear profile with respect to the colour offset from the cluster red-sequence. The resulting number density of background sources was found to be ~ 10 gal arcmin $^{-2}$. They estimate the virial mass and concentration to be $M_{\text{vir}} = 9.11^{+2.54}_{-2.06} \times h_{72}^{-1} 10^{14} M_{\odot}$; $c_{\text{vir}} = 6.88^{+3.42}_{-2.16}$. However, given the quality of their fit to the tangential shear profile (reduced $\chi^2 = 2.26$ and the value of the significance probability Q used to quantify the goodness of the fit of the model), Okabe et al. (2010) concluded that the NFW model does not describe well the data.

6.3 Image reduction

The observations of the cluster Abell 2219 were performed with the Suprime-Cam mounted at the prime focus of the *Subaru* telescope, an 8.2 metre telescope located at the summit of the Mauna Kea. The camera is a mosaic of 10 2048 \times 4096 CCDs which covers a 34 arcmin \times 27 arcmin field of view with a pixel scale of 0.20 arcsec. We analyse data publicly available in the BVR bands, retrieved from the development of the Subaru-Mitaka-Okayama-Kiso Archive System (SMOKA; Baba et al. 2002) Science Archive.

Images are reduced using the VST-TUBE imaging pipeline (Grado et al. 2012), a software designed to work on optical astronomical images taken with different instruments and instrumental setups. The data reduction consisted in the following steps: overscan correction, flat fielding, correction of distortions due to optics and sky background subtraction. The astrometric solution describing the distortions produced by the optics and telescope was computed for each exposure using the ASTROMC code (Radovich et al. 2004). This was then used in the SWARP tool, which was used to resample and co-add all exposures. The absolute photometric calibration of the images in each filter was achieved using standard Stetson fields. The Sloan digital sky survey i -band image of Abell 2219 was taken with the Large Binocular Camera mounted on the prime focus of the 8.4 m Large Binocular Telescope. LBC has a field of 23 arcmin \times 23 arcmin and provides images with a sampling of 0.225'' pixel $^{-1}$. The data were taken in 2010; the data reduction was done by the LBC Data Centre using a pipeline specifically designed for LBC data. Magnitudes in the $BVRi$ filters are in the AB system. The observation nights of the images in each filter, the total exposure time, the photometric zero-point of the final co-added images and the average FWHM for point-like sources are given in Table 4. Finally, galaxy magnitudes are corrected for galactic extinction using the maps produced by Schlegel et al. (1998).

Table 4. Abell 2219 observations summary.

Date	Band	Exp. time (s)	zp	FWHM (arcsec)
25/06/06	B	720	27.54	0.90
18/07/04	V	1080	27.56	0.68
19/07/04, 14/08/07	R	3330	27.59	1.09
07/05/10	SDSS i	3005.8	27.57	0.89

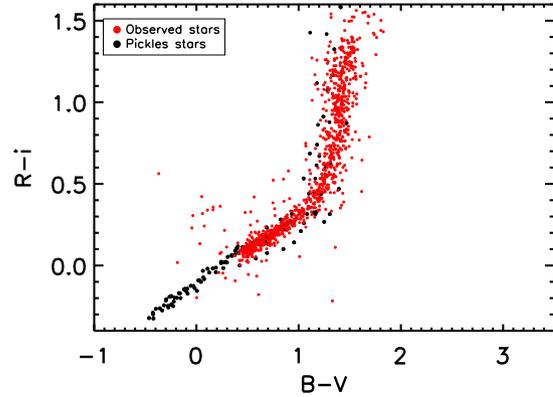


Figure 5. Star observed colours (red dots) and colours obtained by the convolution of the stellar observed spectra from the Pickles' library with the filter transmission curves (black dots), after the offsets given in the text were applied.

We use the deep R -band image to perform the weak-lensing analysis, while photometric catalogues in the B and V bands are derived using SExtractor in dual-mode, with the R -band image used as the detection image. The masking of reflection haloes and diffraction spikes near bright stars is performed with the EXAM code described in Huang et al. (2011). We apply to the B and V magnitudes the offsets derived within the COSMOS survey by Capak et al. (2007): $+0.19$ (B), $+0.04$ (V); according to Huang et al. (2011), no offset is required for the R band. We convolve the stellar spectra from the Pickles' library (Pickles 1985) by the transmission curves. Comparing the colours obtained with those derived for the stars in our catalogue, we derive an offset in i magnitudes of $+0.065$. Fig. 5 shows this comparison after adding the offsets. We verified that the addition of these offsets improves the match of stellar colours in our data with colours obtained from the convolution of the stellar spectra from the Pickles' library by the filter transmission curves.

6.4 Weak-lensing analysis

6.4.1 Shape measurement

The observed shear signal produced by the gravitational field of the cluster Abell 2219 is reconstructed using our new KSB implementation described in Section 3. The selection of stars used to correct the distortions introduced by the PSF components is made in the magnitude MAG AUTO versus $\delta = \text{MU MAX-MAG AUTO}$ space.

