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Numerical simulations of radar echoes rule out basal CO₂ ice deposits at Ultimi Scopuli, Mars.

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Abstract

The principal objective of the radar sounder MARSIS experiment is to look for ice and water in the Martian subsurface. One particular focus of investigations, since 2005, has been the search for basal liquid water in the south polar layered deposits (SPLD). Anomalously strong basal echoes detected from four distinct areas at the base of the deposits at Ultimi Scopuli have been interpreted to indicate the presence of bodies of liquid water in this location, beneath a 1.5 km thick cover of ice and dust. Other explanations for the bright basal reflections have been proposed, however, including the possibility of constructive interference in layered media. Here, we test this mechanism through simulations of MARSIS radar signals propagating in models of CO₂-H₂O ice sequences. We then compare the results to real MARSIS data acquired over Ultimi Scopuli, finding that no CO₂-H₂O ice model sequence reproduces the set of real data. The results of our work have implications in relation to the global CO₂ inventory of Mars.

Keywords: Mars, SPLD, MARSIS, CO₂ ice, radar.

61 1. Introduction

62
63 The radar sounder MARSIS (Mars Advanced Radar for Subsurface and Ionospheric Sounding) on board the
64 European Space Agency orbiter Mars Express, has been probing the Martian subsurface since 2005 (Orosei et al.,
65 2015), exploring the global distribution of rock, soil, ice, and liquid water. The instrument operates within the range
66 1.3-5.5 MHz, transmitting a 1 MHz-bandwidth pulse centered at the frequencies of 1.8, 3, 4 and 5 MHz. Exploration of
67 the Martian subsurface is mostly carried out using the 3 and 4 MHz, or 4 and 5 MHz frequencies, for a nominal depth
68 of investigation of up to 5 km, depending on the electromagnetic properties of the probed material (Picardi et al.,
69 2004).

70 The primary objective for the MARSIS experiment is to look for buried ice and water in the Martian
71 subsurface. It was not until 2018, however, that evidence of bright subsurface reflections originating 1.5 km deep
72 below the surface of Ultimi Scopuli, one of the terminal regions of the South Polar Layered Deposits (SPLD), was
73 obtained (Orosei et al., 2018a) at coordinates 193°E - 81°S. The reflections, much brighter than those detected from
74 the surface, were interpreted as originating from basal ponded water or wet sediments distributed on an area
75 approximately 20 km across. The work reignited the scientific debate about the stability of liquid water at the Martian
76 polar regions, and specifically at the south pole, where average surface temperatures (Mellon et al., 2004) are
77 estimated to be 162 K. A possible explanation for the water being liquid at these conditions is that it contains a high
78 concentration of salts (e.g., perchlorates, chlorides) which are ubiquitous on Mars (Hecht et al., 2009) and have
79 eutectic temperatures as low as 198 K (Pestova et al., 2005). It was subsequently suggested that a high geothermal
80 gradient, such as that produced by a magmatic chamber at depth, is required to reach the eutectic temperatures of
81 concentrated salt solutions at the base of the deposits (Sori and Bramson, 2019). This scenario was disputed,
82 however, on the grounds that there is no evidence of an anomalous geothermal gradient in the region (Lauro et al.,
83 2021). Newly acquired MARSIS orbits of the region further revealed other three previously undetected areas of bright
84 reflections, each approximately 10 km across (Lauro et al., 2021). Because laboratory experiments have shown that
85 hypersaline aqueous solutions can form and persist over time at temperatures as low as 123 K, well below their
86 eutectic temperatures (Toner et al., 2014; Primm et al., 2017), a hypersaline origin of the water was deemed to remain
87 the most likely explanation for the bright reflections (Lauro et al., 2021). It is however difficult to envisage how large
88 bodies of supercooled brines could have persisted in the natural Martian environment over geological timescales
89 (Schroeder and Steinbrügge, 2021; Smith et al., 2021), which remains a significant argument against the liquid brine
90 interpretation. The possibility that bright reflections could be caused by scattering phenomena in a stratified medium
91 was also raised (Hecht et al., 2018). This criticism was however dismissed by Orosei et al. (2018b), who pointed out
92 that they had demonstrated that the strength of the basal echoes at Ultimi Scopuli cannot be explained by multiple
93 reflections. Even though a particular stratigraphic arrangement could produce enhanced basal echoes at one of the
94 MARSIS frequencies of detection, it did not do so at the other observation frequencies. Alternative models suggesting
95 that a range of materials, other than bodies of liquid water, could explain the sources of the strong basal echoes, have
96 recently been published by Bierson et al. (2021), and Smith et al. (2021). The merits of all proposed alternative
97 hypotheses have been reviewed by Schroeder and Steinbrügge (2021), who argued that none of the alternatives is
98 conclusively incompatible with bodies of liquid water as the source of the bright basal reflections, and that each model
99 opens up questions providing new opportunities for future investigations.

100 Distinct layers of CO₂ ice have been detected by radar sounder SHARAD surveys of the SPLD (Phillips et al.,
101 2011; Bierson et al., 2016). The layers, up to 1 km thick, are interlayered with thin (typically 10-40 m thick) water ice,
102 and it has been calculated that sublimation of these massive deposits of CO₂ ice would increase the Martian
103 atmospheric pressure by 100% (Bierson et al., 2016). It has been suggested that, if the conditions leading to formation

104 of the thick CO₂-H₂O ice deposits have occurred at other times, during the approximately 100 Myr history of formation
105 of the SPLD (Manning et al., 2019), then a specific CO₂ ice layering configuration might produce subsurface
106 reflections of higher intensity than those from the surface. In fact, Lalich et al. (2021a,b) obtained strong simulated
107 radar reflections at 4 MHz frequency, mimicking MARSIS waveform, for a model configuration involving 1-100 m thick
108 basal CO₂ ice layers sandwiching a dusty water ice layer of comparable thickness. Such configuration over the extent
109 of MARSIS bright area, would translate in an average ~ 2% of the current CO₂ atmospheric content. While this
110 specific contribution to the atmospheric volatile budget appears negligible, compounding scenarios involving storage
111 of volatiles within the south polar cap can potentially impact estimates of the Martian volatile budget, with
112 consequences on models of Martian climate cycles and atmospheric evolution, and on estimates of the amount of
113 volatiles stored in the Martian crust as part of the geological record.

114 The Martian polar caps are the most voluminous exposed reservoirs of volatiles on Mars. It is estimated that,
115 combined, the polar caps store two thirds of the entire water content of the planet (Catling, 2014; Smith et al., 2020;
116 Jakoski, 2021). While climate models have suggested the possibility that the polar layered deposits may also contain
117 significant amounts of CO₂ (Jakoski et al., 1995), estimates based on thermophysical properties of water and CO₂ ices
118 rejected such interpretation (Mellon, 1996). The thin veneer of seasonal and residual solid CO₂ on the surface of the
119 polar caps is effectively counted as part of the atmosphere (Jakosky, 2019). However, geomorphological (Gulick and
120 Baker, 1989; Baker et al., 1991; Baker, 2001; Ramirez and Craddock, 2018) and geochemical evidence (Bridges et
121 al., 2019) has been interpreted to indicate that early Mars was warmer and wetter than at present, prompting
122 suggestions that a thick CO₂ rich atmosphere existed at the time, to account for greenhouse heating of the planet
123 (Haberle et al., 1994). It has been argued that subsequent loss to space of the CO₂ (and other volatiles) does not fully
124 account for the present-day content of CO₂ in the current thin Martian atmosphere, and that massive quantities of CO₂
125 ought to be sequestered within one or more geological sinks on Mars (Forget et al., 2013). While deep-seated
126 carbonates could in principle accommodate large volumes of CO₂ (Kahn, 1985; Michalski and Niles, 2010; Wray et al.,
127 2016; Jakosky, 2019), as is the case on Earth, there is at present no conclusive evidence for a global deep-crustal
128 reservoir of carbonates on Mars (Edwards and Ehlmann, 2015). The discovery of massive deposits of CO₂ ice (Phillips
129 et al., 2011; Bierson et al., 2016) in the South Polar Layered Deposits (SPLD) by the radar sounder SHARAD, and
130 three-dimensional radar imaging volume calculations suggesting the deposits may contain more than $1.6 \cdot 10^4$ km³ of
131 dry ice (Putzig et al., 2018), have thus paved the way for renewed investigations into the possibility of the polar ice
132 caps as CO₂ sinks.

133 The global abundance of CO₂ on Mars is presently unknown. There are some generally accepted points,
134 however, that hold irrespective of the exact abundance and distribution of CO₂ on Mars. Mars had lost most of its CO₂
135 rich atmosphere by the end of the Noachian (~ 3.7 Ga) (Jakosky, 2019). The rate of input of juvenile gases into the
136 atmosphere through planetary outgassing (i.e., volcanic activity) is not known precisely, but it was higher during the
137 Noachian, waning in the Hesperian, becoming negligible in the Amazonian (Grott et al., 2011), so it would have been
138 insufficient to replenish the lost early atmospheric budget. It is not known how much of the initial Martian atmosphere
139 was left by the end of the Noachian (estimates range from ~ 0.1 bar to 1 bar; Catling, 2014), and how much CO₂ had
140 by then already been stored in crustal sinks, though it is clear that the overall trend since has been one of thinning of
141 the atmosphere to the present-day global value of ~ 6 mbar of pressure. If the massive CO₂ ice deposits (Phillips et
142 al., 2011) were to sublimate, they would add to the atmosphere as much as ~0.006 bar of CO₂ (Bierson et al., 2016).
143 Further compounding the contribution (however small) of putative additional reservoirs of CO₂ ice hidden in the SPLD,
144 could potentially impact estimates of the planetary CO₂ budget.

145 Here, we present the results of our comprehensive testing of model simulations against the database of
146 MARSIS data acquired over more than 10 years of surveys of the south polar region. We modeled the propagation

147 behaviour of electromagnetic waves at MARSIS operating frequencies (1.8-5 MHz) under a variety of conditions
148 involving different configurations of thin layers of solid CO₂, H₂O ice and dust. Within the parameter space we
149 explored, the models do not return values of reflection consistent with the observed MARSIS data. We thus reject the
150 notion that the bright basal reflections detected at Ultimi Scopuli may indicate the presence of thin layers of CO₂-ice at
151 the base of the SPLD.

152 153 **2. Geological Setting**

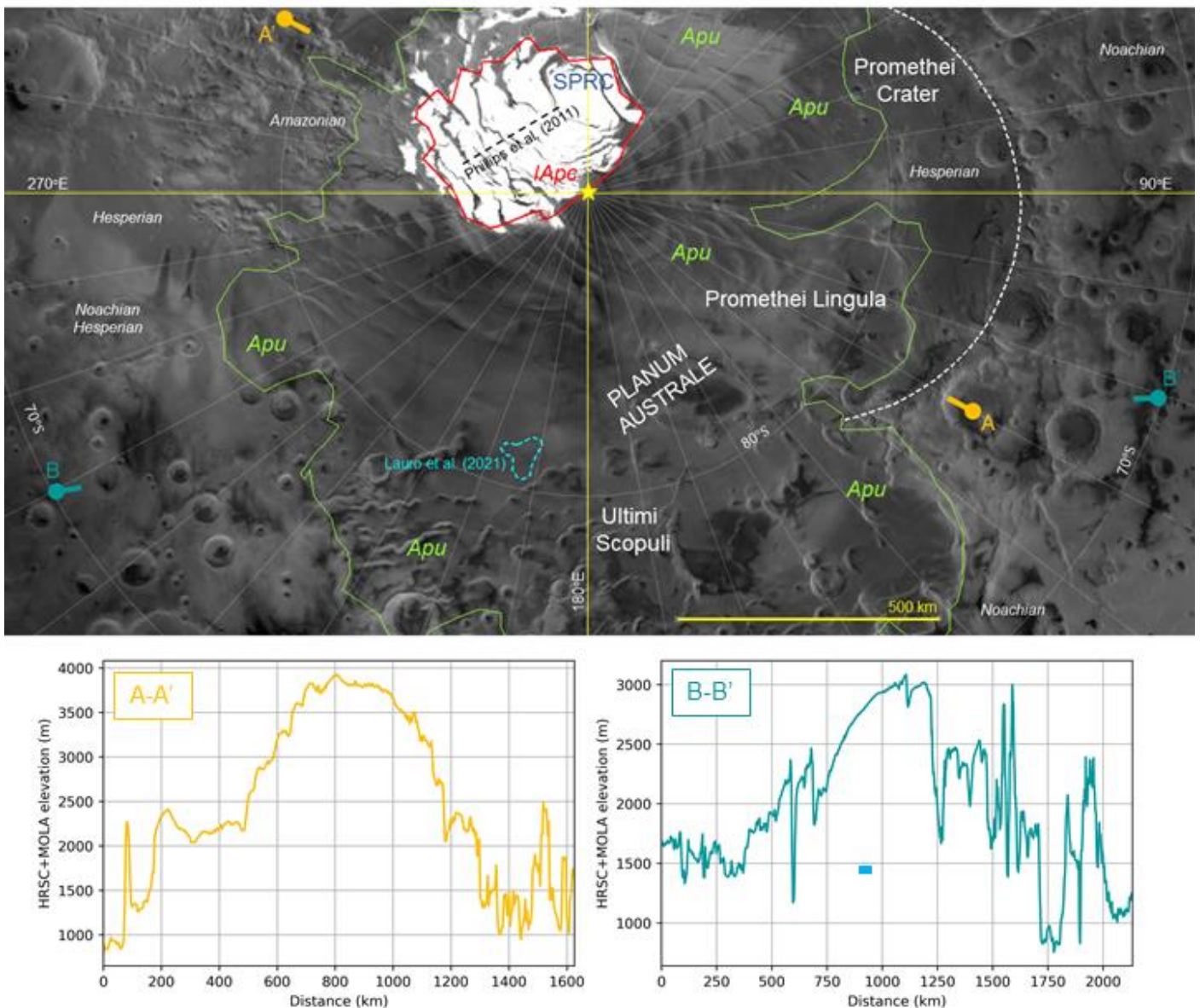
154
155 The South Polar Layered Deposits (SPLD), covering most of the Planum Australe of Mars (Fig. 1), form an up
156 to 3.7 km thick (Plaut et al., 2007) stratified unit, composed of pure water ice interlayered with thin dust-rich deposits.
157 It has been estimated that the content of dust in the SPLD is ~ 10-15% (Plaut et al., 2007; Zuber et al., 2007),
158 although inhomogeneities in the values of inverted density of the deposits indicate that the dust is not uniformly
159 distributed (Li et al., 2012), and that the fraction of dust in the water ice would vary depending on the presence of
160 denser CO₂ ice (Broquet et al., 2021). The layers are directly observed in radar imagery where troughs and scarps cut
161 through the deposits (Milkovich and Plaut, 2008). Owing to the electromagnetic contrast between the values of
162 dielectric permittivity ϵ of ice (3.12; Pettinelli et al., 2015) and dust (> 4; Mattei et al., 2014), layering is also detected
163 by radar sounder observations (Plaut et al., 2006, 2007; Milkovich et al., 2009; Whitten and Campbell, 2018). Layering
164 of the deposits is thought to reflect differential dust and ice accumulation rates influenced by global climate cycles
165 linked to orbital forcing (e.g., Hvidberg et al., 2012). The deposits decrease in elevation and thickness at the margins,
166 eventually waning and revealing underlying Noachian and Hesperian highlands terrains (Tanaka and Kolb, 2001),
167 overlain by a thin mantle of desiccated or water-ice pore-filled soil (Putzig et al., 2005; Jones et al., 2014). The surface
168 age of the SPLD is 7-100 Ma, based on the impact crater record (Herkenhoff and Plaut, 2000; Koutnik et al., 2002).
169 The age of the lowermost strata of the SPLD sequence is presently undetermined. A detailed geological analysis and
170 interpretation of the unconformities observed in Promethei Lingula suggests that the current extent of the SPLD may
171 have started forming no earlier than 250 Ma (Guallini et al., 2018). This estimate is consistent with time-series analysis
172 of HiRISE and CaSSIS stereo-images of the internal structure of the SPLD shown along scarp faces, suggesting that
173 the minimum time of accumulation of the SPLD is 10-30 Myr, resulting in a minimum age of deposition for the basal
174 layers of 17 to 130 Ma (Becerra et al., 2019).

175 The SPLD have been subdivided into three distinct structures (Milkovich and Plaut, 2008): the *Promethei*
176 *Lingula Layer Sequence* (PLL), the *Bench Forming Layer Sequence* (BFL), and the *Inferred Layer Sequence* (IL). The
177 BFL is the top sequence, geographically located in the central west-northwest portion of the SPLD, where the deposits
178 measure their highest elevations (~ 3500-4500 m). The BFL is covered by thin layers of CO₂ ice (part of the South
179 Polar Residual Cap, SPRC). Underlying the BFL is the PLL, while IL is inferred to be the basal sequence in the west
180 of the deposits. Elsewhere, only the PLL exists (ref. to figs. 18-19, in Milkovich and Plaut, 2008), its lateral extent
181 corresponding to that of Tanaka et al.'s (2014) *Apu* (Amazonian polar undivided) unit (Fig. 1).

