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Cosmic Ray Oriented Performance Studies for the JEM-EUSO First Level Trigger

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Abstract

JEM-EUSO is a space mission designed to investigate Ultra-High Energy Cosmic Rays and Neutrinos ($E > 5 \cdot 10^{19}$ eV) from the International Space Station (ISS). Looking down from above its wide angle telescope is able to observe their air showers and collect such data from a very wide area. Highly specific trigger algorithms are needed to drastically reduce the data load in the presence of both atmospheric and human activity related background light, yet retain the rare cosmic ray events recorded in the telescope. We report the performance in offline testing of the first level trigger algorithm on data from JEM-EUSO prototypes and laboratory measurements observing different light sources: data taken during a high altitude balloon flight over Canada, laser pulses observed from the ground traversing the real atmosphere, and model landscapes reproducing realistic aspect ratios and light conditions as would be seen from the ISS itself. The first level trigger logic successfully kept the trigger rate within the permissible bounds when challenged with artificially produced as well as naturally encountered night sky background fluctuations and while retaining events with general air-shower characteristics.

Key words: JEM-EUSO, trigger system, FPGA, nightglow background

1 1. Introduction

Ultra-High Energy Cosmic Rays (UHECR) are observed as Extensive Air 2 Showers (EAS) in the atmosphere surrounding Earth. They are rare events, 3 and the higher their energy, the rarer they get. The greatest mystery sur-4 rounding them is their origin, but also their nature remains contentious. High 5 statistics and high quality data are needed to make progress on both fronts, 6 which means scanning the largest possible volume of atmosphere for EAS. The current ground based experiments run up against natural boundaries 8 limiting their expansion and present difficulties when comparing data ob-9 tained in northern and southern latitudes, as the fraction of common sky is 10 limited. Therefore, space based instruments observing the atmosphere from 11 above with full-sky coverage have long been considered the logical next step 12 in the evolution of UHECR experiments [1]. 13

The International Space Station (ISS) with its existing infrastructure and 14 support systems is a natural first step on this way into space, and JEM-EUSO 15 [2] is a scientific mission under development with the aim of identifying the 16 astrophysical origin and nature of UHECRs from the ISS. JEM-EUSO detects 17 UHECR induced EAS by looking down onto the earth atmosphere. It has a 18 telescope with a large $(\pm 30^{\circ})$ Field of View (FoV) imaging the atmosphere 19 below the ISS onto an array of UV sensitive Multi-Anode Photomultiplier 20 Tubes (MAPMTs) [3]. The MAPMTs (Hamamatsu Photonics R11265-03-21 M64) have 8×8 pixels and for readout purposes 2×2 MAPMTs are grouped 22 into one Elementary Cell (EC). The First Level Trigger (FLT), which is the 23 subject of this article, works at the level of these ECs. Nine ECs form one 24 Photo-Detector Module (PDM), which is the basic unit for the Second Level 25 Trigger (SLT). The Focal Surface (FS) is organised in 137 PDMs. Together 26 these PDMs cover the FS of the telescope with $\sim 3.2 \cdot 10^5$ MAPMT pixels. 27 A detailed description of the electronics and data acquisition for JEM-EUSO 28 can be found in [4], while a sketch of the structure of the FS is shown in 29

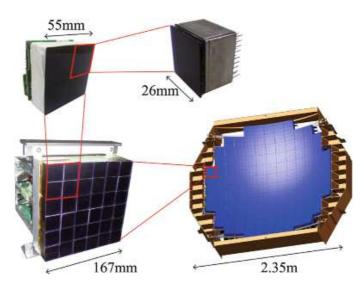


Figure 1: Structure of the Focal Surface. The 2.5 m surface is divided in 137 PDM modules. Each PDM is filled with 9 ECs, with 4 MAPMTs each. The bottom left corner shows the prototype of the mechanical structure with 36 MSPMTs installed. Figure taken from [4].

30 Fig. 1.

The observational concept of JEM-EUSO [5] is based on recording both 31 the fluorescence light emitted during the evolution of EAS as well as the 32 reflected Cerenkov light if the EAS' Cerenkov cone hits a reflective surface 33 as it reaches the ground. EAS from the interaction of UHECRs or neutrinos 34 in the atmosphere will - for 10^{20} eV EAS - typically result in a few thou-35 sand photons detected by the JEM-EUSO detector within a few hundred 36 microseconds. Owing to the large FoV the expected rate of such ultra high 37 energy EAS are approximately one per day. Depending on both the energy 38 and the zenith angle of the EAS, images may be contained inside a single 39 EC or may cross a few PDMs as they are imaged onto across the FS. EAS 40 develop within the lowest 15 km of the Earth atmosphere, so that their dis-41 tance to the ISS, which orbits earth at a height above ground of about 400 42 km, can be considered unchanging whatever an individual EAS' zenith angle 43 happens to be. With that EAS' angular speed across the FS to first order 44 only depends on the EAS' propagation direction relative to the respective FS 45 pixels' direction of view. As EAS traverse the atmosphere at essentially the 46 speed of light, and from the height of the ISS a single square MAPMT pixel's 47 FoV's diagonal measures roughly 750 m on the ground, it takes about 2.5 μ s

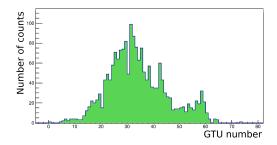


Figure 2: Light curve (number of photoelectron counts) integrated over the shower profile and plotted over time along the shower axis for a simulated UHECR EAS of 2×10^{20} eV (no background is added). Time is measured in GTUs (see Section 2).

for horizontal EAS' image to traverse the diagonal of a MAPMT pixel. As 40 bandwidth for data transmission from the ISS back to Earth is limited, 2.5 50 μ s, the so-called Gate Time Unit (GTU), was adopted as the basic unit for 51 digitization at JEM-EUSO. Given the distance between EAS and ISS, JEM-52 EUSO must be able to detect single photons. The front-end electronics works 53 in single photon-counting mode, which means that HV and electronics gain 54 are adjusted such that after digitization one digital increment corresponds to 55 one photoelectron (PE) count released from the MAPMT's photocathode. 56

In this paper we discuss the FLT algorithm specific to the identification 57 of UHECR and neutrino induced EAS. Fig. 2 shows the temporal evolution 58 in units of GTU of the MAPMT signal for typical simulated proton EAS of 59 energy $2 \cdot 10^{20}$ eV, viewed by the JEM-EUSO telescope under an angle of 60 60°. EAS simulations are performed using the ESAF [6] package adapted 61 to the JEM-EUSO instrument. In Fig. 3, the top panel shows the spacial 62 distribution of simulated EAS scintillation light emission projected back onto 63 the Earth's surface for EAS with a common energy of $E = 10^{20}$ eV traversing 64 the atmosphere under zenith angles of (a) $\theta = 30^{\circ}$, (b) $\theta = 60^{\circ}$ and (c) 65 $\theta = 75^{\circ}$. The inset in the lower left corner of this panel puts these showers 66 into the context of the FoV of the whole FS of JEM-EUSO. The bottom panel 67 presents the image of the EAS in (b) as it would be seen by the JEM-EUSO 68 telescope: the optics inverts the direction of motion, and the photon counts 69 per pixel are integrated over the EAS duration which is of the order of 100 70 $\mu s.$ 71

Looking down from the ISS the FLT has to identify these events in the
 presence of various backgrounds: UV albedo, transient atmospheric phenom ena, and artificial light sources in cities, along transportation networks, and