In Fig. 6, the star locus is defined by objects populating the vertical branch (green dots). Objects with δ lower than stars are classified as spurious detections (black dots) while galaxies are marked with blue dots. Saturated stars are marked with red dots. We use stars with R magnitude in the range [19, 22] to correct for the effects introduced by PSF anisotropy and seeing.

In Fig. 7, the first three panels show stars measured, fitted and residuals ellipticities, while the last one shows stars

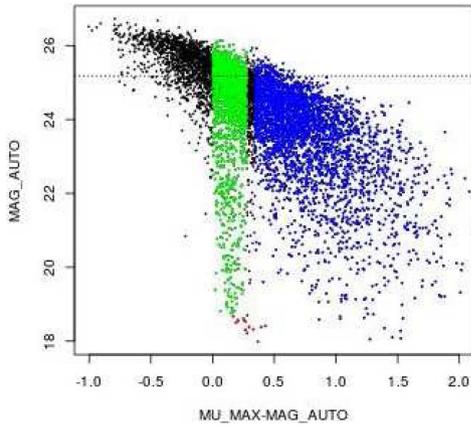


Figure 6. MAG AUTO versus $\delta = \text{MU_MAX} - \text{MAG_AUTO}$ plot. Stars are plotted with green dots, saturated stars with red dots and galaxies with blue dots. Sources with size smaller than stars and sources with ambiguous classification are plotted with black dots.

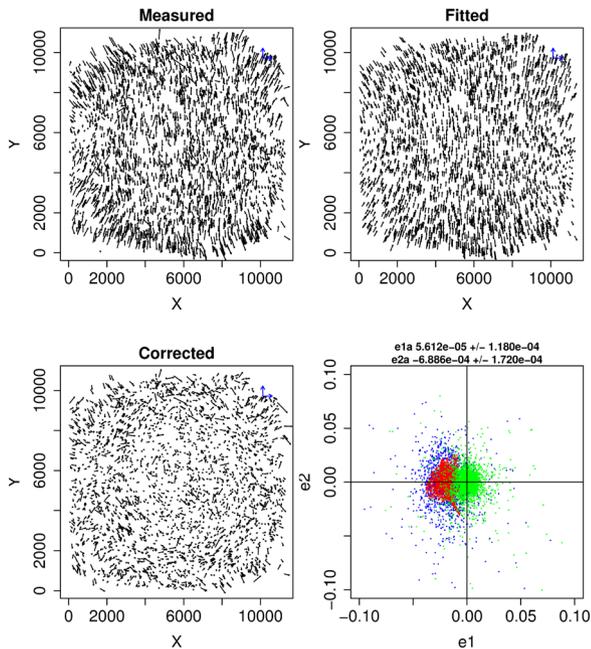


Figure 7. PSF anisotropy correction: the first three panels show the measured, fitted, and residuals ellipticities as function of the position (X and Y are in pixels). In the last panel, blue dots are the measured values, the red dots the modelled values and the green dots the values after the correction.

measured ellipticities (marked with blue dots), stars modelled ellipticities (marked with red dots) and their values after the correction (green dots). We obtain $(e_{\text{aniso},1}) = (0.6 \pm 1.2) \times 10^{-4}$, $(e_{\text{aniso},2}) = (-6.9 \pm 1.7) \times 10^{-4}$.

The tangential and the cross components of the reduced shear, g_T and g_\times , are obtained from the quantities $e_{1\text{iso}}$ and $e_{2\text{iso}}$ used, respectively, as g_1 and g_2 in equation (6) and (7). The radial profiles are obtained averaging both shear components in annular regions centred on the BCG.

We exclude from the catalog galaxies with $\text{SNe} < 5$ (our choice of $\text{SNe} < 5$ corresponds to about $\text{S/N} < 10$), for which the ellipticity measurement is not meaningful, and galaxies with corrected ellipticities $e_{1\text{iso}}$ or $e_{2\text{iso}} > 1$. The final catalogue which will be used for the analysis has a number density of $\sim 20 \text{ gal arcmin}^{-2}$.

6.4.2 Selection of the background galaxies

In this section, we describe the selection of the background galaxies obtained with our method based on colours (Sections 4 and 5.1). For comparison, we also use the alternative methods introduced earlier and based on photometric redshifts, magnitude cuts, and identification of the cluster red-sequence.

6.4.2.1. Selection in colour space

We derive the $V-R$ versus $B-R$ and the $R-i$ versus $B-V$ CC diagrams of the galaxies in the field of Abell 2219. Applying the selection described in the Section 5.1 for the $V-R$ and $B-R$ colours, we obtain a background population with a number density of $\sim 13 \text{ gal arcmin}^{-2}$. The selection in $V-R$ versus $B-R$ space is represented in the upper-left panel of the Fig. 8. In all panels of Fig. 8, galaxy sub-samples are shown using the same colour-coding of Fig. 2.

