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# Elucidating the global distribution of reprocessing gas in NGC 1194

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## ABSTRACT

A joint *XMM–Newton* and *NuSTAR* observation was conducted for the bright, local Seyfert 1.9 galaxy, NGC 1194. The hard spectral form of this active galactic nucleus (AGN) was modelled using the toroidal reprocessor MYTORUS. The decoupled model form provides a good description of the spectrum, with reflection arising from gas with a global average column density  $>4 \times 10^{24} \text{ cm}^{-2}$  and transmission of the continuum through an order-of-magnitude lower column. In this model, the reflection strength is a factor of  $\sim 3$  higher than expected from a simple torus. Such a result may indicate that much of the intrinsic X-ray continuum is hidden from view. An alternative model is that of a patchy torus, where 85 per cent of sightlines are obscured by Compton-thick gas and the remaining 15 per cent by Compton-thin gas. The patchy torus model is based on a solar abundance of Fe and is consistent with X-ray partial-covering results found in other AGN. That a patchy torus model would relieve the issue with the strength of the reflection signature is not an intuitive result: such an insight regarding the geometry of the global reprocessing gas could not have been obtained using ad hoc model components to describe the spectral form.

**Key words:** galaxies:active – galaxies: individual: NGC 1194 – galaxies: Seyfert – X-rays: galaxies.

## 1 INTRODUCTION

Mapping the geometry, mass, and kinematics of circumnuclear material surrounding supermassive black holes in active galactic nucleus (AGN) is critical to understanding the distribution of observed properties in the local population, and, more fundamentally, the way in which material and energy are exchanged between the nuclear black hole and host galaxy.

Recent modelling of a Burst Alert Telescope (BAT)-selected *Suzaku* sample (Tatum et al. 2013) of type 1–1.9 radio-quiet AGN established that, contrary to expectations, such sources generally have very hard X-ray spectral forms at high energies. Even a model where simple disc reflection is the only component in the spectrum is not adequate to account for *NuSTAR* luminosities measured above 10 keV; such models overpredict the Fe line equivalent widths relative to the hardness ratios observed (Tatum et al. 2013). The most direct way to reconcile the distribution of local AGN properties is to allow for a significant component of absorption by Compton-thick gas, partially covering the X-ray source, to different degrees for the type 1–1.9 cases. This can be reconciled with the type 1 nature of some AGN by invoking holes in the gas distribution allowing a time-variable direct view of the nucleus. The high column density of gas required in the Tatum et al. study forms a natural extension of the multizoned X-ray absorber, whose signatures are clear in the 0.5–10-keV band (e.g. Kaspi et al. 2002; Reeves et al. 2004). Kinematic measurements show that the absorber is predominantly seen to be

outflowing (e.g. Blustin et al. 2005), tracing some fraction of the accreting material that is being blown back into the host, driven by radiation pressure, thermal pressure, magnetohydrodynamic forces or a combination of these. X-ray results to date are consistent with AGN having a cloud-filling factor that increases from the pole view down to the plane of the accretion disc, such that systems viewed close to the plane of the disc have a relatively high probability of there being a large number of clouds in the sightline (Nenkova et al. 2008). The gas distribution required is consistent with clumpy torus models (Elitzur & Shlosman 2006).

While the line-of-sight ionized absorption has been usefully constrained by measurement of the depths of spectral features, the establishment of the global covering factor for the gas, and the radial location (and thus the total mass) have proved more challenging. The strengths of observed emission lines can be used in this regard, and in particular, that of the scattered X-ray emission predicted from the reprocessing gas.

The column density corresponding to an X-ray scattering optical depth of unity is commonly used to set the convention of terminology. This relates to the inverse of the Thomson scattering cross-section ( $\sigma_T$ ) such that  $N_H \simeq 1/1.2\sigma_T \simeq 1.25 \times 10^{24} \text{ cm}^{-2}$  – values above this column density being denoted ‘Compton thick’ and those below it ‘Compton thin’. One limitation of work to date has been the widely used assumption of an infinite column density for the reprocessing gas, the so called ‘Compton-thick’ case.

Scattering from Compton-thin gas produces strong fluorescent line emission, most notably from the K shell of Fe, and a rich variety of possible spectra (that depend on geometry, viewing angle, and column density). When the reprocessor has an infinite column

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density the Compton hump peaks at an energy that depends on structure orientation and geometry (e.g. see George & Fabian 1991). Physically motivated models that treat the scattered continuum and Fe K $\alpha$  fluorescent emission-line self-consistently for finite column density material have been recently released in a form suitable for application to X-ray spectral data (e.g. Ikeda, Awaki & Terashima 2009; Murphy & Yaqoob 2009; Brightman & Nandra 2011; Liu & Li 2014). The high-energy bandpass of *Suzaku* (e.g. Yaqoob 2012; Yaqoob et al. 2015) and, later, *NuSTAR* has allowed these toroidal models to be applied to Compton-thick AGN, allowing important insight into the reprocessor details (e.g. Brightman et al. 2015; Baloković et al. 2018; Marchesi et al. 2019), for which scant information was available previously.

The improved methodologies inherent in the MYTORUS model allows one to probe the three-dimensional structure of the X-ray reprocessor in a meaningful way. Initial results offer promise for significant progress. For example, application of the MYTORUS model to the type 2 AGN Mrk 3, produces a solution in which the line-of-sight column density is higher than the global average column (Yaqoob et al. 2015).

Models for the local AGN population provide a recent-epoch data point in the fundamental problem of accounting for the cosmic X-ray background (CXB). This line of investigation gives us a more realistic estimate of the actual contribution of all AGN (type 1, intermediates and type 2) to the integrated absorbed light in the CXB and from that one might draw model possibilities for sources seen at higher redshift. Once this is properly accounted for in AGN population synthesis models, we might not need to find a new population of hitherto undetected Compton-thick AGN to explain the X-ray background.

The hardest AGN offer the most insight into understanding the circumnuclear gas, and one of the hardest sources in the Tatum et al. (2013) sample is NGC 1194. This nearby ( $z = 0.0136$ ) Seyfert 1.9 galaxy hosts a water maser, which has allowed a precise measurement of the mass of the black hole at  $M = 6.5 \pm 0.3 \times 10^7 M_\odot$  (Kuo et al. 2011). NGC 1194 is also a bright X-ray source, detected in the BAT catalogue. NGC 1194 has been classified as X-ray Compton thick and has been discussed by various authors in the literature (e.g. Della Ceca et al. 2008; Georgantopoulos & Akylas 2019; Marchesi et al. 2019, and references therein). Further to this, the source has been classified as Compton-thick on the basis of the X-ray to infrared ratio versus the X-ray hardness ratio (Severgnini, Caccianiga & Della Ceca 2012). Despite the source brightness and interesting properties, NGC 1194 has been poorly studied in the X-ray regime. The source has had only a 16 ks snapshot by *XMM-Newton*, 40 ks with *Suzaku* (Tatum et al. 2013; Tanimoto et al. 2018), and a 32 ks *NuSTAR* exposure available in X-ray archives. Thus, we proposed a joint *XMM-Newton* and *NuSTAR* observation of the target, and those data form the subject of this paper.

Section 2 outlines the details of the new observations. Section 3 examines the time variability, along with a basic spectral parametrization. Section 4 describes the application of a full-covering spherical reprocessor, while Section 5 presents detailed modelling using several variations of the MYTORUS model. Section 6 discusses the results in the context of other information for NGC 1194.