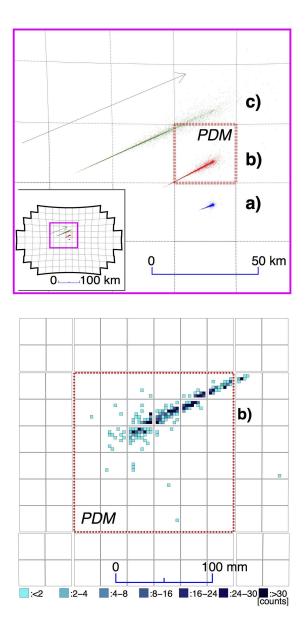


Figure 3: The top panel shows the 3 EAS with the same energy of $E = 10^{20}$ eV but impinging on the atmosphere under different zenith angles of (a) $\theta = 30^{\circ}$, (b) $\theta = 60^{\circ}$ and (c) $\theta = 75^{\circ}$. The most inclined EAS last ~250 μs . The arrow indicates the direction of the EAS transit on the FS. The inset on the bottom left and the grey grid shows how the FoV is imaged on the telescope's FS. The distance scale on this panel refers to the distance the shower develops over as projected onto the Earth surface. The bottom panel shows the image (inverted by the optics) on shower (b) as recorded (integrated over time) by the JEM-EUSO telescope. The distance scale here refers to distance on the FS. The regions enclosed by thick dashed lines in both panels refer to the same PDM. Image taken from [7]. UV background is not added in these plots.

on ships and airplanes. The ISS is moving at about 7.6 km/s so that sta-75 tionary light sources on the ground stay within the FoV of a single pixel for 76 about 70 ms. Such anthropogenic lights, as for example cities, are in the FoV 77 on average during only $\sim 10\%$ of an ISS orbit [7]. Transient Luminous Events 78 (TLEs) within the atmosphere, like electric discharges (Elves, Sprites, Blue 79 Jets and lightning) as well as meteors will have their own triggering schemes 80 to support a separate science program with JEM-EUSO and will be sup-81 pressed by the UHECR and neutrino oriented FLT on the basis of time and 82 light intensity structures. The expected rate of TLEs is $\sim 700/\text{day}$ [8]. The 83 greatest uncertainty is associated with the very slowly varying backgrounds 84 associated with the albedo of the atmosphere: its scattering and reflection 85 of starlight - light reflected from the moon and planets - and airglow. On a 86 clear night the resulting diffuse flux of such photons is expected to amount to 87 between 300 and 1000 photons $m^{-2} ns^{-1} sr^{-1}$, an expectation derived from 88 existing measurements [9, 10]. 89

In this paper we show the results of the offline testing of the current JEM-EUSO FLT concept using data from three main sources rather than simulation: data collected during a high altitude balloon flight over Canada (EUSO-Balloon), laser shots recorded in coincidence with a ground based JEM-EUSO prototype (EUSO-TA), and data taken in a laboratory setup where realistic background scenarios can be explored (EUSO@TurLab).

⁹⁶ 2. Technical requirements

As mentioned before, working on the ISS imposes severe bandwidth con-97 straints on data transfer to the ground. On top of that there is a $\sim 1 \text{ kW}$ 98 limit on power consumption for the whole telescope, including the readout 99 and trigger electronics, high voltage for the MAPMTs, and monitoring. This 100 constraint for example means that triggering cannot be substituted for by 101 massive only computing. Data rate considerations also played into choosing 102 a GTU of 2.5 μ s for digitization in time and the number of 128 GTUs per 103 packet, meaning that each event record will contain the timing evolution of 104 a signal over a time span of 320 μ s. 105

Since the ISS is far above the parts of the atmosphere where EAS develop, the MAPMTs have to be able to detect faint signals using photon-counting. Therefore, 8 bit full scale per pixel is sufficiently large. Under these conditions the total data rate from the telescope before the FLT would be of order $3.2 \cdot 10^5$ pixel/FS × $4 \cdot 10^5$ GTU/s × 8 bit/pixel ≈ 1 Tbps. To achieve the required overall data reduction of $\sim 3 \cdot 10^6$ the FLT will have to reduce the trigger rate to ~ 1 Hz/EC, and the PDM based second level trigger to ~ 0.1 Hz/FS [11, 12].

In the following section, specifics of the FLT logic aimed at detecting UHECR and neutrino induced EAS, which are the main scientific objective of the JEM-EUSO mission, are considered. While TLEs and other bright but slow atmospheric events are part of the exploratory objectives of the mission, we will not consider them here.

¹¹⁹ 3. The 1st Trigger Level: Persistency Tracking Trigger

Persistency is a measure of how long in time a signal "persists" or stays 120 within the FoV of a particular MAPMT pixel. As detailed above a GTU 121 roughly reflects the time horizontal EAS right under the ISS need to travel 122 the diagonal of a MAPMT pixel. To cross the FoV of an EC, where the FLT 123 operates, takes up to 45 μ s for EAS, milliseconds for lightning, hundreds of 124 milliseconds for meteors, and seconds for cities or airplanes. These differences 125 in persistency and the fact that the signal moves from pixel to pixel as EAS 126 pass through the FoV of the telescope were exploited in designing the FLT 127 logic, which is described in [13]. Here we give a summary of how an EAS 128 trigger is formed at the EC/FLT level. 129

¹³⁰ Unless EAS develop along the line of sight of the telescope, the image ¹³¹ of their fluorescence trails in the atmosphere can be tracked across the FS. ¹³² Tracking discriminates EAS images against accidental coincidences of back-¹³³ ground light fluctuations.

To decide if a single MAPMT pixel is seeing signal above a slowly vary-134 ing, non-negligible background, this background level has to be estimated. 135 Two different approaches were pursued: The pixel-based estimate sets one 136 threshold for a whole MAPMT. To obtain this threshold the average over 137 a 128 GTU data packet is calculated for each pixel in the MAPMT, and 138 the maximum of these 64 averages becomes the background estimate and 139 threshold for triggering in the next 128 GTU packet. Stationary or slowly 140 moving anthropogenic light sources within the atmosphere are automatically 141 suppressed by this method. The group-based estimate divides the whole EC 142 into 32 groups of 2×4 pixels, calculates the 128 GTU averages per group, 143 and chooses the maximum of those 32 averages as the threshold for all pixels 144 in the EC during the next 128 GTU packet. The threshold here is a digital 145 value as all calculations are done after digitization. No analog thresholds 146

are used. The results presented in this paper are based on the pixel-basedestimate which turned out to be better performing overall.