The selection based on the colours $R-i$ and $B-V$ according to the criteria in Section 5.1 is represented in Fig. 8 (bottom-left panel). The number density of background galaxies is $\sim 14 \text{ gal arcmin}^{-2}$.

The radial number density profiles of cluster members clearly show overdensities at small radii as expected for cluster member galaxies. The radial number density profiles of background and foreground galaxies do not show clustering at small radii which rises from cluster member galaxies, while they appear almost flat since they are equally distributed on the sky except for a decrease at small radii due to area loss by bright sources and the magnification bias (Gray et al. 2000; Broadhurst et al. 2008) clearly evident in the background samples and with a lower amplitude in the foreground samples indicating the presence of background sources in foreground-selected samples. We check that the decrease at small radii is effectively due to the magnification introduced by lensing measuring its signal in the case of background galaxies selected in the $V-R$ versus $B-R$ space:

$$n_\mu(\theta) = n_0 \mu(\theta)^{2.5s-1}; \quad (17)$$

with $n_0 = dN_0(< m_{\text{cut}})/d\Omega$ the unlensed mean number density of background sources for a given magnitude cutoff, described with a power law with slope, $s = d \log_{10} N_0(< m)/dm > 0$.

The normalization and slope are estimated in the outer region i.e. for $\theta > 10 \text{ arcmin}$. We obtain the following values for $s = 0.02 \pm 0.001$ and $n_0 = 16.8 \pm 0.3 \text{ galaxies arcmin}^{-2}$ that we use to derive the expected magnification for a lens described by an NFW density profile as shown in Fig. 9 (red dot line). For comparison, we derived the expected magnification applied to the sample selected with the magnitude cut $22 \leq R \leq 26$ (back points in Fig. 9). In this case, the rise in the number counts at small radii is a clear evidence of the presence of cluster members in the sample. The model curve (black dot line) is overplotted to data.

The g_T profiles of the cluster and foreground galaxies show a positive value in some bins. This is similar to what found in simulations, where we noticed that the sub-samples of cluster and foreground galaxies are contaminated by background sources. The g_T profiles of background sources nicely rise at small radii. The g_\times profiles for each identified population are consistent with zero, indicating the lack of systematic errors arising from an imperfect PSF correction.

The consistency of the real and simulated data is further verified showing in the Fig. 10 the magnitude distributions of the background sources selected in $V-R$ versus $B-R$ diagram in the two cases; and in the Fig. 11, their size distribution plotting the FLUX RADIUS derived from the R images.

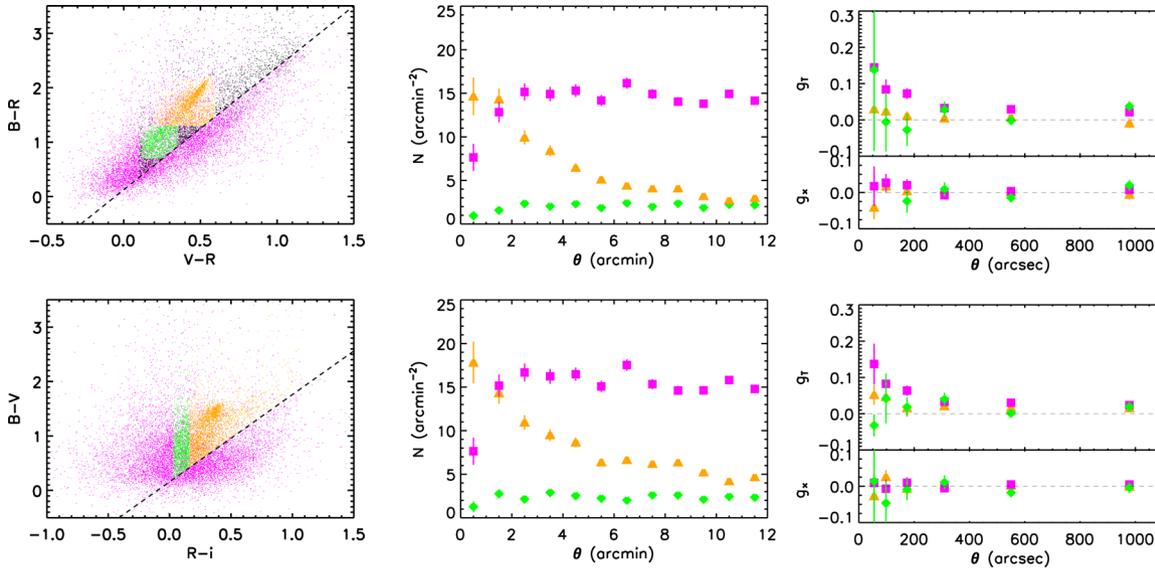


Figure 8. Sample selections in CC space. Left-hand panels: CC diagrams for Abell 2219 derived choosing two combinations of photometric bands. The samples corresponding to cluster member galaxies are displayed with orange dots and the foreground and background samples with green and magenta dots, respectively. Unclassified sources are plotted with black dots. Centre: radial number density profiles of galaxies. Background density profiles (magenta squares) show a decrease in the central region (magnification bias and area loss by bright sources). The cluster density profiles (orange triangles) rise towards the centre indicating a population of cluster members. The foreground density profiles are almost flat and are represented with green diamonds. Right-hand panels: g_r and g_x radial profiles of the corresponding samples.