## 2 OBSERVATIONS

### 2.1 XMM-Newton

*XMM-Newton* conducted an observation of NGC 1194 during 2020 January 16 as part of a coordinated *XMM-Newton*/*NuSTAR* exposure. Table 1 gives a summary of the observations. The *XMM-Newton*

**Table 1.** Observation log.

Date	Observation ID	Total exposure	Flux 2–10 keV
<i>XMM-Newton</i>			
<b>2020-1-16</b>	<b>0852200101</b>	<b>62(MOS), 52(pn)</b>	<b>0.10</b>
2006-2-19	0307000701	16	0.11 <sup>1</sup>
<i>NuSTAR</i>			
<b>2020-1-17</b>	<b>60501011002</b>	<b>56 (per FPM)</b>	<b>2.18</b>
2015-2-28	60061035002	32	1.80 <sup>2</sup>

*Notes.* Observed fluxes are given in units of  $10^{-11}$  erg cm $^{-2}$  s $^{-1}$  based on the *XMM-Newton* pn and the mean FPM data. Bold face type denotes the sequences presented in this paper.

<sup>a</sup>Marchesi et al. (2019).

<sup>b</sup>Extrapolated from Marchesi et al. (2019).

observations were performed in Full Window mode, with the medium filter in place. All data were processed using SAS v18.0.0 and HEASOFT v6.23 software. The EPIC pn and MOS source spectra were extracted using a circular region with a radius of 35 arcsec and background spectra using two circular regions each with a radius of 35 arcsec. After cleaning using standard criteria, i.e. PATTERN  $\leq 12$  for MOS,  $\leq 4$  for pn, pulse invariant channels between 200 and 15000, and FLAG = 0, each observation was filtered to remove a brief period of high background at the end of the exposure. The resulting *XMM-Newton* event files yielded exposures of  $\sim 62$  ks per MOS CCD and  $\sim 52$  ks for the pn.

The count rates over the 0.5–10.0 keV band were 0.02 count s $^{-1}$  MOS $^{-1}$  and 0.10 count s $^{-1}$  for the pn. The background was 7 per cent of the source count rate in that bandpass. We generated the response matrices and the ancillary response files at the source position using the SAS tasks *arfgen* and *rmfgen* and the latest calibration available. The EPIC spectra were binned to a minimum of 50 counts per energy bin: the binned data maintained sampling finer than the spectral resolution of the EPIC CCDs. The *XMM-Newton* RGS1 and RGS2 data were reduced with the standard SAS task RGSPROC. Unfortunately, the source was too faint in the soft band for a useful RGS spectrum to be accumulated.

### 2.2 NuSTAR

*NuSTAR* carries two co-aligned telescopes containing focal plane modules A and B (FPMA, FPMB; Harrison et al. 2013) covering a useful bandpass of  $\sim 3$ –80 keV for AGN. *NuSTAR* observed NGC 1194 on 2020 January 17, overlapping the *XMM-Newton* exposure (Table 1). Event files were created through the NUPipeline task v0.4.6, cleaned applying the standard screening criteria, where passages through the SAA were filtered out, setting the mode to ‘optimised’ in NUCALSA. This yielded a net exposure of  $\sim 111$  ks in the summed data from both focal plane modules. For each of the focal plane modules, source spectra were extracted from a circular region with a radius of 70 arcsec, while background spectra were extracted from a circular region with a 75 arcsec radius located on the same detector. *NuSTAR* source spectra were binned to 100 counts per spectral channel, maintaining a sampling that is finer than the spectral resolution of the instruments. During 2020 January the summed FP modules yielded a source count rate of  $0.067 \pm 0.008$  count s $^{-1}$  over 3–50 keV. The background level was  $\sim 9$  per cent of the total count rate. These rates correspond to observed fluxes  $\sim 1.0 \times 10^{-12}$  and  $\sim 1.6 \times 10^{-11}$  erg cm $^{-2}$  s $^{-1}$  in the 2–10 and 10–50 keV band, respectively.

### 2.3 General considerations

Spectra are analyzed with XSPEC v12.9.1m. We used data over 0.6–10 keV for the pn and MOS and 3–50 keV for *NuSTAR*. This data range was chosen as our main application here is the MYTORUS model, which is not calculated below that energy range. All models included a full covering column of gas at a redshift of zero, that covered all components and whose lower limit was set to the Galactic line-of-sight absorption,  $N_{\text{H,Gal}} = 5.53 \times 10^{20} \text{ cm}^{-2}$  (HI4PI Collaboration 2016). This absorber was parametrized using TBABS (Wilms, Allen & McCray 2000). All other model components were adjusted to be at the redshift of the host galaxy. We assume  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$  throughout.

For the ionized emitter, we generated model tables using version 2.41 of the XSTAR code (Kallman & Bautista 2001; Kallman et al. 2004), allowing abundances that were variable such that they could float in the fit. XSTAR models the reprocessing gas as thin slabs, with parameters of atomic column density and ionization parameter  $\xi$ , defined as

$$\xi = \frac{L_{\text{ion}}}{n_e r^2},$$

which has units  $\text{erg cm s}^{-1}$  and where  $L_{\text{ion}}$  is the ionizing luminosity between 1 and 1000 Rydbergs,  $n_e$  is the gas density in  $\text{cm}^{-3}$ , and  $r$  is the radial distance (cm) of the absorbing gas from the central continuum source. The spectral energy distribution (SED) illuminating this gas was taken to be a simple power law with  $\Gamma = 2$ . The turbulent velocity was taken as  $\sigma = 300 \text{ km s}^{-1}$ .

Parameters are quoted in the rest frame of the source and errors are at the 90 per cent confidence level for one interesting parameter ( $\Delta\chi^2 = 2.706$ ). In the plots, the data are binned more coarsely than allowed in the fit, for visual clarity.

There is a small calibration offset between the different detectors used and so a constant component was allowed in all models, constrained to a range of 0.9–1.1 for the cross-normalization constant between the pn, MOS and the summed *NuSTAR* detectors (Madsen et al. 2017).

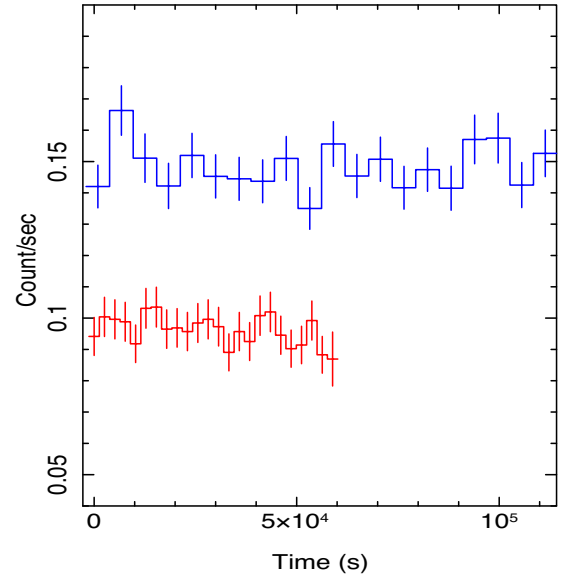
## 3 INITIAL ANALYSIS

### 3.1 Flux stability

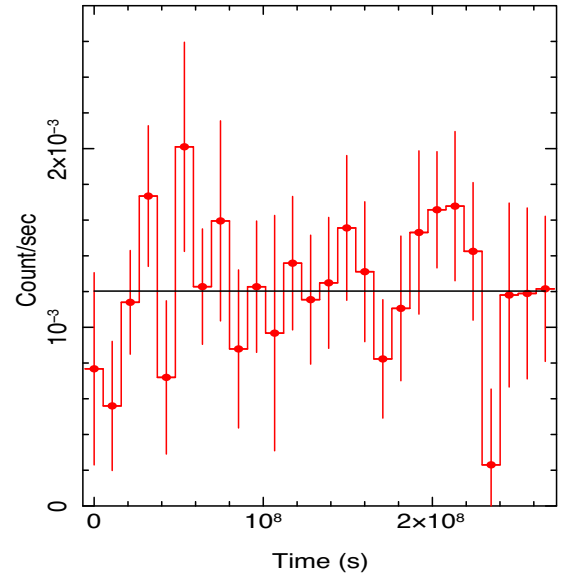
Light curves were constructed for the source, extracting the counts through the source regions detailed in the previous section. The *XMM-Newton* pn data were sampled using 2560 s bins and in the full band of 0.5–10.0 keV. The *NuSTAR* data covering 3.0–70.0 keV were used, sampled on the orbit time-scale of 5814 s and data from the two FP modules were summed. The *NuSTAR* exposure was initiated 920 s after the *XMM-Newton* observation was started. Fig. 1 shows the light curves for the *NuSTAR* and *XMM-Newton* sequences.