182 The South Polar Residual Cap (SPRC), which is part of geological unit *IAPc* (Fig. 1), straddles the prime
183 meridian, approximately between longitudes 250°E-40°E and latitudes 85-90°S. The cap is primarily composed of a
184 high albedo thin (<10 m) layer of CO₂ ice (Byrne and Ingersoll, 2003), with a characteristic «Swiss cheese»
185 appearance (Malin et al., 2001) due to the presence of pits that reveal underlying water ice (Zuber et al., 2007).
186 Observed constant rates of erosion of the pits (Thomas et al., 2013) have been interpreted to point to net loss of CO₂
187 from the SPRC, possibly counterbalanced by observed net gain of seasonal water ice (Brown et al., 2014), suggesting
188 orbital effects are responsible for the CO₂ and H₂O atmosphere-SPRC exchanges. Localized deposits up to 1 km
189 thick, formed by alternating layers of CO₂ ice (each approximately 100 m thick) and H₂O ice (a few m thick) have also

190 been detected (Phillips et al., 2011; Bierson et al., 2016) by the radar sounder SHARAD, in troughs and at the base of
 191 scarps of the BFL structural unit, to the west of the location of the inferred rim of the ~ 870 km diameter Noachian
 192 impact crater Promethei Basin (Fig. 1). This suggests a possible relationship between the basal topography in this
 193 area, and atmospheric conditions that caused deposition of CO₂ ice here. It has been proposed that H₂O ice entombs
 194 the CO₂ ice layers, thus isolating them from the atmosphere and preventing their erosion (Manning et al., 2019). An
 195 alternative hypothesis argues that the deposits are in communication with the atmosphere, and that variations in polar
 196 incident sunlight from Mars's 100 kyr orbital cycles balance the exchange between CO₂ in the deposits and in the
 197 atmosphere. The observed deposits should thus record the last 510 kyr of Martian polar climate history (Buhler et al.,
 198 2021). Observations by SHARAD and MARSIS radars did not detect CO₂ ice layers anywhere else in the SPLD
 199 (Whitten et al., 2017; Whitten and Campbell, 2018; Khuller and Plaut, 2021).

200



201

202 **Fig. 1. Context map of the south polar region.** The map shows the extent of the geological units and geographic features
 203 described in Section 2. The dashed aqua polyline indicates the location of the strong basal reflections detected by Lauro et al.
 204 (2021). The trace of the ground track of the tract of SHARAD orbit 5968-01 traversing the RFZ₃ (Phillips et al., 2011) is shown as a
 205 black dashed line on the south pole residual cap (SPRC). Yellow star: South Pole. Vertical yellow line: prime and anti- meridians.
 206 Dashed white line: approximate position of remnant of rim of Promethei Crater. Lime green polyline: boundary of Amazonian polar
 207 undivided unit (*Apu*). Red polyline: boundary of late Amazonian polar cap unit (*IAPc*). A-A': end points of topographic profile A-A'
 208 (orange). B-B': end points of topographic profile B-B' (green). The short light blue segment in topographic profile B-B' indicates the

209 position and inferred depth of the bright basal reflections. Elevations are blended DEM/DTM (200 m/pixel) from MOLA and HRSC
210 data. Geological units are described by Tanaka et al. (2014). Geographic grid interval: 10°. Mapped with JMARS.

211
212 Ultimi Scopuli is a vast terminal region of the deposits, part of the *Apu* geological unit (Fig. 1) and of the
213 structural PLL sequence. Here, between coordinates 191°E-196°E and 80.3-81.5°S, and at a depth of approximately
214 1.5 km, recent orbiter sounder MARSIS observations (Orosei et al., 2018a; Lauro et al., 2021) revealed the presence
215 of four high permittivity areas, interpreted to be bodies of basal liquid water. The heat flow at the Martian south polar
216 regions is ~ 22-26 mW/m² (Grott et al., 2013; Plesa et al., 2018), and there is presently no evidence of an anomalous
217 geothermal gradient or localized hot spot (such as a magma chamber, as postulated by Sori and Bramson, 2019), that
218 would plausibly increase the basal temperature. The overburden pressure, estimated to be ~ 6 MPa (considering an
219 average density of the SPLD = 1220 kg/m³, from Zuber et al., 2007, which is the same as the best fit value published
220 by Broquet et al., 2021, and also consistent with a local value of ~ 1200 kg/m³, as inferred from the inverted density
221 map of the SPLD published by Li et al., 2012), would also be insufficient for ice to melt. Therefore, to explain the
222 presence of basal liquid water, Orosei et al. (2018a) suggested that the water may be a perchlorate brine.
223 Subsequently, Lauro et al. (2021) further defended this interpretation, on the grounds that supercooled perchlorate
224 brines have been shown to form and persist at sub-eutectic temperatures (Toner et al., 2014; Primm et al., 2017),
225 within the range of those reasonably expected at the base of the SPLD. The reflections at the base of Ultimi Scopuli
226 are located over a gently sloping basement (Lauro et al., 2021).

228 3. Model

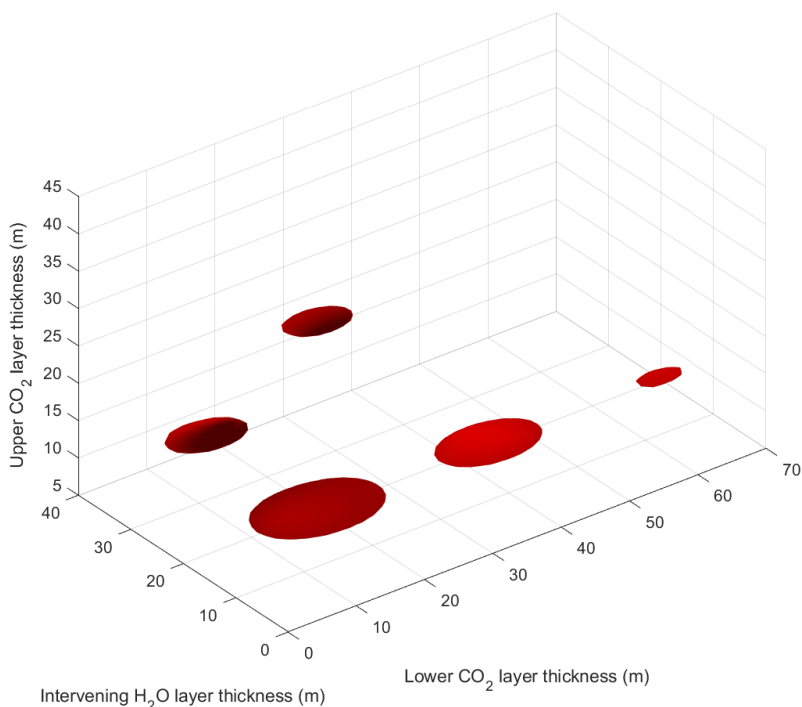
229 3.1. Preliminary Analysis

231
232 The fundamental concept underpinning this study is that the electromagnetic waves locally reflected from
233 interfaces between the layers in a stratified structure, can combine to enhance (constructive interference) or suppress
234 (destructive interference) the overall radar response from a layered structure. As a preliminary analysis, we explored
235 the specific parameter space presented by Lalich et al.'s (2021a) model, consisting of a thin layer of pure water ice
236 sandwiched between two thin layers of CO₂ ice: thus, we modelled the variation of only three parameters, i.e., the
237 thickness of each of the three layers. Layer thicknesses greater than a few tens of m dampen resonance effects,
238 because longer propagation paths decrease the likelihood of simultaneous in-phase sum of sinusoids at different
239 frequencies (Supplementary Figure 5, Orosei et al., 2018a). Therefore, we constrained the simulated thickness of
240 each layer to be ≤ 100 m. In the field of numerical electromagnetic modelling, spatial discretization is usually set to be
241 less than 1/10 of the wavelength. The shortest wavelength transmitted by MARSIS is 30 m in water ice. Thus, we
242 applied simulated thickness increments of 1 m as our sampling interval, to obtain accurate model stratigraphy
243 responses to the MARSIS pulses. We ran more than one million simulations (101 × 101 × 101) at MARSIS operating
244 frequencies of 3, 4 and 5 MHz, obtaining a complete characterization of the properties of Lalich et al.'s (2021a) model.

245 The highest value of subsurface-to-surface echo power ratio we obtained was 4.7 dB. Because the radar
246 response to the stratigraphy is described by sinusoidal functions, we ruled out the eventuality that we might have
247 missed peaks in the subsurface echo: considering that variations in the radar response between consecutive data
248 points are negligible (a fraction of a decibel at most), the model stratigraphy subsurface-to-surface echo power ratios
249 would never exceed ~ 5 dB, even in the case that the position of our maxima were not exact.

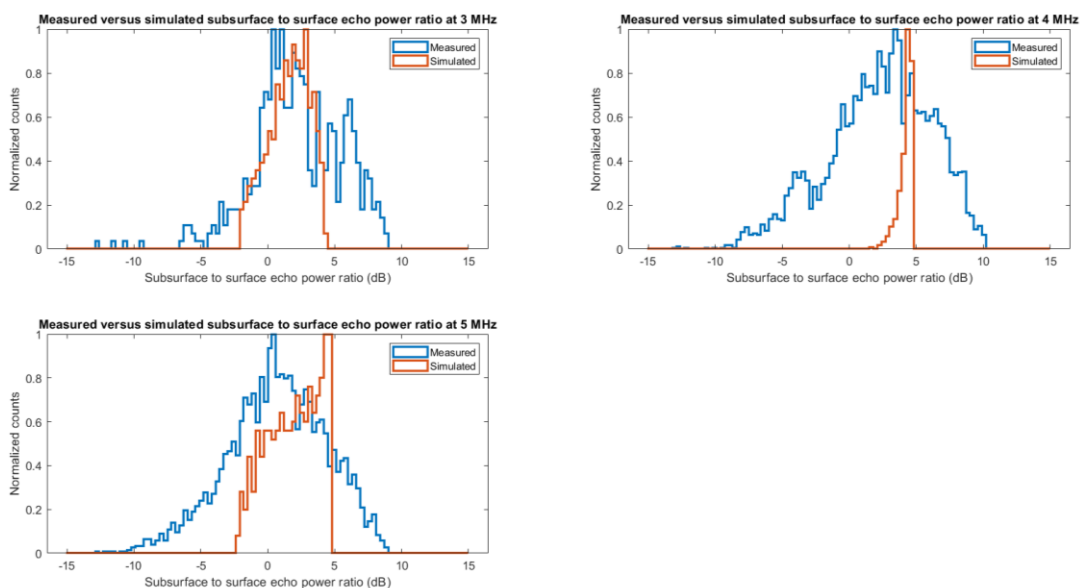
250 For each frequency, we obtained five well-defined regions in parameter space with subsurface-to-surface
251 radar echo power ratio values > 4 dB (Fig. 2). These regions are specific to each frequency, because constructive

252 interference phenomena are primarily controlled by layer thickness / wavelength ratios. Consequently, while basal
 253 echo powers remain high for small variations of thickness around the local maxima at one frequency, the
 254 corresponding echo powers at the other two frequencies are lower and much more widely dispersed (Fig. 3). Thus, if
 255 we take the parameter values corresponding to the maxima found at one frequency as a reference, and calculate the
 256 subsurface-to-surface echo power ratio differences between this frequency and the other two, these will always favour
 257 the reference frequency (Fig. 4).



258
 259 **Fig. 2.** Location of strong basal echoes (> 4 dB greater than surface echoes) produced in simulations at 4 MHz, displayed in the
 260 parameter space of the model by Lalich et al. (2021a), consisting of three quantities, specifically: the thickness of two CO₂ layers,
 261 and of a pure water ice layer sandwiched between them. Simulations at 3 and 5 MHz produce the same pattern, but with local
 262 maxima in different positions.

263



264

265 **Fig. 3.** Histograms of subsurface-to-surface echo power ratios obtained for simulations (orange) with a bottom CO₂ ice layer of
 266 thickness ranging between 8 m and 14 m, a top CO₂ ice layer 9-15 m thick, and an interbedded H₂O ice layer 9-15 m thick. Blue
 267 histograms represent real MARSIS observations.

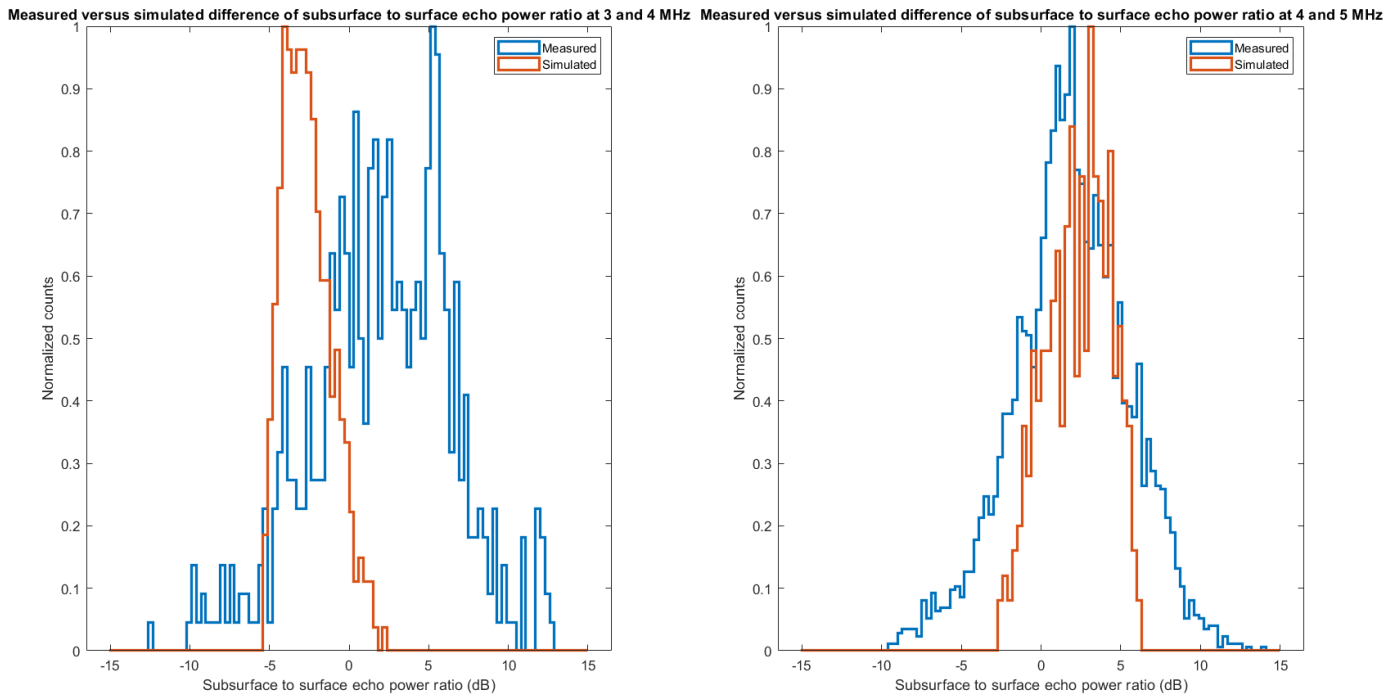


Fig. 4. Histograms of basal-to-surface echo power ratio differences between pairs of MARSIS operating frequencies (3 and 4 MHz, and 4 and 5 MHz). Simulation conditions are the same as for those in Fig. 3. Blue histograms are for real MARSIS observations.

Our preliminary analysis indicates that the distribution and behavior of subsurface-to-surface echo power ratios simulated in the parameter space as described, contrast sharply with that of real MARSIS data. However, idealized models do not consider possible sources of echo power fluctuations that affect real data, and thus some discussion is needed before a comparison between data and model is attempted.

The model used and discussed by Orosei et al. (2018) is computationally the same used by Lalich et al. (2021a, 2021b), that is the one-dimensional solution of Maxwell equations in a plane parallel stratigraphy. The former model had a simpler stratigraphy compared to the latter, with the SPLD represented as a single homogeneous layer composed of a mixture of H₂O ice and dust. The model was run varying basal temperature, dust content and basal permittivity. These three parameters controlled the attenuation within the SPLD and the intensity of the reflection at the bottom of the SPLD. The results showed that even for low amounts of dust and basal temperatures, it was necessary to have a high basal reflectivity to produce the observed strong basal echoes.

The model parameters published by Lalich et al. (2021b) are exactly the same, except for the input of single values for dust fraction, basal temperature and basal permittivity. These authors varied the thickness of the layers at the base of the SPLD.

Because of the similarity of the mathematical formulation of the two models, we compared them directly to one another to assess their capability to reproduce MARSIS observations, following the same approach used by Lalich et al. (2021b). For example, Figure 3 in that paper shows reflections produced by a single CO₂ layer model that are as strong as the ones observed by MARSIS, although the authors neglected to specify that their reference values for MARSIS echo power are the medians of observed values, and that normalized echo power can reach values as high as 10 dB (see Figure 4 and S4 in Orosei et al., 2018), well above any value that can be obtained by the single layer model shown in Figure 3.

