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Authors	FRASCA, Antonio; Montes, D.; ALCALA', JUAN MANUEL; Klutsch, A.; Guillout, P.
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Contemporaneous Broad-Band Photometry and H α Observations of T Tau Stars*

A. Frasca¹, D. Montes², J. M. Alcalá³, A. Klutsch¹ and P. Guillout⁴

¹INAF – Osservatorio Astrofisico di Catania, via S. Sofia, 78, I-95123 Catania, Italy
e-mail: antonio.frasca@oact.inaf.it

²Departamento de Física de la Tierra y Astrofísica & UPARCOS-UCM (Unidad de Física de Partículas y del Cosmos de la UCM), Facultad de Ciencias Físicas, Universidad Complutense de Madrid, E-28040 Madrid, Spain
e-mail: dmontes@ucm.es

³INAF – Osservatorio Astronomico di Capodimonte, via Moiariello 16, I-80131 Napoli, Italy

⁴Université de Strasbourg, CNRS, Observatoire astronomique de Strasbourg, UMR 7550, F-67000 Strasbourg, France

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ABSTRACT

In this paper we present the results of a simultaneous *BVR*I and H α photometric monitoring, contemporaneous to medium-resolution spectroscopy of six T Tau stars in the Taurus-Auriga star forming region. We have characterized the H α photometric system using synthetic templates and the contemporaneous spectra of the targets. We show that we can achieve a precision corresponding to 2–3 Å in the H α equivalent width, in typical observing conditions.

The spectral analysis has allowed us to determine the basic stellar parameters and the values of quantities related to the accretion. In particular, we have measured a significant veiling only for the three targets with the strongest H α emission (T Tau, FM Tau, and DG Tau).

The broad-band photometric variations are found to be in the range of 0.05–0.70 mag and are often paired to variations in the H α intensity, which becomes stronger when the stellar continuum is weaker. In addition, we have mostly observed a redder $V - I$ and a bluer $B - V$ color as the stars become fainter. For most of the targets, the timescales of these variations seem to be longer than the rotation period. One exception is T Tau, for which the broad-band photometry varies with the rotation period. The most plausible interpretation of these photometric and H α variations is that they are due to non-stationary mass accretion onto the stars, but rotational modulation can play a major role in some cases.

Key words: *Stars: pre-main sequence – Stars: activity – Accretion, accretion disks – Open clusters and associations: individual: Taurus*

*Based on observations made with the 0.91-m telescope of Catania Astrophysical Observatory (Italy) and with the INT telescope operated on the island of La Palma by the Isaac Newton Group of Telescopes in the Spanish Observatorio del Roque de los Muchachos of the Instituto de Astrofísica de Canarias.

1. Introduction

Young pre-main sequence (PMS) stars display a rich variety of phenomena in their atmospheres and circumstellar environments, including magnetic activity, mass accretion from circumstellar disks, outflows, and jets, which produce typical signatures in their spectra and energy distributions. However, different phenomena can produce similar effects on the observed properties of PMS stars or can affect the same spectral features. Indeed, most diagnostics of mass accretion from the disk onto the central object for the classical T Tau stars (CTTs) are emission lines that are also diagnostics of chromospheric activity for weak-lined T Tau stars (WTTs). The latter, however, display both much weaker emission lines and significantly lower or no infrared excess than the former, hence undergo much less or no accretion. The boundary between weak accretors and WTTs is not sharp, as chromospheric emission is always present in PMS stars, and may provide a non-negligible contribution to the excess emission with respect to the photospheric flux in weak accretors (see, *e.g.*, Manara *et al.* 2013, Ingleby *et al.* 2011, Frasca *et al.* 2015, Alcalá *et al.* 2017, and references therein).

An important feature of PMS stars is the variability of both their continuum flux and emission lines with both short and long timescales (*e.g.*, Cody *et al.* 2014, Costigan *et al.* 2014, and references therein). The photometric variations in different bands carry important information on the physical mechanisms operating in the PMS stars and in their circumstellar environments. Rotation periods can be determined for both the CTTs and the WTTs. Indeed, due to the rotation of the star, the stellar flux can be modulated by hot spots that are produced by the impact of the accretion flow on the stellar photosphere, for the strongly accreting CTTs, or to photospheric dark spots related to the stellar magnetic fields for both classes of objects. In addition, flux variations on time scales different than the star rotation period can arise from variable accretion and outflows, and obscuration by the circumstellar disk (*e.g.*, Bouvier *et al.* 2007, Rodriguez *et al.* 2017).

High-cadence, long-term photometry, like that obtained using space telescopes, is able to reveal different types of variability with various time scales (periodic variations, dips, bursts, stochastic variations, etc.), which are interpreted as the result of different physical processes and geometric effects (*e.g.*, Cody *et al.* 2014, Sousa *et al.* 2016, Venuti *et al.* 2017). Multiband photometry in optical and near-infrared (NIR) bands is useful, for the periodic variables, to distinguish between hot and cool spots as the source of the modulation, and to determine their mean properties (*e.g.*, Bouvier *et al.* 1993, Frasca *et al.* 2009).

Despite the ever growing number of photometric data of high quality as regards precision and cadence, there are few works in the literature based on multiband photometric observations taken contemporaneously to spectroscopy (*e.g.*, Bouvier *et al.* 2003, Alencar *et al.* 2012, Sousa *et al.* 2016). The study of contemporaneous variations of the continuum flux and emission spectral lines, which are diagnostics of jets, accretion, and/or chromospheric activity, is of great importance for a

better understanding of the phenomena in T Tau stars. Moreover, the acquisition of highly-accurate flux-calibrated spectra is very difficult, hence contemporaneous photometric observations are crucial to overcome this problem, allowing the conversion of emission-line equivalent widths into line fluxes. Furthermore, narrow-band photometry is a very efficient tool to obtain information on line fluxes even in the case of very faint objects, in which spectroscopic observations are normally not possible or would provide very low signal-to-noise useless data.

In this paper we present a study of six T Tau stars in the Taurus-Auriga complex, which is based on simultaneous *BVRI* and narrow-band $H\alpha$ photometry. The photometry, gathered with a 0.9-m telescope, has been contemporaneously acquired to intermediate-resolution spectra that include the $H\alpha$ line. We present and discuss the characteristics and performances of the $H\alpha$ photometric system that can be used for larger observational campaigns.

The spectra are used not only to check and calibrate the $H\alpha$ photometry, but also to derive physical and accretion parameters for these sources.

2. Observations and Data Reduction

We selected a small sample of six bright young stars in the nearby star forming region Taurus-Auriga ($d \simeq 140$ pc, Kenyon *et al.* 1994). Table 1 lists the relevant information reported in the literature regarding their spectral type (SpType), which ranges from G2 to M0, their maximum *V* magnitude (in the range 9–12 mag), and other stellar parameters.

Table 1
Properties of the observed T Tau stars

Name	SpType	V_{\max} [mag]	Period [d]	Distance [pc]	Mass ^(a) [M_{\odot}]	Radius ^(a) [R_{\odot}]
T Tau	K0 IV–Ve	9.8 ^(b)	2.81 ^(f)	139.1 ⁽ⁱ⁾	1.99	3.73
RY Tau	K1 IV–Ve	10.0 ^(c)	5.6 ^(f)	176.7 ⁽ⁱ⁾	2.11	3.10
V773 Tau	K3 Ve	10.7 ^(b)	3.08 ^(g)	138.1 ⁽ⁱ⁾	0.98	3.03
FM Tau	M0–2e	12.0 ^(d)	3.07 ^(h)	140 ⁽ⁱ⁾	0.15	0.97
DG Tau	K6 Ve	11.8 ^(e)	6.3 ^(f)	140 ⁽ⁱ⁾	0.75	1.58
SU Aur	G2IIIine	8.9 ^(c)	2.66 ^(f)	142.4 ⁽ⁱ⁾	2.07	2.72

^(a)Herczeg and Hillenbrand (2014), ^(b)ESA (1997), ^(c)Herbig *et al.* (1988),

^(d)Alfonso-Garzón *et al.* (2012), ^(e)Hauck *et al.* (1990),

^(f)Watson *et al.* (2015), ^(g)Norton *et al.* (2007), ^(h)Davies *et al.* (2014),

⁽ⁱ⁾Gaia Collaboration (2016), ⁽ⁱ⁾Kenyon *et al.* (1994).