For tracking purposes each MAPMT's pixels are grouped into square 3×3 149 cells. Each pixel belongs to more than one cell, but since cells do not span 150 MAPMT borders, pixels along the edges and in particular at the corners of 151 each MAPMT belong to fewer cells than the central pixels. The first panel 152 in Fig. 4 shows three such overlapping cells on a single MAPMT's pixel grid 153 outlined in dark orange. Since it includes neighboring MAPMTs' pixels, the 154 3x3 pixel patch surrounded by the dashed grey line on the other hand does 155 not constitute a cell. A total of 36 cells exists within each MAPMT. Each 156 pixel for which its digitized signal in a certain GTU of a 128 GTU packet 157 surpasses the threshold value $n_{\rm thr}^{\rm pix}$ for that MAPMT (determined from the 158 data in the preceding 128 GTU) contributes to each of the cells it participates 159 in. Apart from the pixel level signal threshold there also is a cell level signal 160 threshold $n_{\rm thr}^{\rm cell}$. For a typical background level of one PE per GTU per pixel these thresholds $n_{\rm thr}^{\rm pix}$ and $n_{\rm thr}^{\rm cell}$ would normally be set to 3 and 31, respectively. 161 162 Persistency at the pixel level is evaluated based on two more parame-163 ters that unlike those introduced above do not depend on the background 164 situation: a pre-determined range of consecutive GTUs N_{pst} over which per-165 sistency is to be evaluated, and a limit $N_{\rm ctd}$ on the number of GTUs within 166 that range for which pixels in the cell are above threshold. Standard values 167 for the pixel related parameters $N_{\rm pst}$ and $N_{\rm ctd}$ are 5 and 3 GTUs, respec-168 tively. These values were determined by means of simulations of EAS signals 169 and of preliminary tests on MAPMT fluctuations with the aim of keeping 170 EAS signals and rejecting background fluctuations. Persistency of a signal 171 at the EC level is also monitored and similarly checked by two parameters 172 for a maximum allowed number of GTUs above threshold N_{GTU}^{thr} in a GTU 173 range N_{GTU} . Too many GTUs with signal indicate high persistency, which 174 is the hallmark of non-EAS induced events like lightning or meteors. This 175 GTU range is started at the GTU in which for the first time a cell threshold 176 is surpassed. Typical values for EAS identification are $N_{GTU} = 73$ and N_{GTU}^{thr} 177 = 72. These two values were decided according to the following considera-178 tions. N_{GTU} is determined by the number of GTUs remaining after the first 179 trigger till the end of the packet. For technical reasons it was decided to have 180 the trigger at GTU 55 of a packet. The N_{GTU}^{thr} indicates that all the events 181 are accepted unless the PDM continues triggering every GTU after the first 182 trigger till the end of the packet. Both values will be fine tuned in future, if 183 needed. In particular N_{GTU}^{thr} can be easily shortened by a few GTUs without 184

impacting the trigger efficiency on EAS. 185

As mentioned above power consumption is a major constraint on the ISS. 186 The current implementation of the FLT was programmed into and tested 187 on a Xilinx Virtex6 model XC6VLX240T [14] FPGA. Given that it required 188 only $\sim 7\%$ of this FPGA's resources to accommodate the logic needed for 1 189 EC, it is expected that one such FPGA can host all 9 ECs belonging to one 190 PDM. The graphic in the following figure reflects the FPGA architecture in 191 its reference to the FPGA's various adders. 192

Fig. 4 uses an event recorded by EUSO@TurLab² to illustrate how a 193 FLT is formed. The background estimate was derived from the the preceding 194 GTU packet as $n_{\text{thr}}^{\text{pix}} = 5$ and $n_{\text{thr}}^{\text{cell}} = 65$; the average background light level 195 for that packet was of order 3.4 PE/GTU/pixel. Using the standard values 196 of $N_{\rm ctd} = 3$ and $N_{\rm pst} = 5$, the graphic follows the cell's pixels' recorded PE 197 counts for an EAS-like event created by a line of LEDs mimicking an almost 198 vertical shower mostly staying in the FoV of the cell's central pixel. In the 199 first panel of Fig. 4 the cell is highlighted in dark orange on the MAPMT's 200 pixel map. The next five sub-panels of the figure after the MAPMT pixel 201 overview show two pixel maps each for that same cell. The five sub-panels 202 represent the five successive GTUs following the cell's first threshold crossing. 203 The pixel map on the left in each sub-panel shows the raw PE counts recorded 204 per MAPMT pixel in the respective GTU. Using the estimated $n_{\rm thr}^{\rm pix}=5$ 205 PE background as threshold then leads to the pattern of threshold-crossing 206 pixels displayed on the right of the sub-panel with each pixel's background 207 subtracted PE signal estimates. The sum of that signal above background 208 is then compared to the cell threshold $n_{\rm thr}^{\rm cell} = 65$ PE. In summary: for the 209 chosen cell and its five GTUs after the cell's threshold crossing, at least one 210 pixel in the cell crosses the pixel threshold for each GTU, and the total signal 211 strength accumulated within this cell in each GTU is enough to contribute 212 to the EC wide evaluation of the event. Therefore, at GTU step 5 the 213 corresponding adder (T) is incremented. The EC-wide check with regard to 214 the GTUs during which the signal passes through all the other EC cells is 215 summarized in the last panel, where the content of the adder T is finally 216 checked before a FLT is issued (or not) to the PDM for second level trigger 217 purposes. 218

219

Persistence is the main concept behind the FLT implementation. In the

 $^{^{2}}$ see Section 4.1 and Fig. 11 for more details.

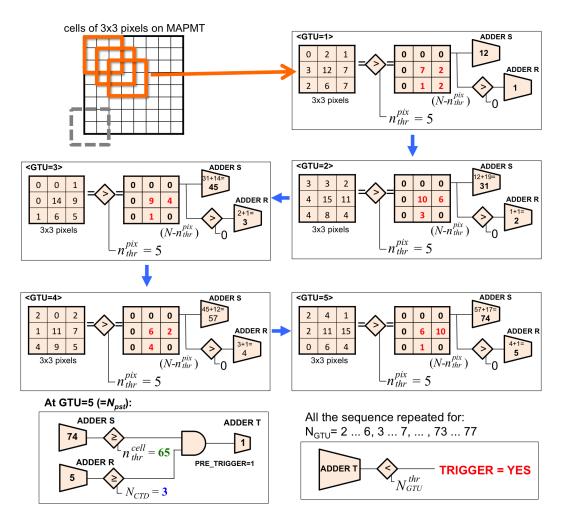


Figure 4: The FLT implementation at the level of the 3×3 cells. See text for details.

following section we will discuss how this current implementation performs in the presence of background, using data recorded with EC modules in dedicated experiments reflecting different aspects of space based EAS observation in the Earth atmosphere.

224 4. FLT Tests using experimental data

The trigger efficiency as a function of EAS energy (commonly referred to as the trigger efficiency curve) captures an important aspect of the experiment's sensitivity. Several publications already discussed the expected JEM-EUSO sensitivity given the current FLT implementation, but so far they were purely based on MC simulations [5, 7, 15]. These simulations all assumed Poisson fluctuations on a UV background intensity that is constant and uniform across the FS.

In this section we report on tests performed offline using data taken with 232 actual ECs in three very different environments, each addressing specific chal-233 lenges JEM-EUSO is expected to face during observation from the ISS: data 234 collected by the EUSO-Balloon flight in 2014 [16], measurements performed 235 by EUSO@TurLab at TurLab [17], and observations in coincidence with a 236 Telescope Array (TA) air fluorescence detector by EUSO-TA [18]. These 237 data sets allow to test the trigger system in very different and complemen-238 tary ways. EUSO@TurLab provides the possibility to control lighting and 230 create realistic event patterns and persistencies, EUSO-Balloon takes data 240 under space-like conditions, and EUSO-TA allows comparison with a well 241 calibrated existing ground-based fluorescence detector. 242

243 4.1. Tests with TurLab measurements

The two main aspects of the FLT that were tested at TurLab, located at 244 the Physics Department of the University of Turin (Italy), were the adequacy 245 of its background estimation and the ability to trigger on EAS while sup-246 pressing other signatures, such as cities, meteors, lightnings, discontinuities 247 in the luminosity due to the presence of clouds, variation in soil condition, 248 moon phase, etc. All these phenomena have variable intensity, duration and 249 extension. Table 1 gives typical values expected for JEM-EUSO for a subset 250 of these conditions which have been reproduced at TurLab to test the trigger 251 logic. 252

Being 15m under ground, the ambient light level in the TurLab laboratory [17] is several orders of magnitude lower than that of the darkest night sky.

light source	intensity	duration	extension	variability
	(cts/pix/GTU)			
UV glow	0.5 - 5	orbit	EC	water, soil, cloud
Urban	3 - 30	seconds	pixel - EC	village - city
EAS	3 - 30	$\sim 100 \mu s$	PMT (track)	EAS energy
Meteor	3 - 100	seconds	EC (track)	magnitude

Table 1: Variability of the signal expected for JEM-EUSO due to different light sources in the FoV of the telescope, ranging from steady UV nightglow to localized and impulsive light bursts such as cities, EAS, meteors. The maximum luminosity of meteors is here defined by the saturation of the front-end electronics.