Comparison of the hard-band flux from our *NuSTAR* observation and the archived data from 2015 indicates there to be no significant variability on time-scales of years. To further check the count rate trend on long-term time-scales, we extracted the light curve from the *Swift* BAT, based on 105 months of sky survey data and available online at <https://swift.gsfc.nasa.gov/results/bs105mon/>.

This BAT light curve, sampled at four months, covers 14–195 keV and includes observations carried out between 2004 December and 2013 August. The source flux shows no significant variability over that  $\sim 9$ -yr baseline (Fig. 2). Fitting the data to a constant flux model produces a  $\chi_r^2 = 22/25$ .



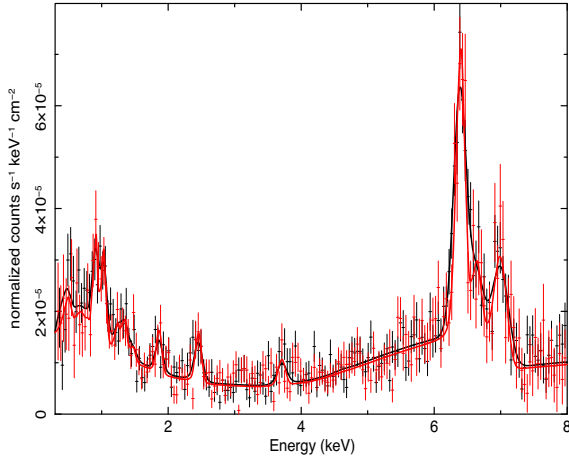
**Figure 1.** The light curves for the *XMM-Newton* pn data (red), over 0.5–10 keV, and sampled at 2560 s and the overlapping *NuSTAR* data (blue), covering 3–70 keV binned at the orbit time-scale (5814 s). The x-axis shows the time from the start of the *XMM-Newton* observation.



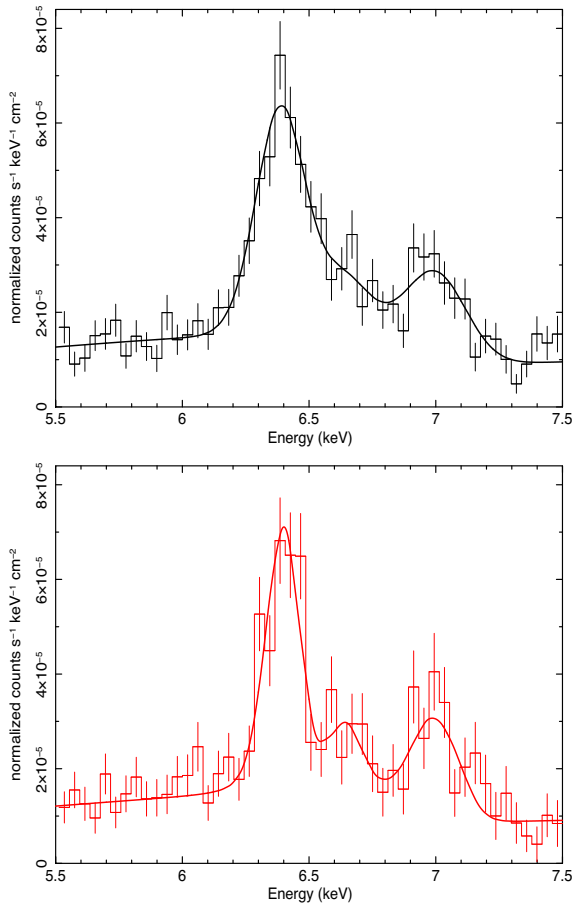
**Figure 2.** The 105 month BAT light curve sampled at four-month intervals. The zero-time point corresponds to 2004 December. The horizontal line represents the best-fitting constant value.

### 3.2 Basic characterization of the X-ray spectrum

Initially just the *XMM-Newton* (pn and MOS) data were fitted to parametrize the soft X-ray continuum and the multiple emission lines that were found to be present both at soft X-rays (Fig. 3) and at Fe K (Fig. 4). The line energies are consistent with blended Ne IX–X in the 0.9–1 keV regime, Mg XI at  $\sim 1.3$  keV, Si XIII at  $\sim 1.8$  keV, and S XV at  $\sim 2.4$  keV. Further to these, a line identified at  $\sim 3.7$  keV is likely from Ca K $\alpha$ . In the Fe K band, emission components include contributions from neutral gas at 6.4 keV, He-like, and H-like Fe at 6.67 and 6.97 keV, plus Fe K $\beta$  emission at 7.06 keV that is modelled as a fixed 13.5 per cent of the fitted Fe



**Figure 3.** Emission lines in the EPIC data from the pn (black) and MOS (red). The model is represented by a solid line (see the text for details).



**Figure 4.** The counts spectrum in the Fe K band, for the pn (top, black) and MOS (bottom, red) data. The model line is shown as for Fig. 3.

$K\alpha$  line flux (Table 2; Fig. 3). The lines are all consistent with being narrow, with  $\sigma < 30$  eV for the Fe  $K\alpha$  component at 6.4 keV. The observed mix of ionized and neutral-line components are not consistent with an origin in a single gas zone. The lines were fitted using Gaussian profiles, and the line parameters are summarized in Table 2. The baseline model over which the lines were fit had the

**Table 2.** X-ray line emission.

$E$ (keV)	ID	Flux $\frac{\text{photons cm}^{-2} \text{s}^{-1}}{10^{-6}}$	EW (eV)
$0.91 \pm 0.01$	Ne IX	$2.77 \pm 0.64$	141
$1.04 \pm 0.01$	Ne X	$2.21 \pm 0.52$	134
$1.23^{+0.04}_{-0.03}$	Mg I $K\alpha$	$0.81 \pm 0.35$	61
$1.34 \pm 0.02$	Mg XI	$1.00 \pm 0.34$	84
$1.47 \pm 0.05$	Mg XII	$0.44 \pm 0.36$	40
$1.88^{+0.01}_{-0.03}$	Si XIII	$1.19^{+0.39}_{-0.22}$	136
$2.45 \pm 0.02$	S XV	$1.27 \pm 0.41$	169
$3.72^{+0.17}_{-0.05}$	Ca $K\alpha$	$0.50 \pm 0.32$	67
$6.39 \pm 0.01$	Fe $K\alpha$	$11.12^{+0.81}_{-0.92}$	550
$6.66 \pm 0.04$	Fe XXV	$2.28^{+0.74}_{-0.65}$	63
$6.96 \pm 0.04$	Fe XXVI	$1.60^{+0.69}_{-0.67}$	78
7.05 (fixed)	Fe $K\beta$	1.50	92