2.1. Photometry

The photometric observations were carried out at the M.G. Fracastoro station (Serra La Nave, Mt. Etna, 1750 m altitude) of the Osservatorio Astrofisico di Catania (OAC – Italy) from 19 to 24 November 2014. We used the the imaging camera with the 91-cm telescope and a set of broad-band filters (Johnson-Cousins *B*, *V*, *R*, *I*) as well as two narrow band filters. The CCD camera adopts a Kodak KAF1001E $1k \times 1k$ chip that, with the focal reducer, covers a field of view of about $13' \times 13'$.[†]

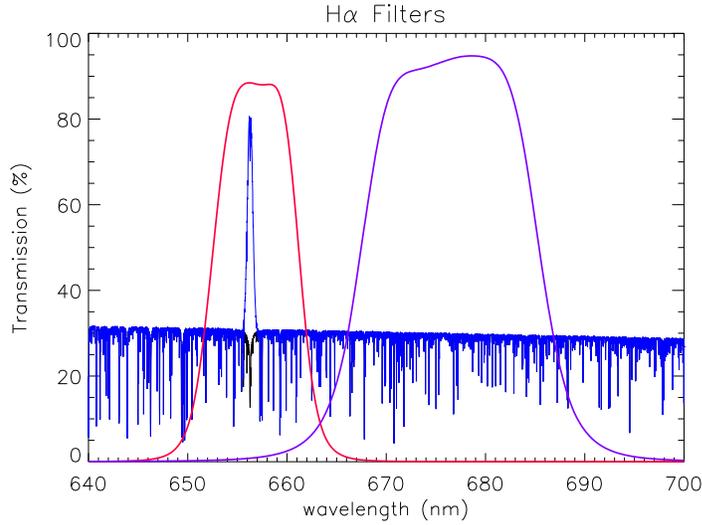


Fig. 1. Passbands of the $H\alpha$ filters (red for $H\alpha_9$ and purple for $H\alpha_{18}$, respectively) superimposed to a synthetic BT-Settl spectrum with $T_{\text{eff}} = 5000$ K and $\log g = 4.0$ (black line). We also overplot a simulated CTT spectrum (blue line) that is obtained adding a Gaussian with FWHM = 6 Å and an area EW = 12 Å to the same synthetic spectrum.

The two narrow-band filters, $H\alpha_9$ and $H\alpha_{18}$, were manufactured by ASAHI SPECTRA[‡]. They have central wavelengths of 656.8 nm and 676.4 nm and full-width at half maximum (FWHM) of 9 nm and 18 nm, respectively. The first filter is centered at the $H\alpha$ wavelength, while the second one collects radiation from a continuum region adjacent to the $H\alpha$. The passbands of these two filters are shown in Fig. 1 overplotted to a BT-Settl synthetic spectrum (Allard *et al.* 2012) with a solar iron abundance, $T_{\text{eff}} = 5000$ K, and $\log g = 4.0$. Fig. 1 also shows a template spectrum that is obtained by summing a Gaussian $H\alpha$ emission profile with FWHM = 6 Å (274 km/s) and equivalent width EW $_{H\alpha} = 12$ Å to the same BT-Settl spectrum.

The exposure times ranged from 2 s to 60 s depending on the filter and the star brightness. For each target, we typically acquired one or two sequences of *BVRI* exposures per night followed by a sequence of three exposures $H\alpha_{18} - H\alpha_9 - H\alpha_{18}$.

[†]http://sln.oact.inaf.it/sln_old/dmdocuments/ccd91rappint2-07.pdf

[‡]<http://www.asahi-spectra.com/>

Data reduction was carried out following standard steps of overscan region subtraction, master-bias subtraction, and division by average twilight flat-field images. The magnitudes were measured on the corrected images through aperture photometry performed with ad-hoc IDL procedures based on DAOPHOT. We extracted the instrumental magnitudes using apertures with a different radius. To minimize the effects of seeing variations from one night to another (in the range of $1''.5$ – $3''.0$) we have chosen the aperture with a radius of 8 pixels ($6''$). The photometric errors related to the photon statistics are in the range of 0.001 – 0.005 mag in all bands for all the targets with the exception of FM Tau, for which they are of about 0.01 mag. Indeed, this target is located in the field of the brighter component V773 Tau that required shorter exposures to avoid saturation. A more reliable evaluation of the precision of our measurements is provided by the standard deviation of the differential magnitudes between two non-variable stars in the target's field. This value is in the range of 0.010 – 0.025 mag for stars with a similar, or slightly fainter, brightness as the targets (see Fig. 2).

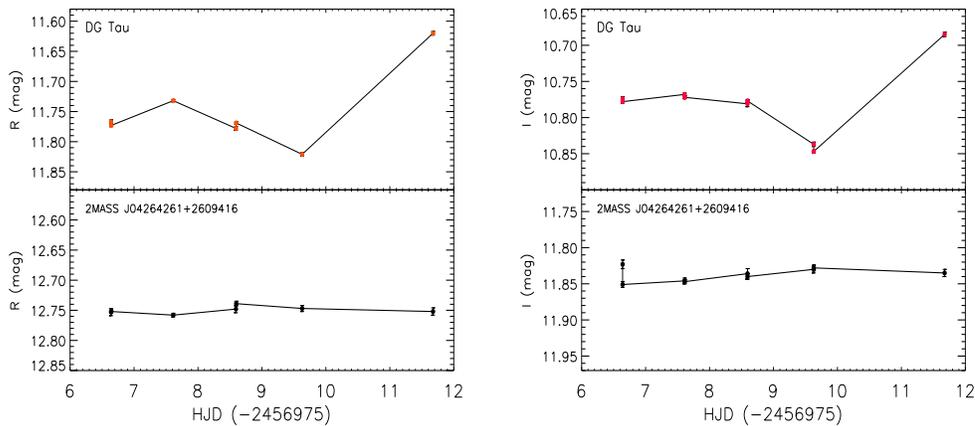


Fig. 2. R and I light curves of DG Tau along with the contemporaneous photometry of a non-variable star (2MASS J04264261+2609416) in the same field of view.

For the broad-band photometry, standard stars in the cluster NGC 7790 (Stetson 2000) were observed to calculate the transformation coefficients to the Johnson-Cousins system and the zero points in the B , V , R , I bands. Local standard stars in the same fields of the targets or near them were also used to derive the zero points. The errors of the broad-band magnitudes include (summed in quadrature) the photon-noise errors and the zero-point errors.

As regards the narrow-band photometry, we measured the magnitude difference $CI_{H\alpha} = H\alpha_{18} - H\alpha_9$, which is related to the intensity of the $H\alpha$ emission (see Section 3.1). This index is analogous to the $(r' - H\alpha)$ color used in the IPHAS survey (Drew *et al.* 2005). Whenever we acquired exposures in the $H\alpha$ continuum ($H\alpha_{18}$) before and after the $H\alpha_9$ frame, we have interpolated the $H\alpha_{18}$ magnitude at the time of the $H\alpha_9$ exposure to calculate $CI_{H\alpha}$. The error of $CI_{H\alpha}$ was calculated,

per each star in the CCD field, by the quadratic sum of the photon-noise errors in the two bands. When we have two $H\alpha_{18}$ exposures before and after the $H\alpha_9$ one we also evaluated the standard deviation of the $H\alpha_{18}$ instrumental magnitudes and took this value as a robust error estimate whenever it was larger than the previous one. The measure of $CI_{H\alpha}$ for the non variable stars in the field of our targets allows us to have an independent estimate of the precision of our measurements, which turns out to be of about 0.02 mag for stars with a brightness comparable to our targets (see Fig. 3 for an example).