Using artificial light sources therefore puts the ambient light levels as well as the distribution of light in the lab under the control of researchers.

At TurLab a rotating tank of 5 m diameter provides the stage on which 257 light emitting as well as light reflecting installations are made. EUSO@Turlab 258 consists of one EC hung off-center from the ceiling above this rotating tank. 259 While in principle the EC can be moved radially it was kept at a radius of 260 roughly 2 m from the center of the tank. The optics imaged 1 cm^2 on the 261 tank's surface onto one pixel 2 m above the tank surface, giving it a FoV of 262 the order of 10^{-5} sr, which is only one order of magnitude larger than a JEM-263 EUSO's pixel. This means that if the adjustable speed of the tank rotation 264 were to be around two minutes, the time it takes a stationary source on the 265 tank surface to cross a pixel would be the same as it will be for JEM-EUSO 266 looking down on Earth from the ISS. 267

As outlined before, the data acquisition (DAQ) in JEM-EUSO will be a 268 seamless sequence of 128 GTU long packets. At EUSO@TurLab the EC's 269 ASIC is read out by a test board which transfers the data to a PC, and 270 this system both limits a data packet to 100 GTUs and imposes a ~ 50 ms 271 deadtime between two consecutive packet acquisitions. In other words, at 272 EUSO@TurLab 100 GTUs = 250 μ s of data are taken every ~50 ms, and a 273 stationary light source on the tank surface would have moved through 50%274 of the FoV of a pixel during that deadtime if the tank rotated with a period 275 of 2 min. It was not possible to synchronize the DAQ with the tank rotation, 276 because it was not foreseen by the control system of the tank. Naturally, the 277 synchronization would have allowed to determine exactly the location in the 278 tank responsible for each trigger. 279

280

Fig. 5 shows the various components of the TurLab setup. It shows the

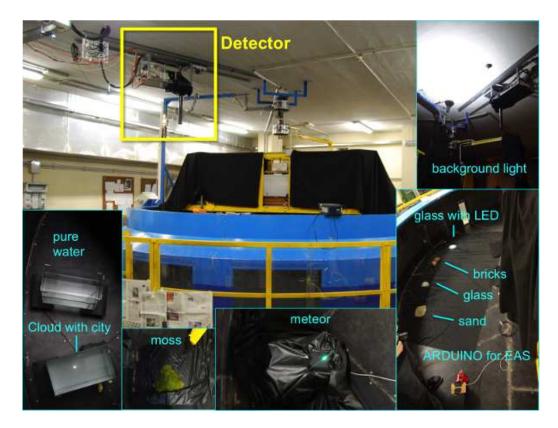


Figure 5: The TurLab rotating tank. The black tube on the ceiling shows the collimator of the experimental setup used to mimic the JEM-EUSO telescope. Light sources and materials used to mimic other phenomena are shown as well.

EC suspended from the ceiling, and various installations on the tank surface 281 designed to emit or reflect light in ways that mimic both anthropogenic and 282 natural lighting situations as they would be seen by JEM-EUSO from the 283 ISS (see Table 1). Light scattered or reflected from the atmosphere back into 284 space - the Earth's albedo - depends on atmospheric conditions like e.g. the 285 presence or absence of clouds, the reflectivity of the Earth's local surface, and 286 lighting conditions like the phase of the moon³ or the presence and density 287 of human habitation. 288

In Fig. 5 one can also see how sand, moss, ground glass, pure water, 280 and a brick were used to mimic the reflection of night-sky light from soils, 290 forests, snow, water, and rocky surfaces, respectively. Water clouded by dis-291 solved particles and illuminated from below is used to mimic clouds, and if 292 illuminated from below cloud cover over e.g. a city. An oscilloscope screen 293 displaying Lissajous traces mimics meteor tracks. As lighting can be con-294 trolled, the TurLab tank allows to verify the performance of the background 295 estimation under realistically varying lighting conditions. Fig. 6 and Fig. 7 296 show two examples of recordings of such features with EUSO@TurLab. 297

In Fig. 6 LED light reflected from ground glass is used to simulate the 298 distributed individual light sources of a city. Each of the four frames in each 299 sub-panel refers to one of the four MAPMTs in the EC. The upper panels 300 show the respective MAPMT's integrated PE counts, in other words the 301 light curve of a city passing through the FoV of the EC. The city entering 302 and exiting the MAPMT's FoV as time progresses is clearly visible in each 303 MAPMT's light curve. The lower row shows 2D pixel maps for the EC's 304 MAPMTs, with the PE counts per pixel for just one GTU on the left, 10 305 GTUs in the center, and 100 GTUs on the right. The red lines in the light 306 curves show the range of GTUs that are used, with the single GTU pixel maps 307 being the first GTU under both red lines; the start times of this integration 308 is the same for all three ranges. Fig. 7 shows the data recorded while passing 309 over the oscilloscope repeating a straight line Lissajous figure taking about 310 one second to complete. While the complete picture emerges after integrating 311 over 1500 GTUs (right panel), the signal still is contained in a single pixel 312 when integrating over only 10 GTUs (center panel). 313

An Arduino board [19] controlling a line of 10 white LEDs was used to mimic a single EAS propagating through the atmosphere at the speed of

³Fluorescence observation of EAS is not possible during daytime.

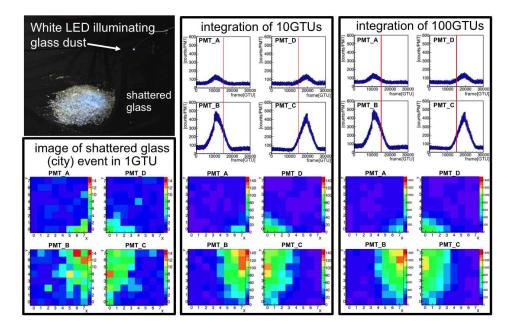


Figure 6: Reproduction of an extended light source, like in case of a city. Top-left plot is a picture of the shattered glass; bottom left plot is the image detected by the MAPMTs in 1 GTU. The right-top plots shows the temporal evolution of the same scene with different time integrations (10 and 100 GTUs). The bottom plot shows one frame per integration taken at the time indicated by the red line in the above plots.

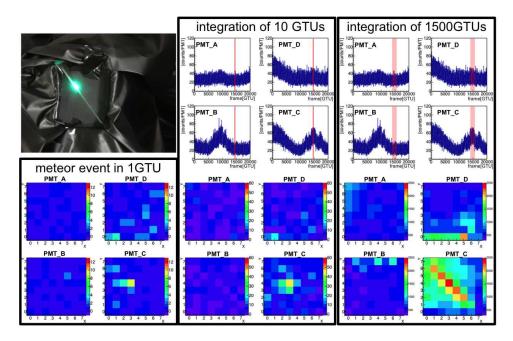


Figure 7: Reproduction of a meteor-like track. See Fig. 6 for details about the meaning of each plot in the figure.

light, resulting in a total duration of about 40 GTUs. As can be seen in
Fig. 8 this signal does no longer stay within one pixel during 10 GTUs, with
the center of light moving visibly between subsequent GTUs.