following form:

$$N_{\text{H,Gal}} \times [N_{\text{H}} \times [(1 - f) \times \text{po}] + f \times [\text{po} + \text{zg}(12)]],$$

where  $N_{\text{H,Gal}}$  is the Galactic column, po is the power law (assumed to fold over at 500 keV).  $f$  is the covering fraction of the absorbing gas,  $N_{\text{H}}$  is an intrinsic absorber that covers only a fraction  $(1 - f)$  of the continuum. The intrinsic absorber is represented by the XSPEC model PHABS, along with the Compton scattering correction for that column density, applied using CABS. zg(12) represents the twelve individual Gaussian-line components. This model fit yielded  $f = 0.078 \pm 0.008$ ,  $\Gamma = 1.08 \pm 0.17$ ,  $N_{\text{H}} = 3.89^{+0.46}_{-0.45} \times 10^{23} \text{ cm}^{-2}$ . This model provided a good fit to the broad X-ray form of the *XMM-Newton* data, with  $\chi^2 = 316/297 \text{ dof}$ .

The absence of O lines in the spectrum is surprising, and so we investigated this further, fitting the pn and MOS spectra jointly, using an emission table generated from XSTAR, allowing variable abundances. This table was assumed to be absorbed by the Galactic column. The host galaxy of NGC 1194 is viewed at high inclination,  $\frac{b}{a} = 0.45$  (Véron-Cetty & Véron 2006). In order to constrain, the column density of gas within the host galaxy, comprising the interstellar medium (ISM) of the host, we retrieved the fluxes of H $\alpha$  and H $\beta$  from SDSS (Ahumada et al., in preparation). The observed ratio of H $\alpha$ /H $\beta$  was  $\approx 17.5$ . Assuming an intrinsic ratio of 2.9 (Osterbrock & Ferland 2006) and the Galactic extinction curve (Savage & Mathis 1979), we calculated a reddening of  $E(B - V) = 1.6$ , which would arise in an H I column density of  $\sim 8.5 \times 10^{21} \text{ cm}^{-2}$  (Shull & van Steenberg 1985). Since, it is likely that some fraction of the gas in the host disc is ionized, we assumed an ISM column density of  $10^{22} \text{ cm}^{-2}$ .

The soft-band lines in NGC 1194 have recently been resolved and found to arise in an extended X-ray region of gas (Braitto et al., in preparation); therefore, we placed the host galaxy absorption in front of the emission table in the model. However, the strength of the soft flux in this source, and the shape of the soft-band spectrum, are not consistent with the entire nuclear spectrum being obscured by an ISM column density as high as  $N_{\text{H,ISM}} \sim 10^{22} \text{ cm}^{-2}$ . This raises the question as to what elements of the observed spectrum are obscured by this ISM gas component. In a similar AGN, NGC 7582, (Braitto et al. 2017) find the dust lane to obscure part of the spatially extended narrow-line region gas, but not the nuclear emission. In NGC 7582, the dust lane is responsible for the suppression of the emission lines below 1 keV, which arise from the AGN NLR.



The detailed morphology of NGC 1194 is under study (Braitto et al., in preparation). It appears that, similar to NGC 7582, the ISM gas lane may obscure part of the NLR and thus the most appropriate application of that absorption in the spectral fit, is to apply it only to the soft-band line-emission gas. We proceed under this assumption in our spectral fits.

The emitter column density was also fixed at  $N_H = 1 \times 10^{22} \text{ cm}^{-2}$ , and the fit returned  $\log \xi = 1.77^{+0.08}_{-0.17}$ . Abundances were allowed to float and found to be  $A_O > 0.57$ ,  $A_{Ne} = 0.84^{+0.65}_{-0.31}$ ,  $A_{Mg} = 0.12^{+0.37}_{-0.12}$ ,  $A_{Si} = 0.70^{+0.66}_{-0.27}$ ,  $A_S = 2.76^{+0.45}_{-1.20}$  of solar. From this fit, it is evident that O line emission is suppressed by the ISM absorber, and once that is taken into account, the data are consistent with solar abundance of O in the emitting gas. Freezing all abundances at 1.0 increased the fit statistic by  $\Delta\chi^2 = 25$ , with the increase driven by  $A_{Mg}$  and  $A_S$ , which are not consistent with solar abundance in this fit. This may indicate either a true deviation from solar abundances, or an inadequacy in the ISM modelling. No further progress can be made on this question, with the data in hand. (This model gas zone cannot account for the Ca and Fe K-shell lines, which are addressed in later fits.)

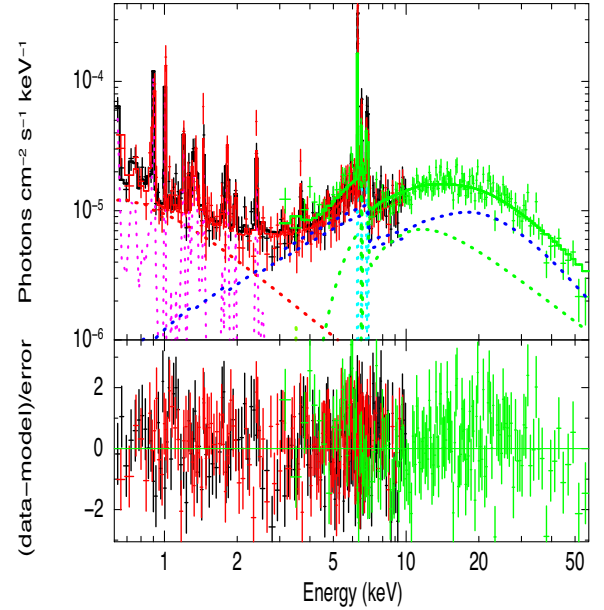
Following this simple parametrization, the *NuSTAR* data were then included in a joint fit with the pn and MOS spectra. The baseline model comprised the following form:

$$N_{H,Gal} \times N_{H1}[N_{H2}[(1-f) \times po] + f \times (po + ref + N_{H,ISM} \times EMIS + zg(5))],$$

where EMIS is the emission profile from the XSTAR table. *Ref* is reflection from gas out of the line of sight – represented here by the XSPEC model PEXRAV, assumed to be viewed orientated at  $60^\circ$ , to be illuminated by the fitted power law and have a cosmic abundance of Fe. Other components are designated as for the previous fit. *zg(5)* represents the five individual Gaussian line components not accounted for by the XSTAR table or by PEXRAV i.e. the two ionized Fe lines, the neutral Fe  $K\alpha$  and  $K\beta$  lines and Ca  $K\alpha$  (Table 2): although these lines might arise in the reflector, they are not included in the PEXRAV model. This model fit yielded  $f = 0.021 \pm 0.005$ ,  $\Gamma = 1.51 \pm 0.04$ ,  $N_{H,1} = 6.91^{+2.19}_{-6.91} \times 10^{20} \text{ cm}^{-2}$  and  $N_{H,2} = 8.74^{+1.36}_{-1.27} \times 10^{23} \text{ cm}^{-2}$ . This model provided a good fit to the broad X-ray form (Fig. 5), with  $\chi^2 = 586/450$  dof. In this fit, the reflection component is relatively strong, yielding  $R = 1.52^{+0.32}_{-0.47}$  against the continuum.  $R=1$  represents a source that subtends  $2\pi$  steradians to the continuum, such as one might expect for a standard geometrically thin accretion disc. High  $R$  values are sometimes explained as being due to a prior episode of higher continuum with time-delays leading to the chance capture of a relatively high reflection signature. That is not likely applicable to NGC 1194, which does not show evidence for continuum variability (Fig. 1 and Table 1). The emission lines are narrow, with an upper limit of  $\sigma < 30 \text{ eV}$ : This suggests an origin from the outer broad-line region (BLR) or torus. The observed line fluxes and equivalent widths are consistent with those detailed in Table 2.