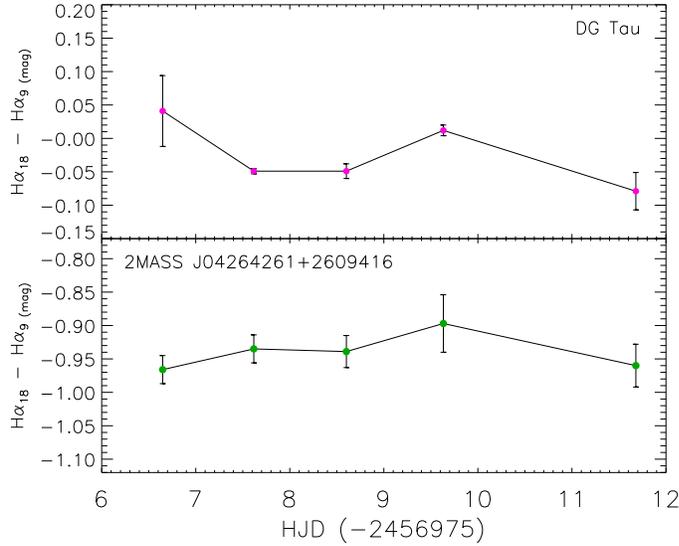


Fig. 3. $H\alpha$ index for DG Tau and the non-variable star 2MASS J04264261+2609416 as a function of time.

2.2. Spectroscopy

Spectroscopic observations were conducted at the 2.5-m Isaac Newton Telescope (INT) with the Intermediate Dispersion Spectrograph (IDS) during three nights, on November 14 and December 6–7, 2014. We collected at least two spectra per target during this observing run. We used the 235-mm camera with the RED+2 CCD detector and the H1800V grating, which provides a dispersion of about $0.34 \text{ \AA}/\text{pix}$. The slit width was 0.220 mm, which corresponds to about $1''.2$ on the sky. The spectra were centered at about 6450 \AA and approximately cover the range $6030\text{--}6870 \text{ \AA}$. The resolution is $R \simeq 9200$, as we verified from the FWHM of the emission lines of the calibration spectra. The exposure times ranged from 90 s to 900 s depending on the star magnitude. The signal-to-noise ratio (S/N) per pixel at the continuum was in the range of 30–200. The reduction of the spectra

was made with IRAF[§], subtracting the overscan and bias, and flat-fielding the images. Optimal extraction of the spectra was performed on the pre-reduced images and the spectra were wavelength calibrated using the emission lines of Cu-Ar arc spectra. Given the small slit width (1.''2) and the variable seeing (in the range of 0.''6–1.''6)[¶], we did not try to flux calibrate the spectra. We instead normalized them to the continuum by fitting low-order polynomials.

3. Data Analysis

3.1. Calibration of the H α Photometry

The H α_9 filter is centered very close to the H α wavelength ($\lambda_c = 656.8$ nm), while the H α_{18} band covers a wavelength range that is not affected by the H α wings and is free from strong photospheric absorption lines (see Fig. 1). Therefore, the magnitude difference in these two bands, $CI_{H\alpha} = H\alpha_{18} - H\alpha_9$, is an index sensitive to the intensity of the H α emission.

However, also in the absence of H α emission, this index depends on other atmospheric parameters, like T_{eff} and $\log g$, because spectral features included in one or both H α bands are sensitive to them (*e.g.*, the wings of the H α line for hot stars or molecular bands for cool stars). To investigate the dependency of $CI_{H\alpha}$ on these parameters we have computed synthetic photometry based on BT-Settl spectra (Allard *et al.* 2012). The synthetic magnitudes have been calculated as:

$$H\alpha_9 = \int F(\lambda)R_9(\lambda) d\lambda, \quad H\alpha_{18} = \int F(\lambda)R_{18}(\lambda) d\lambda \quad (1)$$

where $R_9(\lambda)$ and $R_{18}(\lambda)$ are the response functions of the two H α filters and $F(\lambda)$ is the energy flux distribution of a BT-Settl spectrum (Fig. 1). The results of this synthetic photometry are shown in Fig. 4, which displays a weak dependency (0.05 mag peak-to-peak) on T_{eff} and negligible variations with $\log g$ for FGK stars ($4500 \text{ K} \leq T_{\text{eff}} \leq 7000 \text{ K}$). The index $CI_{H\alpha}$ becomes much more sensitive to $\log g$ for stars hotter than 7000 K due to the increasing dependency of the H α wings on surface gravity of A-type stars. The strong T_{eff} dependence for models cooler than 4500 K is the result of TiO and CaH molecular bands.

To calibrate the index $CI_{H\alpha}$ vs. the equivalent width of the H α line, $EW_{H\alpha}$, we have generated synthetic templates with H α emission mimicking T Tau stars. To this purpose, we have added Gaussian emission profiles centered at 6563 Å with increasing intensity to a BT-Settl spectrum of a given effective temperature (see Fig. 1 for an example). The FWHM of the Gaussian was fixed to 6 Å (274 km/s), which corresponds to a full width at 10% of the maximum of about 500 km/s that

[§]IRAF is distributed by the National Optical Astronomy Observatory, which is operated by the Association of the Universities for Research in Astronomy, inc. (AURA) under cooperative agreement with the National Science Foundation.

[¶]<http://www.ing.iac.es/Astronomy/development/seeing/>

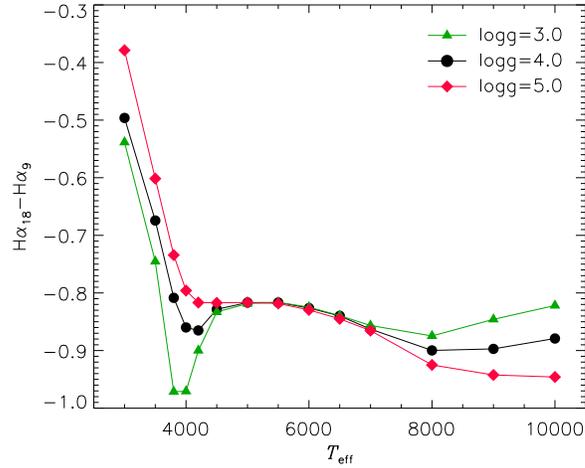


Fig. 4. Dependence of $CI_{H\alpha}$ on T_{eff} and $\log g$ for stars without $H\alpha$ emission, as found from synthetic photometry based on BT-Settl spectra.

is a typical value for the strong accreting CTTs. However, the 90-Å passband of the $H\alpha$ filter makes it almost insensitive to the FWHM of the emission profile, as we have verified with several tests. The calibrations $EW_{H\alpha} - CI_{H\alpha}$, in the T_{eff} range 3000–6000 K, are shown in Fig. 5. We made these calibrations also for different

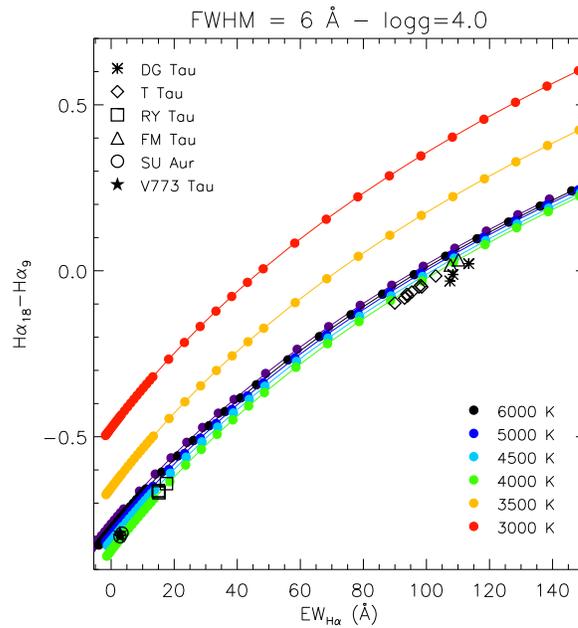


Fig. 5. Calibration relations between the quantity $CI_{H\alpha} = H\alpha_{18} - H\alpha_9$ and the equivalent width, $EW_{H\alpha}$. A different colored dot corresponds to a different value of T_{eff} , as indicated in the legend. The values of $CI_{H\alpha}$ and $EW_{H\alpha}$ measured on the IDS spectra of the targeted T Tau stars (Table 2) are overplotted with different symbols.

values of the gravity, $\log g$, finding that the relations for FGK stars are basically insensitive to $\log g$, as expected on the basis of the results shown in Fig. 4. We note that the calibration does not change appreciably for T_{eff} in the range 4000–6000 K, while the curves at $T_{\text{eff}} = 3500$ K and 3000 K are offset by about +0.2 mag and +0.4 mag, respectively. As discussed above, this T_{eff} dependency is caused by molecular bands encompassed by the $H\alpha_{18}$ filter for low values of T_{eff} . This means that, if we want to convert a measure of $CI_{H\alpha}$ into $EW_{H\alpha}$, we need an independent estimate of T_{eff} that should be accurate for cool objects ($T_{\text{eff}} < 4500$ K).