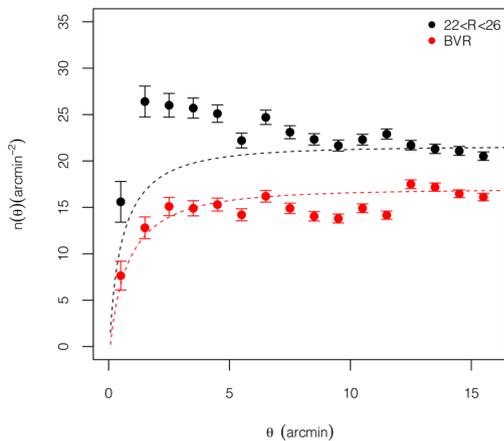


Figure 9. Number counts profile of the background objects selected in $V-R$ versus $B-R$ diagram of Abell 2219 (red points) and with a magnitude cut (black points). Model curves are overplotted for comparison.

6.4.2.2 Alternative selection methods

We compare results obtained with our selection method to those obtained with other techniques to identify lensed galaxies: the identification based on a magnitude cut, that was based on the identification of the cluster red-sequence and finally the use of photometric redshifts.

The selection based on magnitude cuts consists in considering as background galaxies those having magnitude in the range $22 \leq R \leq 26$.

The selection of background galaxies on the base of the colours respect to the cluster red-sequence is done as described in Section 5.3 (Okabe’s method).

In the following, we reported our analysis with photometric redshifts. We use the photometry in four bands to derive photometric redshifts using the ZEBRA code in maximum-likelihood mode as described in Section 5.3. As done analysing the simulations, we do not

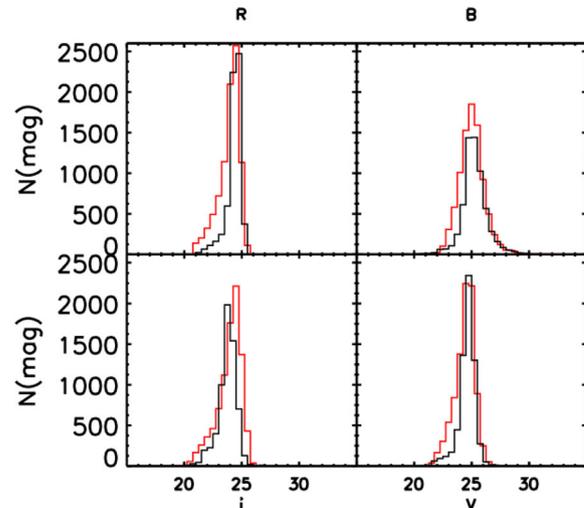


Figure 10. Magnitude distributions of objects selected in $V-R$ versus $B-R$ diagram of the simulated images (black lines) and of the images of Abell 2219 (red lines).

use galaxies with $\text{photo-}z_s > 2$ and $\sigma_z/(1+z) > 0.1$. For the galaxies with spectroscopic redshift given in Boschin et al. (2004), we find $\langle z_{\text{phot}} - z_{\text{spec}} \rangle = -0.017 \pm 0.003$. The sample given in Boschin et al. (2004) is characterized by galaxies with $z \leq 0.4$ so it is not possible to verify the accuracy of sources which have higher redshift value. The accuracy of the photometric redshifts derived for $z > 0.4$ is checked from the comparison with the Sloan Digital Sky Survey photometric redshifts of galaxies in the Data Release 9, although the latter are obtained from shallower images. We find $\langle z_{\text{phot}} - z_{\text{sloan}} \rangle = -0.066 \pm 0.007$. We select as background galaxies those having $z_{\text{phot}} > 0.27$ (Fig. 12, both panels, magenta squares), as foreground galaxies those having $z_{\text{phot}} < 0.19$ (Fig. 12, both panels, green diamonds) and as cluster galaxies those having 0.19

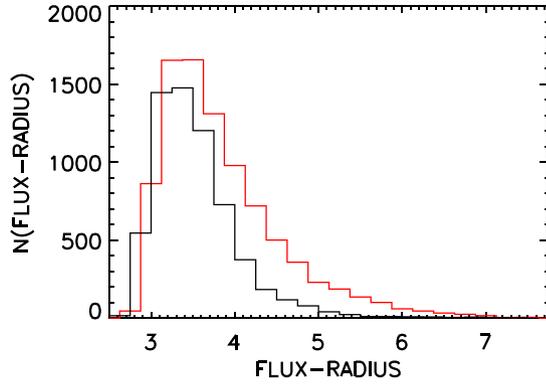


Figure 11. Size distribution for galaxies selected in $V-R$ versus $B-R$ diagram in the simulations (black line) compared to that derived for galaxies in the images of Abell 2219 (red line).