#### 4 SPHERICAL REPROCESSOR MODEL

Next, we investigated more physically motivated models for the X-ray reprocessor around NGC 1194. First, we compared the data to a model utilizing reprocessing in a uniform full-covering sphere of gas. The reprocessor was implemented using an additive table named SPHERE0708.FITS that is a publicly available table originating from the calculations of Brightman & Nandra (2011). The sphere is assumed to be illuminated by a power-law continuum with a high



**Figure 5.** The 2020 January broad-band spectrum of NGC 1194, fitting the *XMM-Newton* pn (black), MOS (red), and *NuSTAR* (green) data together. The model components are an absorbed power law (blue), an XSTAR emission line table (magenta), PEXRAV for a reflection component (green dotted line), a scattered fraction of the unabsorbed power law (red dotted line) and Fe K emission (aqua).

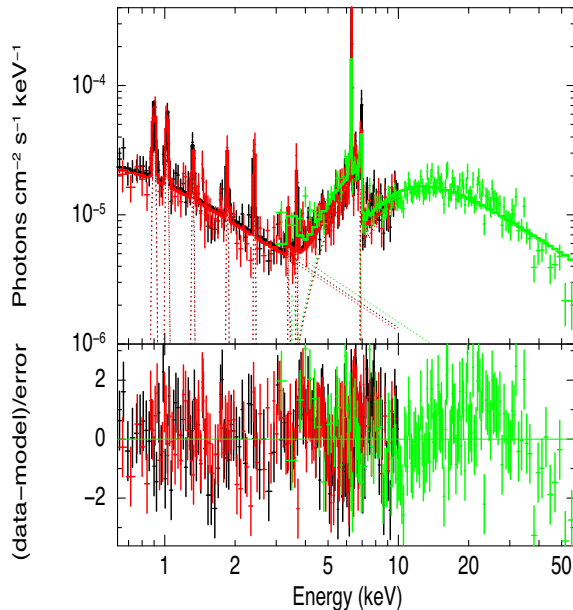
energy cut-off at 500 keV. The Fe abundance was allowed to be free and the other metal abundances (that are controlled by a single scaling parameter) were set to solar, using Anders & Grevesse (1989). The photoelectric absorption cross-sections used for the sphere are those of Verner et al. (1996). In generating the table, Brightman & Nandra (2011) self-consistently calculate the Fe  $K\alpha$  and  $K\beta$  emission lines from the gas, and incorporate  $K\alpha$  emission from other elements, including Mg and Ca. In this model, the Fe line emission, scattered X-rays, and absorption opacity cannot be decoupled. The other emission lines in the data (from ionized gas) are modelled using eight individual Gaussian components and fitted line parameters are consistent with those tabulated previously (Table 2). Galactic absorption was added to the model, and the final model form was

$$N_{H,Gal} \times [pow + sphere0708 + zg(8)].$$

The model yielded  $\chi^2 = 598/437$  dof. The local column was found to be in excess of the Galactic value, with  $N_{H,Gal} = 9.07^{+4.66}_{-4.25} \times 10^{20} \text{ cm}^{-2}$ , covering a continuum with  $\Gamma = 1.21^{+0.05}_{-0.04}$ , with a column density of  $N_H = 4.76^{+0.25}_{-0.32} \times 10^{23} \text{ cm}^{-2}$  and Fe abundance  $A_{Fe} = 1.71^{+0.13}_{-0.12}$ . The fit does not provide adequate curvature above 10 keV to explain the shape of the high energy hump (Fig. 6), and also yields an unrealistically flat power-law index for the illuminating continuum. With these two limitations, we do not consider the spherical reprocessor further.

#### 5 MYTORUS

The toroidal X-ray reprocessor model, MYTORUS (Murphy & Yaqoob 2009; Yaqoob et al. 2010), is illuminated by a central X-ray continuum source. The global covering factor of the reprocessor is 0.5, corresponding to a solid angle subtended by the structure at the central X-ray source of  $2\pi$  (an opening half-angle of  $60^\circ$ ). MYTORUS self-consistently calculates the Fe  $K\alpha$  and Fe  $K\beta$  fluorescent line



**Figure 6.** The spherical reprocessor of Brightman & Nandra (2011), compared to the data. Colours as for Fig. 5.

emission and includes the effects of absorption and Compton scattering on both the continuum and lines. MYTORUS offers the ability to leverage constraints on the Fe K line emission and spectral curvature together, to obtain both the global and line-of-sight column density. The model is therefore one of the most powerful available for fitting X-ray spectra. Although the nominal covering factor is fixed at 50 per cent of the sky as seen by the continuum source, one can model other physical possibilities using the ‘decoupled’ mode that allows for a different line-of-sight column to the gas responsible for reflection (the Fe line emission strength can still be linked self-consistently to the column of the reflector). Such a decoupling is often required in spectral analyses of high-quality broad-band X-ray spectra, because the column density in and out of the line-of-sight column density are not always the same (e.g. Risaliti et al. 2009; Marchese et al. 2012; Yaqoob et al. 2015).

In our fits, a single power law was assumed to be seen in transmitted and scattered light, and this was implemented by linking the photon indices of these components in all fits. The general model form was

$$N_{\text{H}} \times [(MYTZ \times \text{pow}) + \text{pow} + MYTS + \text{gs} \times MYTL + \text{zg}(10)],$$

where *MYTZ* designates the line-of-sight absorption, *MYTS* the scattered X-rays, and *MYTL* the line emission (both from material out of the line of sight, hereafter referred to as the ‘global’ gas), from MYTORUS. *zg*(10) represents 10 individual Gaussian lines (Table 2) present in the data that are not accounted for by the MYTORUS model. *gs* is a Gaussian smoothing component convolved with the torus neutral Fe line profile, using a kernel of  $\sigma = 30$  eV at 6 keV. The power-law continuum is split over two sightlines, one is absorbed by *MYTZ* and the other is absorbed only by  $N_{\text{H}}$ , a full-covering column of gas whose lower limit is set at the Galactic line-of-sight absorption, but whose upper limit is not constrained.

### 5.1 Coupled model fits

Taking the ‘Coupled Model’ to mean one for which the line of sight and global column densities are tied together, we performed two variations on this fit. First, the angle of the global gas was

fixed at  $\theta = 90^\circ$  and the normalization of the global-scattered light and line emission were directly tied to the illuminating continuum. This model did not provide an acceptable fit to the data, yielding  $\chi^2 = 709/438$  dof with  $N_{\text{H,MYT}} = 7.2^{+0.2}_{-0.4} \times 10^{23} \text{ cm}^{-2}$ , while the full covering gas component gave  $N_{\text{H}} = 2.04^{+0.49}_{-0.44} \times 10^{21} \text{ cm}^{-2}$ , exceeding the Galactic line-of-sight value. 2.9 per cent of the total intrinsic continuum flux escaped without absorption by the torus (Table 3).