The capability of the model in Orosei et al. (2018) to produce strong echoes is thus not due to a more complex mathematical implementation or to a greater number of parameters, but to the fact that basal permittivity can have

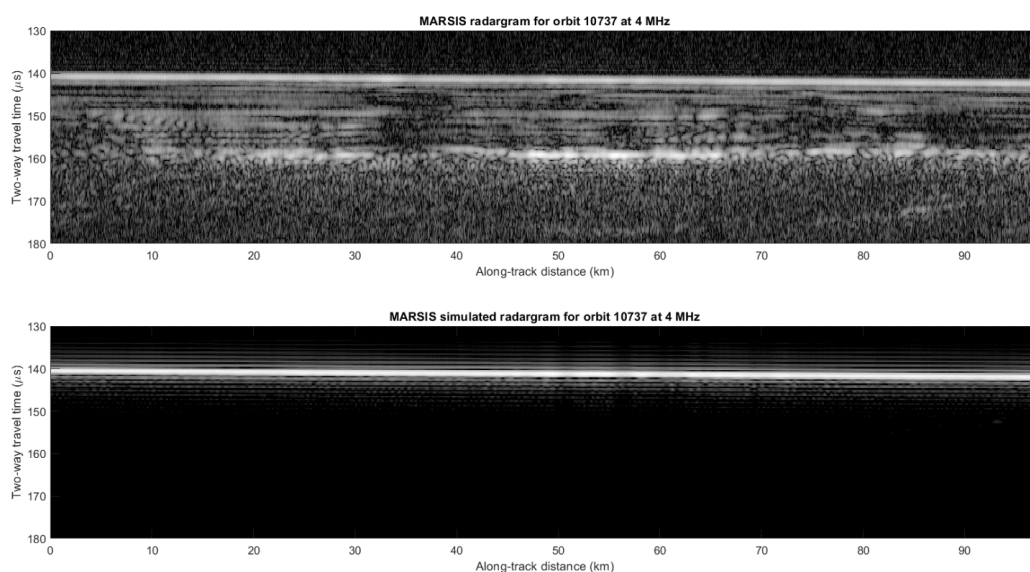
295 higher values than those considered by Lalich et al. (2021a, 2021b). In fact, Orosei et al.'s (2018) model could
296 produce basal echoes spanning the full range of observed values (see their Figures 4 and S4), whereas Lalich et al.'s
297 (2021a, 2021b) resulted only in a few instances of basal echo power comparable to observations. The purpose of the
298 present work is therefore to establish the conditions (if these exist), by which Lalich et al.'s (2021) model reproduces
299 the strong radar echoes reported by Orosei et al. (2018). We thus performed comprehensive simulations based on the
300 same mathematics, using a range of possible parameters beyond that explored in that earlier work.

301 In the following sections, we discuss the sources of variability and error affecting MARSIS observations,
302 evaluating their effect to derive a model of the uncertainties that can be applied to model results, for a more realistic
303 comparison between simulations and data.

305 *3.2 Idealized vs. actual SPLD surface geometry*

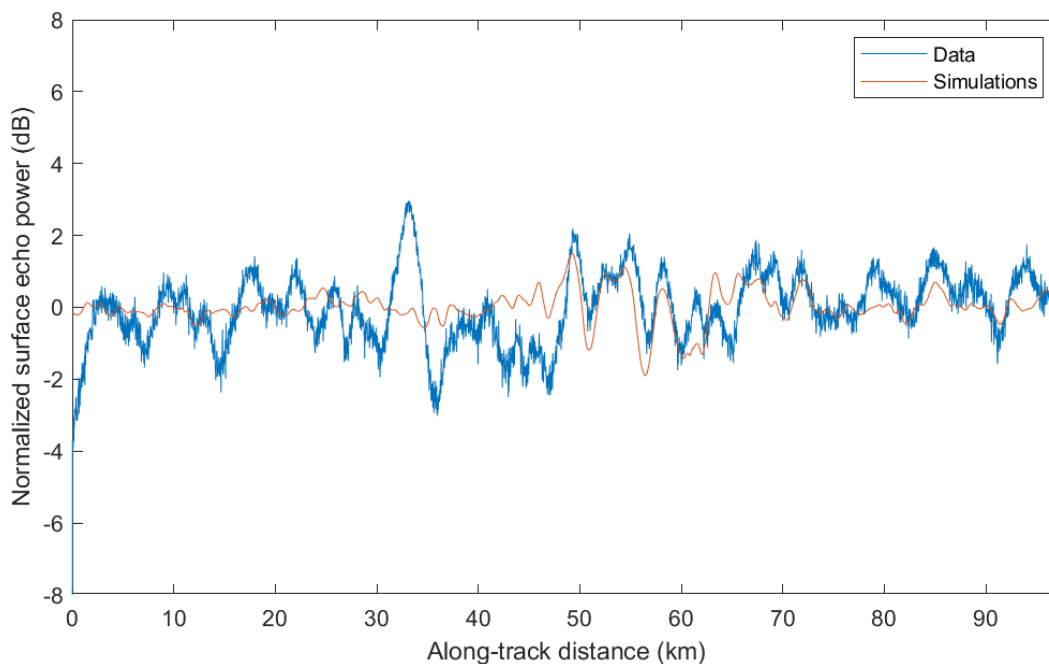
306
307 The assumption of a perfectly flat surface does not prevent the possibility of a comparison with real data, as
308 already shown by Orosei et al. (2018). Furthermore, the subsurface echo power for a given orbit is always normalized
309 to its median value, in order to prevent local minima of surface echo power from producing unrealistically high
310 normalized subsurface echo power values. The choice of using surface median power implies the assumption that the
311 SPLD surface is locally homogeneous, and that surface echo power fluctuations are due to small perturbations of this
312 assumed homogeneity.

313 To verify if this assumption holds, we need to examine the properties of echo power measured by MARSIS.
314 To this end, we retrieved the same data employed to produce Figure 2 in Orosei et al. (2018) and made publicly
315 available at the time of publication of that paper (Orosei and Cicchetti, 2018). We also produced a simulation of
316 surface echo power similar to those presented and discussed in Orosei et al. (2015). Simulations are used to validate
317 subsurface echo detection by visual comparison: all features present in both the real and simulated radargrams are
318 interpreted as surface clutter, and are therefore excluded from interpretation. A comparison between measurements
319 and simulations for orbit 10737 is shown in Figure 5.



321
322 **Fig. 5.** Top: MARSIS radargram for orbit 10737 at 4 MHz, redrawn following Orosei et al. (2018), versus the corresponding
323 simulation of surface scattering produced using the MOLA topographic dataset (bottom).

325 The simulated radargram allows also to estimate surface power fluctuations due to surface roughness alone,
326 sampled at the resolution of the MOLA topographic dataset. A comparison with the measured surface power (both
327 arbitrarily normalized to a 0 dB median power for ease of comparison) is shown in Figure 6: it is clear that surface
328 roughness alone accounts for the surface echo power fluctuations at distances between 50 and 60 km along track,
329 which correspond to the location of the bright subsurface reflector.
330



331
332 **Fig. 6.** Measured vs. simulated surface echo power for orbit 10737 at 4 MHz, both arbitrarily normalized to a 0 dB median power for
333 ease of comparison.

334
335 The residual difference between data and simulations can be attributed to several factors, such as: surface
336 roughness at scales below MOLA resolution or variations in near-surface density and composition, but the exact
337 cause is difficult to ascertain with available data. Residual power fluctuations are found to have a standard deviation of
338 less than 1 dB, corresponding to a fluctuation of 0.5 in the value of the relative dielectric constant in the assumption
339 that its median value is comprised between 3 and 4 (typical of SPLD materials). Because of the small uncertainties
340 resulting from this analysis, we consider the ratioing with the median surface power as a reasonable calibration
341 method for subsurface echo power.

342 Furthermore, a previous analysis of SHARAD echoes by Grima et al. (2012) over this area (called Reference
343 Zone or RZ in the paper because of its flatness) showed that surface reflections are predominantly coherent, implying
344 that scattering is near-specular and thus approximates that of a plane surface. This conclusion is even truer at
345 MARSIS wavelength, as the effect of roughness on scattering does not depend on RMS topographic height, but rather
346 on the ratio between RMS topographic height and the wavelength.

347 It could be argued that even a small roughness causes a decrease (however small) of the mean power of
348 surface echoes, leading to an overestimation of the subsurface-to-surface echo power ratio. In fact, the opposite is
349 true: electromagnetic models and three-dimensional propagation simulations (e.g., Jonard et al., 2019) demonstrate
350 that the wavefront of the radar pulse penetrating through the subsurface is disrupted more than the echo
351 backscattered at the surface, because it crosses the rough surface twice. Thus, the subsurface to surface echo power
352 ratio calculated assuming no surface roughness effect, should be considered a minimum value, with the real ratio
353 likely higher.

354

355 3.3 Signal distortion and noise

356

357 The MARSIS pulse is not an ideal chirp. Its actual shape and power are unknown: in fact, the spectra reported
358 by Jordan et al. (2009) do not include the frequency-dependent effects of antenna gain, because these could not be
359 measured on ground owing to the length of the antenna. A test comparing the results obtained using Jordan et al.'s
360 (2009) chirp spectra with those obtained using the ideal chirp, does not demonstrate any significant difference.

361 Therefore, the MARSIS data have been range-compressed using an ideal chirp since the beginning of the mission.
362 The range compression is not optimal, with both resolution and signal-to-noise ratio affected. This, however, does not
363 affect the ratio of subsurface to surface echo power. In an abstract mathematical representation, radar sounding can
364 be described as the convolution of the radar waveform with the ideal electromagnetic response of the observed scene.
365 In Fourier transform terms, this convolution is the multiplication of the stratigraphy spectral response with a band-
366 limited signal spectrum, which clips the spectral response bandwidth and produces a limited resolution time-domain
367 response. All these mathematical operations are linear, so the shape of the pulse spectrum will change the power
368 level of echoes from two different interfaces, but not their ratio, which is the quantity used in this analysis.

369 A few available data also point to a possible dependence of gain on temperature, although a full
370 characterization of this effect was never performed, because it was considered to be negligible under the conditions at
371 which MARSIS acquires its observations. Furthermore, MARSIS raw data takes are at most 25 seconds long, which is
372 too short a time to produce any significant gain variation. This is further corroborated by the plot of surface echo power
373 shown in Figure 6, which does not show any trend of varying gain across the entire observation. Thus, within the limits
374 of what it is technically possible to ascertain about the causes and extent of non-ideality of range compression in
375 MARSIS, there is no effect that should be considered as introducing a bias in the use of median surface echo power
376 as a calibration reference for subsurface echo power.

377 Plots of the noise level in the observations (Orosei and Cicchetti, 2018) analysed by Orosei et al. (2018) show
378 constant noise level across a single observation, and nearly-constant noise levels across observations at the same
379 frequency (around -5 dB at 3 and 4 MHz and -11 dB at 5 MHz when expressed in uncalibrated digital numbers). The
380 median signal-to-noise ratio is > 20 dB for all observations, exceeding 30 dB.

381 Interferences caused by the electronics of the Mars Express spacecraft were characterized in the
382 commissioning phase of MARSIS. Limited available documentation shows a correlation with the activity of the HRSC
383 camera, though no systematic analysis has ever been published. Overall, data acquired over the lifetime of the
384 mission show that the combined electromagnetic interferences (EMI) caused by orbiter navigation and experiment
385 instrumentations, have a much smaller effect on the functioning of MARSIS than (for example) on its Mars
386 Reconnaissance Orbiter companion SHARAD. Thus, the operating assumption is that EMI have a negligible effect on
387 the subsurface to surface echo power ratio.

388 The MARSIS data analyzed by Orosei et al. (2018) consist of near-simultaneous observations at two
389 frequencies, namely 3 and 4 MHz, or 4 and 5 MHz. Thus, it is possible to compute the subsurface to surface echo
390 power ratio at two frequencies for the same observation. Doing so for data in which surface echo power has a
391 standard deviation below 1 dB (as noted in Table S1 in the Supplementary Materials of Orosei et al., 2018), one is left
392 only with data at 4 and 5 MHz. The difference between normalized power at those two frequencies has a median
393 value of 1.36 dB, which can be used for a back-of the envelope computation based on equations found in Porcello et
394 al. (1974) to obtain an estimate of the loss tangent (see also Orosei et al., 2020 for the effect of such loss tangent
395 value on SHARAD measurements). The resulting value is approximately 0.03, which is about fifteen times greater
396 than the one used by Lalich et al. (2021b). Because this value is affected by a large error (the standard deviation of

397 the difference between normalized power at the two frequencies is almost 4 dB), we did not use the estimated loss
398 tangent in simulations. Instead, we considered a variable attenuation as one of the factors decreasing the measured
399 ratio between subsurface and surface power produced by a model.

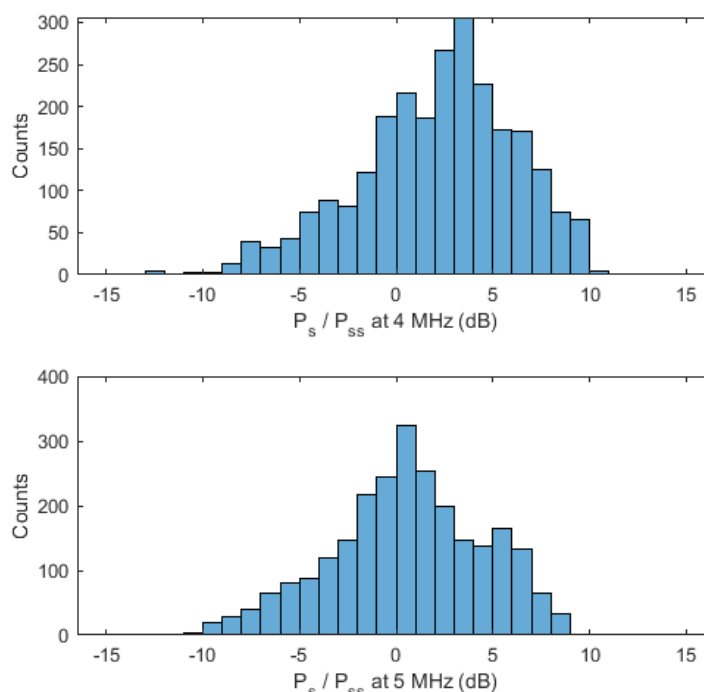
400

401 3.4 Other factors affecting subsurface to surface echo power ratio

402

403 Factors affecting subsurface echo power not included in Orosei et al.'s (2018) models are: surface roughness
404 (as discussed above); subsurface roughness (affecting basal echo intensity only); fraction of the radar footprint
405 covered by the basal reflector (see for example Haynes et al., 2018, for a discussion); signal attenuation within the
406 SPLD due to dielectric losses (as discussed above); signal attenuation within the SPLD due to reflection losses in the
407 thinly layered stratigraphy (see for example Courville et al., 2021, for a detailed modelling of propagation in a layered
408 medium at SHARAD frequencies). Their combined effect is demonstrated by the large dispersion of values observed
409 (Figure 7).

410



411

412 **Fig. 7.** Histograms of measured values of subsurface to surface echo power ratio over the bright reflector analysed in Orosei et al.
413 (2018) at 4 MHz (top) and 5 MHz (bottom).

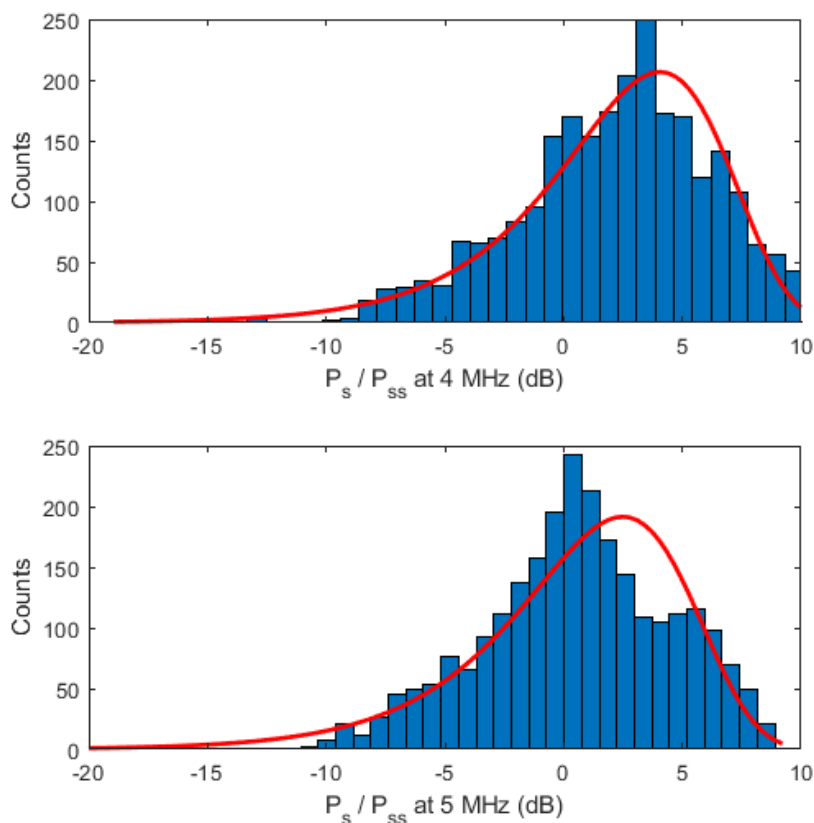
414

415 The histograms are asymmetric, with a cut-off at high values and a long tail towards low ones. This
416 characteristic matches the intuitive notion that the intensity of a basal reflection cannot increase indefinitely, and would
417 thus have a maximum possible value obtained when all the factors contributing to decreasing basal echo intensity as
418 discussed above are minimized. Thus, a comparison between simulations and data requires the introduction of a
419 model for the random dispersion of measurements. As it is not possible to quantify separately the importance of each
420 factor contributing to basal echo power attenuation, we derived a model probability density function directly from the
421 data (Fig. 8).