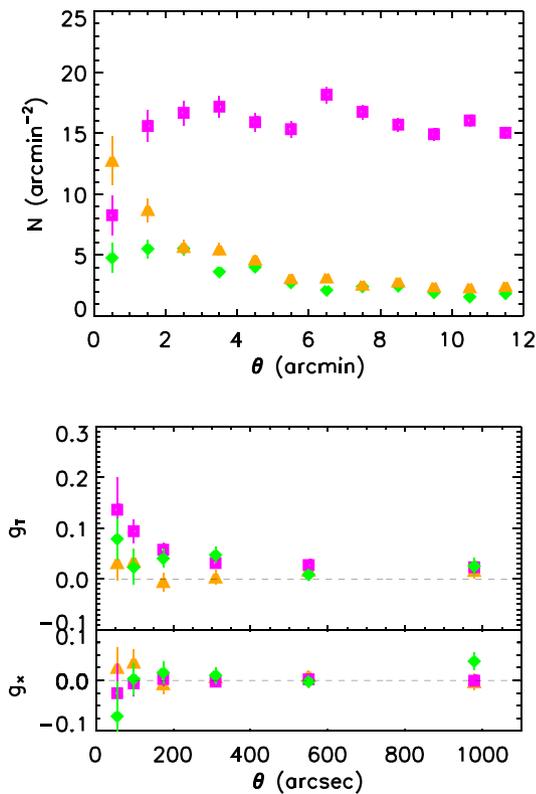


Figure 12. Selection with photometric redshifts. Background galaxies having $z_{\text{phot}} > 0.27$ are marked with magenta squares, cluster galaxies having $0.19 \leq z_{\text{phot}} \leq 0.27$ with orange triangles and foreground galaxies having $z_{\text{phot}} < 0.19$ with green diamonds. The rise in the radial number density profile of foreground galaxies is due to the low accuracy of photometric redshifts derived caused by a limited spectral coverage.

$\leq z_{\text{phot}} \leq 0.27$ (Fig. 12, both panels, orange triangles). The number density of background galaxies selected in this way is ~ 14 gal arcmin $^{-2}$. The radial number density profile of foreground galaxies shows an overdensity at small radii indicating the presence of cluster galaxies in the sample (Fig. 12, upper panel). This result reflects the low accuracy of the photometric redshifts obtained, due to the limited number of bands available and in particular to the lack of photometry in the u passband.

For comparison, we compute the fraction of background sources common to both the BVR and the photo- z selections. We find that

the two samples overlap at the level of 68 per cent. In the case of the $BVRi$ and to the photo- z selections, we find that the 72 per cent of the sources are in common.

6.4.3 Mass estimate

Mass estimates of Abell 2219 are derived by fitting the reduced tangential shear profile with an NFW lens model, as described in the Section 5.4. In this case, we perform the fit in two different ways in order to compare our results with those presented in literature. In the first case, we use c_{vir} and M_{vir} as free parameters. In the second, we adopt the mass–concentration relation proposed by Bullock et al. (2001) to get rid of one of the two parameters. We denote these two methods to estimate the mass with the acronyms *nfw* and *mnfw*, respectively.

In Table 5, we summarize the results of our analysis. The first column describes the selection method applied. The second column reports the number densities of background sources selected in each case. The recovered masses and concentrations are given in columns 4–8. For comparison, we add the results in the case of a magnitude cut selection and in the case of colour–magnitude selection as described in Okabe et al. (2010). We perform all the fits in a range of $30 \text{ arcsec} < \theta < 1000 \text{ arcsec}$ since the lower limit for the radius ($\theta = 30 \text{ arcsec}$) is chosen sufficiently large so that the weak-lensing limit $g_{\text{T}} \approx \gamma_{\text{T}}$ holds. The shear profiles obtained from the selection in CC spaces and with photo- z s and the corresponding fitted models are displayed in Fig. 13. In the following, we compare our estimates of the mass of Abell 2219 with previous analyses found in literature. The weak-lensing analysis made by Bardeau et al. (2007) gives a value of

$$M_{200} = 2094 \pm 435 \times h_{70}^{-1} 10^{12} M_{\odot} \text{ and } c_{200} = 3.84 \pm 0.99.$$

The comparison of our results with those found in the paper of Bardeau et al. (2007) is done performing the NFW fit with the fixed value of $c_{\text{vir}} = 4.83$, corresponding to $c_{200} = 3.84$. The values of M_{200} found in our work are smaller by a factor of 2.4–3 with respect to that found in Bardeau et al. (2007).