The FLT estimates the background for the current data packet from the 319 data collected in the preceding data packet. To mitigate possible adverse 320 effects of the DAQ-imposed deadtime between the acquisition of consecutive 321 packets on background estimation the tank rotation was slowed to complete 322 one rotation in 9 min, reducing the offset of a stationary light source between 323 consecutive acquisition packets from 50% of a pixel to roughly 10% of a pixel; 324 with the JEM-EUSO DAQ a 128 GTU offset at ISS speed would correspond 325 to 0.5% of the pixel size projected onto the ground. Given the deadtime 326 between the 100 GTU acquisition packets a total of ~ 3 seconds of data is 327 collected during one 9 minute rotation. 328

The DAQ at TurLab collects that data "as is": it simply reads out the PE counts for each MAPMT pixel in each GTU from the EC's ASIC and writes them to disk. The subsequent trigger simulation is then implemented

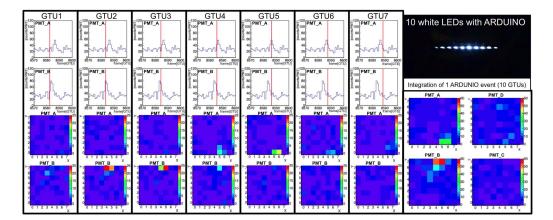


Figure 8: Reproduction of a cosmic ray-like track. Top-right picture shows the integrated light sequence reproducing a cosmic-ray track. Bottom-right plot shows the integrated number of counts during the light sequence. The left part of the figure displays 7 frames of 1 GTU each taken during the reproduction of the track. The time at which the frames are taken is shown in the above corresponding plots which present the time evolution of the total number of counts recorded by the MAPMT.

in VHDL⁴ according to the schema described in [13].

The light collected on one of the EC's MAPMTs during a complete 9 333 minute rotation of the tank is shown in the upper panel of Fig. 9. Changes 334 in the background light level are clearly seen, and the various contraptions 335 that precipitate them are labeled in the figure. The two bridges refer to the 336 footbridges to cross the tank which, despite being covered by some black 337 fabric, are a source of quite variable light reflection. In general the black 338 fabric was used to make as dark as possible specific regions of the tank to 330 help increase the dynamic range of the light intensity seen by the MAPMTs 340 during the tank rotation. Pure water in a little transparent tank was used 341 to mimic a mirror-like condition which induces much higher reflection. The 342 vellow bar is a pole on the rotating tank which passes a few cm below the 343 collimator, thus filling a significant portion of the FoV of the detector for a 344 short time. The second panel of Fig. 9 shows the average PE count in the 345 preceding 100 GTU packet for the pixel with the maximal average count in 346 that same MAPMT, which is the value used for the threhold setting in the 347 current packet. The final panel shows how the trigger simulation reacted 348

⁴Very high speed integrated circuit Hardware Description Language

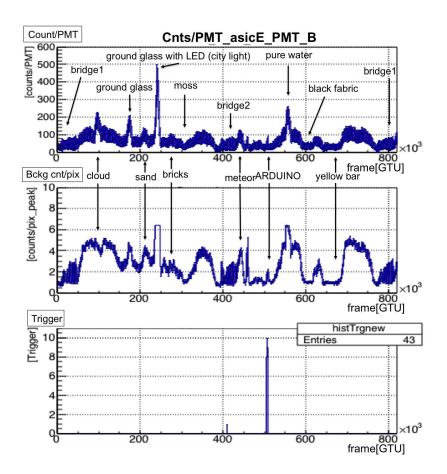


Figure 9: Reproduction of a full TurLab rotation with many types of light. Top plot shows the sum of the light collected by 1 MAPMT as a function of time. Middle plot shows the light intensity monitored by the pixel responsible to set the trigger thresholds of the MAPMT. Bottom plots show the triggered events. Except for two spurious cases due to quite variable background conditions (see middle panel) which could not be properly followed with the 50 ms dead time between packets, all the triggers coincide with the cosmic ray-like events generated by Arduino.

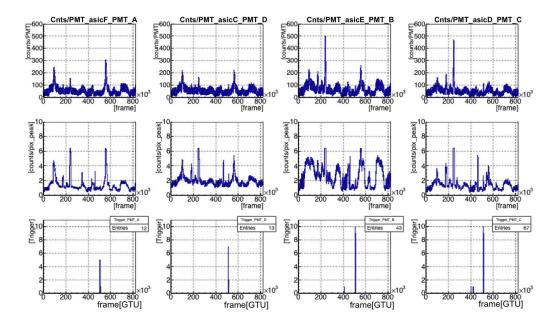


Figure 10: Reproduction of a full TurLab rotation with many types of light sources as shown in fig. 9 for the entire EC.

to this input. It shows when FLTs were issued based on signals in that MAPMT. Almost all triggers coincide with passing over the Arduino driven LED chain as it should be; the one that is not is due to a specific location near one of the two bridges crossing the tank where the variations of light reflection were still too fast to be compensated by the slower rotation of the tank. Fig. 10 shows top and bottom plots of Fig. 9 for all four MAPMTs. A similar response is obtained for all four MAPMTs all along the rotation.

To directly assess the impact of background on the FLT trigger scheme for 356 EAS dedicated measurements were made at TurLab with the tank rotation 357 stopped and the EC stationary above the Arduino driven white LED strip 358 simulating EAS. Ambient light levels then controlled the background to the 359 LED induced signal. These ambient light levels were varied between 0.1 360 and 2.0 PE per pixel and per GTU, reflecting expectations for typical ISS 361 observation background. The Arduino EAS were generated 1 ms apart in 362 order to reduce the probability of recording such Arduino EAS in consecutive 363 data packets, in which case the first EAS would set the background level for 364 the second EAS. As the DAQ for EUSO@TurLab was not synchronized with 365 track timing in the Arduino, extracting the packets containing a complete 366

³⁶⁷ Arduino track required some event selection.

This selection started from a 4×4 pixel box in that MAPMT which contained the brightest part of the Arduino LED simulated EAS. The stationary tank was oriented such that the Arduino LEDs were all within the field of view of a single MAPMT and the Arduino EAS were always crossing the same MAPMT pixels.

The LED sequence for these Arduino EAS was kept stable with ~ 30 PE at maximum, which corresponds to recording a $\sim 6 \times 10^{19}$ eV EAS in JEM-EUSO. A mask above the LEDs was used as an aperture to avoid unwanted reflections of LED light from nearby structures above the tank. The voltages supplied to the LEDs were also adjusted to dim the LEDs that were closer to the ends of the strip in an effort to provide a realistic EAS profile when the Arduino board sequentially lights up the LEDs in the strip.

If the PE count in the 4×4 pixel box smoothed over 5 GTUs exceeded the corresponding background estimate by more than 4σ, the data were considered an Arduino EAS candidate. Such a candidate would subsequently be rejected if the excess occurred only in the first or last five GTUs of a 100 GTU data packet, or if the preceding data packet also contained an Arduino EAS candidate. Fig. 11 shows PE counts for the relevant MAPMT over

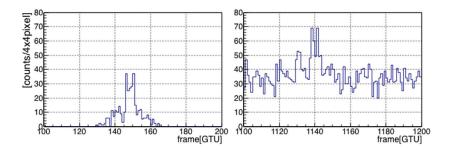


Figure 11: Example of light curves of two extracted Arduino events without background counts and with a background condition of ~ 2 counts GTU^{-1} pixel⁻¹, respectively.