422

423 The long left-sided tails of the histograms in Figure 7 are characteristic of generalized extreme value
distributions, often used to model the maxima of sequences of random variables. Although this is a heuristic method

424 for the selection of the model distribution, it has the merit of matching data histograms better than other parametric
425 distributions.
426

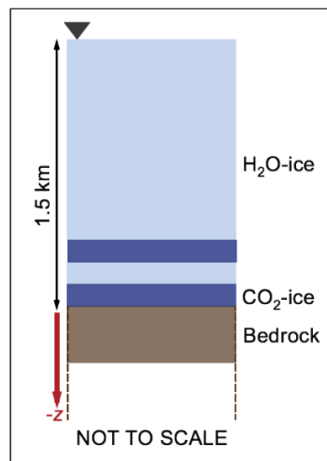


427
428 **Fig. 8.** Best-fit extreme value distributions plotted over the histograms shown in Figure 7.

429
430 Through the remainder of this paper, we use these best-fit distributions as the template for randomly
431 generating synthetic measurements from models.

432
433 *3.5. Electromagnetic propagation model*

434
435 The propagation of radar waves across stratified media is described by non-linear equations dependent on
436 the electromagnetic properties and thickness of the layers, as well as the probing frequency. Here, we model
437 electromagnetic behaviour of stratified CO₂ ice and dust-laden H₂O ice, by assuming a 1D geometry (Lombardo et al.,
438 2000; Pettinelli et al., 2003; Kartashov et al., 2004; Nunes and Phillips, 2006; Courville et al., 2021). Specifically, we
439 simulate the propagation of radar pulses through twelve idealized models of the SPLD, consisting of sets of perfectly
440 parallel planar interstratified basal layers composed of CO₂ ice and H₂O ice, where the latter has variable contents of
441 dust. The total thickness of the simulated deposits was set at the value derived through MARSIS radargrams over the
442 study area (1.5 km). The material underlying the deposits was set to be basaltic rock. The simulations were run by
443 code adapted from MATLABTM scripts publicly available through a public research data repository (Orosei and
444 Cicchetti, 2018). The code computes the complex reflectivity of the layer stack across the range of frequencies within
445 the pulse bandwidth, multiplying it by the signal spectrum and Inverse-Fourier transforming to the time domain.
446 Simulations were run for one, two or three CO₂ ice layers, the lowermost of which is in direct contact with the bedrock,
447 while the remaining ones are separated by layers of H₂O ice of variable thickness (Fig. 9). CO₂ ice layers are



449

450 **Fig. 9. Schematic stratigraphy used in numerical simulations.** The model stratigraphy is here shown as a set of planar parallel
 451 layers of CO₂ ice (dark blue) and H₂O ice (light blue), bounded by an upper half space representing the Martian atmosphere (as
 452 semi-infinite vacuum) and a lower half space representing the Martian bedrock (brown). Simulations are run for one, two or three
 453 lossless CO₂ ice layers, the lowermost of which is in direct contact with the bedrock, while the remaining ones are separated by
 454 layers of H₂O ice of variable thickness. The topmost layer of H₂O ice represents the bulk of the stratigraphy. Layer thicknesses of
 455 CO₂ ice and interlayer dusty water ice range from 0 to up to 100 m, with total deposit thickness = 1.5 km. The red arrow indicates
 456 the depth direction (-z). The inverted triangle on top of the sequence indicates the surface of the deposits. The sketch is drawn not
 457 to scale.

458

459 Typically, complex relative dielectric permittivities are strongly dependent on the frequency of the propagating
 460 electromagnetic wave and the temperature of the medium. Attenuation of the signal is independent of frequency at
 461 MARSIS wavelength for pure water ice: we used the empirical formulae reported by Mätzler (1998) for the average
 462 surface temperature value of 160 K (Clifford, 1987) taken as a constant across the entire model stratigraphy. We find
 463 that the real part of the dielectric permittivity of water ice is 3.11, with loss tangent values between $2.5 \cdot 10^{-8}$ (3 MHz)
 464 and $3.7 \cdot 10^{-8}$ (5 MHz). The loss tangent of CO₂ ice is not well constrained at the temperature of interest, with
 465 extrapolations and upper limits only being reported in the published literature. In our computations therefore, we set it
 466 to be = 0, to produce the highest possible subsurface-to-surface echo power ratio values. We did not consider the
 467 effect of the complex part of the permittivity of CO₂ on the reflection coefficient, because it is generally assumed to be
 468 low ($\leq 10^{-3}$). Computing the Fresnel reflection coefficient at normal incidence for an interface between a layer of CO₂
 469 ice and an underlying bedrock with permittivity values such as those used in the simulations, the results, with and
 470 without the inclusion of the imaginary part, differ by < 0.00001 .

471

472 3.6. Simulations

473

474 To calculate the simulated radar waveforms, in our model stratigraphies we assumed perfectly smooth
 475 surfaces, an assumption that holds also for small roughness of the illuminated surface, where “small roughness” is
 476 defined (Ulaby et al., 1981) by a height $< \lambda / 32$, which gives a value of 2.35 m at 4MHz. We considered the
 477 electromagnetic wave transmitted by the radar having a normal incidence with respect to the model stratified structure
 478 (ref. to sketch in Fig. 9), and assumed a 1D-model where the only variation of the electromagnetic properties occurs
 479 along the z-axis (depth). Furthermore, we assumed spatially homogenous layers with a constant complex dielectric
 480 permittivity (ϵ_i) and thickness (h_i). The stratified structure is bounded by two half-spaces representative of the
 481 atmosphere and the bedrock. This model is the same adopted in previous work (Orosei et al., 2018a; Lauro et al.,
 482 2019, Lalich et al., 2021a, 2021b). The received signal $y(t)$ is thus computed as:

483

$$y(t) = \mathcal{F}^{-1}[X(f)R(f)], \quad (1)$$

where \mathcal{F}^{-1} is the inverse Fourier transform, $X(f)$ is the Fourier transform of the signal transmitted by MARSIS antenna, $R(f)$ is the frequency response ($\Gamma_1(f)$) of the layered structure computed from the recursive scattering function $\Gamma_i(f)$:

$$\begin{aligned} R(f) &= \Gamma_1(f) \\ \Gamma_i(f) &= \frac{\rho_i + \Gamma_{i+1} e^{-2jk_{i+1}h_{i+1}}}{1 + \rho_i \Gamma_{i+1} e^{-2jk_{i+1}h_{i+1}}} \end{aligned} \quad (3)$$

with k_i the complex-valued wavenumber of the i -th layer, and ρ_i the Fresnel reflection coefficient at the boundary between layer $i-1$ and i , given by:

$$\rho_i = \frac{\sqrt{\varepsilon_{i-1}} - \sqrt{\varepsilon_i}}{\sqrt{\varepsilon_{i-1}} + \sqrt{\varepsilon_i}} \quad (4)$$

The geometry of the problem is modelled as a layered structure, where the upper and lower half space are the same for each simulation, whereas the assumed stratified layer between the atmosphere and the bedrock changes on the basis of the number of CO₂ ice layers, and the thickness of the layers. The ratio between basal and surface power is determined by considering the ratio between the maximum of the square amplitude of $y(t)$ at a time interval containing the echo return from the surface, and that occurring at a time interval containing the echo return from the upper side of the shallower CO₂ ice layer.

Using a multiprocessor machine, we simulated the three CO₂ layer model for every layer thickness between 0 and 50 m at 1-m intervals for three CO₂ layers and two intervening H₂O layers, for a total of 51⁵ simulations. This allowed us to sample completely the parameter space of the models with two- and three-basal CO₂ layers, and thus to locate and explore the properties of local maxima embedded in the three- and five-parameter spaces respectively. We identified local maxima of subsurface to surface power ratio, both at 4 and 5 MHz, and for each of them we extracted the basal echo power ratio for all simulations whose layers are within 2 m of the thickness of those corresponding to the local maximum, thus simulating the effect of random but limited variation of the stratigraphy.

4. Results

To implement a quantitative comparison between data and simulation results, for each local maximum of subsurface to surface echo power ratio we produced the histogram of power values extracted from its neighbourhood, and convolved it with the template distributions shown in Figure 8, obtaining a simulated histogram which represents the global effect of the factors affecting propagation within the SPLD on model results. An example of a synthetic distribution is shown in Figure 10.

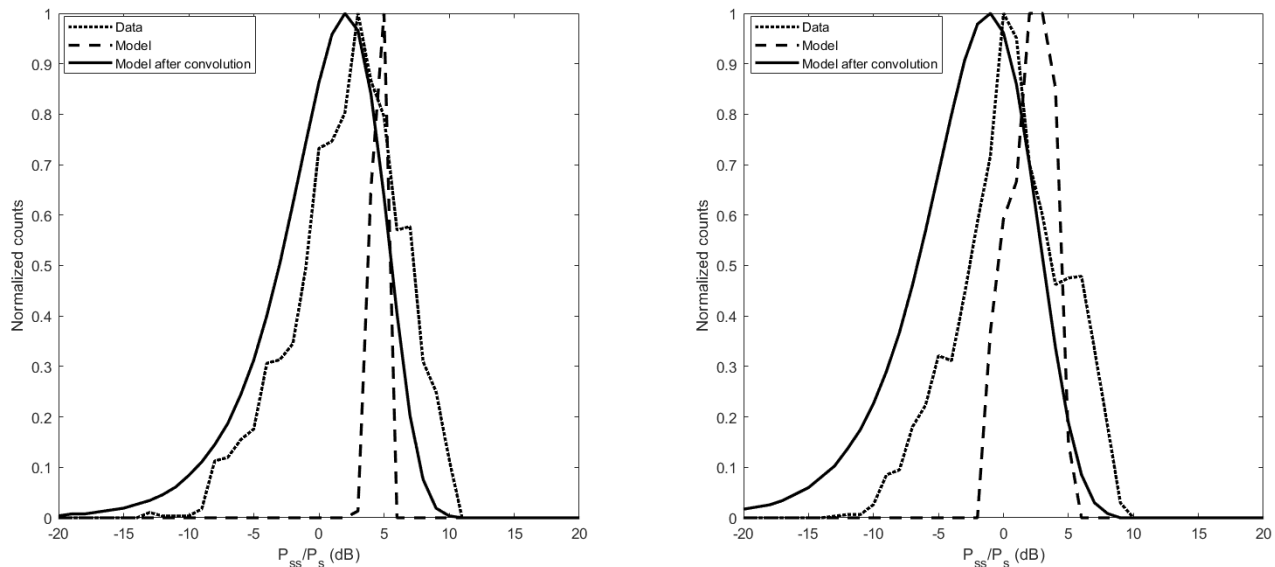


Fig. 10. Comparison between the histogram of measured values of subsurface to surface echo power ratio at 4 (left) and 5 MHz (right) with the histogram of model results in the neighbourhood of a given local maximum (see text for details), and with the synthetic histogram of model results produced through the convolution of the histogram of model results with the distribution template shown in Figure 8. The histograms are computed both for model results at 4 and 5 MHz.

To make a quantitative comparison of model results with data, the synthetic histograms were used to generate a set of random values representing measurements. This was done by taking each histogram bin, and generating a set of uniformly distributed random numbers in an interval corresponding to the bin, whose quantity is proportional to the bin height. Finally, the set of synthetic measurements was compared to actual measurements through a statistical test to determine if the two sets of samples could be produced by the same distribution. We used the two-sample Kolmogorov-Smirnov test (Kolmogorov, 1933; Smirnov, 1939), which is a nonparametric test for the null hypothesis that the data in two sets of samples are from the same continuous distribution. It is widely employed in statistics for its robustness and its sensitivity to both location and shape of distributions. The test was performed using the implementation available in the MATLAB^(TM) proprietary programming language, and setting the significance level for the rejection of the null hypothesis at 5%.

A configuration of layers producing a local maximum at one frequency does not result in a local maximum at the other frequency. For this reason, both local maxima at 4 MHz and at 5 MHz had to be considered. Thus, for a given model (either two- or three-CO₂ ice layers) we simultaneously tested samples of model results at 4 and 5 MHz extracted for a set of layer thicknesses corresponding to a local peak of basal echo power at either 4 or 5 MHz, and for all layer thicknesses within two meters of the one producing the peak. This resulted in four possible combinations of number of layers and frequency of the echo power maximum. We found and tested 487 local maxima at 4 MHz and 687 at 5 MHz for the two-layer model, and 2981 local maxima at 4 MHz and 5232 at 5 MHz for the three-layer model. Results are shown in Figure 11.

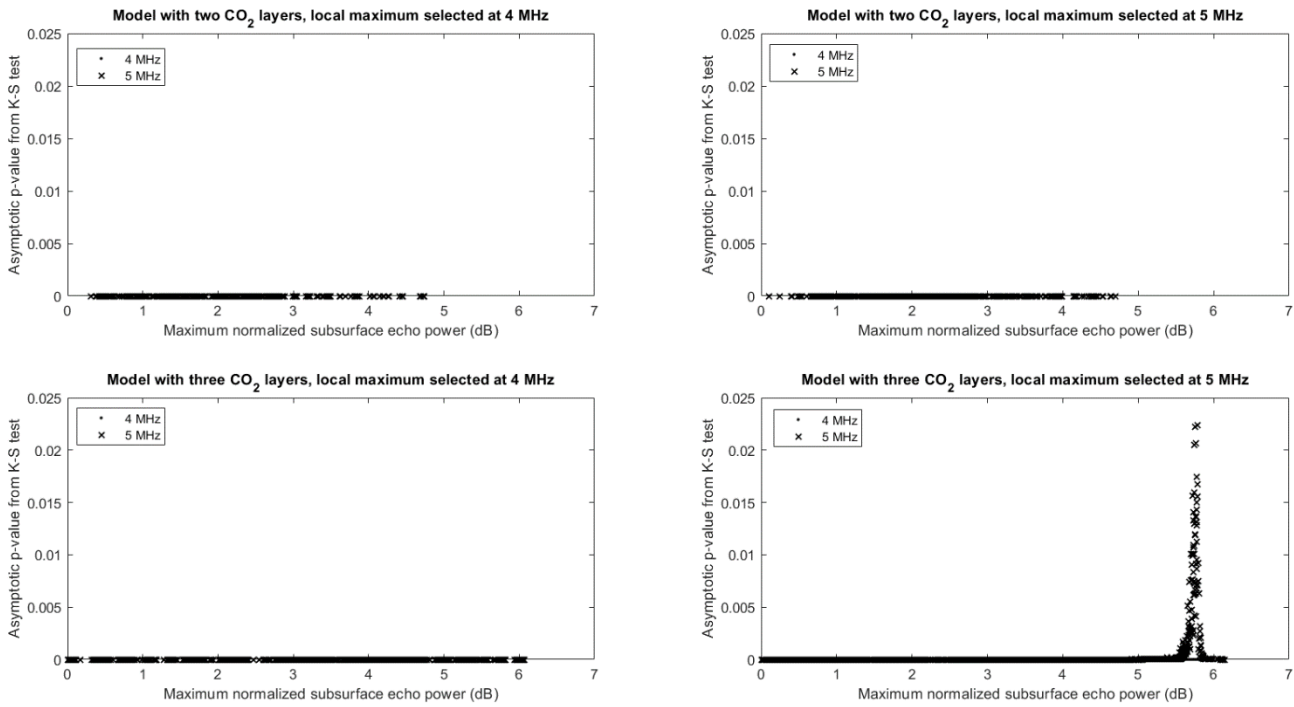


Fig. 11. Plots of the asymptotic p-values produced by the Kolmogorov-Smirnov test versus the maximum subsurface to surface power ratio of simulated echoes, for models with two- and three-CO₂ ice layers and for local maxima at 4 and 5 MHz. Each point represents the probability that model results and MARSIS measurements belong to the same continuous distribution. For each tested set of simulations, probabilities are computed both at 4 MHz (dots) and at 5 MHz (crosses).

No set of two- and three-layer models resulted in simulated values exceeding the 5% significance level threshold. However, it can be seen in the bottom right plot in Figure 11 that a few layer configurations resulting in a local maximum of basal echo power at 5 MHz approach the 5% threshold. This happens only at 5 MHz (crosses in Fig. 11), while at 4 MHz (dots) the probability remains negligible. This quantitatively confirms the observation that, in a model based on the coherent interference between layers, a configuration capable of producing a constructive interference at one frequency is not capable of doing the same at another frequency.

None of the modelled scenarios involving CO₂ ice layers reproduce the distribution of basal echo intensities observed by MARSIS over the bright areas in Ultimi Scopuli. Median values of the simulated basal echo power are lower than those observed in real MARSIS data, and the spread of the simulated values does not match the range of the measured data. We thus conclude that the interpretation that best reconciles all radar observations in Ultimi Scopuli remains that of a strong basal reflector overlain by a thick dust-laden SPLD.

5. Conclusions

The results of our simulations do not support claims that CO₂ ice layers interbedded with water ice can produce the bright reflections detected by MARSIS from the base of the SPLD in Ultimi Scopuli. Therefore, we rule out the presence of hidden deposits of basal CO₂ ice in this area. Based on these results, and the lack of detection of sequestered CO₂ ice deposits elsewhere in the SPLD (Whitten et al., 2017; Whitten and Campbell, 2018; Khuller and Plaut, 2021) beyond those reported by Phillips et al. (2011) and Bierson et al. (2016), we posit that it is unlikely that hidden CO₂ was stored in the late Amazonian (< 250 Ma) south polar deposits. This conclusion is consistent with the geological record of diminishing volcanic activity on Mars during the Hesperian and throughout the Amazonian, with very small amounts of CO₂ released into the atmosphere by volcanic degassing in post-Noachian times.