The weak-lensing analysis of Hoekstra (2007) gives a virial mass of $11.3^{+3.2}_{-2.7} \times h^{-1} 10^{14} M_{\odot}$ derived using the relation of Bullock et al. (2001), in agreement with our results.

Okabe et al. (2010) derived the weak-lensing mass on the base of *Subaru* data in only two photometric bands (R and V). They found a value of $M_{\text{vir}} = 9.11^{+2.54}_{-2.06} \times h_{72}^{-1} 10^{14} M_{\odot}$ and a concentration of $c_{\text{vir}} = 6.88^{+3.42}_{-2.16}$.

In order to get a comparison with the result obtained by Okabe et al. (2010), we performed all the fits using the same values of the cosmological parameters of their paper and fixing c_{vir} to the value found by the authors. The value of the virial mass found by Okabe et al. (2010) is in agreement with our results except for that derived selecting in the $R-i$ versus $B-V$ space. However, the authors concluded that the fit of the NFW model to data does not give an acceptable result on the base of the significance probability Q used to quantify the goodness of the fit of the model, unlike the results of this work. This can be due to the different selections of background sources in both works. This is further supported by results on simulations in which the Okabe’s method is applied. In this case, the fitted mass is lower than those derived by the selection based on colours (Table 3). Okabe et al. (2013) presented a new analysis with a significantly revision of their colour-cut technique used to select background sources. The paper shows the results of the stacked analysis of 50 galaxy clusters but the individual cluster masses are not reported. They claim that the new masses are

Table 5. Summary of the fit parameters derived for the cluster A2219. Column 1: selection criterium; Column 2: number density of the sample; Column 3: fitting method; Column 4: M_{vir} ; Column 5: c_{vir} ; Column 6: M_{200} ; Column 7: c_{200} ; Column 8: M_{500} ; Column 9: $\langle\beta\rangle$. For each selection in the first row, we give the results obtained using the relation found by Bullock et al. (2001) (labelled mnfw) while in the second row (labelled nfw) we keep both M_{vir} and c_{vir} as free parameters.

Selection	Density (arcmin^{-2})	Method	M_{vir} ($10^{15} M_{\odot}$)	c_{vir}	M_{200} ($10^{15} M_{\odot}$)	c_{200}	M_{500} ($10^{15} M_{\odot}$)	$\langle\beta\rangle$
$V-R$ versus $B-R$	13	mnfw	$1.48^{+0.14}_{-0.12}$	4.27	$1.23^{+0.11}_{-0.09}$	3.39	$0.82^{+0.07}_{-0.06}$	0.737
		nfw	$1.65^{+0.21}_{-0.19}$	$3.16^{+0.47}_{-0.40}$	$1.32^{+0.16}_{-0.13}$	$2.47^{+0.39}_{-0.32}$	$0.81^{+0.07}_{-0.07}$	
$R-i$ versus $B-V$	14	mnfw	$1.46^{+0.12}_{-0.11}$	4.30	$1.21^{+0.10}_{-0.09}$	3.40	$0.81^{+0.06}_{-0.06}$	0.742
		nfw	$1.93^{+0.32}_{-0.24}$	$2.48^{+0.41}_{-0.40}$	$1.50^{+0.21}_{-0.17}$	$1.92^{+0.33}_{-0.32}$	$0.85^{+0.07}_{-0.07}$	
z_{phot}	14	mnfw	$1.91^{+0.16}_{-0.15}$	4.15	$1.58^{+0.14}_{-0.12}$	3.29	$1.04^{+0.09}_{-0.08}$	0.622
		nfw	$1.93^{+0.28}_{-0.26}$	$3.50^{+0.60}_{-0.54}$	$1.56^{+0.19}_{-0.19}$	$2.75^{+0.49}_{-0.44}$	$0.98^{+0.09}_{-0.09}$	
$22 \leq R \leq 26$	18	mnfw	$1.10^{+0.08}_{-0.08}$	4.44	$0.91^{+0.07}_{-0.07}$	3.52	$0.61^{+0.05}_{-0.05}$	0.720
		nfw	$1.71^{+0.24}_{-0.20}$	$2.14^{+0.36}_{-0.32}$	$1.29^{+0.17}_{-0.14}$	$1.64^{+0.31}_{-0.29}$	$0.71^{+0.08}_{-0.07}$	
VR	10	mnfw	$1.33^{+0.13}_{-0.11}$	4.33	$1.11^{+0.10}_{-0.10}$	3.43	$0.74^{+0.06}_{-0.07}$	0.735
		nfw	$1.48^{+0.27}_{-0.22}$	$3.30^{+0.68}_{-0.63}$	$1.20^{+0.17}_{-0.17}$	$2.58^{+0.56}_{-0.51}$	$0.73^{+0.08}_{-0.08}$	