385

time in GTU units. The left panel shows a typical event produced by the
LED strip without background. This highlights the event's original shape.
The right one is a similar event produced under high background.

Events selected by this procedure were then fed into the VHDL trigger simulation. Table 2 shows that all the selected Arduino EAS also triggered in the FLT simulation. This together with the fact that when observing the rotating tank with its various implementations of atmospheric as well as

$N_{PEave./4 \times 4pix}$	N _{triggered}	N _{extracted}
0.21	24	24
0.36	13	13
0.56	15	15
0.78	21	21
1.00	26	26
1.23	16	16
1.47	26	26
1.69	22	22
1.93	20	20
2.11	31	31

Table 2: Number of triggered and extracted cosmic-ray-like-track events in various background photon level conditions. $N_{PEave./4 \times 4pix}$ indicates the average background level expressed in counts per pixel per GTU evaluated on a 4×4 pixel-box during the preceding packet of data, where no Arduino event was extracted.

ground-based light sources, albedo effects and generally varying background light levels gave rise to only a few spurious triggers under very specific conditions gives confidence that the current FLT implementation is ready for deployment in JEM-EUSO.

397 4.2. Tests with EUSO-Balloon data

The EUSO-Balloon [20, 21] data taken during a 5 hour flight at 38 km 398 altitude in the vicinity of Timmins in Canada provides another testbed for 399 the FLT. Again the adequacy of the newly adopted background estimation 400 method with respect to keeping the trigger rate within the permissible bounds 401 in the presence of artificially and naturally encountered fluctuations in the 402 background lighting conditions as well as the FLT's ability to trigger on 403 relevant optical phenomena was studied. While at TurLab the optics and 404 speed could be adjusted to match event duration and persistence in a pixel's 405 FoV for the various phenomena recreated there, the EUSO-Balloon's speed 406 and trajectory could not be controlled to that extent. On the other hand 407 EUSO-Balloon looked down on a real Earth environment just as JEM-EUSO 408 will, albeit from a much closer distance than the ISS. Thus were TurLab 400 strove to be realistic in an artificial environment, EUSO-Balloon was looking 410 at realistic settings compromising on perspective. Thus they each capture 411 different aspects of the challenges presented to an FLT operating at JEM-412

413 EUSO.

In flight the EUSO-Balloon optics imaged a 60 km^2 surface area onto one 414 full PDM with its 9 ECs and a total of 36 MAPMTs. Just as at TurLab 415 the data acquisition did not allow to record data continuously, but took 128 416 GTU data packets at 18 Hz, translating into 320 μ s of data recorded every 417 \sim 55 ms. Data taking on the balloon was organized into runs. A new run 418 would be started after either 200 or 2000 data packets had been recorded. 419 In total about $4 \cdot 10^7$ GTUs were recorded with the camera looking down on 420 natural backgrounds like forests, lakes, and clouds, as well as city lights. 421 Fig. 12 traces the time variation of a typical one of the 2304 pixels in the

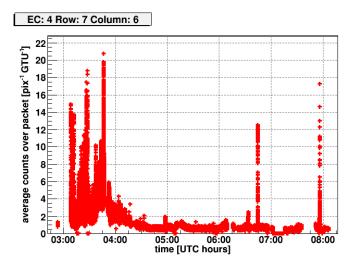


Figure 12: PE counts averaged over each recorded data packet for one typical pixel, covering the entire balloon flight.

422

camera. The city of Timmins for example crossed the field of view of that
pixel between 03:00 and 04:00 UTC. The low counts at the center of the data
taking period correspond to times when the balloon was passing over forests,
lakes, and clouds. The highest PE counts were recorded when the balloon
passed over an active mining operation.

For about two hours a helicopter was flying under the balloon. The helicopter carried three different light sources: a UV LED, a Xe-Flashlamp and a UV laser. This sources had the purpose of calibrating the detector response and simulating EAS-like event patterns that in this paper are used offline to test the FLT. Data analysis focused on about one hour of balloon flight during which the helicopter was shooting underneath the balloon and in the FoV of the camera. During this hour no other system tests interfered
with the measurements, and the flight path crossed over dark as well as bright
areas. This is essential to test the FLT logic under extreme conditions.

The data was divided into two data blocks: One block with so-called "nominal" background levels as are expected for operation on the ISS with $\sim 0.6 \text{ PE/pixel/GTU}$ averaged over a packet, and one block with a "high" background of 5–10 PE/pixel/GTU average over a packet in some parts of the FoV. This latter block of data was collected while passing over the mine or a city. About 15% of MAPMT pixels that were not working properly were excluded from the analysis.

The light sources on the helicopter were set up to emit signature patterns 444 that each served a distinct purpose [22]. First in the sequence was a UV-LED 445 (375 nm wavelength), the light output of which steadily increased with time 446 over 12 GTUs. From the balloon this signal appears as a stationary source 447 typically contained in a single MAPMT pixel. The UV-LED light output 448 was kept stable throughout the night and designed to raise the signal level 449 from ~ 1 to ~ 50 PE over the 12 GTUs in that pixel. This signal provides 450 a normalization for the distance between helicopter and EUSO-Balloon and 451 allows to determine an effective threshold for the FLT. 452

Next in the sequence was a laser pulse shot horizontally away from the 453 helicopter. This laser shot was fired about 25 GTU after the end of the 454 UV-LED signal, delivering ~ 5 mJ over 7.5 ns at a wavelength of 355 nm. 455 Depending on where in the balloon's FoV the helicopter happened to be 456 at that time, it could take a maximum of 10 GTUs before the laser pulse 457 would leave the balloon's FoV. The number of photons scattered out of such 458 a laser pulse roughly corresponds to the fluorescence light emitted at shower 459 maximum from a $\sim 10^{20}$ eV EAS according to ESAF simulations [23]. 460

The balloon's altitude being low compared to the ISS however meant that 461 the $\sim 400 \text{ m} \times 400 \text{ m}$ of a 3×3 pixel cell on the ground was crossed by the 462 laser pulse in ~ 1 GTU, while the FLT is integrating over 5 GTU to establish 463 a threshold crossing. To retain the ability for the FLT to trigger, the last 464 light source in the sequence of light sources operated on the helicopter is a 465 xenon (Xe) discharge lamp (wavelength 337 nm) emitting its light over ~ 8 466 GTU, hereafter referred to as Xe flasher. This Xe flasher is triggered $\sim 5 \ \mu s$ 467 after the laser shot, and its light curve reaches its maximum three GTU from 468 its start, decreasing thereafter. This is reminiscent of the light curve along 469 EAS, and four different flash intensities were used to mimic different EAS 470 energies. A total helicopter light sequence therefore extends over ~ 50 GTU. 471

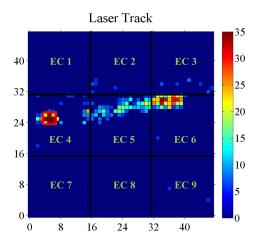


Figure 13: Image of a helicopter event obtained by integrating the counts in each pixel for the whole packet=1960 of RUN=043202 (128 GTUs). A threshold is applied to the minimum signal level to emphasise the location of the track. The UV-LED and Xe-flasher signals are centred around pixel at axis of abscissae X=5 and axis of ordinates Y=25.