A key parameter is how sensitive the index $CI_{H\alpha}$ is to $EW_{H\alpha}$. This can be expressed as the derivative of the calibration relations. We find that, for all T_{eff} values, the index $CI_{H\alpha}$ changes by about 0.01 mag for a variation of equivalent width $\Delta EW_{H\alpha} = 1 \text{ \AA}$, for $EW_{H\alpha} < 20 \text{ \AA}$, while its variation decreases to $\approx 0.005 \text{ mag/\AA}$ for $EW_{H\alpha} \approx 100 \text{ \AA}$. With the typical accuracy on $CI_{H\alpha}$ of 0.02 mag (Section 2.1), we should be able to detect $EW_{H\alpha}$ variations at the level of 2–3 \AA .

3.2. Analysis of the IDS Spectra

The IDS spectra cover the wavelength range 6000–7000 \AA almost entirely. This range includes useful diagnostics of chromospheric activity, magnetospheric accretion and jets (*e.g.*, [OI] $\lambda 6300$, $\lambda 6364$, $H\alpha$, HeI $\lambda 6678$, [SII] $\lambda 6716$, $\lambda 6731$) and the lithium absorption line at 6708 \AA . This spectral range is also suitable to determine basic stellar parameters, such as T_{eff} , $\log g$, and $v \sin i$, thanks to the numerous photospheric absorption lines. For this task we used the code ROTFIT, adopting a grid of synthetic BT-Settl spectra (Allard *et al.* 2012) as templates (see Frasca *et al.* 2017 for more details). We analyzed nearly all the spectral range covered by IDS

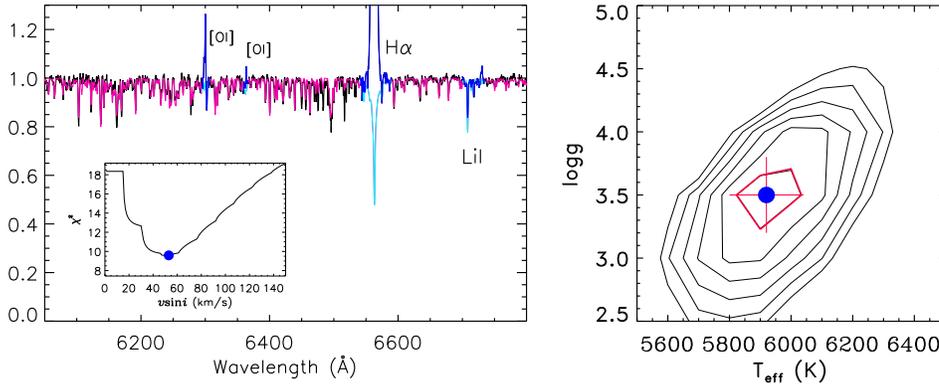


Fig. 6. *Left panel:* Continuum-normalized IDS spectrum of RY Tau (black line) with the best fitting template overlotted (red line). The blue and cyan lines denote the masked spectral regions of the observed and template spectrum, respectively. The inset shows the χ^2 as a function of $v \sin i$, which clearly shows the effect of the spectral sampling ($\approx 15 \text{ km/s}$). *Right panel:* χ^2 contour maps in the $T_{\text{eff}} - \log g$ plane. The blue dot indicates the best values of T_{eff} and $\log g$, while the 1σ confidence level is denoted by the thick red contour. The errorbars on T_{eff} and $\log g$ are also indicated.

with the setup selected by us. We only excluded the low-signal edges, and masked the strongest emission lines and the LiI $\lambda 6708 \text{ \AA}$ absorption. As shown in Frasca *et al.* (2017), this spectral region is well suited for the determination of T_{eff} , $\log g$, and $v \sin i$ for K-type and early M-type stars.

An example of the results of ROTFIT for RY Tau is shown in Fig. 6.

As a result of the strong mass accretion, the spectra of T Tau stars can be affected by veiling, which is better observed at near-UV or optical wavelengths, and sometimes also in the near-IR bands (*e.g.*, Fischer *et al.* 2011, Frasca *et al.* 2017). For this reason, we also left the veiling free to vary in the ROTFIT processing. This was accomplished by adding to the templates a wavelength-independent veiling in steps of 0.1 per each iteration. Noticeable values of the veiling, r_{6500} , were detected for T Tau, DG Tau, and FM Tau, which are also the objects with the strongest H α emission. The χ^2 profile as a function of r_{6500} is shown in Fig. 7 for T Tau and DG Tau.

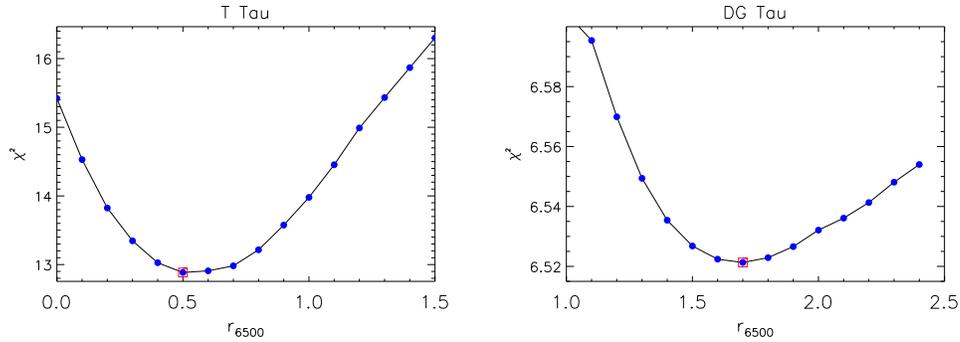


Fig. 7. χ^2 minimization of veiling for T Tau (*left panel*) and DG Tau (*right panel*).

Table 2

Stellar parameters derived from the IDS spectra

Name	T_{eff} [K]	Err	$\log g$ [dex]	Err	$v \sin i$ [km/s]	Err	r_{6500}
T Tau	4570	70	3.20	0.30	19	3	0.5
RY Tau	5920	120	3.50	0.30	52	6	0.0
V773 Tau ^(a)	4460	40	3.30	0.30	70	4	0.0
FM Tau ^(b)	4070	190	4.50	0.90	3.5:
DG Tau	4350	90	3.10	0.60	16	19	1.7
SU Aur	5870	110	3.50	0.20	61	7	0.0

^(a)Multiple system. Uncertain stellar parameters.

^(b)Uncertain stellar parameters and no $v \sin i$ determination because of the low S/N ratio at the continuum.

Table 3

Equivalent widths, $H\alpha$ color index, and accretion parameters derived from the IDS spectra