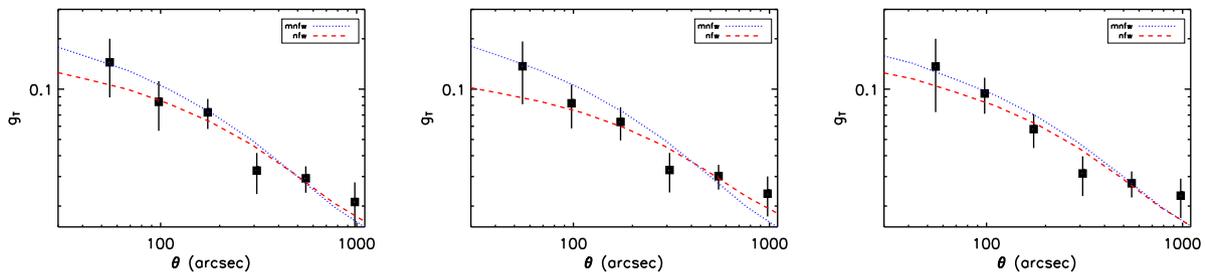


Figure 13. Shear profiles obtained with the different selection methods. Left-hand panel: selection in $V-R$ versus $B-R$ space. Central panel: selection in $R-i$ versus $B-V$ space. Right-hand panel: selection with photo- z s. The two models nfw and mnfw are overlapped.

higher by 14 per cent–20 per cent compared to the previous analysis, suggesting that such difference could arise from systematics in the shape measurement methods and contamination by unlensed sources.

7 CONCLUSION

Dilution of the lensing signal due to unlensed sources can crucially affect the estimation of the cluster mass. We developed a new selection method for the background lensed sources based on the simultaneous analysis of the shear signal and of the colours of the galaxies, with the photometry from the COSMOS used as training set. Our selection method can be applied to the weak-lensing analysis of clusters, under different cluster redshifts, masses and filter combinations. Its performance in the selection of the background sources and its effect on the dilution bias are tested using simulations. To this end, we produced realistic *Subaru* and LBT images containing lensing signal with the SKYLENS code under different cluster redshifts ($z = 0.23-0.45$), masses ($M_{\text{vir}} = 0.5-1.55 \times 10^{15} M_{\odot}$) and filter combinations. We showed with simulations that three photometric bands are enough to obtain a safe sample of background sources for the weak-lensing analysis. The values of M_{vir} obtained from the fit of the expected NFW profile to the shear for the simulated clusters support a low contamination by unlensed sources in the samples selected by cuts in the CC space.

We find that the fitted values of M_{vir} obtained with our selection agree with the true masses within $\lesssim 10-15$ per cent. On average, the ratio between estimated and true masses is $\sim 0.98 \pm 0.09$.

The simulations also allowed us to estimate the residual fraction of unlensed galaxies in the samples of background galaxies, and the fraction of background galaxies mis-identified as foreground and cluster members. Colour selections and photometric redshifts give less than the 11 per cent of unlensed sources contaminating the background-selected samples, even if the photo- z s have a low accuracy due to the limited spectral coverage. A fraction of background lensed sources is lost since they are misidentified as foreground and cluster galaxies in a percentage varying for each kind of selection done. However, the low accuracy of the derived photo- z s could lead to an additional bias in the cluster mass reconstruction due to the wrong identification of the value of $\langle\beta\rangle$ when estimated from the photometric redshifts themselves. Its value is smaller than the true value resulting in an overestimate of the mass of the lens by ~ 11 per cent with respect to the input value.

Selections based on a simple magnitude cut or on the analysis of the colour–magnitude diagram following the technique described in Okabe et al. (2010) lead to a lower mass due to a residual contamination by unlensed sources.

We further apply the developed selection method to Abell 2219. In our work, we used a new KSB implementation based on the PSFEX code (Radovich et al. 2015) to measure the observed shear signal produced by the gravitational field of the cluster.

As shown by simulations the high purity of the selected samples obtained allows us to derive fiducial samples of background sources on which an accurate cluster mass reconstruction can be performed.

The virial mass of Abell 2219 was obtained for each kind of selection by fitting the NFW model to data. Our results are in

agreement with the previous analysis of Hoekstra (2007) and Okabe et al. (2010) (it makes exception the selection in the $R-i$ versus $B-V$ space). However differently from our results, Okabe et al. (2010) found that the NFW shear model does not fit the observed shear. This can arise from different selection criteria and source redshifts identification used in the analyses, which can origin the other different results found in literature.

To conclude, in this paper, we showed how using colours from three photometric bands can improve the selection of background galaxies for the weak-lensing analysis of galaxy clusters: the selection is driven both by an external photometric catalogue with known redshifts (the COSMOS catalogue), and the maximization of the amplitude of the shear in the observed cluster. The procedure proposed in this paper can be easily automated and applied to large data sets.