Fig. 13 shows the integrated PE count per pixel for a 128 GTU data 472 packet containing a whole helicopter light sequence. An offset was subtracted 473 throughout to highlight the excess along the laser track. The helicopter with 474 the UV-LED and the Xe flasher was in the pixel at x=5, y=25. The left panel 475 of Fig. 14 shows the PE sum for each GTU in that same data packet for the 476 3×3 pixel cell that is centered on the helicopter position at x=5, y=25. The 477 UV-LED ramp can be seen to start from GTU 19. The laser shot is seen at 478 GTU 55 and 56, and the Xe flash lamp is seen between GTU 58 and 65. The 479 signal peaking at GTU 72 is attributed to an afterpulse in the Xe flasher. 480

Fig. 13 also identifies the ECs that the various signals appear in. For the 481 UV-LED and the Xe flasher this is EC number 4, and the laser track starts 482 in EC number 4 and then moves out through ECs 5 and 6. The right hand 483 panel of Fig. 14 shows the number of ensuing FLT triggers per GTU for all 484 ECs in that same data packet. The red line refers to triggers in EC number 485 4, green to EC number 5, and blue to EC number 6. All signals are detected 486 by the FLT algorithm, and the delay between the signal and trigger timings 487 simply reflect the 5 GTU integration. 488

As the EUSO-Balloon DAQ and the helicopter light sequence each run on their own respective clocks and the DAQ recorded only $2.5 \times 128 = 320 \ \mu s$ every ~55 ms (~0.6%), the vast majority of helicopter light sequences were not recorded. There will also be events where only a part of the helicopter
light sequence overlapped with a DAQ data packet. Running the offline FLT
simulation through the data, 274 events were found in which at least two
ECs triggered the FLT algorithm.

Another peculiarity of EUSO-Balloon was that the optical module under 496 the balloon spun, constantly changing the FoV's alignment with respect to 497 both surface features and laser direction, and did so at a varying rate. This 498 meant that especially at the edges of the FoV stationary and other light 499 sources would often enter or exit the FoV during the \sim 55 ms dead time be-500 tween data packets. This complication should clearly be more relevant where 501 stationary light sources on the ground play a significant role. As described 502 above, the data was sorted into two blocks: one with nominal and the other 503 with high background over locations lit up by human activity. The first block 504 contains a total integrated time corresponding to ~ 8.5 seconds, and the lat-505 ter corresponding to ~ 6.5 seconds, with about one order of magnitude more 506 background light in this latter block's data on some parts of the FoV. 507

As expected the trigger rate under the more severe background conditions 508 is higher: In the high background block of data the FLT algorithm triggered 509 on 148 laser events and 59 others, while it triggered 126 laser and 17 other 510 events in the nominal background data block. Assuming all other events are 511 background, this puts the background rates for the current FLT trigger logic 512 at 2.0 Hz per 9 ECs for the nominal background data block and 9.1 Hz per 513 9 ECs for the high background data block. Under both conditions the rate 514 requirement of ~ 1 Hz per EC is met. In particular this means that despite 515 the particular challenge posed by the combination of balloon spinning and 516 DAQ deadtime the background estimation using the preceding data packet 517 works well. 518

The event shown in Fig. 13 and Fig. 14 and triggered on by the FLT 519 algorithm can be used to estimate an energy threshold for EAS that would 520 pass the FLT. Averaging over the seven lowest PE/GTU values that raise 521 trigger alerts in EC number 4 the average signal excess becomes 81 ± 13 and 522 the average background 39 ± 1 PE/GTU. This is a signal over background 523 ratio of 2.1 \pm 0.3 for the 3 \times 3 pixel cell raising the trigger. Comparing 524 this to ESAF simulations for EUSO-Balloon [23] under nominal background 525 conditions this ratio is reached for vertical EAS initiated by a $\sim 5 \times 10^{18}$ eV 526 proton. As the simulation also shows that showers at higher zenith angle 527 provide higher signal/GTU, this value should be considered an upper limit 528 for the energy threshold of FLT-triggered events recorded by EUSO-Balloon. 529

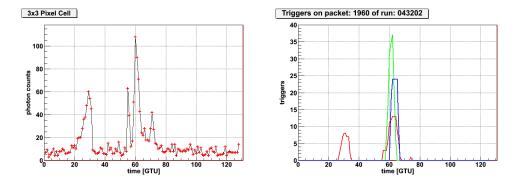


Figure 14: Left: Number of counts recorded in the 3×3 pixel-cell centred around (X=5,Y=25) during the entire packet. See text for details. Right: Sequence of trigger alerts in the different ECs crossed by the laser track during the entire packet. See text for details.

Given its FoV and the measured CR rate at this energy this means that the FLT should trigger one event for every 24 hours of EUSO-Balloon livetime.

532 4.3. Tests with EUSO-TA data

TurLab measurements and EUSO-Balloon data were used to verify that 533 the FLT and in particular its background estimation perform and meet the 534 requirements under various realistic or even challenging background condi-535 tions. While comparing the lowest light level in EUSO-Balloon events that 536 raised a FLT with simulation produced an estimate for the energy thresh-537 old in detecting cosmic ray particles, this is still a far cry from obtaining 538 an efficiency curve for the FLT. This problem is addressed with data from 539 EUSO-TA. 540

The EUSO-TA [18] telescope is a prototype of the JEM-EUSO space telescope with two 1 m^2 square Fresnel lenses. Just as for EUSO-Balloon its electronics comprise a full PDM with 9 ECs and 36 MAPMTs.

It is located right in front of the Black Rock Mesa (BRM) fluorescence 544 detector (FD) site of the TA experiment in the Utah West Desert, USA [24]. 545 EUSO-TA's FoV of $11^{\circ} \times 11^{\circ}$ is contained within that of the BRM's FD and 546 aligned such that it contains the vertical tracks from the pulsed 355 nm laser 547 at TA's central laser facility (CLF). During TA data taking on the moonless 548 parts of nights with amenable weather the CLF fires 300 vertical laser pulses 549 of 3 mJ at 10 Hz every half hour. Providing atmospheric and calibration data 550 for all three of TA's FD sites it is located centrally at an equal distance of 551 21 km from each of the TA FDs, and therewith also 21 km from EUSO-TA. 552

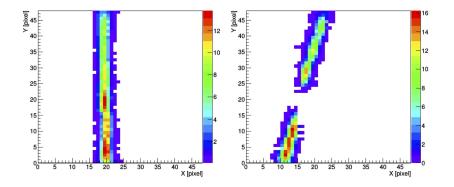


Figure 15: Left: an average of ~ 250 shots of CLF laser; right: an average of ~ 150 inclined shots of the Colorado School of Mines laser, located at 40 km from TA-EUSO, the missing part due to a non- functioning MAPMT in the center of the focal surface. The color scale on both pictures denotes the detector counts. Figure taken from [25].

⁵⁵³ Depending on the offset between GTU boundaries and laser shot, laser tracks ⁵⁵⁴ took 6 to 8 GTUs for their image to cross the PDM at EUSO-TA. The left ⁵⁵⁵ panel of Fig. 15 shows an average over ~ 250 such CLF shots as recorded by ⁵⁵⁶ EUSO-TA.

The inclined laser track shown in the right panel of Fig. 15 is from a set of laser events produced with the help of a mobile UV laser belonging to the Colorado School of Mines. The missing piece in this laser track average was due to a defective MAPMT in EUSO-TA.

Also using a 355 nm laser the pulses from this mobile laser can be adjusted in intensity within a range of 1 to 86 mJ. As the laser itself is steerable, the geometry of the laser track can be varied more freely, and for the average over the \sim 150 laser pulses shown here the laser was shot at a distance of 40 km with a pulse energy of 62 mJ.