Name	HJD 2456900+	10% $W_{H\alpha}$ [km/s]	$EW_{H\alpha}$ [Å]	$CI_{H\alpha}$ [mag]	EW_{HeI} [Å]	A_V [mag]	f_{R_0} [erg/cm ² sÅ]	$\log L_{acc}^{H\alpha}$ [L _⊙]	$\log L_{acc}^{HeI}$ [L _⊙]	$\log \dot{M}_{acc}^{H\alpha}$ [M _⊙ /yr]	$\log \dot{M}_{acc}^{HeI}$ [M _⊙ /yr]
T Tau	76.7366	482 ± 23	100.60 ± 0.19	-0.016	0.28 ± 0.04	1.7	1.28(9)e-12	0.49 ± 0.46	0.12 ± 0.46	-6.64	-7.01
T Tau	98.4723	377 ± 8	95.89 ± 0.16	-0.049	0.19 ± 0.05	1.6	1.18(8)e-12	0.42 ± 0.46	-0.14 ± 0.52	-6.70	-7.26
T Tau	98.4761	376 ± 15	95.03 ± 0.23	-0.044	0.24 ± 0.08	1.6	1.18(8)e-12	0.42 ± 0.46	-0.01 ± 0.56	-6.71	-7.14
T Tau	98.4790	377 ± 15	95.27 ± 0.28	-0.049	0.24 ± 0.09	1.6	1.18(8)e-12	0.42 ± 0.46	-0.01 ± 0.58	-6.71	-7.14
T Tau	98.7006	392 ± 8	92.33 ± 0.32	-0.064	0.11 ± 0.03	1.6	1.18(8)e-12	0.40 ± 0.46	-0.43 ± 0.48	-6.72	-7.56
T Tau	98.7057	390 ± 8	89.22 ± 0.54	-0.076	0.15 ± 0.03	1.6	1.18(8)e-12	0.39 ± 0.46	-0.26 ± 0.49	-6.74	-7.39
T Tau	98.7088	377 ± 22	87.11 ± 0.22	-0.097	0.19 ± 0.05	1.6	1.18(8)e-12	0.38 ± 0.46	-0.14 ± 0.52	-6.75	-7.26
T Tau	99.6538	392 ± 8	90.53 ± 0.29	-0.082	0.14 ± 0.02	1.6	1.18(8)e-12	0.39 ± 0.46	-0.30 ± 0.46	-6.73	-7.43
T Tau	99.6594	390 ± 45	91.01 ± 0.26	-0.070	0.11 ± 0.03	1.6	1.18(8)e-12	0.40 ± 0.46	-0.43 ± 0.52	-6.73	-7.56
RY Tau	76.7259	602 ± 8	17.84 ± 0.12	-0.641	0.09 ± 0.02	2.3	1.8(5)e-12	0.05 ± 0.63	-0.04 ± 0.64	-7.18	-7.27
RY Tau	98.5170	633 ± 15	14.74 ± 0.36	-0.669	0.10 ± 0.02	2.4	2.1(6)e-12	0.03 ± 0.63	0.09 ± 0.64	-7.21	-7.14
RY Tau	98.5218	619 ± 15	14.57 ± 0.38	-0.664	0.05 ± 0.04	2.4	2.1(6)e-12	0.02 ± 0.63	-0.28 ± 1.18	-7.21	-7.52
RY Tau	98.5297	634 ± 15	14.74 ± 0.28	-0.662	0.10 ± 0.01	2.4	2.1(6)e-12	0.03 ± 0.63	0.09 ± 0.62	-7.21	-7.14
V773 Tau	99.6959	393 ± 15	3.24 ± 0.10	-0.800	0.04 ± 0.02	2.9	5.0(3)e-13	-1.67 ± 0.48	-1.46 ± 0.67	-8.58	-8.37
V773 Tau	99.7085	410 ± 15	3.33 ± 0.09	-0.790	0.08 ± 0.02	2.9	5.0(3)e-13	-1.66 ± 0.48	-1.08 ± 0.52	-8.57	-7.99
FM Tau	99.6959	376 ± 8	103.81 ± 0.66	0.018	1.53 ± 0.20	0.7	1.20(7)e-14	-1.80 ± 0.48	-1.50 ± 0.47	-8.39	-8.09
FM Tau	99.7085	391 ± 15	103.70 ± 0.88	0.033	1.69 ± 0.18	0.7	1.20(7)e-14	-1.80 ± 0.48	-1.45 ± 0.46	-8.39	-8.04
DG Tau	76.7318	512 ± 15	101.91 ± 0.39	-0.011	0.76 ± 0.09	2.4	3.0(4)e-13	-0.21 ± 0.57	-0.12 ± 0.55	-7.29	-7.19
DG Tau	98.6024	499 ± 15	89.07 ± 1.26	-0.031	1.20 ± 0.28	2.1	2.7(4)e-13	-0.32 ± 0.57	0.08 ± 0.58	-7.40	-6.99
DG Tau	98.6077	515 ± 9	103.60 ± 0.51	0.021	0.95 ± 0.20	2.1	2.7(4)e-13	-0.25 ± 0.57	-0.04 ± 0.57	-7.32	-7.12
SU Aur	76.7415	527 ± 15	3.44 ± 0.06	-0.801	0.07 ± 0.01	2.4	2.4(4)e-12	-0.84 ± 0.53	-0.27 ± 0.52	-8.12	-7.55
SU Aur	98.5475	528 ± 15	3.64 ± 0.16	-0.790	0.01 ± 0.01	1.8	2.3(3)e-12	-0.83 ± 0.52	...	-8.11	...

The stellar parameters were derived with ROTFIT using the average of the best spectra for each target and are listed in Table 2 along with the corresponding errors. We measured the equivalent width of the $H\alpha$ line, $EW_{H\alpha}$, and the full width at 10% of peak height ($10\% W_{H\alpha}$), which are both related to the mass accretion rate. Another useful diagnostic of accretion included in our IDS spectra is the HeI line λ 6678 Å, whose equivalent width, EW_{HeI} , has been measured after the subtraction of the synthetic templates, to remove a nearby iron line. The values of these parameters for each individual spectrum are reported in Table 3 along with the accretion luminosity and mass accretion rate, calculated as described in Section 4.2.

We have also measured the photometric index $CI_{H\alpha}$ on each spectrum by integrating it over the $H\alpha_{18}$ and $H\alpha_9$ passbands, as we did with the synthetic templates (Section 3.1). The $EW_{H\alpha}$ and $CI_{H\alpha}$ values measured on the IDS spectra are overplotted, with a different symbol for each star, to the $EW_{H\alpha}-CI_{H\alpha}$ calibration curves in Fig. 5. We note that these values are superimposed on the curves obtained with the synthetic spectra.

4. Results

4.1. Stellar Parameters

The determination of stellar parameters from the spectra of CTTs is hampered by the presence of strong emission lines and continuum emission arising from the matter accreting onto the central star that overlaps the photospheric spectrum giving rise to the veiling. For this reason one can find very different determinations of effective temperature and spectral type in the literature. This can also be the result of the different spectral ranges and resolutions of the used spectra. Moreover, in some cases the parameters are not directly measured from spectra, but they are derived from photometric indices. For the $v \sin i$, instead, it is normally found a better consistency among the various values in the literature.

As an example, among the catalogs available in the VizieR Service for Astronomical Catalogues at Centre de Données astronomiques de Strasbourg (CDS), a SpType ranging from F8 V to K1 IV and T_{eff} in the range 2600–6890 K are listed for T Tau. The values $T_{\text{eff}} = 4750$ K derived by Wright *et al.* (2003) from a SpType– T_{eff} calibration and $T_{\text{eff}} = 4817$ K reported by Huber *et al.* (2016) are in marginal agreement with our determination of $T_{\text{eff}} = 4570 \pm 70$ K. We note that the inclusion of veiling among the parameters derived with our code simultaneously to T_{eff} , $\log g$, and $v \sin i$ allows an unbiased determination of these parameters. The effect of the veiling is to fill the core of the absorption lines, mimicking a higher photospheric temperature, so that, when it is taken into account, lower T_{eff} values are found. The only determination of surface gravity available in the CDS is $\log g = 2.73$ by Huber *et al.* (2016). We note that T Tau has two companions (T Tau Sa and Sb) with a separation of about $0.''7$ (Dyck *et al.* 1982, Kasper *et al.* 2016) which are so weak ($\Delta J > 10$ mag, $\Delta H \approx 5.5$ mag, $\Delta K \approx 2.2$ mag

for T Tau Sa, Kasper *et al.* 2016) that they do not affect either the determination of its atmospheric parameters or its optical photometry. The values of $v \sin i$ of 23.0 ± 1.2 km/s (Nguyen *et al.* 2012), 20.1 ± 1.2 km/s (Hartmann *et al.* 1986), and 20.7 km/s (Głębocki and Gnaciński 2005) are all in fairly good agreement with each other and with our value of 19 ± 3 km/s. A low value of veiling at 7500 \AA , $r_{7500} = 0.1$, is reported by Herczeg and Hillenbrand (2014), which is in line with our slightly larger value, $r_{6500} = 0.5$, at a shorter wavelength.

For RY Tau, our determination of $T_{\text{eff}} = 5920 \pm 120$ K is in very good agreement with the value of $T_{\text{eff}} = 5945$ K reported by Kim *et al.* (2013), which is based on Spitzer low-resolution IR spectroscopy. It is also in line with the spectral classification G1–2 IV by Petrov *et al.* (1999) and G0 by Herczeg and Hillenbrand (2014). However, very different values of T_{eff} , ranging from 4920 K to 5080 K and based on SpType– T_{eff} calibrations, are also reported in previous works (*e.g.*, Kenyon and Hartmann 1995, Davies *et al.* 2014). Our $v \sin i = 52 \pm 6$ km/s agrees very well with the literature values of 55 ± 3 km/s (Mora *et al.* 2001), 52 ± 2 km/s (Petrov *et al.* 1999), and 48.7 ± 3.8 km/s (Hartmann *et al.* 1986).