567 With the polar caps being the most voluminous reservoir of volatiles on Mars, and their role in regulating the
568 planetary climate, records of the composition and physical properties of the polar layered deposits are fundamental to
569 constrain the entire geological and climatological history of Mars. The ground penetrating radar MARSIS continuing
570 acquisition of observations over the polar regions is an essential element of our investigations into the origin,
571 evolution, and inventory of Martian volatiles and of the planet evolution. This is why interpretations of MARSIS data
572 must be critically assessed against the framework of the physical theory of the propagation of radar waves through
573 stratified media, as we have demonstrated in this paper. Rigorous application of this approach to all investigations
574 based on MARSIS data will advance our knowledge of the nature and properties of Mars's polar materials.

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577
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Numerical simulations of radar echoes rule out basal CO₂ ice deposits at Ultimi Scopuli, Mars.

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Abstract

The principal objective of the radar sounder MARSIS experiment is to look for ice and water in the Martian subsurface. One particular focus of investigations, since 2005, has been the search for basal liquid water in the south polar layered deposits (SPLD). Anomalously strong basal echoes detected from four distinct areas at the base of the deposits at Ultimi Scopuli have been interpreted to indicate the presence of bodies of liquid water in this location, beneath a 1.5 km thick cover of ice and dust. Other explanations for the bright basal reflections have been proposed, however, including the possibility of constructive interference in layered media. Here, we test this mechanism through simulations of MARSIS radar signals propagating in models of CO₂-H₂O ice sequences. We then compare the results to real MARSIS data acquired over Ultimi Scopuli, finding that no CO₂-H₂O ice model sequence reproduces the set of real data. The results of our work have implications in relation to the global CO₂ inventory of Mars.

Keywords: Mars, SPLD, MARSIS, CO₂ ice, radar.

61 1. Introduction

62
63 The radar sounder MARSIS (Mars Advanced Radar for Subsurface and Ionospheric Sounding) on board the
64 European Space Agency orbiter Mars Express, has been probing the Martian subsurface since 2005 (Orosei et al.,
65 2015), exploring the global distribution of rock, soil, ice, and liquid water. The instrument operates within the range
66 1.3-5.5 MHz, transmitting a 1 MHz-bandwidth pulse centered at the frequencies of 1.8, 3, 4 and 5 MHz. Exploration of
67 the Martian subsurface is mostly carried out using the 3 and 4 MHz, or 4 and 5 MHz frequencies, for a nominal depth
68 of investigation of up to 5 km, depending on the electromagnetic properties of the probed material (Picardi et al.,
69 2004).

70 The primary objective for the MARSIS experiment is to look for buried ice and water in the Martian
71 subsurface. It was not until 2018, however, that evidence of bright subsurface reflections originating 1.5 km deep
72 below the surface of Ultimi Scopuli, one of the terminal regions of the South Polar Layered Deposits (SPLD), was
73 obtained (Orosei et al., 2018a) at coordinates 193°E - 81°S. The reflections, much brighter than those detected from
74 the surface, were interpreted as originating from basal ponded water or wet sediments distributed on an area
75 approximately 20 km across. The work reignited the scientific debate about the stability of liquid water at the Martian
76 polar regions, and specifically at the south pole, where average surface temperatures (Mellon et al., 2004) are
77 estimated to be 162 K. A possible explanation for the water being liquid at these conditions is that it contains a high
78 concentration of salts (e.g., perchlorates, chlorides) which are ubiquitous on Mars (Hecht et al., 2009) and have
79 eutectic temperatures as low as 198 K (Pestova et al., 2005). It was subsequently suggested that a high geothermal
80 gradient, such as that produced by a magmatic chamber at depth, is required to reach the eutectic temperatures of
81 concentrated salt solutions at the base of the deposits (Sori and Bramson, 2019). This scenario was disputed,
82 however, on the grounds that there is no evidence of an anomalous geothermal gradient in the region (Lauro et al.,
83 2021). **Newly acquired MARSIS orbits of the region further revealed other three previously undetected areas of bright
84 reflections, each approximately 10 km across (Lauro et al., 2021).** Because **laboratory experiments have shown that
85 hypersaline aqueous solutions can form and persist over time at temperatures as low as 123 K, well below their
86 eutectic temperatures (Toner et al., 2014; Primm et al., 2017), a hypersaline origin of the water was deemed to remain
87 the most likely explanation for the bright reflections (Lauro et al., 2021). It is however difficult to envisage how large
88 bodies of supercooled brines could have persisted in the natural Martian environment over geological timescales
89 (Schroeder and Steinbrügge, 2021; Smith et al., 2021), which remains a significant argument against the liquid brine
90 interpretation.** The possibility that bright reflections could be caused by scattering phenomena in a stratified medium
91 was also raised (Hecht et al., 2018). This criticism was however dismissed by Orosei et al. (2018b), who pointed out
92 that they had demonstrated that the strength of the basal echoes at Ultimi Scopuli cannot be explained by multiple
93 reflections. Even though a particular stratigraphic arrangement could produce enhanced basal echoes at one of the
94 MARSIS frequencies of detection, it did not do so at the other observation frequencies. Alternative models suggesting
95 that a range of materials, other than bodies of liquid water, could explain the sources of the strong basal echoes, have
96 recently been published by Bierson et al. (2021), and Smith et al. (2021). The merits of all proposed alternative
97 hypotheses have been reviewed by Schroeder and Steinbrügge (2021), who argued that none of the alternatives is
98 conclusively incompatible with bodies of liquid water as the source of the bright basal reflections, and that each model
99 opens up questions providing new opportunities for future investigations.

100 Distinct layers of CO₂ ice have been detected by radar sounder SHARAD surveys of the SPLD (Phillips et al.,
101 2011; Bierson et al., 2016). The layers, up to 1 km thick, are interlayered with thin (typically 10-40 m thick) water ice,
102 and it has been calculated that sublimation of these massive deposits of CO₂ ice would increase the Martian
103 atmospheric pressure by 100% (Bierson et al., 2016). It has been suggested that, if the conditions leading to formation

of the thick CO₂-H₂O ice deposits have occurred at other times, during the approximately 100 Myr history of formation of the SPLD (Manning et al., 2019), then a specific CO₂ ice layering configuration might produce subsurface reflections of higher intensity than those from the surface. In fact, Lalich et al. (2021a,b) obtained strong simulated radar reflections at 4 MHz frequency, mimicking MARSIS waveform, for a model configuration involving 1-100 m thick basal CO₂ ice layers sandwiching a dusty water ice layer of comparable thickness. Such configuration over the extent of MARSIS bright area, would translate in an average ~ 2% of the current CO₂ atmospheric content. While this specific contribution to the atmospheric volatile budget appears negligible, compounding scenarios involving storage of volatiles within the south polar cap can potentially impact estimates of the Martian volatile budget, with consequences on models of Martian climate cycles and atmospheric evolution, and on estimates of the amount of volatiles stored in the Martian crust as part of the geological record.

The Martian polar caps are the most voluminous exposed reservoirs of volatiles on Mars. It is estimated that, combined, the polar caps store two thirds of the entire water content of the planet (Catling, 2014; Smith et al., 2020; Jakoski, 2021). While climate models have suggested the possibility that the polar layered deposits may also contain significant amounts of CO₂ (Jakoski et al., 1995), estimates based on thermophysical properties of water and CO₂ ices rejected such interpretation (Mellon, 1996). The thin veneer of seasonal and residual solid CO₂ on the surface of the polar caps is effectively counted as part of the atmosphere (Jakosky, 2019). However, geomorphological (Gulick and Baker, 1989; Baker et al., 1991; Baker, 2001; Ramirez and Craddock, 2018) and geochemical evidence (Bridges et al., 2019) has been interpreted to indicate that early Mars was warmer and wetter than at present, prompting suggestions that a thick CO₂ rich atmosphere existed at the time, to account for greenhouse heating of the planet (Haberle et al., 1994). It has been argued that subsequent loss to space of the CO₂ (and other volatiles) does not fully account for the present-day content of CO₂ in the current thin Martian atmosphere, and that massive quantities of CO₂ ought to be sequestered within one or more geological sinks on Mars (Forget et al., 2013). While deep-seated carbonates could in principle accommodate large volumes of CO₂ (Kahn, 1985; Michalski and Niles, 2010; Wray et al., 2016; Jakosky, 2019), as is the case on Earth, there is at present no conclusive evidence for a global deep-crustal reservoir of carbonates on Mars (Edwards and Ehlmann, 2015). The discovery of massive deposits of CO₂ ice (Phillips et al., 2011; Bierson et al., 2016) in the South Polar Layered Deposits (SPLD) by the radar sounder SHARAD, and three-dimensional radar imaging volume calculations suggesting the deposits may contain more than $1.6 \cdot 10^4$ km³ of dry ice (Putzig et al., 2018), have thus paved the way for renewed investigations into the possibility of the polar ice caps as CO₂ sinks.

The global abundance of CO₂ on Mars is presently unknown. There are some generally accepted points, however, that hold irrespective of the exact abundance and distribution of CO₂ on Mars. Mars had lost most of its CO₂ rich atmosphere by the end of the Noachian (~ 3.7 Ga) (Jakosky, 2019). The rate of input of juvenile gases into the atmosphere through planetary outgassing (i.e., volcanic activity) is not known precisely, but it was higher during the Noachian, waning in the Hesperian, becoming negligible in the Amazonian (Grott et al., 2011), so it would have been insufficient to replenish the lost early atmospheric budget. It is not known how much of the initial Martian atmosphere was left by the end of the Noachian (estimates range from ~ 0.1 bar to 1 bar; Catling, 2014), and how much CO₂ had by then already been stored in crustal sinks, though it is clear that the overall trend since has been one of thinning of the atmosphere to the present-day global value of ~ 6 mbar of pressure. If the massive CO₂ ice deposits (Phillips et al., 2011) were to sublimate, they would add to the atmosphere as much as ~0.006 bar of CO₂ (Bierson et al., 2016). Further compounding the contribution (however small) of putative additional reservoirs of CO₂ ice hidden in the SPLD, could potentially impact estimates of the planetary CO₂ budget.

Here, we present the results of our comprehensive testing of model simulations against the database of MARSIS data acquired over more than 10 years of surveys of the south polar region. We modeled the propagation

behaviour of electromagnetic waves at MARSIS operating frequencies (1.8-5 MHz) under a variety of conditions involving different configurations of thin layers of solid CO₂, H₂O ice and dust. Within the parameter space we explored, the models do not return values of reflection consistent with the observed MARSIS data. We thus reject the notion that the bright basal reflections detected at Ultimi Scopuli may indicate the presence of thin layers of CO₂-ice at the base of the SPLD.

2. Geological Setting

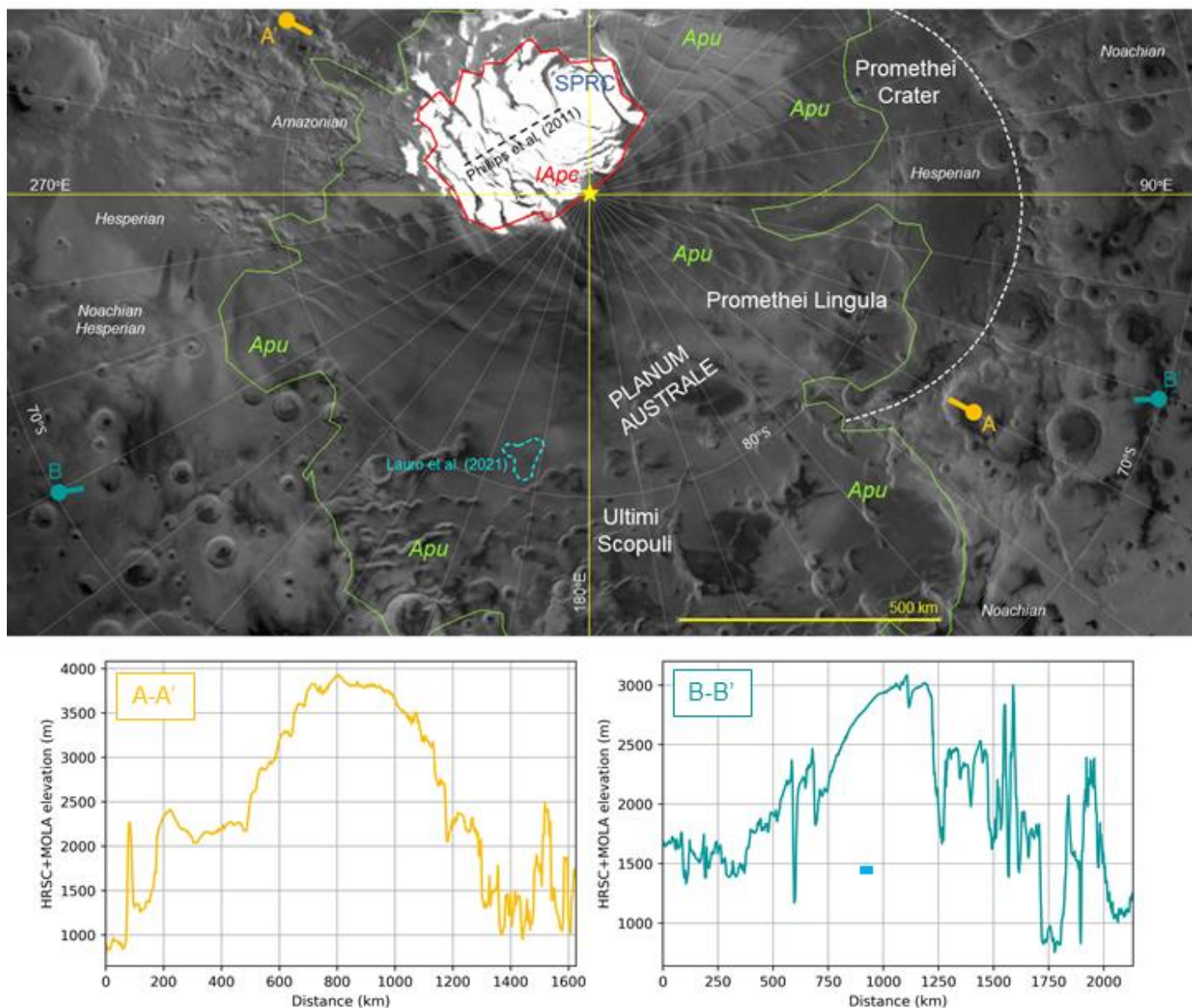
The South Polar Layered Deposits (SPLD), covering most of the Planum Australe of Mars (Fig. 1), form an up to 3.7 km thick (Plaut et al., 2007) stratified unit, composed of pure water ice interlayered with thin dust-rich deposits. It has been estimated that the content of dust in the SPLD is ~ 10-15% (Plaut et al., 2007; Zuber et al., 2007), although inhomogeneities in the values of inverted density of the deposits indicate that the dust is not uniformly distributed (Li et al., 2012), and that the fraction of dust in the water ice would vary depending on the presence of denser CO₂ ice (Broquet et al., 2021). The layers are directly observed in radar imagery where troughs and scarps cut through the deposits (Milkovich and Plaut, 2008). Owing to the electromagnetic contrast between the values of dielectric permittivity ϵ of ice (3.12; Pettinelli et al., 2015) and dust (> 4; Mattei et al., 2014), layering is also detected by radar sounder observations (Plaut et al., 2006, 2007; Milkovich et al., 2009; Whitten and Campbell, 2018). Layering of the deposits is thought to reflect differential dust and ice accumulation rates influenced by global climate cycles linked to orbital forcing (e.g., Hvidberg et al., 2012). The deposits decrease in elevation and thickness at the margins, eventually waning and revealing underlying Noachian and Hesperian highlands terrains (Tanaka and Kolb, 2001), overlain by a thin mantle of desiccated or water-ice pore-filled soil (Putzig et al., 2005; Jones et al., 2014). The surface age of the SPLD is 7-100 Ma, based on the impact crater record (Herkenhoff and Plaut, 2000; Koutnik et al., 2002). The age of the lowermost strata of the SPLD sequence is presently undetermined. A detailed geological analysis and interpretation of the unconformities observed in Promethei Lingula suggests that the current extent of the SPLD may have started forming no earlier than 250 Ma (Guallini et al., 2018). This estimate is consistent with time-series analysis of HiRISE and CaSSIS stereo-images of the internal structure of the SPLD shown along scarp faces, suggesting that the minimum time of accumulation of the SPLD is 10-30 Myr, resulting in a minimum age of deposition for the basal layers of 17 to 130 Ma (Becerra et al., 2019).

The SPLD have been subdivided into three distinct structures (Milkovich and Plaut, 2008): the *Promethei Lingula Layer Sequence* (PLL), the *Bench Forming Layer Sequence* (BFL), and the *Inferred Layer Sequence* (IL). The BFL is the top sequence, geographically located in the central west-northwest portion of the SPLD, where the deposits measure their highest elevations (~ 3500-4500 m). The BFL is covered by thin layers of CO₂ ice (part of the South Polar Residual Cap, SPRC). Underlying the BFL is the PLL, while IL is inferred to be the basal sequence in the west of the deposits. Elsewhere, only the PLL exists (ref. to figs. 18-19, in Milkovich and Plaut, 2008), its lateral extent corresponding to that of Tanaka et al.'s (2014) *Apu* (Amazonian polar undivided) unit (Fig. 1).

