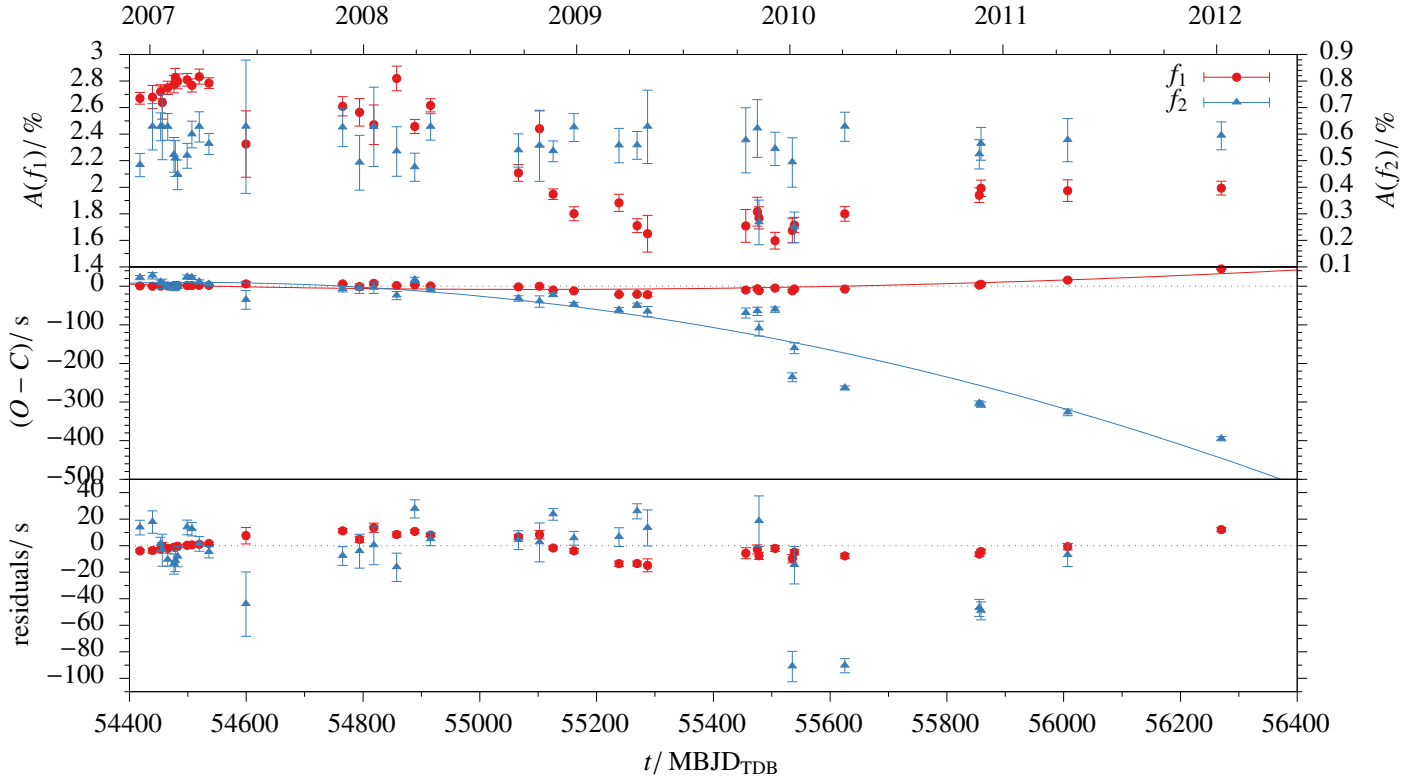
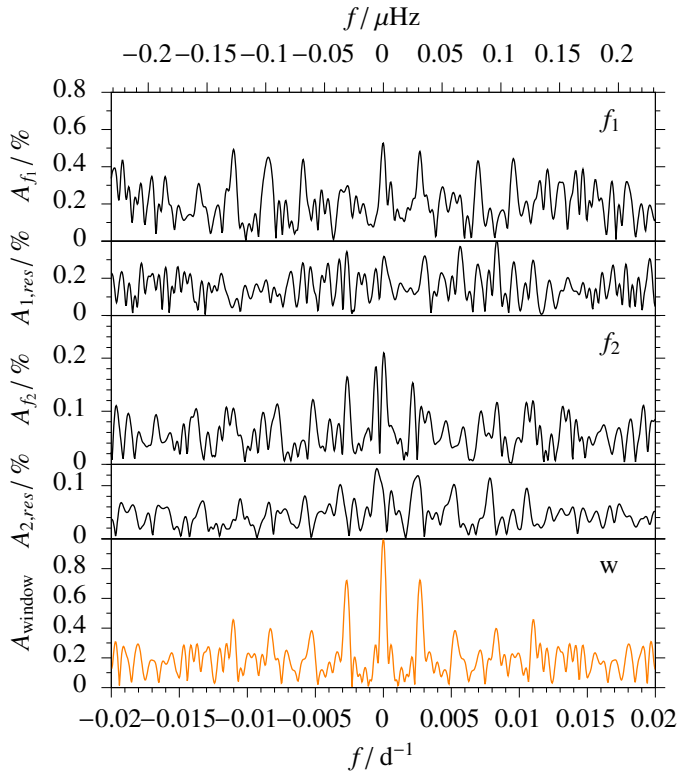




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<b>Authors</b>	Mackebrandt, F.; Schuh, S.; SILVOTTI, Roberto; Kim, S. -L.; Kilkenny, D.; et al.
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**Fig. 8.** Results for the two main pulsations of DW Lyn. *Top panel:* amplitudes. *Middle panel:* fits of the O–C data with second order polynomials in time. *Lower panel:* residuals.



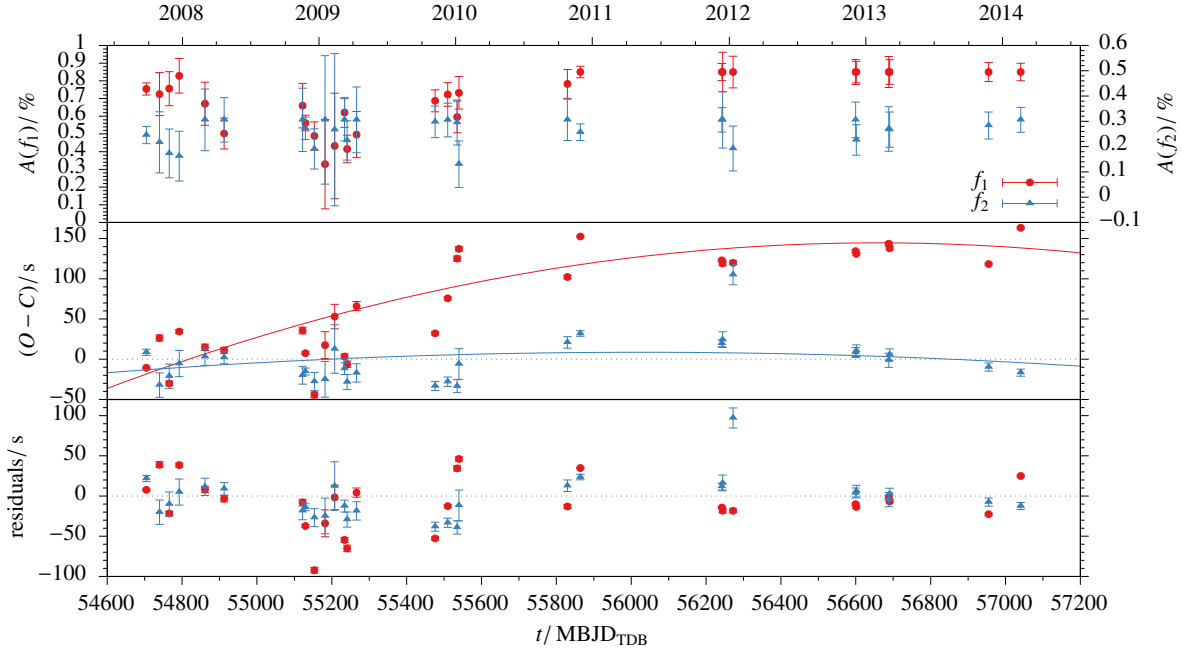
**Fig. 9.** Amplitude spectrum of V1636 Ori of the main pulsation frequency  $f_1 = 631.7346 \text{ d}^{-1}$  (top),  $f_2 = 509.9780 \text{ d}^{-1}$  (middle) with the respective residuals after the pre-whitening below and the normalised window-function (bottom).