A useful variation on the strictly coupled model is to allow a modest offset of orientation angle for the global gas, and we investigated this possibility, setting  $\theta = 70^\circ$  and again, the normalization of the global-scattered light and lines were tied directly to the illuminating continuum. This fit yielded  $\chi^2 = 698/438$  dof. In this case, the torus column density was  $N_{\text{H,MYT}} = 8.4^{+0.4}_{-0.3} \times 10^{23} \text{ cm}^{-2}$  while the full-covering layer had  $N_{\text{H}} = 1.72^{+0.47}_{-0.43} \times 10^{21} \text{ cm}^{-2}$ . Again we require a small escape of flux without absorption by the torus, in this case,  $\sim 2$  per cent. While the fit is superior to the previous coupled model, the model has too shallow of a high-energy hump to account for the data (Fig. 7).

Previous work (Georgantopoulos & Akylas 2019; Marchesi et al. 2019) had found the coupled MYTORUS model to provide a good fit to NGC 1194. Marchesi et al. (2019) found a column density that was consistent with that in our coupled MYTORUS fit. To make a direct comparison, we reduced the archived *NuSTAR* data from OBSID 60061035002 taken during 2015 Feb 28 (Table 1). Those data were reduced following the same method as for our new observation. We then applied the coupled MYTORUS model derived from our 2020 data and fit it to the 2015 *NuSTAR* data. We confirmed the adequacy of the coupled model for this short exposure using *NuSTAR* alone (cf. Georgantopoulos & Akylas 2019; Marchesi et al. 2019), this fit yielded  $\chi^2 = 115/92$  dof. We then took the coupled model fit to the archived 2015 data and re-applied it to our new joint *XMM* and *NuSTAR* data. This yielded  $\chi^2 = 3356/435$  dof. The coupled model fails to account for the data below 3 keV, where the observed flux is much higher than the model prediction. We return to this point in the discussion.

In conclusion, the coupled models provide a statistically unacceptable fit to the combined *XMM-Newton* and *NuSTAR* data, and a visual inspection of the fits shows this to be due to the fact that the curvature below 10 keV requires a lower column density than that which is required to explain the observed Fe K emission line and overall reflection signature strength.

### 5.2 Decoupled Model Fits

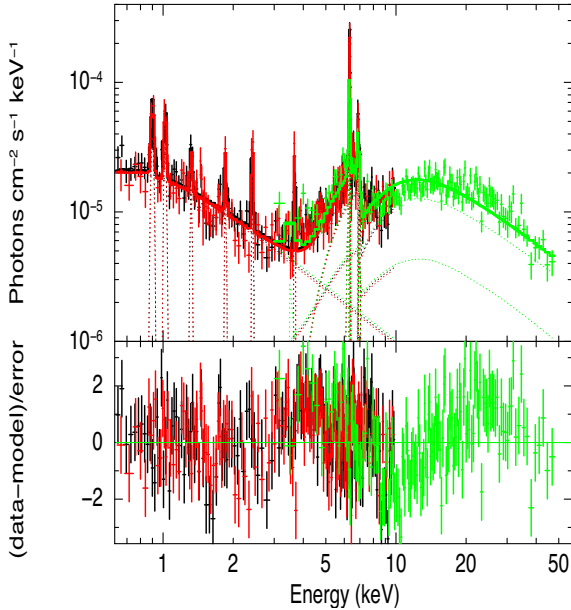
The standard, coupled configuration representing a toroidal reprocessor with an opening angle of  $60^\circ$  fails to account for the relative strength of the hard spectral excess. So-called decoupled models allow for the components in and out of the line of sight to be separated, such that one may view the transmitted continuum through a different column density than that from which the X-rays are scattered. So these column densities are independent in the fit. The toroidal global scattered gas  $-(MYTS + MYTL)$  that is out of the line of sight, is fixed at an inclination angle of  $\theta = 0^\circ$  and is viewed directly without attenuation. The direct or transmitted component (*MYTZ*) is fixed at  $90^\circ$  inclination and is attenuated by the full column of gas in the MYTORUS model.

The column densities of the line of sight and global gas were decoupled. As the strength of the globally-scattered component has been established to be strong in this source, we went directly to a fit that allowed the global gas component to be viewed at  $\theta = 0^\circ$ . This fit gave  $\chi^2 = 596/437$  dof. (Table 3 column 4, Decoupled 1). A photon

**Table 3.** Torus model fits.

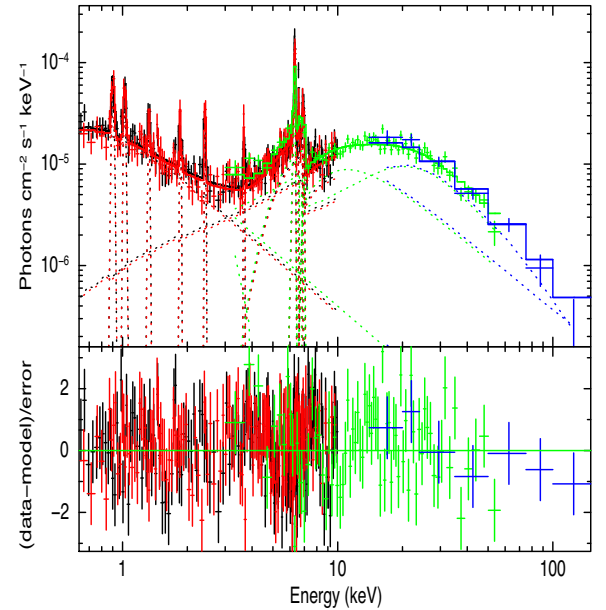
Par	COUPLED I	COUPLED II	DECOUPLED I	DECOUPLED II	Patchy torus: COUPLED III
$\Gamma$	1.40 (<1.41)	1.40 (<1.41)	$2.06^{+0.05}_{-0.05}$	$1.87^{+0.06}_{-0.08}$	$1.42^{+0.06}_{-0.0p}$
$-H_{\text{abs}}$	$0.72^{+0.02}_{-0.04}$	$0.84^{+0.04}_{-0.03}$	$1.20^{+0.08}_{-0.05}$	$0.81^{+0.10}_{-0.12}$	$1.54^{+0.50}_{-0.25}$ (85%) $0.34^{+0.11}_{-0.10}$ (15%)
$/10^{24} \text{ cm}^{-2}$					
$N_{\text{H, scatt}}$	0.72(l)	0.84(l)	9.91(>7.63)	9.02 (>4.18)	1.54(l) 0.34(l)
$/10^{24} \text{ cm}^{-2}$					
$e_{\text{scatt}}$	90°(f)	70°(f)	0(f)	0(f)	0(f)/0(f)
$W$	1.0(f)	1.0(f)	1.0(f)	$2.61^{+0.65}_{-0.45}$	1.0(f)/1.0(f)
$TO_{\text{unabs}}$	$2.91^{+0.25}_{-0.18}\%$	$2.00^{+0.20}_{-0.19}$	$0.21^{+0.05}_{-0.05}$	$0.92^{+0.51}_{-0.30}$	$1.51^{+0.30}_{-0.29}$
%					
$N_{\text{H,Gal}}$	$2.04^{+0.49}_{-0.44}$	$1.72^{+0.47}_{-0.43}$	$2.30^{+0.50}_{-0.0p}$	$1.95^{+0.49}_{-0.46}$	$1.02^{+0.48}_{-0.42}$
$/10^{21} \text{ cm}^{-2}$					
$\chi^2$	709/438	698/438	596/437	520/436	526/436

Notes. *p* indicates the parameter pegged at the preset limit. *f* indicates the parameter was fixed. *l* indicates the parameter was linked (see the text for details).