The South Polar Residual Cap (SPRC), which is part of geological unit *IAPc* (Fig. 1), straddles the prime meridian, approximately between longitudes 250°E-40°E and latitudes 85-90°S. The cap is primarily composed of a high albedo thin (<10 m) layer of CO₂ ice (Byrne and Ingersoll, 2003), with a characteristic «Swiss cheese» appearance (Malin et al., 2001) due to the presence of pits that reveal underlying water ice (Zuber et al., 2007). Observed constant rates of erosion of the pits (Thomas et al., 2013) have been interpreted to point to net loss of CO₂ from the SPRC, possibly counterbalanced by observed net gain of seasonal water ice (Brown et al., 2014), suggesting orbital effects are responsible for the CO₂ and H₂O atmosphere-SPRC exchanges. Localized deposits up to 1 km thick, formed by alternating layers of CO₂ ice (each approximately 100 m thick) and H₂O ice (a few m thick) have also

190 been detected (Phillips et al., 2011; Bierson et al., 2016) by the radar sounder SHARAD, in troughs and at the base of
 191 scarps of the BFL structural unit, to the west of the location of the inferred rim of the ~ 870 km diameter Noachian
 192 impact crater Promethei Basin (Fig. 1). This suggests a possible relationship between the basal topography in this
 193 area, and atmospheric conditions that caused deposition of CO₂ ice here. It has been proposed that H₂O ice entombs
 194 the CO₂ ice layers, thus isolating them from the atmosphere and preventing their erosion (Manning et al., 2019). An
 195 alternative hypothesis argues that the deposits are in communication with the atmosphere, and that variations in polar
 196 incident sunlight from Mars's 100 kyr orbital cycles balance the exchange between CO₂ in the deposits and in the
 197 atmosphere. The observed deposits should thus record the last 510 kyr of Martian polar climate history (Buhler et al.,
 198 2021). Observations by SHARAD and MARSIS radars did not detect CO₂ ice layers anywhere else in the SPLD
 199 (Whitten et al., 2017; Whitten and Campbell, 2018; Khuller and Plaut, 2021).

200



201

202 **Fig. 1. Context map of the south polar region.** The map shows the extent of the geological units and geographic features
 203 described in Section 2. The dashed aqua polyline indicates the location of the strong basal reflections detected by Lauro et al.
 204 (2021). The trace of the ground track of the tract of SHARAD orbit 5968-01 traversing the RFZ₃ (Phillips et al., 2011) is shown as a
 205 black dashed line on the south pole residual cap (SPRC). Yellow star: South Pole. Vertical yellow line: prime and anti- meridians.
 206 Dashed white line: approximate position of remnant of rim of Promethei Crater. Lime green polyline: boundary of Amazonian polar
 207 undivided unit (*Apu*). Red polyline: boundary of late Amazonian polar cap unit (*IAPc*). A-A': end points of topographic profile A-A'
 208 (orange). B-B': end points of topographic profile B-B' (green). The short light blue segment in topographic profile B-B' indicates the

209 position and inferred depth of the bright basal reflections. Elevations are blended DEM/DTM (200 m/pixel) from MOLA and HRSC
210 data. Geological units are described by Tanaka et al. (2014). Geographic grid interval: 10°. Mapped with JMARS.

211
212 Ultimi Scopuli is a vast terminal region of the deposits, part of the *Apu* geological unit (Fig. 1) and of the
213 structural PLL sequence. Here, between coordinates 191°E-196°E and 80.3-81.5°S, and at a depth of approximately
214 1.5 km, recent orbiter sounder MARSIS observations (Orosei et al., 2018a; Lauro et al., 2021) revealed the presence
215 of four high permittivity areas, interpreted to be bodies of basal liquid water. The heat flow at the Martian south polar
216 regions is ~ 22-26 mW/m² (Grott et al., 2013; Plesa et al., 2018), and there is presently no evidence of an anomalous
217 geothermal gradient or localized hot spot (such as a magma chamber, as postulated by Sori and Bramson, 2019), that
218 would plausibly increase the basal temperature. The overburden pressure, estimated to be ~ 6 MPa (considering an
219 average density of the SPLD = 1220 kg/m³, from Zuber et al., 2007, which is the same as the best fit value published
220 by Broquet et al., 2021, and also consistent with a local value of ~ 1200 kg/m³, as inferred from the inverted density
221 map of the SPLD published by Li et al., 2012), would also be insufficient for ice to melt. Therefore, to explain the
222 presence of basal liquid water, Orosei et al. (2018a) suggested that the water may be a perchlorate brine.
223 Subsequently, Lauro et al. (2021) further defended this interpretation, on the grounds that supercooled perchlorate
224 brines have been shown to form and persist at sub-eutectic temperatures (Toner et al., 2014; Primm et al., 2017),
225 within the range of those reasonably expected at the base of the SPLD. The reflections at the base of Ultimi Scopuli
226 are located over a gently sloping basement (Lauro et al., 2021).

228 3. Model

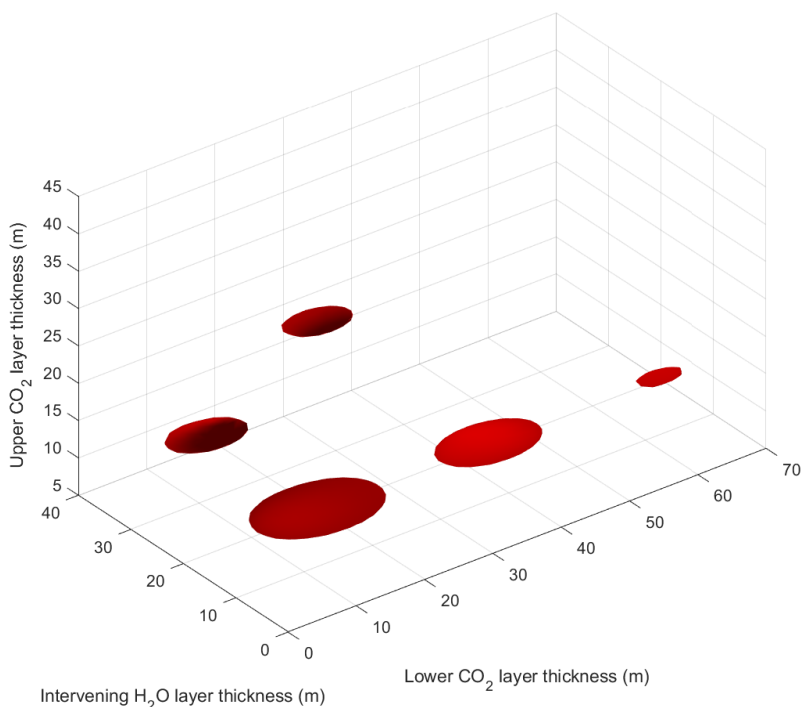
229 3.1. Preliminary Analysis

231
232 The fundamental concept underpinning this study is that the electromagnetic waves locally reflected from
233 interfaces between the layers in a stratified structure, can combine to enhance (constructive interference) or suppress
234 (destructive interference) the overall radar response from a layered structure. As a preliminary analysis, we explored
235 the specific parameter space presented by Lalich et al.'s (2021a) model, consisting of a thin layer of pure water ice
236 sandwiched between two thin layers of CO₂ ice: thus, we modelled the variation of only three parameters, i.e., the
237 thickness of each of the three layers. Layer thicknesses greater than a few tens of m dampen resonance effects,
238 because longer propagation paths decrease the likelihood of simultaneous in-phase sum of sinusoids at different
239 frequencies (Supplementary Figure 5, Orosei et al., 2018a). Therefore, we constrained the simulated thickness of
240 each layer to be ≤ 100 m. In the field of numerical electromagnetic modelling, spatial discretization is usually set to be
241 less than 1/10 of the wavelength. The shortest wavelength transmitted by MARSIS is 30 m in water ice. Thus, we
242 applied simulated thickness increments of 1 m as our sampling interval, to obtain accurate model stratigraphy
243 responses to the MARSIS pulses. We ran more than one million simulations (101 × 101 × 101) at MARSIS operating
244 frequencies of 3, 4 and 5 MHz, obtaining a complete characterization of the properties of Lalich et al.'s (2021a) model.

245 The highest value of subsurface-to-surface echo power ratio we obtained was 4.7 dB. Because the radar
246 response to the stratigraphy is described by sinusoidal functions, we ruled out the eventuality that we might have
247 missed peaks in the subsurface echo: considering that variations in the radar response between consecutive data
248 points are negligible (a fraction of a decibel at most), the model stratigraphy subsurface-to-surface echo power ratios
249 would never exceed ~ 5 dB, even in the case that the position of our maxima were not exact.

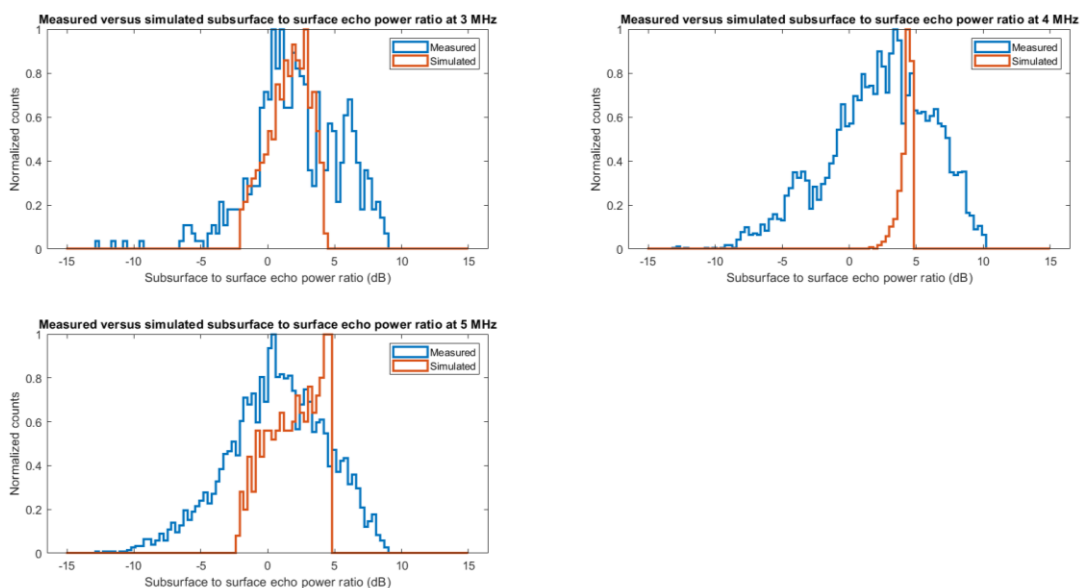
250 For each frequency, we obtained five well-defined regions in parameter space with subsurface-to-surface
251 radar echo power ratio values > 4 dB (Fig. 2). These regions are specific to each frequency, because constructive

252 interference phenomena are primarily controlled by layer thickness / wavelength ratios. Consequently, while basal
 253 echo powers remain high for small variations of thickness around the local maxima at one frequency, the
 254 corresponding echo powers at the other two frequencies are lower and much more widely dispersed (Fig. 3). Thus, if
 255 we take the parameter values corresponding to the maxima found at one frequency as a reference, and calculate the
 256 subsurface-to-surface echo power ratio differences between this frequency and the other two, these will always favour
 257 the reference frequency (Fig. 4).



258
 259 **Fig. 2.** Location of strong basal echoes (> 4 dB greater than surface echoes) produced in simulations at 4 MHz, displayed in the
 260 parameter space of the model by Lalich et al. (2021a), consisting of three quantities, specifically: the thickness of two CO₂ layers,
 261 and of a pure water ice layer sandwiched between them. Simulations at 3 and 5 MHz produce the same pattern, but with local
 262 maxima in different positions.

263



264

265 **Fig. 3.** Histograms of subsurface-to-surface echo power ratios obtained for simulations (orange) with a bottom CO₂ ice layer of
 266 thickness ranging between 8 m and 14 m, a top CO₂ ice layer 9-15 m thick, and an interbedded H₂O ice layer 9-15 m thick. Blue
 267 histograms represent real MARSIS observations.

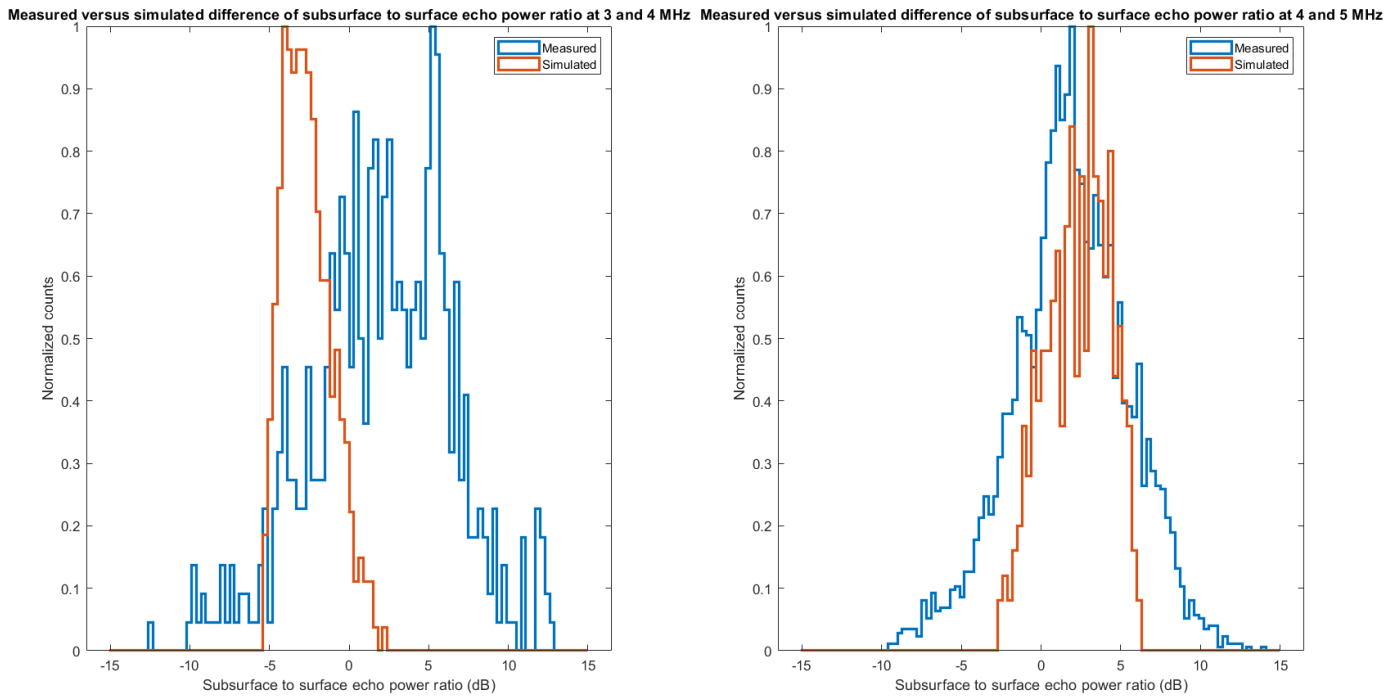


Fig. 4. Histograms of basal-to-surface echo power ratio differences between pairs of MARSIS operating frequencies (3 and 4 MHz, and 4 and 5 MHz). Simulation conditions are the same as for those in Fig. 3. Blue histograms are for real MARSIS observations.

Our preliminary analysis indicates that the distribution and behavior of subsurface-to-surface echo power ratios simulated in the parameter space as described, contrast sharply with that of real MARSIS data. However, idealized models do not consider possible sources of echo power fluctuations that affect real data, and thus some discussion is needed before a comparison between data and model is attempted.

The model used and discussed by Orosei et al. (2018) is computationally the same used by Lalich et al. (2021a, 2021b), that is the one-dimensional solution of Maxwell equations in a plane parallel stratigraphy. The former model had a simpler stratigraphy compared to the latter, with the SPLD represented as a single homogeneous layer composed of a mixture of H₂O ice and dust. The model was run varying basal temperature, dust content and basal permittivity. These three parameters controlled the attenuation within the SPLD and the intensity of the reflection at the bottom of the SPLD. The results showed that even for low amounts of dust and basal temperatures, it was necessary to have a high basal reflectivity to produce the observed strong basal echoes.

The model parameters published by Lalich et al. (2021b) are exactly the same, except for the input of single values for dust fraction, basal temperature and basal permittivity. These authors varied the thickness of the layers at the base of the SPLD.

Because of the similarity of the mathematical formulation of the two models, we compared them directly to one another to assess their capability to reproduce MARSIS observations, following the same approach used by Lalich et al. (2021b). For example, Figure 3 in that paper shows reflections produced by a single CO₂ layer model that are as strong as the ones observed by MARSIS, although the authors neglected to specify that their reference values for MARSIS echo power are the medians of observed values, and that normalized echo power can reach values as high as 10 dB (see Figure 4 and S4 in Orosei et al., 2018), well above any value that can be obtained by the single layer model shown in Figure 3.