$f_3 = 642.0516 \text{ d}^{-1}$  are presented next to  $f_1$  in Fig. 11. Another peak at about  $665 \text{ d}^{-1}$  consists of at least two frequencies at  $664.488549 \text{ d}^{-1}$  and  $665.478133 \text{ d}^{-1}$ , but they are not sufficiently resolvable within the individual epochs, and lead to uncertainties in the O–C analysis that are too large.

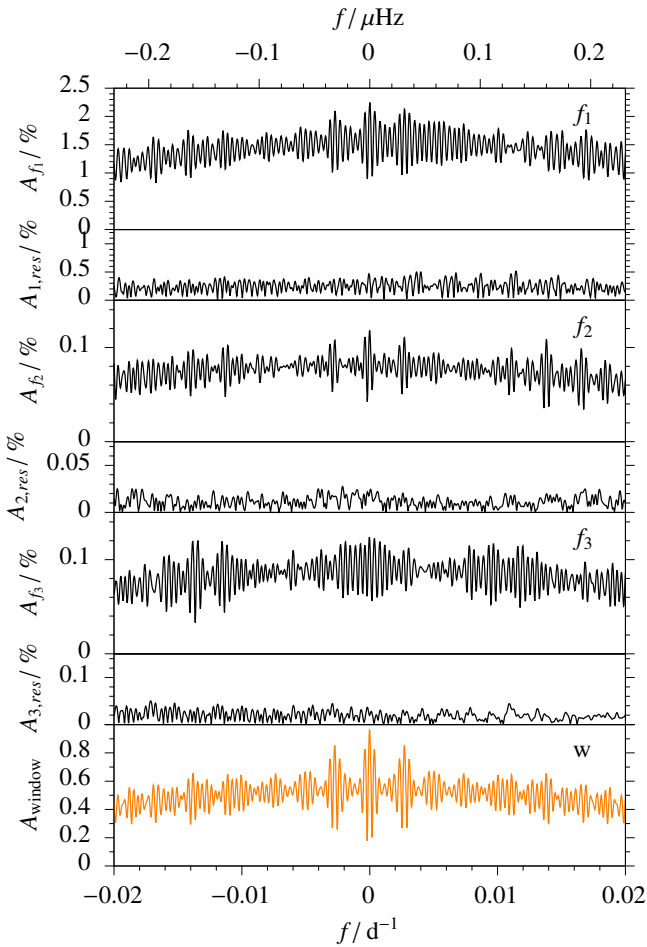
Figure 12 shows the resulting O–C diagram and the amplitudes at different epochs. Due to the large observational gap from 2003 to 2008 with only one block of observations in between, we had difficulties avoiding errors in cycle count. In order to avoid a phase jump, we increased the averaging window for initial phase values to  $q = 6$ . With this set up, the changes in pulsation frequencies read as follows:  $\dot{P}/P_{f1} = (1.7 \pm 1.6) \times 10^{-7} \text{ d}^{-1}$ ,  $\dot{P}/P_{f2} = (2.4 \pm 0.4) \times 10^{-5} \text{ d}^{-1}$  and  $\dot{P}/P_{f3} = (4.0 \pm 0.5) \times 10^{-6} \text{ d}^{-1}$ . While  $f_2$  and  $f_3$  show no significant variation of pulsation amplitude,  $f_1$  varies by 1.5 per cent (amplitude) or 50 per cent (relative). Thus, the corresponding phase changes should be interpreted with caution. Charpinet et al. (2006) identified the radial order  $k$  and degree  $l$  from asteroseismic modelling to be  $f_1: l = 2, k = 2$ ;  $f_2: l = 4, k = 1$ ;  $f_3: l = 3, k = 2$ . These combinations do not allow a direct comparison of our  $\dot{P}$  measurements to the model calculations from Charpinet et al. (2002), but the sign of  $\dot{P}$  indicates QQ Vir to be in the stage of He burning.

#### 4.4. V541 Hya

The amplitude spectrum in Fig. B.4 shows two pulsation modes with frequencies at  $f_1 = 635.32218 \text{ d}^{-1}$  and at  $f_2 = 571.28556 \text{ d}^{-1}$ . Both of them show a complex behaviour (Fig. 13), indicating unresolved multiplets and/or frequency changes that we see also in the O–C diagrams (Fig. 14). The



**Fig. 10.** Results for the two main pulsations of V1636 Ori. *Top panel:* amplitudes. *Middle panel:* fits of the O–C data with second order polynomials in time. *Lower panel:* residuals.



**Fig. 11.** Amplitude spectrum of QQ Vir of the main pulsation frequency  $f_1 = 626.877628 \text{ d}^{-1}$  (top),  $f_2 = 552.00713 \text{ d}^{-1}$  (top middle),  $f_3 = 642.0516 \text{ d}^{-1}$  (bottom middle) with the respective residuals after the pre-whitening below and the normalised window-function (bottom).