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APPENDIX A: BIAS IN THE ESTIMATED ELLIPTICITIES

The accuracy of the weak-lensing analysis crucially depends on the performance of the method used to correct for the distortions introduced by the PSF components. Tests on mock images were

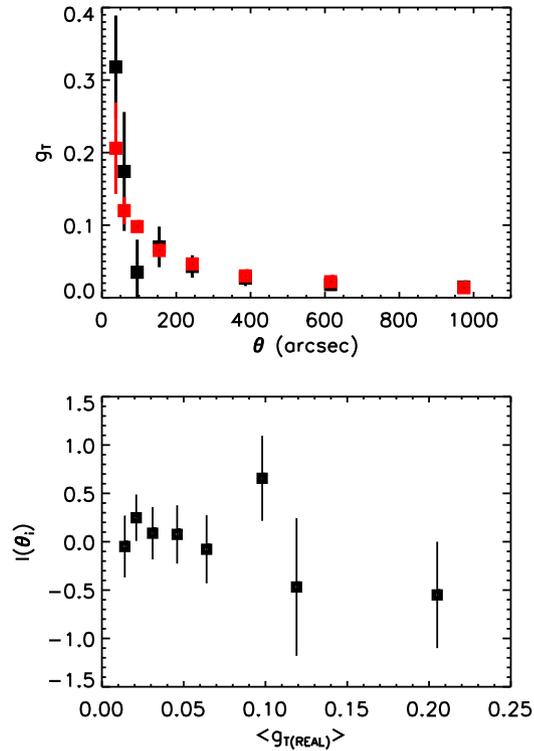


Figure A1. Top panel: g_T radial profiles for the simulated lensed galaxies obtained using the measured KSB ellipticities (black squares) and their input values (red squares). Background galaxies are selected on the base of their input value of the redshift in the simulation. Bottom panel: $l(\theta_i)$ as function of the amplitude of the mean true reduced shear value $\langle g_{T(\text{REAL})} \rangle$.

launched in the past to exploit this issue (Erben et al. 2001; Heymans et al. 2006; Massey et al. 2007; Bridle et al. 2010; Kitching et al. 2012). Ellipticities derived by different methods for the shape measurement are affected by a multiplicative (m) and an additive (c) bias defined as: $e_{\text{obs}} - e_{\text{true}} = me_{\text{true}} + c$ (Heymans et al. 2006). We try to quantify this multiplicative bias measuring the shear from the simulated images and comparing it to the true input value. To this end, we use the set of simulations described in the Section 2 tailored on Abell 2219 observations in the $BVRi$ bands with exposure times as listed in Table 4. Simulated images are characterized by a spatially constant Gaussian PSF with FWHM 0.6 arcsec. We derive radial profiles of the tangential component of the reduced shear, obtained using the measured ellipticities (Fig. A1 top panel, black

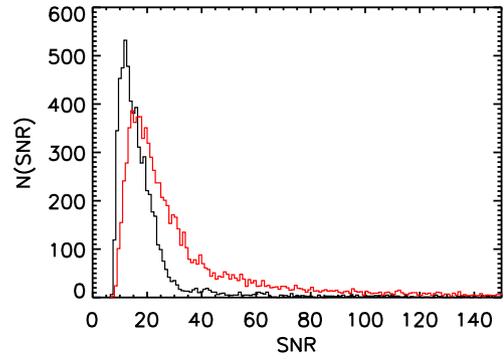


Figure A2. Distribution of the signal-to-noise ratio of the background galaxies selected in the $V-R$ versus $B-R$ CC diagram in the simulations (black line) and in the images of Abell 2219 (red line).

squares) and their input values (Fig. A1 top panel, red squares) for background galaxies selected on the base of the input value of the redshift in the simulations. In each radial bin θ_i , we compute the difference between the mean tangential reduced shear obtained using the measured KSB ellipticities ($g_{T(\text{KSB})}$) and that derived using their true input values ($g_{T(\text{REAL})}$):

$$l(\theta_i) = \frac{\langle g_{T(\text{REAL})} \rangle - \langle g_{T(\text{KSB})} \rangle}{\langle g_{T(\text{REAL})} \rangle}. \quad (\text{A1})$$

We derive the mean value of l avoiding the three inner bins characterized by a low statistic. In Fig. A1 (bottom panel), we show how $l(\theta_i)$ changes with the amplitude of the mean true reduced shear value $\langle g_{T(\text{REAL})} \rangle$. We find that, for these simulations, our KSB pipeline on average underestimates the measured shear of about 5 per cent with respect to the input values. Since the background sources derived in the analyses of A2219 and of the simulated data are both dominated by an $\text{SNR} \leq 40$ (Fig. A2), we apply the correction factor of 1.05 to the measured ellipticities and use the so-corrected ellipticities to derive the cluster mass. A dependence of the amplitude of the multiplicative shear calibration factor on the amplitude of the shear signal could arise in intermediate lensing regimes. We will further investigate this aspect in a future work.

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