The capability of the model in Orosei et al. (2018) to produce strong echoes is thus not due to a more complex mathematical implementation or to a greater number of parameters, but to the fact that basal permittivity can have

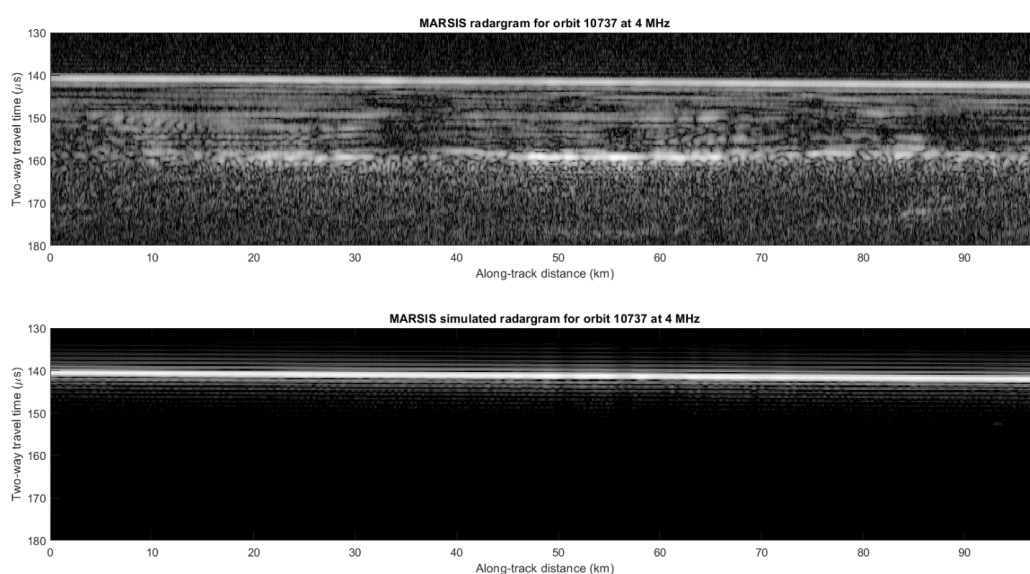
295 higher values than those considered by Lalich et al. (2021a, 2021b). In fact, Orosei et al.'s (2018) model could
296 produce basal echoes spanning the full range of observed values (see their Figures 4 and S4), whereas Lalich et al.'s
297 (2021a, 2021b) resulted only in a few instances of basal echo power comparable to observations. The purpose of the
298 present work is therefore to establish the conditions (if these exist), by which Lalich et al.'s (2021) model reproduces
299 the strong radar echoes reported by Orosei et al. (2018). We thus performed comprehensive simulations based on the
300 same mathematics, using a range of possible parameters beyond that explored in that earlier work.

301 In the following sections, we discuss the sources of variability and error affecting MARSIS observations,
302 evaluating their effect to derive a model of the uncertainties that can be applied to model results, for a more realistic
303 comparison between simulations and data.
304

305 *3.2 Idealized vs. actual SPLD surface geometry*

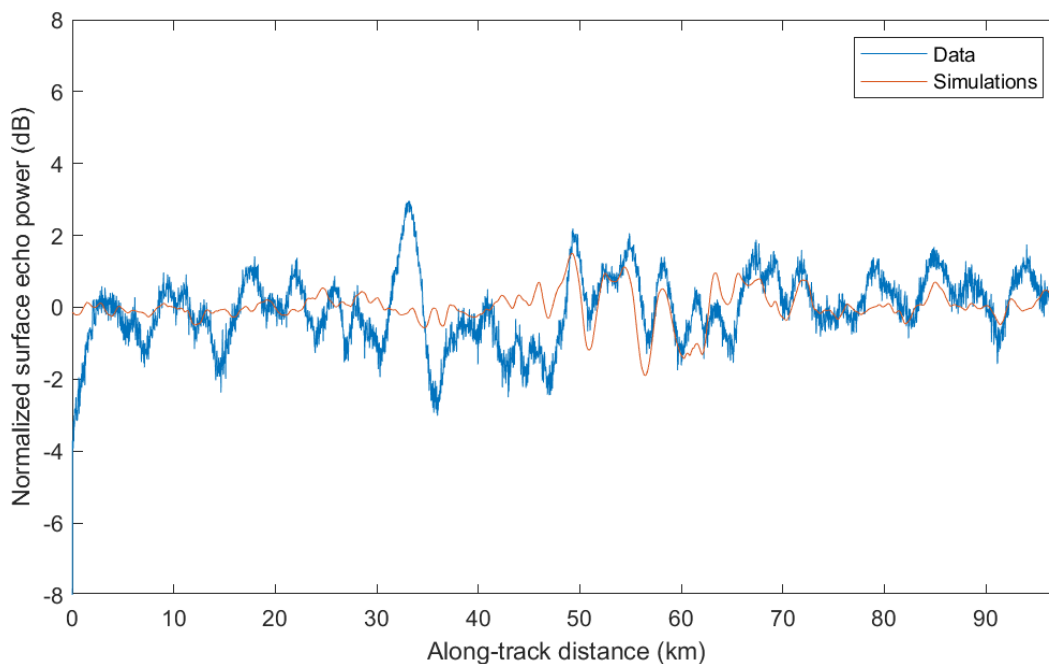
306
307 The assumption of a perfectly flat surface does not prevent the possibility of a comparison with real data, as
308 already shown by Orosei et al. (2018). Furthermore, the subsurface echo power for a given orbit is always normalized
309 to its median value, in order to prevent local minima of surface echo power from producing unrealistically high
310 normalized subsurface echo power values. The choice of using surface median power implies the assumption that the
311 SPLD surface is locally homogeneous, and that surface echo power fluctuations are due to small perturbations of this
312 assumed homogeneity.

313 To verify if this assumption holds, we need to examine the properties of echo power measured by MARSIS.
314 To this end, we retrieved the same data employed to produce Figure 2 in Orosei et al. (2018) and made publicly
315 available at the time of publication of that paper (Orosei and Cicchetti, 2018). We also produced a simulation of
316 surface echo power similar to those presented and discussed in Orosei et al. (2015). Simulations are used to validate
317 subsurface echo detection by visual comparison: all features present in both the real and simulated radargrams are
318 interpreted as surface clutter, and are therefore excluded from interpretation. A comparison between measurements
319 and simulations for orbit 10737 is shown in Figure 5.
320



321
322 **Fig. 5.** Top: MARSIS radargram for orbit 10737 at 4 MHz, redrawn following Orosei et al. (2018), versus the corresponding
323 simulation of surface scattering produced using the MOLA topographic dataset (bottom).
324

325 The simulated radargram allows also to estimate surface power fluctuations due to surface roughness alone,
326 sampled at the resolution of the MOLA topographic dataset. A comparison with the measured surface power (both
327 arbitrarily normalized to a 0 dB median power for ease of comparison) is shown in Figure 6: it is clear that surface
328 roughness alone accounts for the surface echo power fluctuations at distances between 50 and 60 km along track,
329 which correspond to the location of the bright subsurface reflector.
330



331
332 **Fig. 6.** Measured vs. simulated surface echo power for orbit 10737 at 4 MHz, both arbitrarily normalized to a 0 dB median power for
333 ease of comparison.

334
335 The residual difference between data and simulations can be attributed to several factors, such as: surface
336 roughness at scales below MOLA resolution or variations in near-surface density and composition, but the exact
337 cause is difficult to ascertain with available data. Residual power fluctuations are found to have a standard deviation of
338 less than 1 dB, corresponding to a fluctuation of 0.5 in the value of the relative dielectric constant in the assumption
339 that its median value is comprised between 3 and 4 (typical of SPLD materials). Because of the small uncertainties
340 resulting from this analysis, we consider the ratioing with the median surface power as a reasonable calibration
341 method for subsurface echo power.

342 Furthermore, a previous analysis of SHARAD echoes by Grima et al. (2012) over this area (called Reference
343 Zone or RZ in the paper because of its flatness) showed that surface reflections are predominantly coherent, implying
344 that scattering is near-specular and thus approximates that of a plane surface. This conclusion is even truer at
345 MARSIS wavelength, as the effect of roughness on scattering does not depend on RMS topographic height, but rather
346 on the ratio between RMS topographic height and the wavelength.

347 It could be argued that even a small roughness causes a decrease (however small) of the mean power of
348 surface echoes, leading to an overestimation of the subsurface-to-surface echo power ratio. In fact, the opposite is
349 true: electromagnetic models and three-dimensional propagation simulations (e.g., Jonard et al., 2019) demonstrate
350 that the wavefront of the radar pulse penetrating through the subsurface is disrupted more than the echo
351 backscattered at the surface, because it crosses the rough surface twice. Thus, the subsurface to surface echo power
352 ratio calculated assuming no surface roughness effect, should be considered a minimum value, with the real ratio
353 likely higher.

3.3 Signal distortion and noise

The MARSIS pulse is not an ideal chirp. Its actual shape and power are unknown: in fact, the spectra reported by Jordan et al. (2009) do not include the frequency-dependent effects of antenna gain, because these could not be measured on ground owing to the length of the antenna. A test comparing the results obtained using Jordan et al.'s (2009) chirp spectra with those obtained using the ideal chirp, does not demonstrate any significant difference. Therefore, the MARSIS data have been range-compressed using an ideal chirp since the beginning of the mission. The range compression is not optimal, with both resolution and signal-to-noise ratio affected. This, however, does not affect the ratio of subsurface to surface echo power. In an abstract mathematical representation, radar sounding can be described as the convolution of the radar waveform with the ideal electromagnetic response of the observed scene. In Fourier transform terms, this convolution is the multiplication of the stratigraphy spectral response with a band-limited signal spectrum, which clips the spectral response bandwidth and produces a limited resolution time-domain response. All these mathematical operations are linear, so the shape of the pulse spectrum will change the power level of echoes from two different interfaces, but not their ratio, which is the quantity used in this analysis.

A few available data also point to a possible dependence of gain on temperature, although a full characterization of this effect was never performed, because it was considered to be negligible under the conditions at which MARSIS acquires its observations. Furthermore, MARSIS raw data takes are at most 25 seconds long, which is too short a time to produce any significant gain variation. This is further corroborated by the plot of surface echo power shown in Figure 6, which does not show any trend of varying gain across the entire observation. Thus, within the limits of what it is technically possible to ascertain about the causes and extent of non-ideality of range compression in MARSIS, there is no effect that should be considered as introducing a bias in the use of median surface echo power as a calibration reference for subsurface echo power.

Plots of the noise level in the observations (Orosei and Cicchetti, 2018) analysed by Orosei et al. (2018) show constant noise level across a single observation, and nearly-constant noise levels across observations at the same frequency (around -5 dB at 3 and 4 MHz and -11 dB at 5 MHz when expressed in uncalibrated digital numbers). The median signal-to-noise ratio is > 20 dB for all observations, exceeding 30 dB.

Interferences caused by the electronics of the Mars Express spacecraft were characterized in the commissioning phase of MARSIS. Limited available documentation shows a correlation with the activity of the HRSC camera, though no systematic analysis has ever been published. Overall, data acquired over the lifetime of the mission show that the combined electromagnetic interferences (EMI) caused by orbiter navigation and experiment instrumentations, have a much smaller effect on the functioning of MARSIS than (for example) on its Mars Reconnaissance Orbiter companion SHARAD. Thus, the operating assumption is that EMI have a negligible effect on the subsurface to surface echo power ratio.

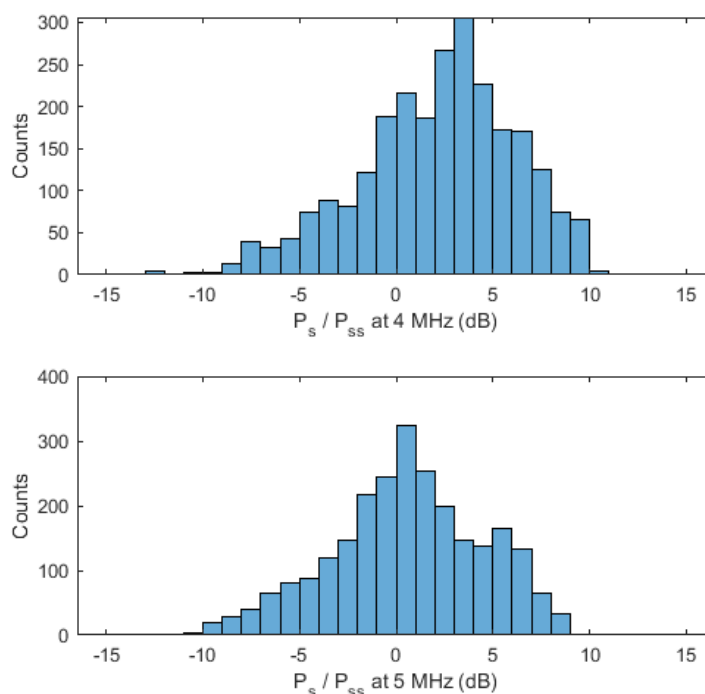
The MARSIS data analyzed by Orosei et al. (2018) consist of near-simultaneous observations at two frequencies, namely 3 and 4 MHz, or 4 and 5 MHz. Thus, it is possible to compute the subsurface to surface echo power ratio at two frequencies for the same observation. Doing so for data in which surface echo power has a standard deviation below 1 dB (as noted in Table S1 in the Supplementary Materials of Orosei et al., 2018), one is left only with data at 4 and 5 MHz. The difference between normalized power at those two frequencies has a median value of 1.36 dB, which can be used for a back-of-the-envelope computation based on equations found in Porcello et al. (1974) to obtain an estimate of the loss tangent (see also Orosei et al., 2020 for the effect of such loss tangent value on SHARAD measurements). The resulting value is approximately 0.03, which is about fifteen times greater than the one used by Lalich et al. (2021b). Because this value is affected by a large error (the standard deviation of

397 the difference between normalized power at the two frequencies is almost 4 dB), we did not use the estimated loss
398 tangent in simulations. Instead, we considered a variable attenuation as one of the factors decreasing the measured
399 ratio between subsurface and surface power produced by a model.

401 3.4 Other factors affecting subsurface to surface echo power ratio

402
403 Factors affecting subsurface echo power not included in Orosei et al.'s (2018) models are: surface roughness
404 (as discussed above); subsurface roughness (affecting basal echo intensity only); fraction of the radar footprint
405 covered by the basal reflector (see for example Haynes et al., 2018, for a discussion); signal attenuation within the
406 SPLD due to dielectric losses (as discussed above); signal attenuation within the SPLD due to reflection losses in the
407 thinly layered stratigraphy (see for example Courville et al., 2021, for a detailed modelling of propagation in a layered
408 medium at SHARAD frequencies). Their combined effect is demonstrated by the large dispersion of values observed
409 (Figure 7).

410



411

412 **Fig. 7.** Histograms of measured values of subsurface to surface echo power ratio over the bright reflector analysed in Orosei et al.
413 (2018) at 4 MHz (top) and 5 MHz (bottom).

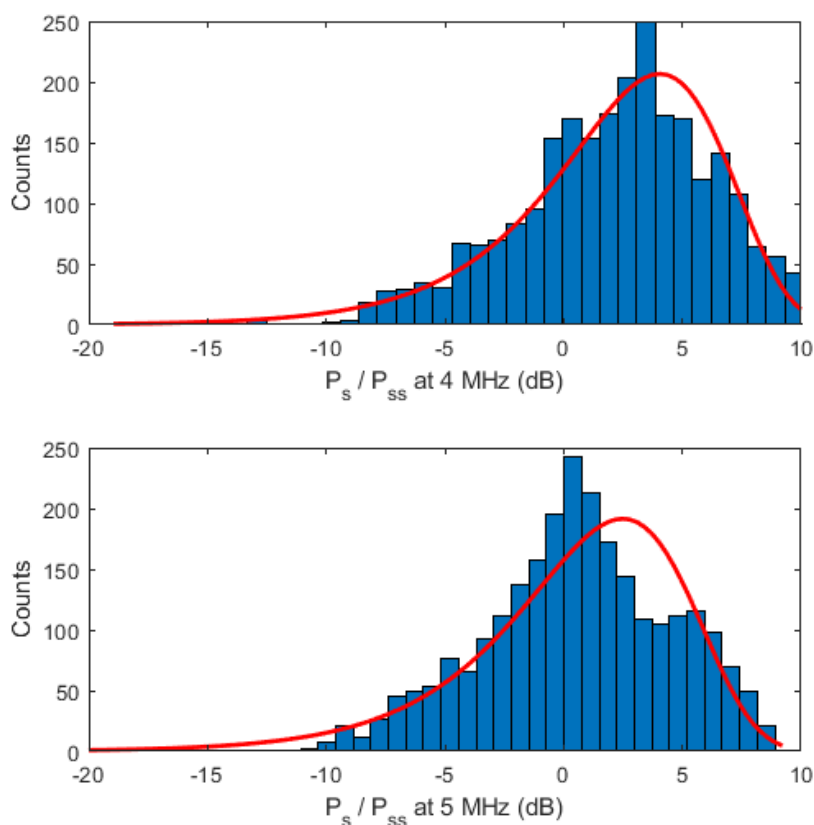
414

415 The histograms are asymmetric, with a cut-off at high values and a long tail towards low ones. This
416 characteristic matches the intuitive notion that the intensity of a basal reflection cannot increase indefinitely, and would
417 thus have a maximum possible value obtained when all the factors contributing to decreasing basal echo intensity as
418 discussed above are minimized. Thus, a comparison between simulations and data requires the introduction of a
419 model for the random dispersion of measurements. As it is not possible to quantify separately the importance of each
420 factor contributing to basal echo power attenuation, we derived a model probability density function directly from the
421 data (Fig. 8).

422

423 The long left-sided tails of the histograms in Figure 7 are characteristic of generalized extreme value
distributions, often used to model the maxima of sequences of random variables. Although this is a heuristic method

424 for the selection of the model distribution, it has the merit of matching data histograms better than other parametric
425 distributions.
426

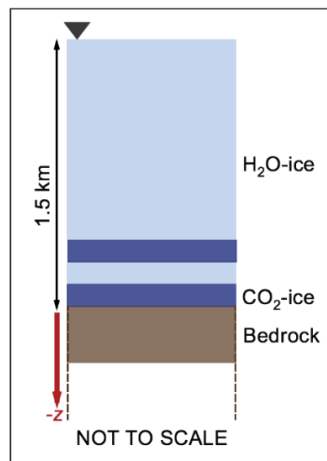


427
428 **Fig. 8.** Best-fit extreme value distributions plotted over the histograms shown in Figure 7.