**Figure 7.** The *XMM-Newton*/*NuSTAR* data from 2020 with fit and residuals for Coupled Model II. Colours as for Fig. 5.

index  $\Gamma \sim 2$  was recovered. The global scattering gas has a column density  $>7.6 \times 10^{24} \text{ cm}^{-2}$ , while the line-of-sight absorption was  $N_{\text{H,MYT}} \sim 1.20^{+0.08}_{-0.05} \times 10^{24} \text{ cm}^{-2}$ . The relative flatness of the model compared to the data around 20 keV suggests that this configuration of model components underpredicts the strength of the reflected light component.

All fits thus far have had the normalization of the global-scattered light and lines tied directly to the illuminating continuum. Once a scaling constant is introduced for the reflected light and line emission, one can account for the inadequacy of the previous model. The best-fitting scalar was a factor of  $\sim 3$  above the nominal reflection strength (Table 3, column 5, Decoupled II), and this addition to the model yielded a significant improvement to the fit, with  $\chi^2 = 520/436$  dof (Fig. 8). In this case, a photon index  $\Gamma \sim 1.9$  was recovered. The

**Figure 8.** The BAT 105 month spectrum (blue points) compared to the *XMM-Newton*/*NuSTAR* data from 2020, and Decoupled Model II. Colours as for Fig. 5.

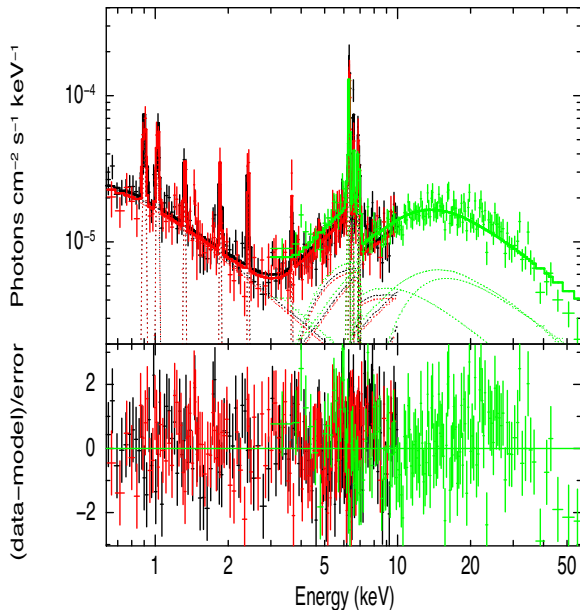
global reprocessing gas has a column density  $>4 \times 10^{24} \text{ cm}^{-2}$ , an order-of-magnitude higher than the line-of-sight absorption.

To check the source spectral constancy over long time-scales, we overlaid the average 105-month BAT spectrum on to the best-fitting MYTORUS decoupled Model II. The BAT data are consistent with the current epoch *NuSTAR* data, with no significant instrument cross-normalization required (Fig. 8). The intrinsic nuclear luminosities for this model are  $L_{2-10} = 4.1 \times 10^{42} \text{ erg s}^{-1}$  and  $L_{10-50} = 1.2 \times 10^{43} \text{ erg s}^{-1}$ .

### 5.3 Patchy torus model

The factor of 3 scaling factor for the scattered spectrum indicates the need for a more complex version of this model. One possibility is





**Figure 9.** *XMM-Newton/NuSTAR* data fit to the double torus model, Coupled III. Colours as for Fig. 5.

that a substantial fraction of the direct emission is hidden from our sightline by a Compton-thick absorber. In such a case, the scattered component results from a luminosity a factor of  $\sim 3$  higher than the observed X-ray continuum.

We therefore constructed a model that allows the intrinsic continuum to pass through two different sightlines, of independent column densities, to mimic a more complex case, i.e. a patchy torus reprocessor. The model therefore has two line-of-sight absorbing zones, each with global scattering gas that is coupled in column density to each of the line-of-sight absorbers. The global/scattering components are inclined at  $\theta = 0^\circ$ , as for some previous models.

This provided a good fit to the data (Table 3). A small amount (1.5 per cent) of continuum is allowed to leak through unabsorbed to explain the soft excess, as found for previous models. Of the remainder of the continuum, 85 per cent is viewed through a column density of  $\sim 1.5 \times 10^{24} \text{ cm}^{-2}$ , 15 per cent through a much lower column of  $\sim 3 \times 10^{23} \text{ cm}^{-2}$ . Each of these line-of-sight absorbers is allowed an associated global scattering component with a scaling factor frozen at 1.0. This fit works very well ( $\chi^2 = 526/436$ ), relieving the need for arbitrary scaling of the global scattered component.

## 6 DISCUSSION

The presence of Compton-thick absorption appears to be common in maser AGN (Greenhill, Tilak & Madejski 2008). The maser source NGC 1194 is no exception, with Compton-thick reprocessing gas revealed in *NuSTAR* (Masini et al. 2016) and *Suzaku* (Tanimoto et al. 2018) data. As also expected for maser AGN, NGC 1194 is highly inclined to our line of sight. There is significant reddening of the host galaxy continuum and the narrow-line region emission lines. The optical line details have led to a Seyfert 1.9 classification, indicating that a small amount of light can be seen directly from the optical BLR and, in turn, that the nuclear obscuring gas might not be homogeneous.

New data from simultaneous 2020 *XMM-Newton* and *NuSTAR* observations of NGC 1194 span the energy range 0.6–50 keV. The *XMM-Newton* spectrum shows many strong soft-band emission lines

that are a superposition of contributions from ionized extended gas mapped in *Chandra* images (Braitto et al., in preparation) and neutral lines, likely dominated by contributions from a toroidal nuclear reprocessor and the narrow-line region gas.

The coupled MYTORUS model was applied to account for the broad X-ray form of the combined 2020 data. The single coupled model configuration does not provide a good fit to the data as the curvature below 10 keV requires a lower line-of-sight column density than that which is required to explain the observed Fe K emission line and overall reflection signature strength from global gas out of the line of sight. Further to this, the single zone coupled models are not found to be a statistically acceptable fit to the data. While the coupled model provided an acceptable fit to the shorter, archived *NuSTAR* observation from 2015, we found that the model fit to that earlier epoch underpredicted the X-ray flux below 3 keV that was sampled by our *XMM-Newton/NuSTAR* campaign. The availability of simultaneous data down to lower X-ray energies drives the need for a more complex model than was previously published for this AGN.

Decoupling the column density for reflection and absorption greatly improves the MYTORUS fit, such that the shape of the model spectrum is well matched to the data. However, while decoupled II provides a good statistical fit, the X-ray continuum required to provide the strength of the global reflection observed, is a factor of  $2.61^{+0.65}_{-0.45}$  higher than the continuum observed (Table 3, column 5, fit parameter C). As NGC 1194 does not show a variable continuum, this enhanced reflection cannot be attributed to a previous, unobserved epoch of higher continuum. It may be that this higher continuum is present, but that Compton-thick gas is responsible for hiding the full continuum from direct view.

If, indeed, a large fraction of the continuum is hidden from view, then examination of the SED for NGC 1194 might help confirm that hypothesis. In a study of SEDs for a sample of Compton-thick AGN, Brightman et al. (2017) calculate the bolometric correction factors  $\kappa_{\text{BOL}} = \log \frac{L_{\text{BOL}}}{L_X}$  where  $L_{\text{BOL}}$  and  $L_X$  are the bolometric and the 2–10 keV luminosity. Taking the intrinsic X-ray luminosity to be  $L_{2-10} = 4.1 \times 10^{42} \text{ erg s}^{-1}$ , from the DECOUPLED II fit, and the bolometric luminosity from Brightman et al. (2017) to be  $L_{\text{BOL}} = 5.5 \times 10^{44} \text{ erg s}^{-1}$ , we estimate  $\kappa_{\text{BOL}} \sim 2.13$ . This value is similar to that found by Brightman et al. (2017) and higher than expected compared to the rest of their Compton-thick sample.