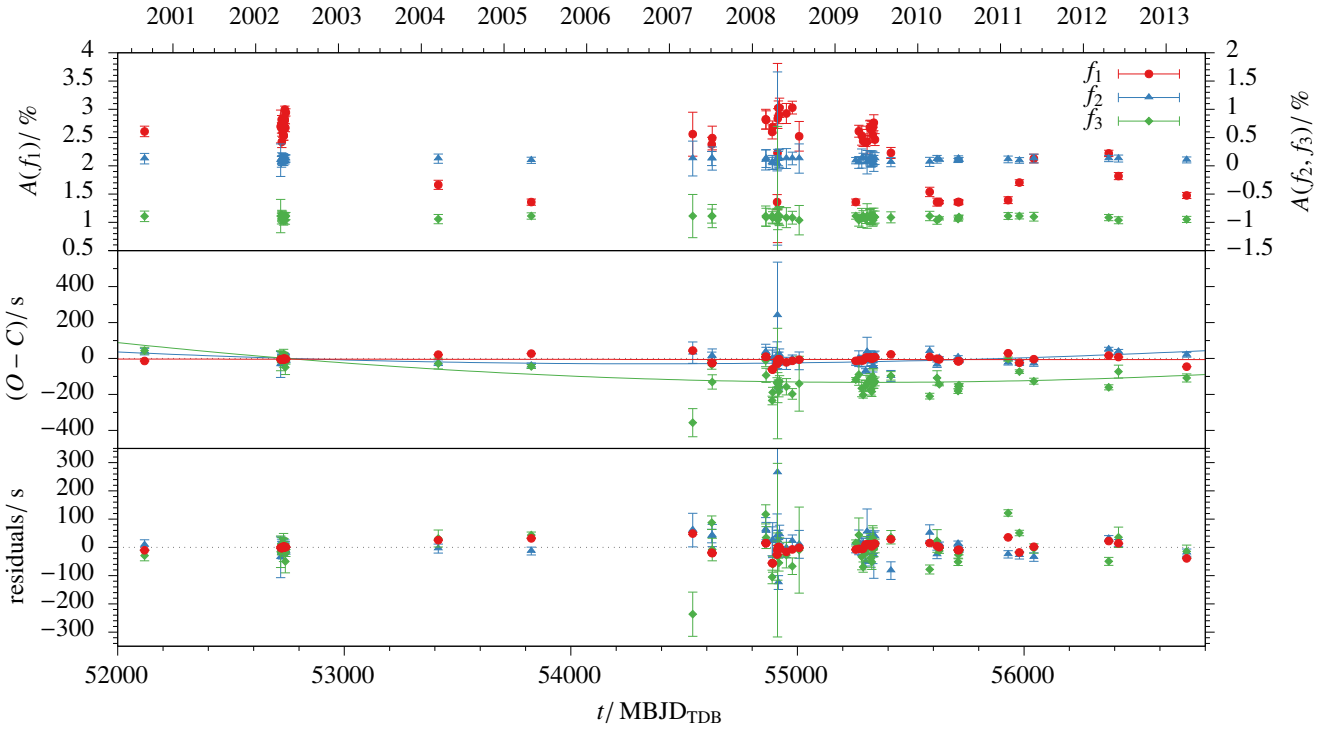
S/N for a third frequency at  $603.88741 \text{ d}^{-1}$  is not sufficient for the O–C analysis. Similar to V1636 Ori, the amplitude spectrum obtained from the TESS light curve in Fig. A.2 shows no evidence for  $g$ -mode pulsations with amplitudes greater than 0.4 per cent.

Randall et al. (2009) speculated about rotational mode splitting for  $f_3$  with  $\Delta f_{3,-} = 5.12 \mu\text{Hz}$  and  $\Delta f_{3,+} = 3.68 \mu\text{Hz}$ . The asteroseismic modelling associates  $f_1$  with a  $l = 0$  mode and  $f_2$  with  $l = 0$  or 1 mode (depending on the favoured model).  $f_3$  corresponds to a  $l = 2$  mode. They caution this interpretation due to their limited resolution in frequency space, the mode splitting could be an unresolved quintuplet. Our data set shows no clear evidence for a mode splitting with  $\Delta f_{3,-} = 5.12 \mu\text{Hz}$  or  $\Delta f_{3,+} = 3.68 \mu\text{Hz}$  (see Fig. 15) but rather a mode splitting for  $f_1$  and  $f_2$  with about  $\Delta f = 0.08 \mu\text{Hz}$  (Fig. 13). Assuming these modes are of degree  $l = 1$ , this could be interpreted as a triplet. But Randall et al. (2009) model these modes with a degree of  $l = 0$ , which does not support a mode splitting into triplets.

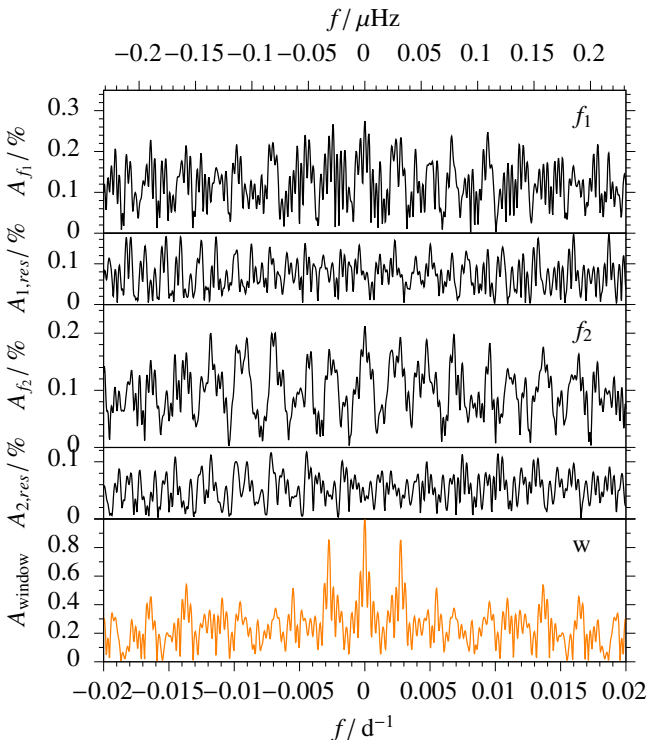
The O–C diagram in Fig. 14 shows the analysis of the two main pulsation modes and the variation of the pulsation amplitudes. The second order fits in time correspond to changes in period of  $\dot{P}/P_{f_1} = (-1.49 \pm 0.11) \times 10^{-5} \text{ d}^{-1}$  and  $\dot{P}/P_{f_2} = (-0.7 \pm 1.5) \times 10^{-5} \text{ d}^{-1}$ . For  $f_2$ , the change in period does not significantly differ from the null hypothesis. Assuming these changes origin from stellar evolution, V541 Hya might just have passed the point of sign change in  $\dot{P}$  and at the beginning of the contraction phase. While the arrival times scatter widely, the amplitudes of both pulsations remain almost constant within the uncertainties. If V541 Hya is in its evolution close to starting the contraction phase, as indicated by a  $\dot{P}$  close to zero, the changes in stellar structure may cancel the strict phase coherence.

#### 4.5. Testing the sub-stellar companion hypothesis

In order to set upper limits to the mass of a companion, we computed a series of synthetic O–C curves for different orbital



**Fig. 12.** Results for the three main pulsations of QQ Vir. *Top panel:* amplitudes.  $f_3$  has a vertical offset of  $-1$  for clarity. *Middle panel:* fits of the O–C data with second order polynomials in time. *Lower panel:* residuals.



**Fig. 13.** Amplitude spectrum of V541 Hya of the main pulsation frequency  $f_1 = 635.32218 \text{ d}^{-1}$  (top),  $f_2 = 571.28556 \text{ d}^{-1}$  (middle) with the respective residuals after the pre-whitening below and the normalised window-function (bottom).

periods and companion masses, assuming circular orbits, and compared these curves with the O–C measurements after subtracting the long-term variations.