The stellar parameters of V773 Tau must be taken with care, because this object is indeed a multiple system composed of three stars, separated by 50–110 mas (A–B) and about 250 mas (A–C) (*e.g.*, Boden *et al.* 2012, and references therein). The brightest star (A) is also a double-lined spectroscopic binary with components of SpType K2 and K5 (Welty 1995) and an orbital period of about 51 d. Our value of $v \sin i = 70$ km/s is larger than the values of about 41 km/s reported by Welty (1995) for each component. It is likely an effect of the blended lines of the individual components, which are not resolved with the intermediate resolution of the IDS spectrograph.

For FW Tau the accuracy of the stellar parameters is low, mainly because of the low signal-to-noise ratio at the continuum, $S/N \approx 30$, and the strong veiling.

In addition to T Tau, the objects with a veiling significantly larger than zero are FM Tau and DG Tau. Our value of $r_{6500} = 3.5$ for FM Tau is quite uncertain, due to the low S/N ratio of our IDS spectra. Values of r_{7500} in the range 0.45–0.72 are reported by Herczeg and Hillenbrand (2014). For DG Tau they quote $r_{7500} = 0.43$ –0.66 and Fischer *et al.* (2011) measured the veiling at different wavelengths with $r_{6500} \approx 1.0$, which is smaller than the value of 1.7 measured by us. They derived larger veiling values, up to about 2.0, only at wavelengths shorter than 5000 \AA . These differences can either arise from the different methods used to measure the veiling or from a real variation due to a variable mass accretion rate.

A noteworthy result is the determination of surface gravity that, excluding FM Tau, is in the range $3.1 < \log g < 3.5$. As shown in the $\log g$ – T_{eff} diagram (Fig. 8), where the pre-main sequence (PMS) isochrones of Siess *et al.* (2000) are displayed with blue lines, the position of our targets (red dots) suggests ages in the range of 0.3–3 Myr, which is in line with the values from the literature (*e.g.*, Kenyon *et al.* 1990, Briceño *et al.* 1999).

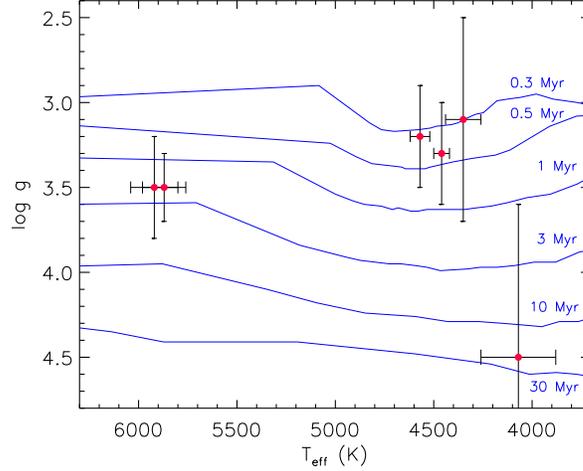


Fig. 8. $T_{\text{eff}}-\log g$ diagram. The solid blue lines are the PMS isochrones with a solar metallicity from Siess *et al.* (2000) at 0.3 Myr, 0.5 Myr, 1 Myr, 3 Myr, 10 Myr, and 30 Myr.

4.2. Accretion Parameters

The availability of optical photometry nearly contemporaneous to the IDS spectra allowed us to obtain an absolute flux scale for the IDS spectra. This permits to convert the equivalent widths of the emission lines into fluxes and luminosities. In particular, from the spectral energy distribution in the *BVRI* bands and the stellar parameters reported in Table 2, we evaluated the extinction, A_V , and the flux at Earth in the *R* band corrected for the extinction, f_{R_0} , which are listed in Table 4. We found for all sources values of A_V larger than those reported by Herczeg and Hilenbrabd (2014) but smaller or equal to those derived from near-IR analyses (*e.g.*, Rebull *et al.* 2010, Furlan *et al.* 2011). Our values agree reasonably well with those reported by López-Martínez *et al.* (2015).

Then the line luminosity, L_{line} , can be calculated as

$$L_{\text{line}} = 4\pi d^2 f_{R_0} EW_{\text{line}} \quad (2)$$

where d is the distance to the star (Table 1). We adopted the f_{R_0} value from the OAC photometry closest in time to the IDS spectra.

We used the relations between the accretion luminosity, L_{acc} , and L_{line} proposed by Alcalá *et al.* (2017, Table B.1) to derive L_{acc} from the two accretion diagnostics in the IDS spectra, namely $\text{H}\alpha$ and $\text{HeI } \lambda 6678 \text{ \AA}$. These two sets of L_{acc} values are reported in Table 3. The L_{acc} errors include the errors on L_{line} , which are dominated by the f_{R_0} error (from 7 to 30%), the error of the distance (from 4% to 15%), and the standard deviation of the $L_{\text{acc}}-L_{\text{line}}$ relations reported by Alcalá *et al.* (2017). We find a marginal agreement between the L_{acc} values derived from $\text{H}\alpha$ and HeI lines, with differences of 0.5 dex, at most. These differences can be ascribed to both the low intensity of the HeI line, which may yield

less accurate results, and to the contribution of different phenomena giving rise to the H α emission. It has been shown, indeed, that in accreting stars the Balmer emission often displays a relevant optically-thin component which arises from the accretion funnels and overlaps the emission produced by the accretion shock over the stellar photosphere (*e.g.*, Alcalá *et al.* 2017, Frasca *et al.* 2017, and references therein).

The accretion luminosities were converted into mass accretion rates, \dot{M}_{acc} , using the relation

$$\dot{M}_{\text{acc}} = \left(1 - \frac{R_{\star}}{R_{\text{in}}}\right)^{-1} \frac{L_{\text{acc}} R_{\star}}{GM_{\star}} \approx 1.25 \frac{L_{\text{acc}} R_{\star}}{GM_{\star}} \quad (3)$$

where R_{\star} and R_{in} are the stellar radius and inner-disk radius, respectively, and M_{\star} is the star mass (Gullbring *et al.* 1998, Hartmann *et al.* 1998). Here we assumed R_{in} to be $5R_{\star}$, as in Alcalá *et al.* (2017) and previous studies. The R_{\star} and M_{\star} values were taken from Herczeg and Hillenbrand (2014) and are also reported in Table 1. The \dot{M}_{acc} values, calculated both from H α and HeI, are also listed in Table 3. We estimate an average uncertainty in $\dot{M}_{\text{acc}}^{\text{H}\alpha}$ of about 0.5 dex.

Our values of L_{acc} and \dot{M}_{acc} agree reasonably well with the literature values, with the exception of T Tau for which we found values of $\log \dot{M}_{\text{acc}}^{\text{H}\alpha} \approx 0.5-0.8$ dex larger than those from the literature. The accretion rates we derive for T Tau from the HeI $\lambda 6678$ line are, on the other hand, in good agreement with the values of $\log \dot{M}_{\text{acc}}$ of $[-7.5, -7.2]$ and -7.34 reported by Calvet *et al.* (2004) and White and Ghez (2001), respectively. For RY Tau, the mass accretion rates we derive are in agreement with the range of values $[-7.2, -7.0]$ quoted by Calvet *et al.* (2004). The same holds for SU Aur, for which Calvet *et al.* (2004) report $\log \dot{M}_{\text{acc}} \simeq -8.2$, and for DG Tau ($\log \dot{M}_{\text{acc}} \simeq -7.34$, White and Ghez 2001). The value of $\log \dot{M}_{\text{acc}}^{\text{H}\alpha} \simeq -8.4$ that we found for FM Tau is in between the values of -8.1 and -8.9 reported by Valenti *et al.* (1993) and White and Ghez (2001), respectively.