Brightman et al. (2017) explore the range of  $\kappa_{\text{BOL}}$  values exhibited, as a function of Eddington ratio ( $\lambda_{\text{Edd}}$ ). Based on the sample distribution and a value of  $\lambda_{\text{Edd}} = 0.08$  for NGC 1194 (Brightman et al. 2017), the value of  $\kappa_{\text{BOL}}$  should lie within 0.8–1.5 for NGC 1194. The initial estimation  $\kappa_{\text{BOL}} \sim 2.13$  is significantly higher than the expected range. If we apply a correction to the X-ray continuum luminosity, based upon the DECOUPLED II fit, then (propagating the uncertainties on that scalar, Table 3, column 5) we estimate a corrected  $\kappa_{\text{BOL}}$  in the range 1.61–1.79, much closer to the expected value for NGC 1194, based on the Brightman et al. (2017) sample study.

An alternative model using two zones of reprocessing gas allows representation of a more complex scenario, such as a patchy torus. In the context of such, 85 per cent of sightlines to the nucleus are through Compton-thick gas and 15 per cent through Compton-thin gas. Such a structure effectively gives partial-covering of the X-ray continuum, which is a common observational phenomenon in local AGN (e.g. Turner et al. 2011; Risaliti et al. 2009, and references therein). This result is consistent with the torus itself being intrinsically patchy, with a range of column densities (from a patchwork of clouds) viewed over different sightlines. Alternatively, the Compton-thin gas could arise

from the BLR region, located closer in and viewed through holes in a Compton-thick torus.

## 7 CONCLUSIONS

We have performed broad-band X-ray spectroscopy of a joint *XMM-Newton* and *NuSTAR* observation of the Seyfert 1.9 galaxy, NGC 1194. The source flux is steady and the spectrum shows strong-line emission and a very hard form. The spectrum can be described by torus models with solutions indicating the global column density to be an order of magnitude higher than the line-of-sight absorption. In such models, the reflection strength is a factor of  $\sim 3$  higher than would be nominally expected. Such a result may indicate that a large fraction of continuum photons are hidden from view in the *XMM-Newton*/*NuSTAR* bandpass. This idea is supported by the atypical SED of the source.

Alternatively, the spectral form can be modelled using two toroidal reprocessor components that could represent a more complex structure, such as a patchy torus. In this case, 85 per cent of the sightlines through this patchy torus are viewed through Compton-thick gas and the remaining 15 per cent through Compton-thin gas.

Both of these torus model solutions are based upon an assumption of a solar abundance of Fe. It is of particular interest that the application of the patchy torus model alleviates the need for a significantly greater intrinsic luminosity to be seen by the reprocessor, than is viewed by the observer. This is not an intuitive result and this insight could not have been obtained by the application of ad hoc models to these data.

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## DATA AVAILABILITY

This research has made use of data obtained from the High Energy Astrophysics Science Archive Research Center, provided by NASA's Goddard Space Flight Center.

## REFERENCES

Anders E., Grevesse N., 1989, *Geochim. Cosmochim. Acta*, 53, 197  
 Baloković M. et al., 2018, *ApJ*, 854, 42  
 Blustin A. J., Page M. J., Fuerst S. V., Branduardi-Raymont G., Ashton C. E., 2005, *A&A*, 431, 111

Braito V., Reeves J. N., Bianchi S., Nardini E., Piconcelli E., 2017, *A&A*, 600, A135  
 Brightman M. et al., 2015, *ApJ*, 805, 41  
 Brightman M. et al., 2017, *ApJ*, 844, 10  
 Brightman M., Nandra K., 2011, *MNRAS*, 414, 3084  
 Della Ceca R. et al., 2008, *A&A*, 487, 119  
 Elitzur M., Shlosman I., 2006, *ApJ*, 648, L101  
 Georgantopoulos I., Akylas A., 2019, *A&A*, 621, A28  
 George I. M., Fabian A. C., 1991, *MNRAS*, 249, 352  
 Greenhill L. J., Tilak A., Madejski G., 2008, *ApJ*, 686, L13  
 Harrison F. A. et al., 2013, *ApJ*, 770, 103  
 HI4PI Collaboration, 2016, *A&A*, 594, A116  
 Ikeda S., Awaki H., Terashima Y., 2009, *ApJ*, 692, 608  
 Kallman T., Bautista M., 2001, *ApJS*, 133, 221  
 Kallman T. R., Palmeri P., Bautista M. A., Mendoza C., Krolik J. H., 2004, *ApJS*, 155, 675  
 Kaspi S. et al., 2002, *ApJ*, 574, 643  
 Kuo C. Y. et al., 2011, *ApJ*, 727, 20  
 Liu Y., Li X., 2014, *ApJ*, 787, 52  
 Madsen K. K., Beardmore A. P., Forster K., Guainazzi M., Marshall H. L., Miller E. D., Page K. L., Stuhlinger M., 2017, *AJ*, 153, 2  
 Marchese E., Braito V., Della Ceca R., Caccianiga A., Severgnini P., 2012, *MNRAS*, 421, 1803  
 Marchesi S. et al., 2019, *ApJ*, 872, 8  
 Masini A. et al., 2016, *A&A*, 589, A59  
 Murphy K. D., Yaqoob T., 2009, *MNRAS*, 397, 1549  
 Nenkova M., Sirocky M. M., Nikutta R., Ivezić Ž., Elitzur M., 2008, *ApJ*, 685, 160  
 Osterbrock D. E., Ferland G. J., 2006, in Osterbrock D. E., Ferland G. J., eds, *Astrophysics of Gaseous Nebulae and Active Galactic Nuclei. Astrophysics of Gaseous Nebulae and Active Galactic Nuclei*. 2nd. ed. Univ. Science Books, Sausalito, CA  
 Reeves J. N., Nandra K., George I. M., Pounds K. A., Turner T. J., Yaqoob T., 2004, *ApJ*, 602, 648  
 Risaliti G. et al., 2009, *MNRAS*, 393, L1  
 Savage B. D., Mathis J. S., 1979, *ARA&A*, 17, 73  
 Severgnini P., Caccianiga A., Della Ceca R., 2012, *A&A*, 542, A46  
 Shull J. M., van Steenberg M. E., 1985, *ApJ*, 294, 599  
 Tanimoto A., Ueda Y., Kawamuro T., Ricci C., Awaki H., Terashima Y., 2018, *ApJ*, 853, 146  
 Tatum M. M., Turner T. J., Miller L., Reeves J. N., 2013, *ApJ*, 762, 80  
 Turner T. J., Miller L., Kraemer S. B., Reeves J. N., 2011, *ApJ*, 733, 48  
 Verner D. A., Ferland G. J., Korista K. T., Yakovlev D. G., 1996, *ApJ*, 465, 487  
 Véron-Cetty M.-P., Véron P., 2006, *A&A*, 455, 773  
 Wilms J., Allen A., McCray R., 2000, *ApJ*, 542, 914  
 Yaqoob T., 2012, *MNRAS*, 423, 3360  
 Yaqoob T., Murphy K. D., Miller L., Turner T. J., 2010, *MNRAS*, 401, 411  
 Yaqoob T., Tatum M. M., Scholtes A., Gottlieb A., Turner T. J., 2015, *MNRAS*, 454, 973

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