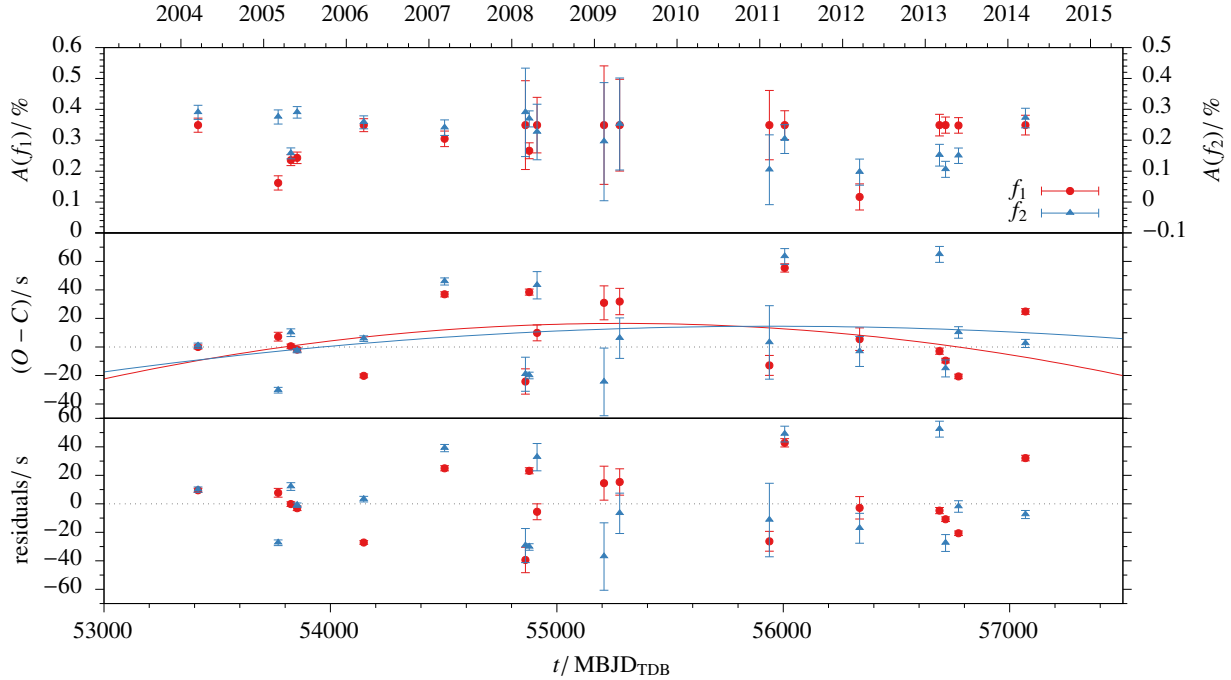
For each synthetic O–C curve, we selected the phase that gives the best fit to the data using a weighted least squares algorithm. For each observational point, we computed the difference, in absolute value and in  $\sigma$  units (where  $\sigma$  is the O–C error), between O–C and the synthetic value. The greyscale in Figs. 16 and 17 corresponds to the mean value of this difference in  $\sigma$  units, which means that the presence of a companion is indicated by a minimum (bright areas) of this parameter. We see that in V1636 Ori, QQ Vir, and V541 Hya, the mean difference for  $f_1$  is always very high, implying that the data are not compatible with a companion. However, these results are limited by the fact that the O–C diagrams of these stars are “contaminated” by other irregular variations, presumably due to other reasons like non-linear interactions between different pulsation modes, for example, and therefore these constraints to the orbital period and mass of a companion must be taken with some caution. For the  $f_2$  and  $f_3$  measurements, the mean difference to the synthetic data is smaller in sigma units (because of the larger uncertainties) and very uniform. The uncertainties of the O–C measurements are not small enough to favour a set of models in the period-mass parameter space.

For  $f_1$  of DW Lyn, there is a significant minimum at about 1450 d ( $\sim 4$  yr) and  $\sim 5 M_{\text{Jup}}^2$ , which is also well visible in the O–C diagram of Fig. 8. This periodicity is not visible in the second frequency  $f_2$  which, however, has much larger error bars due to the much lower amplitude of  $f_2$  with respect to  $f_1$ .

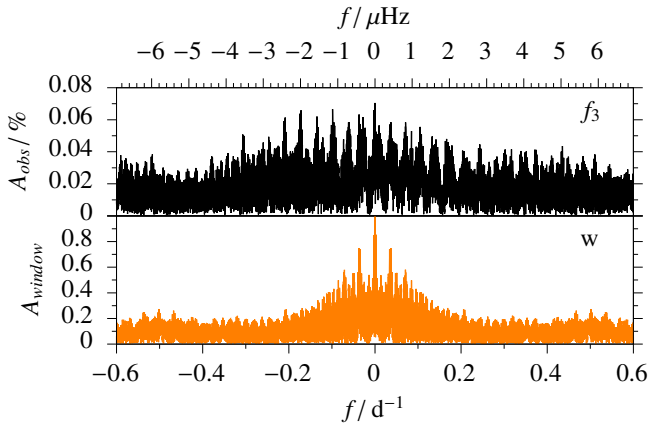
Lutz et al. (2011) described a periodicity at 80 days, detected for  $f_2$ . We can recover this signal, however, with a low significance. This would correspond to a light-travel time amplitude of 4 s (for  $m \sin i \approx 15 M_{\text{Jup}}$ ), which is smaller than the 15 s measured by Lutz et al. (2011). Nevertheless, this signal is not confirmed by  $f_1$ . Thus, we rule out a companion induced signal in the arrival times due to the lack of simultaneous signals in  $f_1$

<sup>2</sup>  $1 M_{\text{Jup}}$  (Jupiter mass) =  $1.899 \times 10^{27}$  kg.





**Fig. 14.** Results for the two main pulsations of V541 Hya. *Top panel:* amplitudes. *Middle panel:* fits of the O–C data with second order polynomials in time. *Lower panel:* residuals.



**Fig. 15.** Amplitude  $A$  spectrum with respect to the pulsation frequency  $f_3 = 603.88741 \text{ d}^{-1}$  of V541 Hya (*top*) and the normalised window function (*bottom*).

and  $f_2$  with similar amplitude. The tentative signal in  $f_2$  is better explained by mode beating, as already described in Sect. 4.1. The variations seen in the first 200 days of the O–C diagram in Fig. 8 correspond to a periodicity of about 80 days and are accompanied by variations in the amplitude of the pulsation.

For V1636 Ori, Lutz (2011) predicted a period at 160 d and amplitude of 12 s. This can not be confirmed as a companion-induced signal. A periodic signal with an amplitude of 6.5 s (for  $m \sin i \approx 15 M_{\odot}$ ) is indicated in the analysis of  $f_1$ , but at a low significance and accompanied by many other signals of similar significance. This periodicity is not confirmed by a significant signal in the measurements of  $f_2$ .

## 5. Summary and conclusion

In this work, we present ground-based multi-site observations for the four sdBs, DW Lyn, V1636 Ori, QQ Vir, and V541 Hya.