429
430 Through the remainder of this paper, we use these best-fit distributions as the template for randomly
431 generating synthetic measurements from models.

432 433 *3.5. Electromagnetic propagation model*

434
435 The propagation of radar waves across stratified media is described by non-linear equations dependent on
436 the electromagnetic properties and thickness of the layers, as well as the probing frequency. Here, we model
437 electromagnetic behaviour of stratified CO₂ ice and dust-laden H₂O ice, by assuming a 1D geometry (Lombardo et al.,
438 2000; Pettinelli et al., 2003; Kartashov et al., 2004; Nunes and Phillips, 2006; Courville et al., 2021). Specifically, we
439 simulate the propagation of radar pulses through twelve idealized models of the SPLD, consisting of sets of perfectly
440 parallel planar interstratified basal layers composed of CO₂ ice and H₂O ice, where the latter has variable contents of
441 dust. The total thickness of the simulated deposits was set at the value derived through MARSIS radargrams over the
442 study area (1.5 km). The material underlying the deposits was set to be basaltic rock. The simulations were run by
443 code adapted from MATLABTM scripts publicly available through a public research data repository (Orosei and
444 Cicchetti, 2018). The code computes the complex reflectivity of the layer stack across the range of frequencies within
445 the pulse bandwidth, multiplying it by the signal spectrum and Inverse-Fourier transforming to the time domain.
446 Simulations were run for one, two or three CO₂ ice layers, the lowermost of which is in direct contact with the bedrock,
447 while the remaining ones are separated by layers of H₂O ice of variable thickness (Fig. 9). CO₂ ice layers are



449

450 **Fig. 9. Schematic stratigraphy used in numerical simulations.** The model stratigraphy is here shown as a set of planar parallel
 451 layers of CO₂ ice (dark blue) and H₂O ice (light blue), bounded by an upper half space representing the Martian atmosphere (as
 452 semi-infinite vacuum) and a lower half space representing the Martian bedrock (brown). Simulations are run for one, two or three
 453 lossless CO₂ ice layers, the lowermost of which is in direct contact with the bedrock, while the remaining ones are separated by
 454 layers of H₂O ice of variable thickness. The topmost layer of H₂O ice represents the bulk of the stratigraphy. Layer thicknesses of
 455 CO₂ ice and interlayer dusty water ice range from 0 to up to 100 m, with total deposit thickness = 1.5 km. The red arrow indicates
 456 the depth direction (-z). The inverted triangle on top of the sequence indicates the surface of the deposits. The sketch is drawn not
 457 to scale.

458

459 Typically, complex relative dielectric permittivities are strongly dependent on the frequency of the propagating
 460 electromagnetic wave and the temperature of the medium. Attenuation of the signal is independent of frequency at
 461 MARSIS wavelength for pure water ice: we used the empirical formulae reported by Mätzler (1998) for the average
 462 surface temperature value of 160 K (Clifford, 1987) taken as a constant across the entire model stratigraphy. We find
 463 that the real part of the dielectric permittivity of water ice is 3.11, with loss tangent values between $2.5 \cdot 10^{-8}$ (3 MHz)
 464 and $3.7 \cdot 10^{-8}$ (5 MHz). **The loss tangent of CO₂ ice is not well constrained at the temperature of interest, with
 465 extrapolations and upper limits only being reported in the published literature. In our computations therefore, we set it
 466 to be = 0, to produce the highest possible subsurface-to-surface echo power ratio values. We did not consider the
 467 effect of the complex part of the permittivity of CO₂ on the reflection coefficient, because it is generally assumed to be
 468 low ($\leq 10^{-3}$). Computing the Fresnel reflection coefficient at normal incidence for an interface between a layer of CO₂
 469 ice and an underlying bedrock with permittivity values such as those used in the simulations, the results, with and
 470 without the inclusion of the imaginary part, differ by < 0.00001 .**

471

472 3.6. Simulations

473

474 To calculate the simulated radar waveforms, in our model stratigraphies we assumed perfectly smooth
 475 surfaces, an assumption that holds also for small roughness of the illuminated surface, where “small roughness” is
 476 defined (Ulaby et al., 1981) by a height $< \lambda / 32$, which gives a value of 2.35 m at 4MHz. We considered the
 477 electromagnetic wave transmitted by the radar having a normal incidence with respect to the model stratified structure
 478 (ref. to sketch in Fig. 9), and assumed a 1D-model where the only variation of the electromagnetic properties occurs
 479 along the z-axis (depth). Furthermore, we assumed spatially homogenous layers with a constant complex dielectric
 480 permittivity (ϵ_i) and thickness (h_i). The stratified structure is bounded by two half-spaces representative of the
 481 atmosphere and the bedrock. This model is the same adopted in previous work (Orosei et al., 2018a; Lauro et al.,
 482 2019, Lalich et al., 2021a, 2021b). The received signal $y(t)$ is thus computed as:

483

$$y(t) = \mathcal{F}^{-1}[X(f)R(f)], \quad (1)$$

where \mathcal{F}^{-1} is the inverse Fourier transform, $X(f)$ is the Fourier transform of the signal transmitted by MARSIS antenna, $R(f)$ is the frequency response ($\Gamma_1(f)$) of the layered structure computed from the recursive scattering function $\Gamma_i(f)$:

$$\begin{aligned} R(f) &= \Gamma_1(f) \\ \Gamma_i(f) &= \frac{\rho_i + \Gamma_{i+1} e^{-2jk_{i+1}h_{i+1}}}{1 + \rho_i \Gamma_{i+1} e^{-2jk_{i+1}h_{i+1}}} \end{aligned} \quad (3)$$

with k_i the complex-valued wavenumber of the i -th layer, and ρ_i the Fresnel reflection coefficient at the boundary between layer $i-1$ and i , given by:

$$\rho_i = \frac{\sqrt{\varepsilon_{i-1}} - \sqrt{\varepsilon_i}}{\sqrt{\varepsilon_{i-1}} + \sqrt{\varepsilon_i}} \quad (4)$$

The geometry of the problem is modelled as a layered structure, where the upper and lower half space are the same for each simulation, whereas the assumed stratified layer between the atmosphere and the bedrock changes on the basis of the number of CO₂ ice layers, and the thickness of the layers. The ratio between basal and surface power is determined by considering the ratio between the maximum of the square amplitude of $y(t)$ at a time interval containing the echo return from the surface, and that occurring at a time interval containing the echo return from the upper side of the shallower CO₂ ice layer.

Using a multiprocessor machine, we simulated the three CO₂ layer model for every layer thickness between 0 and 50 m at 1-m intervals for three CO₂ layers and two intervening H₂O layers, for a total of 51⁵ simulations. This allowed us to sample completely the parameter space of the models with two- and three-basal CO₂ layers, and thus to locate and explore the properties of local maxima embedded in the three- and five-parameter spaces respectively. We identified local maxima of subsurface to surface power ratio, both at 4 and 5 MHz, and for each of them we extracted the basal echo power ratio for all simulations whose layers are within 2 m of the thickness of those corresponding to the local maximum, thus simulating the effect of random but limited variation of the stratigraphy.

4. Results

To implement a quantitative comparison between data and simulation results, for each local maximum of subsurface to surface echo power ratio we produced the histogram of power values extracted from its neighbourhood, and convolved it with the template distributions shown in Figure 8, obtaining a simulated histogram which represents the global effect of the factors affecting propagation within the SPLD on model results. An example of a synthetic distribution is shown in Figure 10.

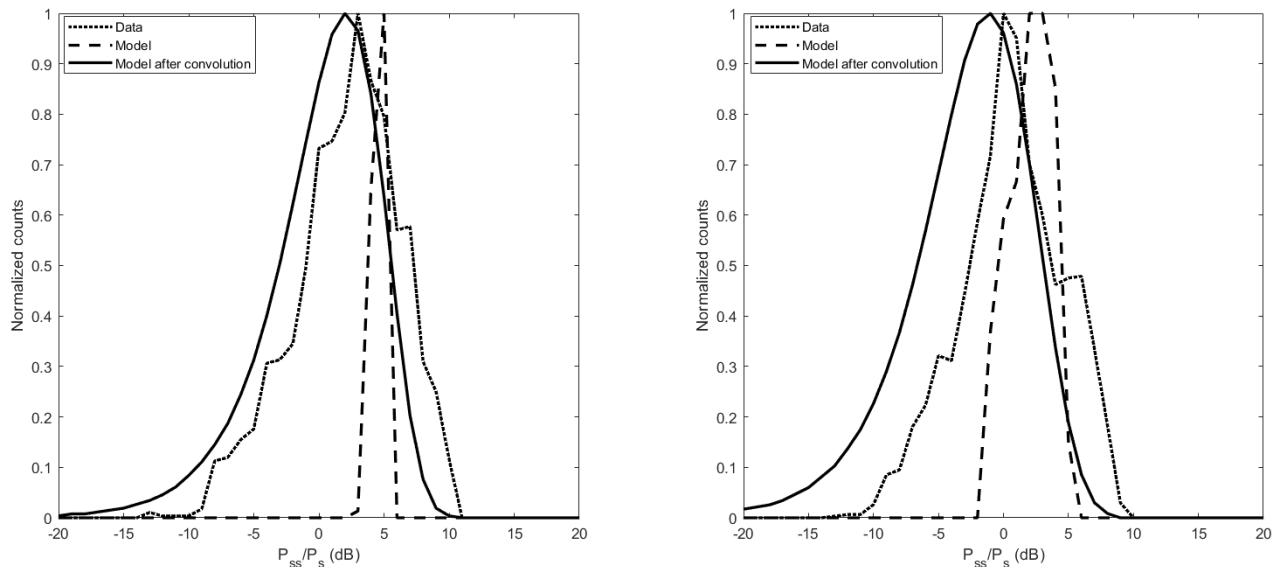


Fig. 10. Comparison between the histogram of measured values of subsurface to surface echo power ratio at 4 (left) and 5 MHz (right) with the histogram of model results in the neighbourhood of a given local maximum (see text for details), and with the synthetic histogram of model results produced through the convolution of the histogram of model results with the distribution template shown in Figure 8. The histograms are computed both for model results at 4 and 5 MHz.

To make a quantitative comparison of model results with data, the synthetic histograms were used to generate a set of random values representing measurements. This was done by taking each histogram bin, and generating a set of uniformly distributed random numbers in an interval corresponding to the bin, whose quantity is proportional to the bin height. Finally, the set of synthetic measurements was compared to actual measurements through a statistical test to determine if the two sets of samples could be produced by the same distribution. We used the two-sample Kolmogorov-Smirnov test (Kolmogorov, 1933; Smirnov, 1939), which is a nonparametric test for the null hypothesis that the data in two sets of samples are from the same continuous distribution. It is widely employed in statistics for its robustness and its sensitivity to both location and shape of distributions. The test was performed using the implementation available in the MATLAB^(TM) proprietary programming language, and setting the significance level for the rejection of the null hypothesis at 5%.

A configuration of layers producing a local maximum at one frequency does not result in a local maximum at the other frequency. For this reason, both local maxima at 4 MHz and at 5 MHz had to be considered. Thus, for a given model (either two- or three-CO₂ ice layers) we simultaneously tested samples of model results at 4 and 5 MHz extracted for a set of layer thicknesses corresponding to a local peak of basal echo power at either 4 or 5 MHz, and for all layer thicknesses within two meters of the one producing the peak. This resulted in four possible combinations of number of layers and frequency of the echo power maximum. We found and tested 487 local maxima at 4 MHz and 687 at 5 MHz for the two-layer model, and 2981 local maxima at 4 MHz and 5232 at 5 MHz for the three-layer model. Results are shown in Figure 11.

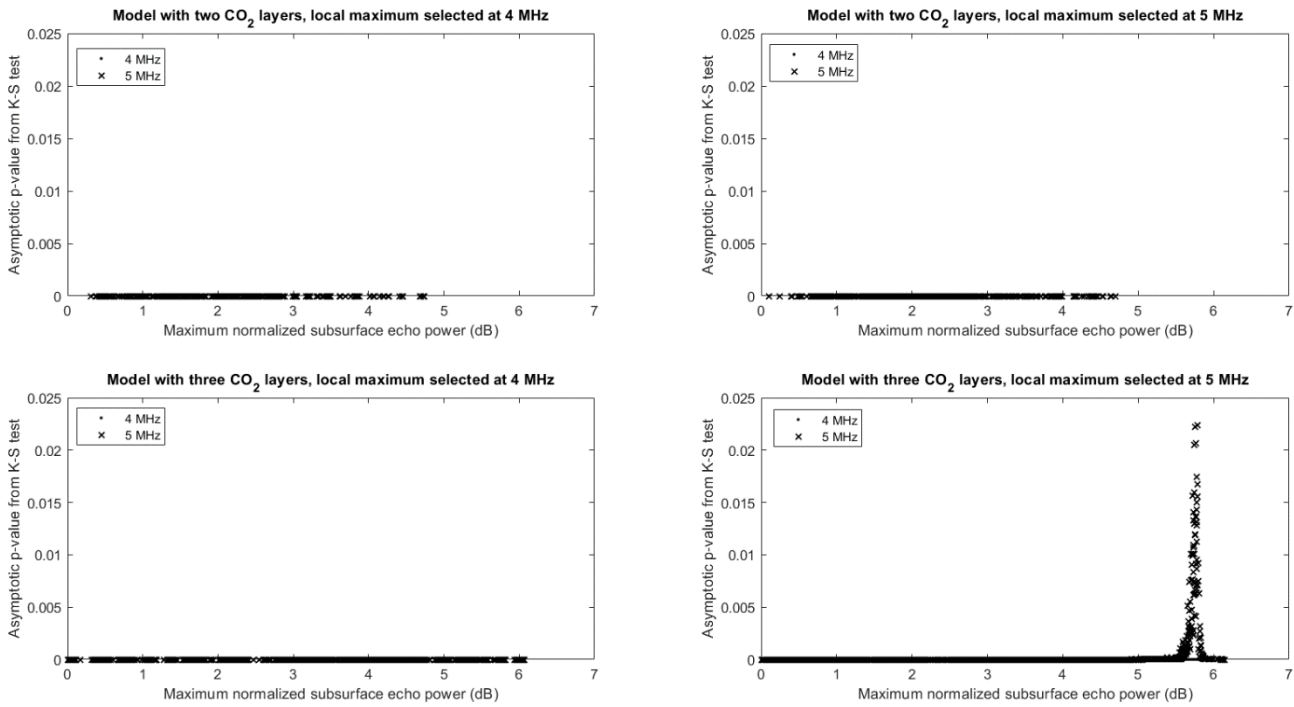


Fig. 11. Plots of the asymptotic p-values produced by the Kolmogorov-Smirnov test versus the maximum subsurface to surface power ratio of simulated echoes, for models with two- and three-CO₂ ice layers and for local maxima at 4 and 5 MHz. Each point represents the probability that model results and MARSIS measurements belong to the same continuous distribution. For each tested set of simulations, probabilities are computed both at 4 MHz (dots) and at 5 MHz (crosses).

No set of two- and three-layer models resulted in simulated values exceeding the 5% significance level threshold. However, it can be seen in the bottom right plot in Figure 11 that a few layer configurations resulting in a local maximum of basal echo power at 5 MHz approach the 5% threshold. This happens only at 5 MHz (crosses in Fig. 11), while at 4 MHz (dots) the probability remains negligible. This quantitatively confirms the observation that, in a model based on the coherent interference between layers, a configuration capable of producing a constructive interference at one frequency is not capable of doing the same at another frequency.

None of the modelled scenarios involving CO₂ ice layers reproduce the distribution of basal echo intensities observed by MARSIS over the bright areas in Ultimi Scopuli. Median values of the simulated basal echo power are lower than those observed in real MARSIS data, and the spread of the simulated values does not match the range of the measured data. We thus conclude that the interpretation that best reconciles all radar observations in Ultimi Scopuli remains that of a strong basal reflector overlain by a thick dust-laden SPLD.

5. Conclusions

The results of our simulations do not support claims that CO₂ ice layers interbedded with water ice can produce the bright reflections detected by MARSIS from the base of the SPLD in Ultimi Scopuli. Therefore, we rule out the presence of hidden deposits of basal CO₂ ice in this area. Based on these results, and the lack of detection of sequestered CO₂ ice deposits elsewhere in the SPLD (Whitten et al., 2017; Whitten and Campbell, 2018; Khuller and Plaut, 2021) beyond those reported by Phillips et al. (2011) and Bierson et al. (2016), we posit that it is unlikely that hidden CO₂ was stored in the late Amazonian (< 250 Ma) south polar deposits. This conclusion is consistent with the geological record of diminishing volcanic activity on Mars during the Hesperian and throughout the Amazonian, with very small amounts of CO₂ released into the atmosphere by volcanic degassing in post-Noachian times.

567 With the polar caps being the most voluminous reservoir of volatiles on Mars, and their role in regulating the
568 planetary climate, records of the composition and physical properties of the polar layered deposits are fundamental to
569 constrain the entire geological and climatological history of Mars. The ground penetrating radar MARSIS continuing
570 acquisition of observations over the polar regions is an essential element of our investigations into the origin,
571 evolution, and inventory of Martian volatiles and of the planet evolution. This is why interpretations of MARSIS data
572 must be critically assessed against the framework of the physical theory of the propagation of radar waves through
573 stratified media, as we have demonstrated in this paper. Rigorous application of this approach to all investigations
574 based on MARSIS data will advance our knowledge of the nature and properties of Mars's polar materials.

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