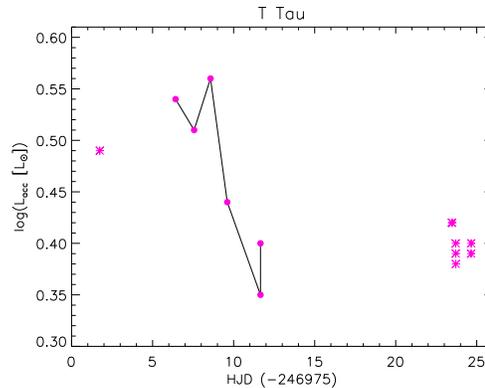


Fig. 9. Accretion luminosity for T Tau calculated from IDS spectra (asterisks) and from OAC photometry (dots) as a function of time.

We point out that using the calibrations between $CI_{H\alpha}$ and $EW_{H\alpha}$ of Section 3.1 and the simultaneous broad-band photometry we can calculate L_{acc} from the photometric data alone, if the distance to the object is known. As an example we show in Fig. 9 the accretion luminosity as a function of time for T Tau.

4.3. Broad-Band and $H\alpha$ Variations

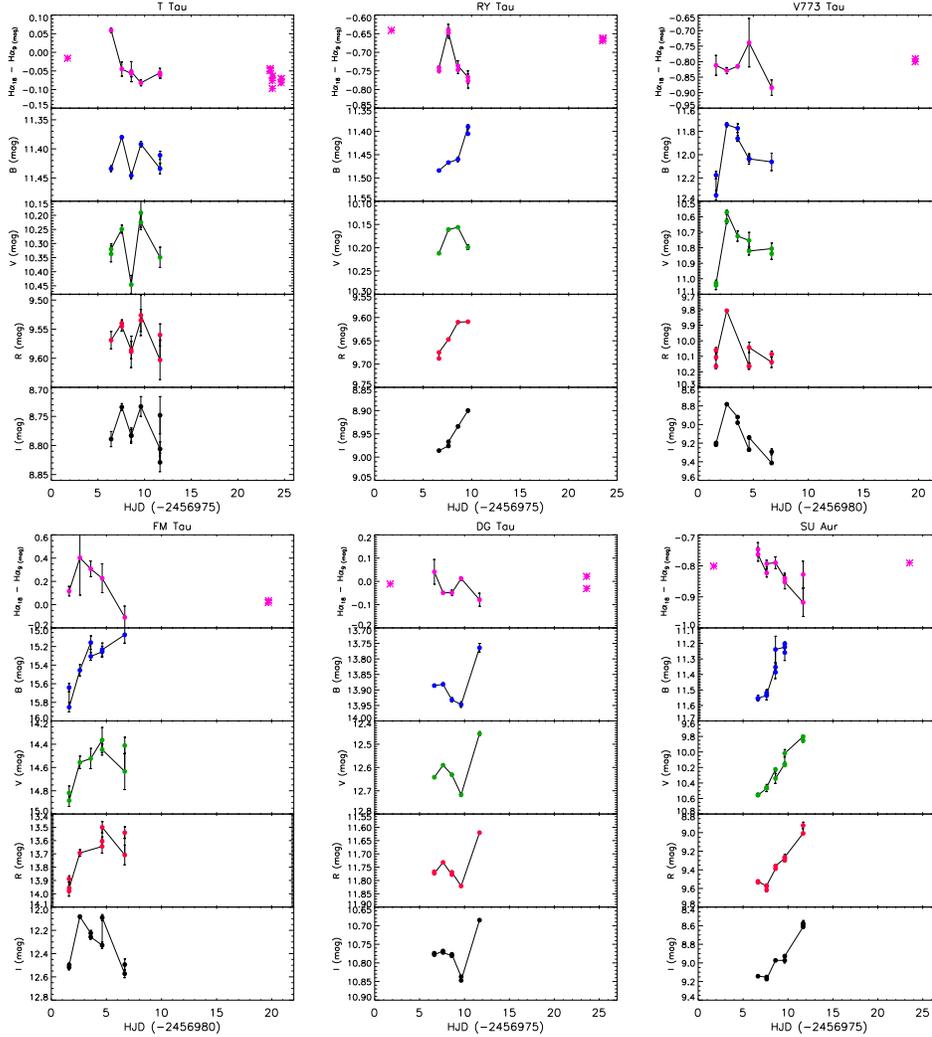


Fig. 10. *From left to right and from top to bottom:* Light curves of T Tau, RY Tau, V773 Tau, FM Tau, DG Tau, and SU Aur. In each box, from the *top* to the *bottom panel*, we plot the color index $CI_{H\alpha}$, B , V , R , and I magnitudes, respectively, as a function of the heliocentric Julian day. In the *upper panel* of each box, the dots represent the color index measured at OAC with the narrow-band filters, while the asterisks refer to the values measured on the IDS spectra.

Despite the short time baseline, we observed significant variations on both broad-band and $H\alpha$ photometry for all the targets (Fig. 10).

T Tau – The broad-band brightness shows low-amplitude (0.05–0.15 mag) variations that could be related to the rotation of the central object ($P_{\text{rot}} \simeq 2.8$ d) with hot (impact shocks) or cool (magnetic activity) spots on its surface. The $\text{H}\alpha$ index does not follow this trend, but displays a decay by about 0.08 mag ($\approx 10 \text{ \AA}$) from the first (JD = 2456976) to the last (JD = 2456999) day of IDS observations.

The rotation period of about 3.0 d was also found by Percy *et al.* (2010) by the analysis of long-term photometric time series. Ismailov *et al.* (2010) found periodicity in the equivalent widths of optical/UV emission lines, especially $\text{H}\beta$ and CaII lines, on a time scale of 33 ± 1.5 d. They note that this is not the rotation period of the star (2.8 d) and the line variation must arise somewhere in the magnetosphere, disk or outflow region.

RY Tau – The brightness increases by about 0.1 mag in *BRI* bands while the $\text{H}\alpha$ index shows an enhancement in the second night of photometric observations (JD = 2456982) by about 0.1 mag ($\Delta EW_{\text{H}\alpha} \approx 10 \text{ \AA}$) followed by a smooth decrease. The values of $CI_{\text{H}\alpha}$ measured on the IDS spectra before and after the photometric observations are close to the maximum photometric value.

The *V* magnitude of about 10.2 mag is near the median value between 1985 and 1995 visible in Fig. 1 of Petrov *et al.* (1999). These authors report also on an intensification of about one magnitude in the *V*-band in October–November 1996, ascribed to a change of opacity of the dusty circumstellar environment.

V773 Tau – This object displays larger *BVRI* variations than the other targets above, from 0.3 mag up to about 0.5 mag peak-to-peak, with a possible variation of $CI_{\text{H}\alpha}$ that does not seem to be correlated with the broad-band variations.

Xiao *et al.* (2012) found a period of 3.075 d, based on TrEs photometry which is the same as that reported by Norton *et al.* (2007) and listed in Table 1. Large night-to-night X-ray variations were found by Guenther *et al.* (2000) during a multi-wavelength campaign of about 10 d. They also found correlation between the strength of $\text{H}\alpha$ and X-ray brightness.

The multiple nature of this object can be the cause of the complex behavior of variations observed in the integrated light of the components. A large photometric variability is reported by Boden *et al.* (2012), who observed variations of more than 2 magnitudes on timescales of hundreds/thousands days in the relative A–B *K*-band photometry. They ascribe these variations to circum-B material resulting in variable line-of-sight extinction.

FM Tau, *DG Tau*, and *SU Aurigae* – Large photometric variations (in the range of 0.3–0.7 mag) have been observed for the remaining three objects, namely *FM Tau*, *DG Tau*, and *SU Aur*. In these cases an anticorrelation between the broad-band and $\text{H}\alpha$ photometry, *i.e.*, the $\text{H}\alpha$ increases as the star becomes fainter, is clearly visible. This is testified by the high values of Spearman’s rank correlation coefficients ρ (Press *et al.* 1992) between these quantities, which are reported in Table 4.

Table 4
Rank correlation coefficients between H α and light curves

Name	H α -B		H α -V		H α -R	
	ρ	σ	ρ	σ	ρ	σ
T Tau	0.000	1.000	0.036	0.938	-0.109	0.816
RY Tau	0.488	0.220	-0.439	0.276	0.487	0.218
V773 Tau	0.200	0.747	0.195	0.750	0.564	0.322
FM Tau	0.300	0.623	0.200	0.747	0.359	0.553
DG Tau	0.616	0.269	0.872	0.054	0.875	0.050
SU Aur	0.896	0.001	0.865	0.003	0.661	0.053

ρ = Rank correlation coefficient.