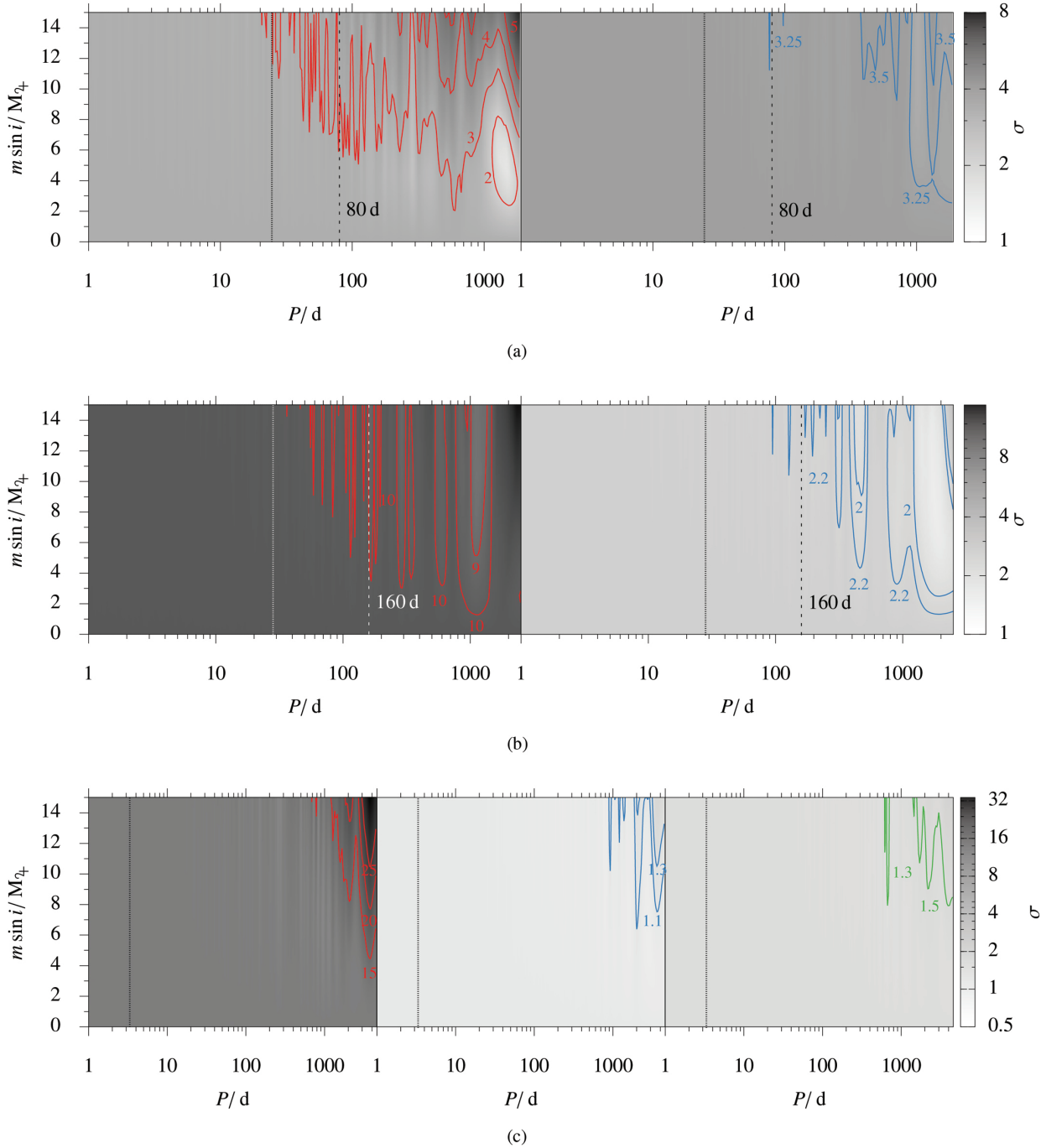
We investigated variations in the arrival times of their dominant stellar pulsation modes to draw conclusions about secular period drifts and possible sub-stellar companions. All light curves are analysed homogeneously.

From the O–C measurements, we derive an evolutionary timescale from the change in period  $\dot{P}$ . Comparing to model calculations from Charpinet et al. (2002), we infer the evolutionary phase of the target. Although some  $\dot{P}$  measurements are influenced by mode splitting, we can tell from the sign of  $\dot{P}_1$  of DW Lyn that the star is likely still in the stage of central He burning. We can draw a similar conclusion from the sign of  $\dot{P}$  of QQ Vir. The  $\dot{P}$  measurements of V1636 Ori are likely affected by mode splitting, making it difficult to interpret the results in the context of stellar evolution. V541 Hya shows  $\dot{P}$  measurements close to zero, which indicates the star being at the transition phase between He burning and contraction due to the depletion of He in the core.

Comparing the atmospheric properties from Table 1 with the evolutionary tracks for different models from Fig. 1 in Charpinet et al. (2002), we can confirm the hypothesis that DW Lyn and QQ Vir are in their He burning phase. V541 Hya agrees within  $2\sigma$  of the  $\log g$  measurement with one model at the turning point between the two evolutionary stages.

However, we can not exclude frequency and amplitude variations on smaller timescales than resolvable by our data set. Using temporally higher resolved *Kepler*-data of KIC 3527751, Zong et al. (2018) cautioned about long-term frequency or phase evolutions ascribing to non-linear amplitude and frequency modulations in pulsating sdBs. We see such effects already in our data set, even with a low temporal resolution compared to the *Kepler* sampling with a duty cycle of more than 90 per cent.

Observations on DW Lyn and V1636 Ori were published by Lutz et al. (2008a, 2011); Schuh et al. (2010); Lutz (2011). Our analysis of these observations, including extended data sets, do not confirm the tentative companion periods of 80 and 160 days,

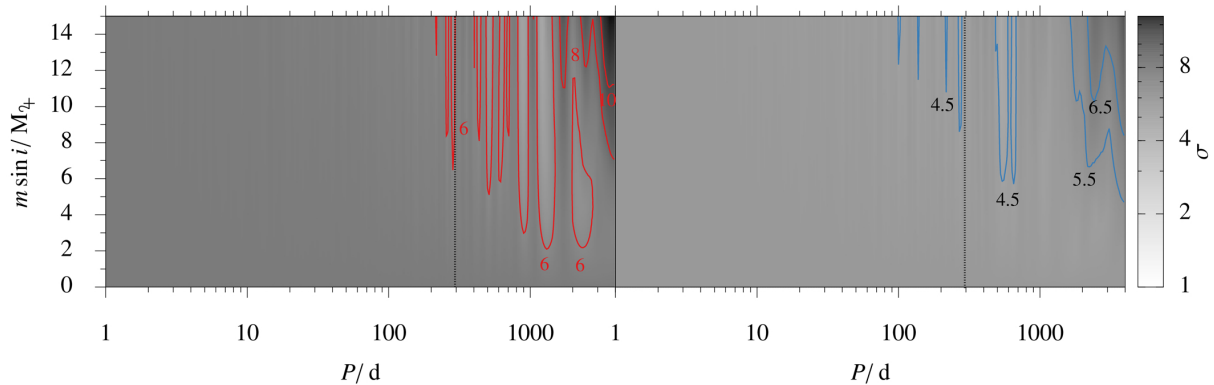


**Fig. 16.** Minimum companion mass as a function of orbital period. Greyscale shows the difference between the O–C measurements and artificial O–C data generated for a given combination of companion mass and orbit. We note that at this stage, the phase optimisation of the artificial data is done independently for each pulsation frequency. The median of gaps in between the epochs is indicated by a vertical dotted line. See text for more details. (a) DW Lyn. Contour lines for  $f_1$  are placed at 2, 3, 4, and 5  $\sigma$  (left panel), and for  $f_2$  at 3.25 and 3.5  $\sigma$  (right panel), as indicated by their labels. The planetary signal proposed by Lutz et al. (2011) at a period of 80 d is indicated as dashed line. (b) V1636 Ori. Contour lines for  $f_1$  are placed at 9 and 10  $\sigma$  (left panel), and for  $f_2$  at 2 and 2.2  $\sigma$  (right panel), as indicated by their labels. The planetary signal proposed by Lutz et al. (2011) at a period of 160 d is indicated as dashed line. (c) QQ Vir. Contour lines for  $f_1$  are placed at 15, 20, and 25  $\sigma$  (left panel), for  $f_2$  at 1.1 and 1.3  $\sigma$  (middle panel), and for  $f_3$  at 1.3 and 1.5  $\sigma$  (right panel), as indicated by their labels.

respectively. These signals more likely arise due to mode beating indicated by partly unresolved frequency multiplets and amplitude modulations.

Almost all analysed pulsation modes show formal significant changes in arrival times, but the amplitudes of these periodic signals do not correlate with frequencies, excluding the light-travel time effect due to orbital reflex motions for such variations and

thus giving upper limits on companion masses. Only DW Lyn might have a planetary companion on a long orbital period, as indicated by one arrival time measurement. But this can not be confirmed with a second measurement, due to larger uncertainties. Additionally, more studies question the presence of already proposed companions, for example, Krzesinski (2015); Hutchens et al. (2017). Our unique sample of long-term observations shows



**Fig. 17.** Continuation of Fig. 16. V541 Hya. Contour lines for  $f_1$  are placed at 6, 8, and  $10\sigma$  (left panel), and for  $f_2$  at 4.5, 5.5, and  $6.5\sigma$  (right panel), as indicated by their labels.