Varying the laser pulse energies with this mobile laser at 34 km from 566 EUSO-TA produced the trigger efficiency curve for the FLT that is shown 567 in Fig. 16. As at these distances the laser pulses typically cross a few 568 pixel/GTU, the FLT logic was adapted by setting $N_{pst} = 1$, while n_{thr}^{pix} and 569 $n_{\rm thr}^{\rm cell}$ were modified accordingly to keep the FLT trigger rates below the 1 570 Hz/EC requirement. To determine the trigger efficiency, an external trig-571 ger, synchronized with the laser shooting, was supplied by the TA-FD to the 572 EUSO-TA DAQ to always have the laser track inside a 128 GTU packet. 573 The efficiency can then be determined by running the adapted FLT algo-574

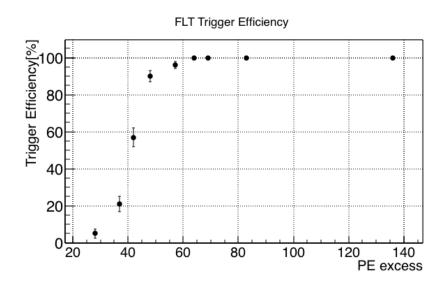


Figure 16: Trigger efficiency curve of FLT as a function of the signal excess recorded by TA-EUSO.

rithm over these data packets and counting the packets that raise an FLT. Laser pulse energies between 3 and 5 mJ were used for the first four points with signal excesses above background < 50 PE over all pixels in that GTU. Above 50 PE overall signal excess, which corresponds to ~ 25 PE in the relevant 3×3 pixel cell and 6 mJ pulse energy for this geometry, the trigger efficiency becomes 90% and higher. In a corresponding analyis for the 21 km CLF geometry the CLF's 3 mJ pulses were seen with 94% efficiency.

582 5. Conclusions

The FLT logic for use in JEM-EUSO as described above and implemented in VHDL was shown here a.) to work well in the presence of artificially produced as well as naturally encountered fluctuations in the background lighting conditions and b.) to keep the FLT rate within the permissible bounds while c.) being efficient at identifying event types with general EAS characteristics.

The FLT trigger as presented here is working at the MAPMT level and is based on the local persistency of a signal excess in a 3×3 pixel area, persisting a few GTUs. To achieve this an automatic evaluation of the average background level is derived from the preceding data package, as strategy that has proven successful even when individual data packages were separated by ⁵⁹⁴ up to a few hundred μ s. Rejection for events with time duration too large ⁵⁹⁵ for an EAS signal, namely longer than 72 GTUs on the ISS, is also imple-⁵⁹⁶ mented. This implementation for one EC requires only a few per cent of the ⁵⁹⁷ resources of commercial FPGAs, which allows to implement it within the ⁵⁹⁸ power constraints imposed on the ISS.

Tests performed with EUSO-Balloon and EUSO-TA data, as well as mea-599 surements performed at the TurLab facility, allowed validating the main 600 functions of the algorithm. The system automatically adjusts the thresh-601 olds to keep the rate of triggers on background fluctuations below 1 Hz/EC 602 even in the case of slow background variations. The FLT level trigger detects 603 EAS-like events with light intensities comparable to those JEM-EUSO would 604 observe in the energy range $E > 5 \cdot 10^{19}$ eV and in the presence of expected 605 night sky background. These results strenghten those obtained in [7] and 606 successive publications as they show that the trigger concept developed from 607 simulation can be effectively implemented in hardware and performs well on 608 real data. 609

The FLT has shown to be quite effective in rejecting city-like and other slow but bright events such as meteors. Of the few spurious triggers that occurred most were artefacts of discontinuities introduced by the available equipment.

The examples shown in this paper are only a sub-sample of all tests performed on the data reported here and the ongoing activities at the TurLab facility and EUSO-TA.

The VHDL logic of the FLT is currently being implemented on the FPGA of the PDM board. EUSO-SPB [26], the next stratospheric balloon flight, is expected to host this trigger logic on-board to verify its performance on real EAS.

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647 References

- [1] R. Benson and J. Linsley, Proc. 17th Int. Cosmic Ray Conf. (Paris),
 (1981) 8.
- ⁶⁵⁰ [2] J.H. Adams et al. (JEM-EUSO Coll.), Exp. Astronomy 40 (2015) 3.
- ⁶⁵¹ [3] F. Kajino et al. (JEM-EUSO Coll.), Nucl. Inst. Meth. A 623 (2010) 422.
- ⁶⁵² [4] J.H. Adams et al. (JEM-EUSO Coll.), Exp. Astronomy 40 (2015) 19.
- ⁶⁵³ [5] M. Bertaina et al. (JEM-EUSO Coll.), Adv. Space Res. 53 (2014) 1515.
- ⁶⁵⁴ [6] C. Berat et al., Astrop. Phys. 33 (2010) 221.
- ⁶⁵⁵ [7] J.H. Adams Jr. et al. (JEM-EUSO Coll.), Astrop. Phys. 44 (2013) 76.
- ⁶⁵⁶ [8] M. Sato et al., Int. Journal of Mod. Phys. A 20/29 (2005) 6903.
- [9] N. Sakaki et al., Proc. 30th Int. Cosmic Ray Conf. (Merida), HE 5 (2007)
 1037.
- ⁶⁵⁹ [10] G.K. Garipov et al. (Tatiana Coll.), JETP Letters 4 (2005) 185.

- [11] J. Bayer et al. (JEM-EUSO Coll.), Proc. 32th Int. Cosmic Ray Conf.
 (Beijing), 3 (2011) 168; arXiv:1204.5065.
- [12] J. Bayer et al. (JEM-EUSO Coll.), Proc. 33th Int. Cosmic Ray Conf.
 (Rio de Janeiro), #0432 (2013); arXiv:1307.7071.
- [13] M. Bertaina et al. (JEM-EUSO Coll.), Nucl. Instr. & Meth. A 824 (2016)
 253.
- ⁶⁶⁶ [14] http://www.xilinx.com/support/documentation/data_sheets/ds150.pdf
- ⁶⁶⁷ [15] M. Bertaina et al. (JEM-EUSO Coll.), Exp. Astronomy 40 (2015) 117.
- ⁶⁶⁸ [16] J.H. Adams et al. (JEM-EUSO Coll.), Exp. Astronomy 40 (2015) 281.
- [17] M. Bertaina et al. (JEM-EUSO Coll.), EPJ Web of Conferences 89
 (2015) 03003.
- ⁶⁷¹ [18] J.H. Adams et al. (JEM-EUSO Coll.), Exp. Astronomy 40 (2015) 301.
- 672 [19] http://www.arduino.cc
- [20] P. von Ballmoos et al. (JEM-EUSO Coll.), Proc. 34th Int. Cosmic Ray
 Conf. (Den Haag), #0725 (2015).
- [21] M. Bertaina et al. (JEM-EUSO Coll.), Proc. 34th Int. Cosmic Ray Conf.
 (Den Haag), #0890 (2015).
- [22] J. Eser et al. (JEM-EUSO Coll.), Proc. 34th Int. Cosmic Ray Conf. (Den Haag), #0860 (2015).
- [23] F. Fenu et al. (JEM-EUSO Coll.), Proc. 34th Int. Cosmic Ray Conf.
 (Den Haag), #0639 (2015).
- [24] T. Abu-Zayyad, et al. (Telescope Array Coll.), Nucl. Instrum. Meth.
 A689 (2012) 87.
- [25] M. Casolino et al. (JEM-EUSO Coll.), Proc. 34th Int. Cosmic Ray Conf.
 (Den Haag), #0636 (2015).
- [26] L. Wiencke et al. (JEM-EUSO Coll.), Proc. 34th Int. Cosmic Ray Conf.
 (Den Haag), #0165 (2015).