σ = Two-sided significance of the deviation from zero.

Anticorrelation between the intensity of H α emission and light curves has been frequently observed in chromospherically active stars and it is ascribed to a close spatial association between dark spots in the photosphere and bright active regions in the chromosphere (*e.g.*, Catalano *et al.* 1996, Biazzo *et al.* 2007, Frasca *et al.* 2008). However, the four CTTs that show a high or marginal correlation between H α and broad-band luminosity (RY Tau, FM Tau, DG Tau, and SU Aur) all have shorter rotational periods than the timescales of the variations observed by us and their H α emission profiles are very strong and broad, suggesting mass accretion as their origin. Therefore, the observed anticorrelation is very unlikely to be the result of chromospheric/photospheric surface features, but it could rather arise from a variable accretion. In this hypothesis, the increase in the accretion rate indicated by the H α line is accompanied by a weakening of the source in *BVRI*, caused by the extinction produced by the material falling on the star. The opposite behavior would be observed during a decrease of the accretion activity.

This idea is supported by the correlations between the *V* magnitude and colors, which are shown in Fig. 11 for three sources. It is interesting to note that when the star gets fainter it becomes redder in the *V-I* color, but at the same time it becomes bluer in the *B-V* color. This could mean that when the accretion is stronger there is a greater extinction that shows up as a reddening of the *V-I* color, while the *B-V* color is dominated by the flux excess produced by the impact shock. However, this does not necessarily imply an intrinsic variation of the accretion, because the rotational modulation of accretion-related features (impact shock and accretion funnels) could produce similar color variations. This is probably the case for T Tau, which displays a nice *B-V* and *V-I* anticorrelation even if the broad-band flux varies with the rotational period.

For SU Aur the period of 2.66 d reported by Watson *et al.* (2015) is the same as that reported by Cody *et al.* (2013) on the basis of a 25-d MOST campaign.

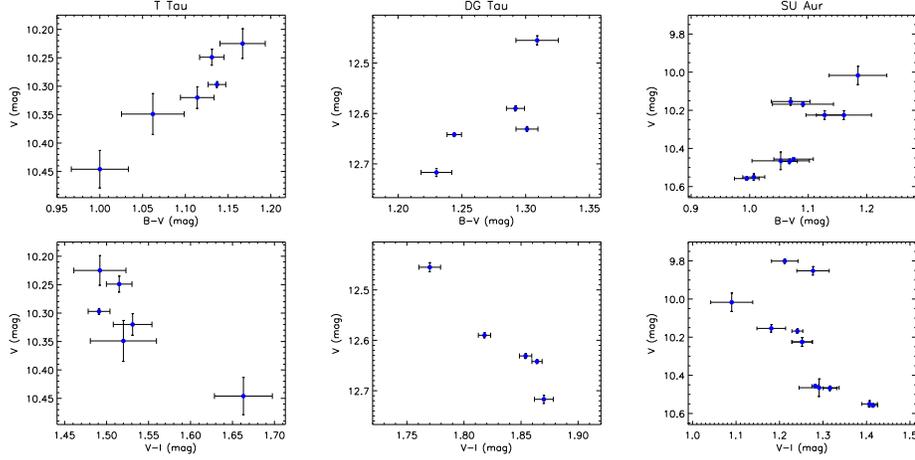


Fig. 11. V magnitude as a function of the $B-V$ (top panels) and $V-I$ (lower panels) for T Tau, DG Tau, and SU Aur.

The MOST data also display a dimming of about 0.4 mag during the last 5–6 d of that observing campaign. Our photometry shows a brightening of ≈ 0.6 mag in a 6-d interval with no hint of rotational modulation. This event is accompanied by a mild simultaneous decrease of the $H\alpha$ emission. The $H\alpha$ emission observed about twenty days later with the IDS seems to be back to the initial high value. Unfortunately, we do not have photometric data up to this date.

5. Discussion and Conclusions

We have presented the results of simultaneous $BVRI$ and narrow-band $H\alpha$ photometry of six T Tau stars in the Taurus-Aurigae complex gathered with the 0.9-m telescope of OAC. Intermediate resolution spectra were also taken with the IDS at the INT telescope nearly contemporaneously (before and/or after) the photometric data and were used both to trace the $H\alpha$ variations on a longer time scale and to derive basic stellar and accretion parameters.

Using synthetic BT-Settl and IDS spectra we carried out a characterization of the narrow-band $H\alpha$ photometric system and obtained calibration relations between the equivalent width $EW_{H\alpha}$ and the color index $CI_{H\alpha}$, also highlighting the dependencies on other parameters such as T_{eff} and $\log g$. In particular, we found that, with the typical photometric accuracy of 0.02 mag for $CI_{H\alpha}$, we can resolve $EW_{H\alpha}$ variations at a level of 2–3 Å.

The analysis of the spectra allowed us to detect a significant veiling only for the sources with the largest $H\alpha$ intensities (T Tau, FM Tau, and DG Tau), as found in previous works (*e.g.*, Fischer *et al.* 2011, Herczeg and Hillenbrand 2014). We also inferred surface gravities in the range of $3.1 < \log g < 3.5$, with the exception of the faintest target, namely FM Tau, for which we found a rather uncertain value

of 4.5, likely due to the low signal-to-noise spectra. These values of gravity are consistent with the very young age of the targets (Fig. 8).

We detected broad-band photometric variations ranging from 0.05 mag to 0.7 mag over the six nights of observations for all the targets. These objects mostly display a redder $V - I$ and a bluer $B - V$ color when they become fainter. Moreover, the sources with the largest variation amplitudes display an anticorrelation between the broad-band magnitudes and $CI_{H\alpha}$, *i.e.*, when the source gets fainter the $H\alpha$ intensity enhances, and *vice versa*.

Despite the rather short time baseline of our photometric and spectroscopic observations (≤ 25 d), the timescales of broad-band and $H\alpha$ variations appear to be longer than the rotational period, with the exception of DG Tau and T Tau. In particular, for the latter source, the broad-band luminosity displays a clear modulation with the rotation period, while the behavior of the $B - V$ and $V - I$ colors seems to be related to accretion features. Since the strong and broad $H\alpha$ profiles of these sources are basically due to accretion, we think that chromospheric plages can be ruled out as the source of the observed anticorrelation between broad-band photometry and $H\alpha$ intensity. The most plausible explanation is an intrinsic variation in the accretion rate, at least for the objects with timescales of variations longer than the rotation period (FM Tau and SU Aur), while a geometric effect can play a major role when the variation timescales are close to the rotation period.

Dynamical instability in the magnetosphere was invoked by Bouvier *et al.* (2003) to explain the photometric and spectral-line variations observed on AA Tau, while the line and continuum variations of V2129 Oph were ascribed by Alencar *et al.* (2012) to rotational modulation of hot and cold spots plus an episodic increase of the mass accretion rate. The changes of mass accretion rates observed spectroscopically by Costigan *et al.* (2014) on a sample of Herbig Ae and T Tau stars (including four of the sources studied in the present paper) are mainly ascribed to a rotational modulation of the accretion-related emission, since the timescales of the variations are comparable with their rotation periods. Periodic $H\alpha$ variations were observed by Sousa *et al.* (2016) for 8 out of 58 CTTs in the young cluster NGC 2264, while the remaining do not display any clear periodicity in the line variations. Three of these sources are AA Tau-like objects (dippers), for which the dips in the CoRoT light curves are ascribed to periodic occultations of the central star by a wrapped disk and the $H\alpha$ variability is likely arising from the geometry of the accretion columns that project over the impact region at particular rotational phases. Other sources display instead CoRoT light curves dominated by a spot-like rotational modulation. These cases represent a “stable” accretion regime. Anyway, several sources do not show any clear periodicity of the light or $H\alpha$ variations and are characterized by stochastic variations and/or bursts. These sources are likely in an “unstable” regime.

Our data show variation timescales longer than the rotation period for at least two sources (FM Tau and SU Aur), suggesting an intrinsic variation of the accre-

tion, while for the remaining sources no clear behavior appears or the variations occur with the rotation period (as for T Tau). Simultaneous photometric/H α observations on a longer time baseline can help to shed light on the different processes responsible for the observed variations and to investigate the relative importance of them.

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