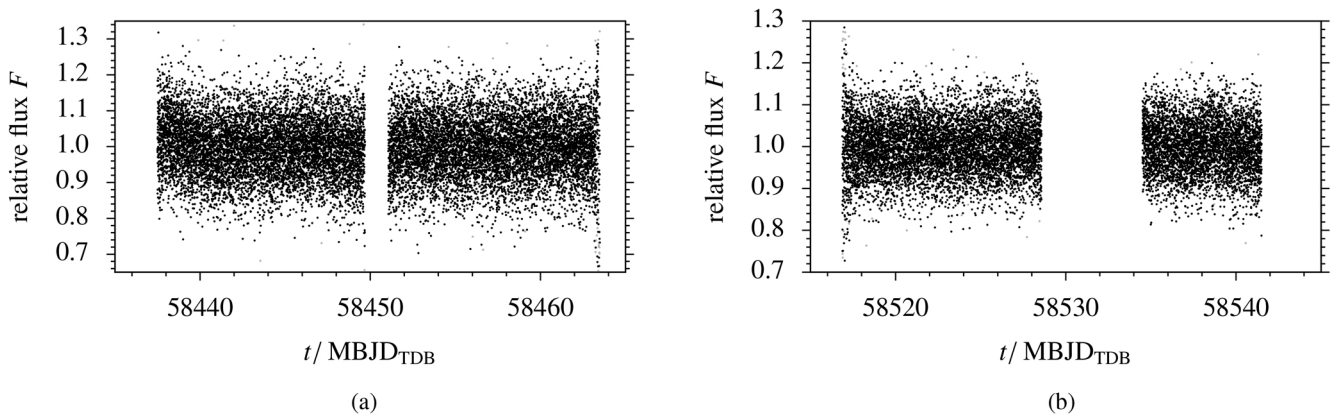
a complex behaviour of mode- and amplitude interactions in sdBs which should be addressed in further studies. Until this has been addressed, caution is advised when interpreting O–C pulse arrival times in terms of companions.

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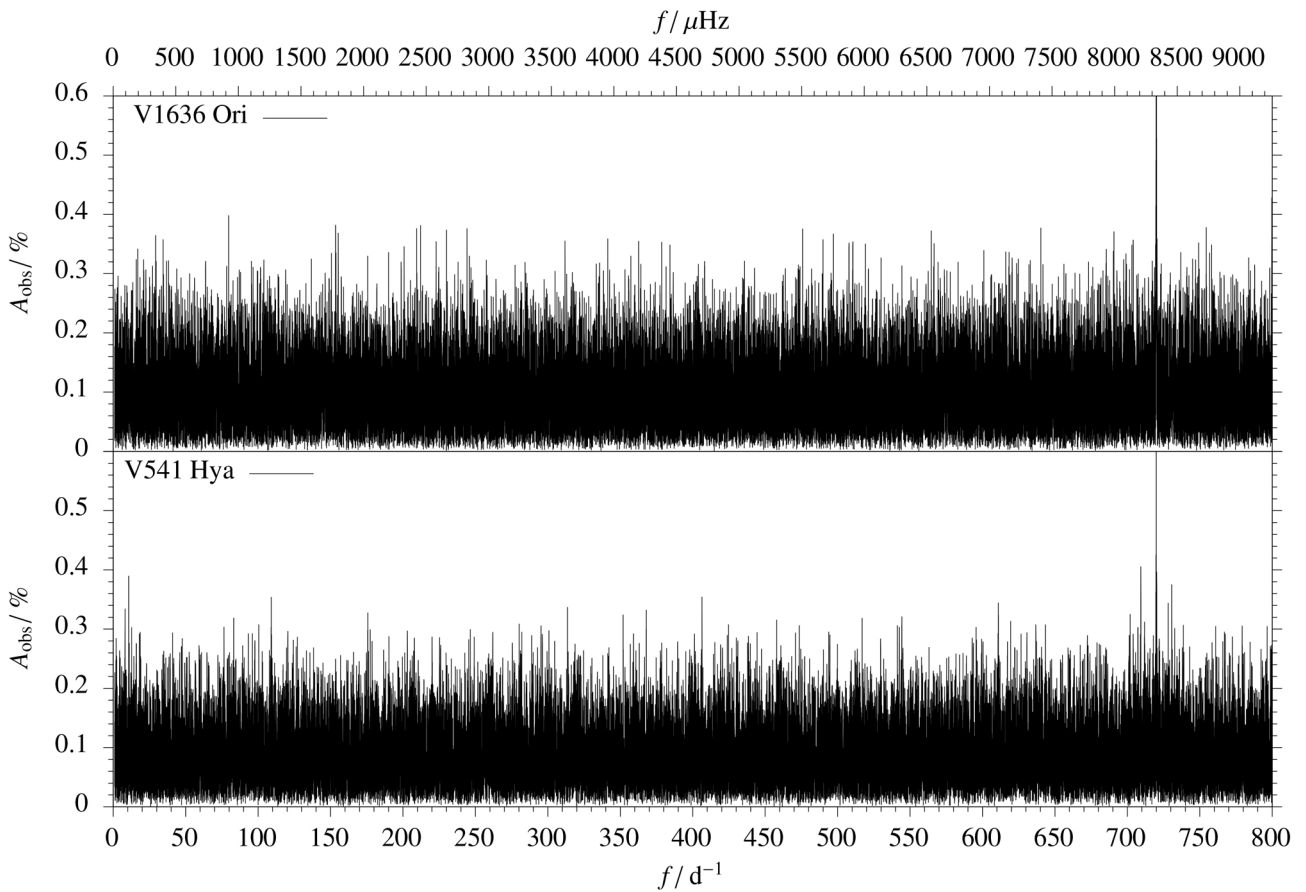
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## Appendix A: TESS data



**Fig. A.1.** Light curves of the TESS observations. Grey points are considered outliers and partially exceeding the plotting range. (a) V1636 Ori. (b) V541 Hya.

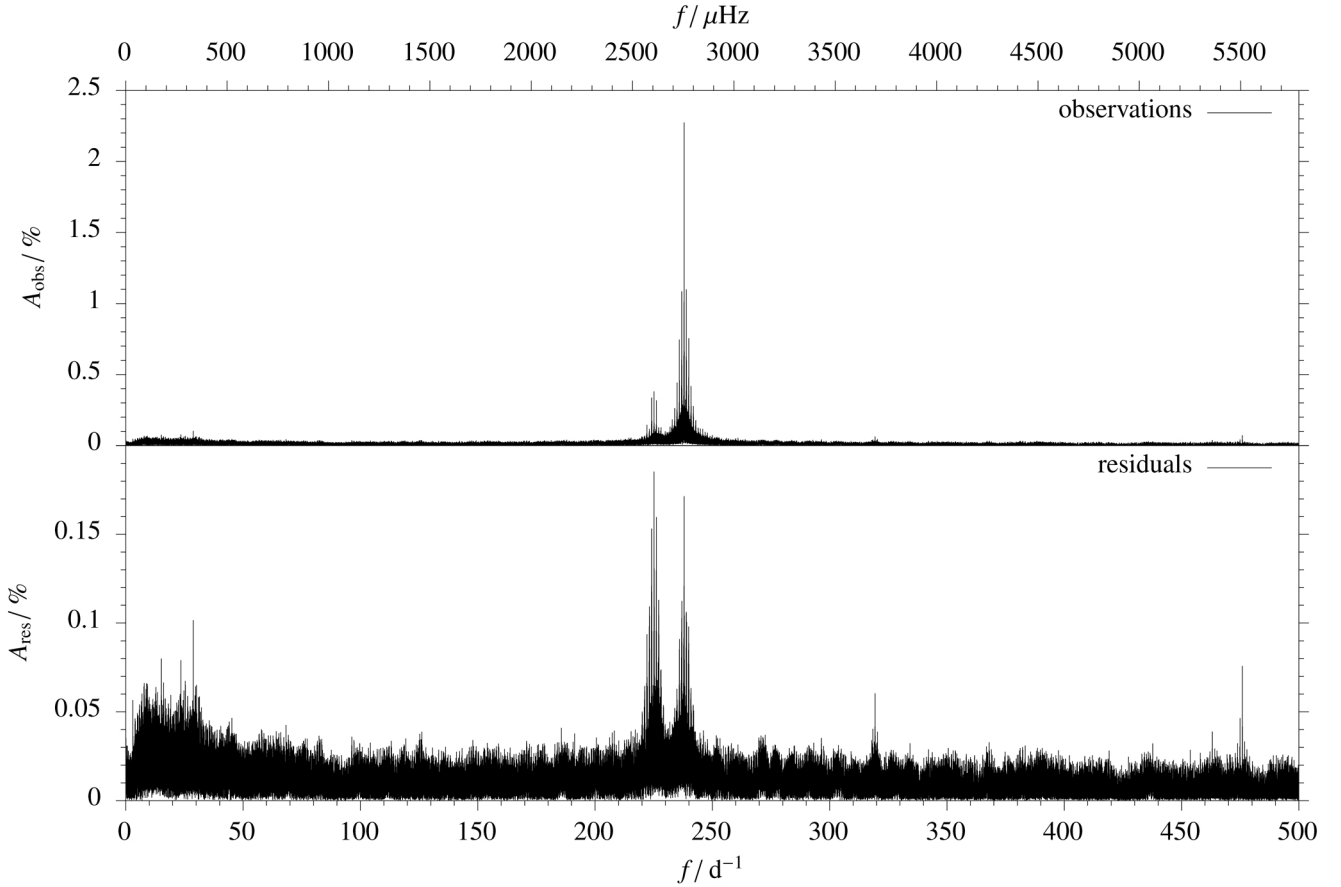


**Fig. A.2.** Amplitude  $A$  spectrum of the TESS observations. *Upper panel:* spectrum of V1636 Ori. *Lower panel:* spectrum of V541 Hya. The only peak above the noise level is the 120 s alias due to the cadence of the observations.

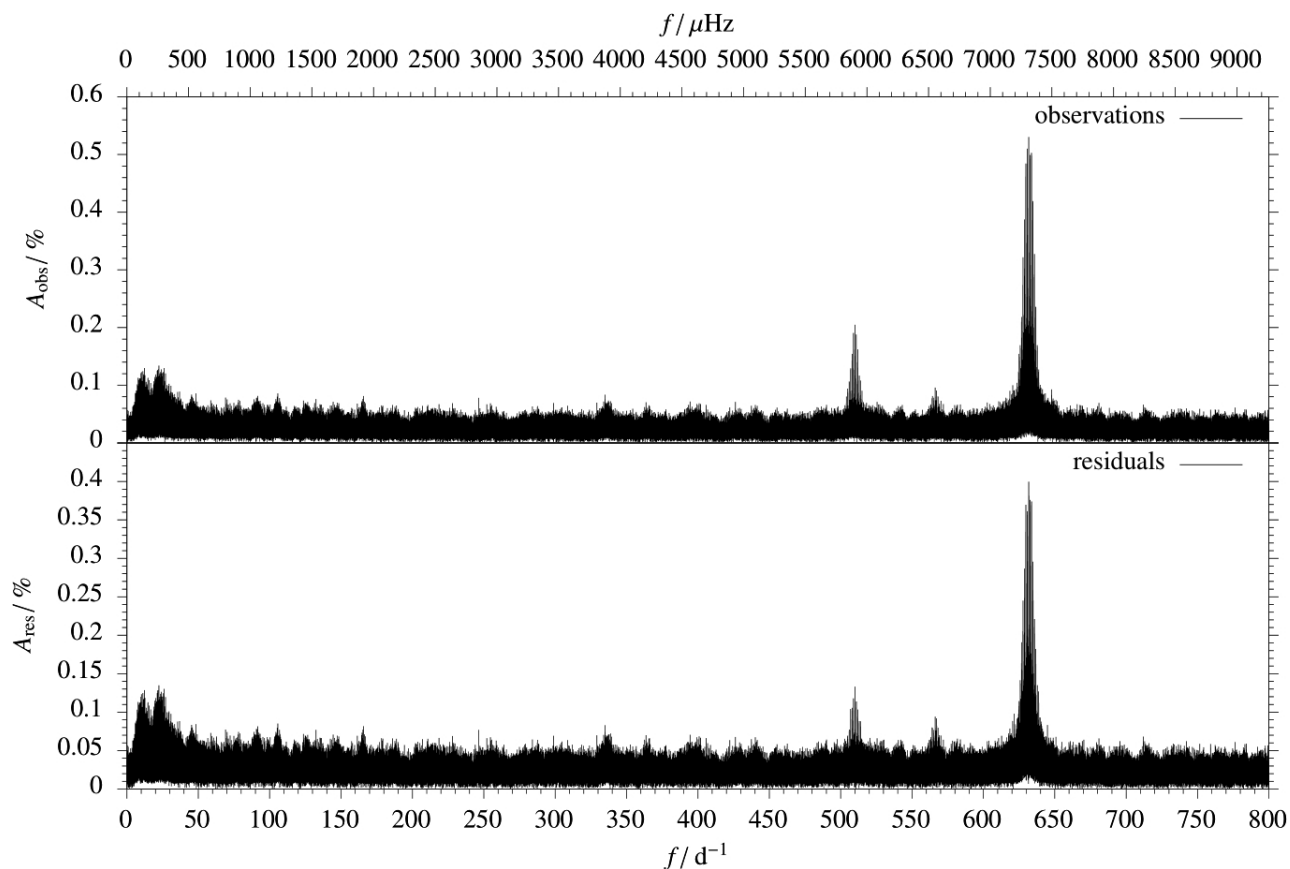
## Appendix B: Amplitude spectra

**Table B.1.** Additional pulsation modes identified for our targets not used in the O–C analysis due to their low S/N.

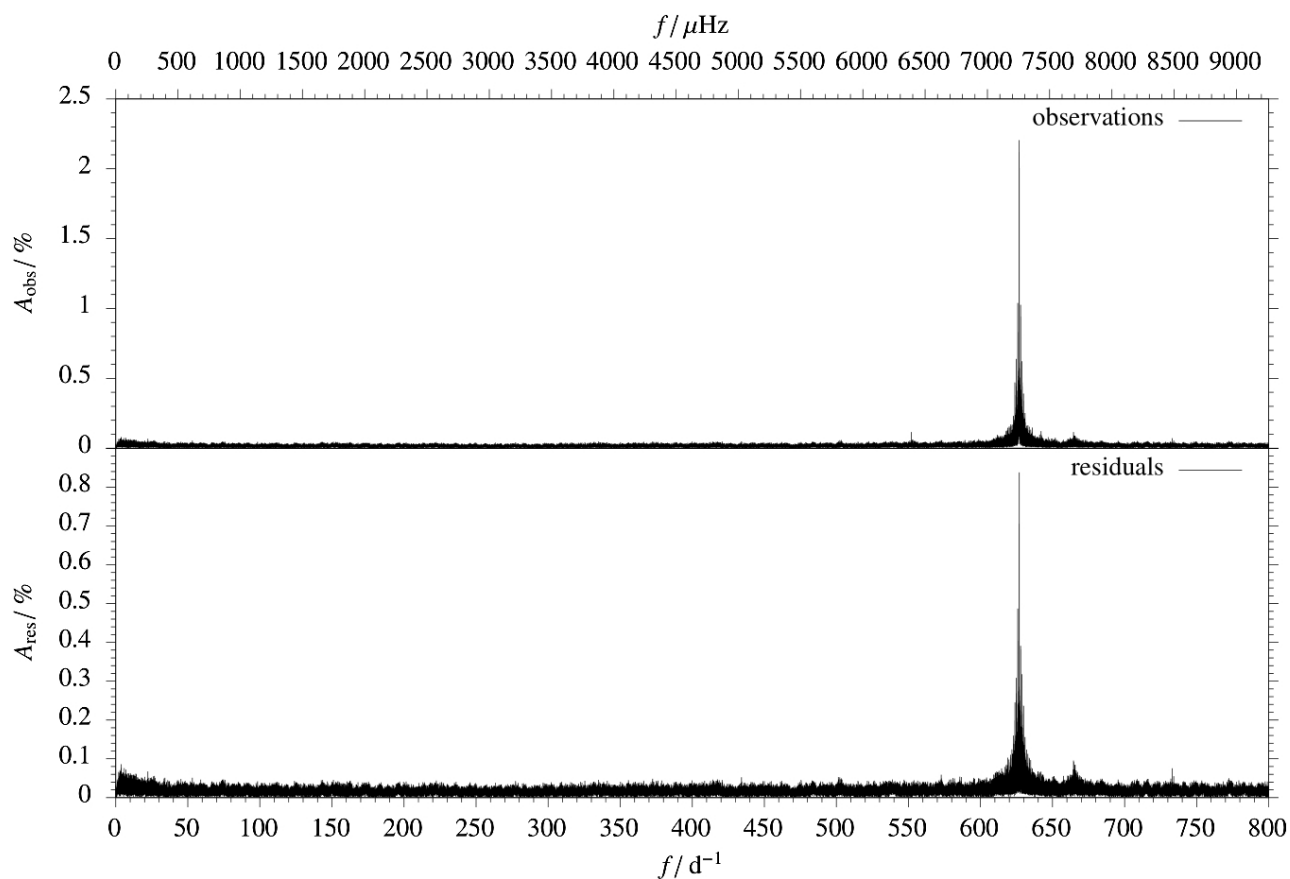
Target	$f/d^{-1}$	$A/\%$
DW Lyn	475.8231(2)	0.09(18)
	319.4042(3)	0.06(12)
	463.0100(6)	0.03(18)
V1636 Ori	566.24031(3)	0.6(3)
QQ Vir	733.0704(1)	0.3(1)
	664.4886(1)	0.2(1)
	572.73611(5)	0.19(9)
	664.7122(1)	0.1(1)
	434.1522(6)	0.01(7)
	502.410(2)	0.01(9)
V541 Hya	531.16759(16)	0.03(7)
	603.88741(6)	0.03(8)



**Fig. B.1.** Amplitude spectrum of DW Lyn. *Upper panel:* observations  $A_{\text{obs}}$ . *Lower panel:* residuals  $A_{\text{res}}$  after subtracting the light curve models from the observations.

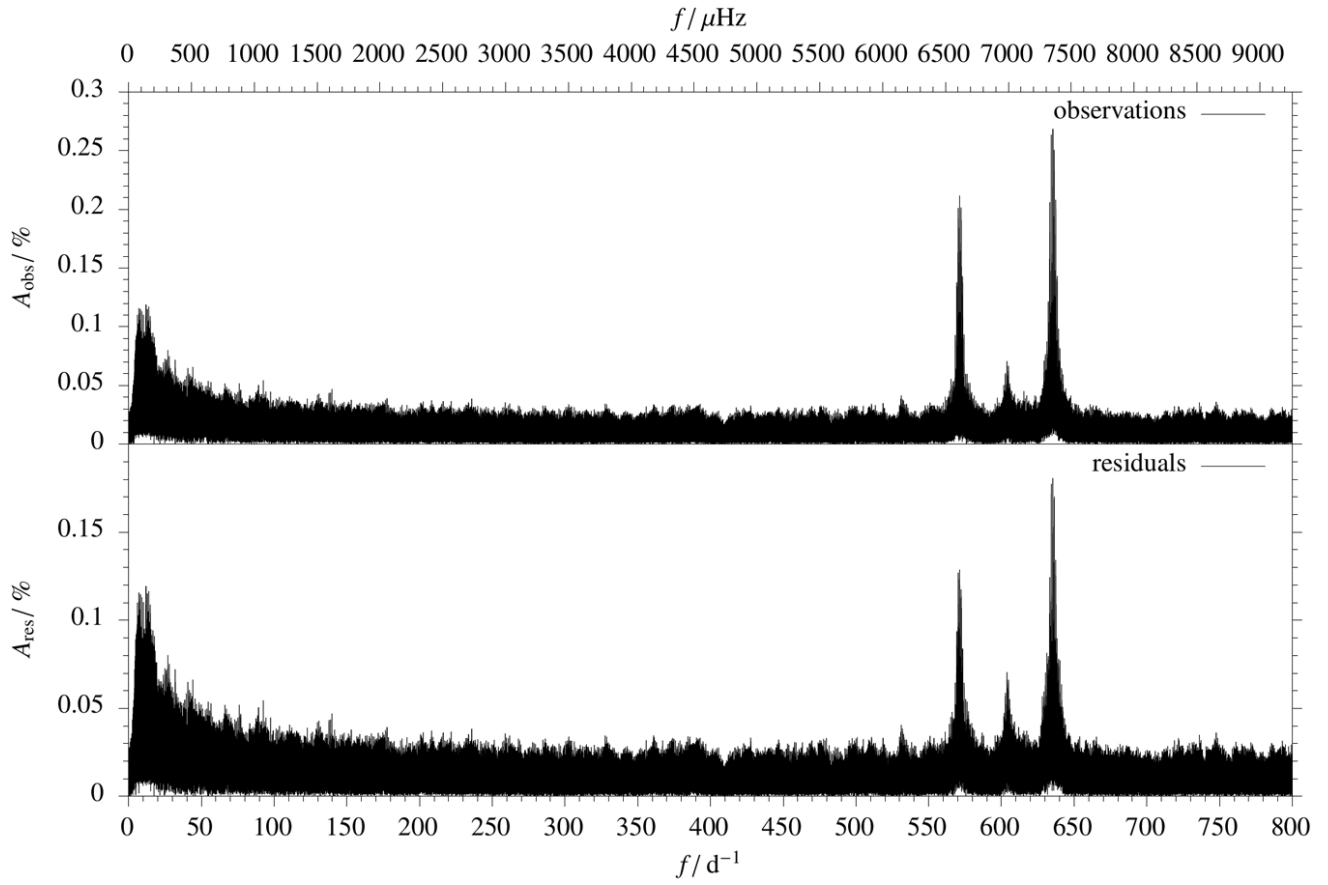


**Fig. B.2.** Same as Fig. B.1 but for V1636 Ori.



**Fig. B.3.** Same as Fig. B.1 but for QQ Vir.





**Fig. B.4.** Same as Fig. B.1 but for V541 Hya.