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Title	Discovery of TOI-1260d and the characterization of the multiplanet system
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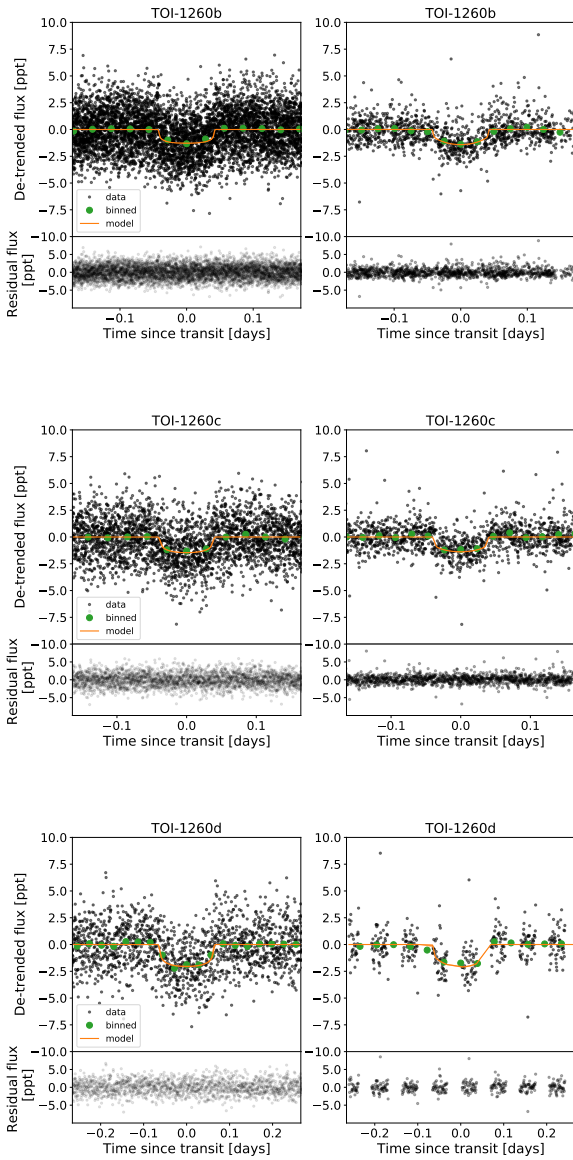


Figure 2. Phase folded light curves of TOI-1260b (top), TOI-1260c (middle), TOI-1260d (bottom). The TESS data are shown in the left panels and the CHEOPS data are shown in the right panels. Residuals of each transit are shown below each phase-folded light curves. The phase binned data are denoted by green points and the orange line shows the best-fitted transit models for each planet.

mass determination of planets b and c are consistent within 1-sigma with values derived in G21.

The mass precision of the planets is the main source of uncertainty in the determination of the planetary bulk densities in the system. This work highlights the need to strategically obtain more RVs for the system in order to understand the effect of stellar activity on the RVs of the system and better constrain the planetary masses.

6 DISCUSSION

The follow-up photometric observations of TOI-1260 allows the precise characterisation of the two inner transiting planets and confirms the planetary nature of the transiting outer planetary companion. Figure 5 shows the mass-radius diagram of known exoplanets with masses below $30 M_{\oplus}$ and radii less than $4 R_{\oplus}$. We proceed now with the discussion of the interior composition and atmospheric evolution of the planetary system.

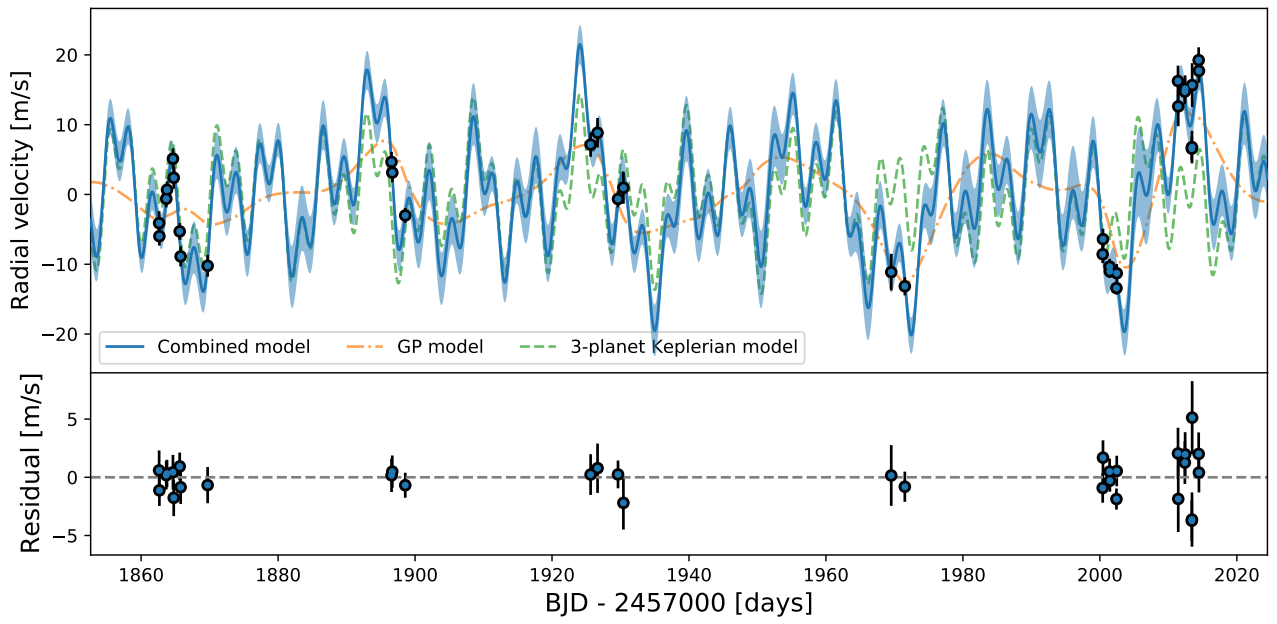


Figure 3. *Top:* Time-series HARPS-N RVs of TOI-1260. The RVs were modelled using a three-planet Keplerian RV model and a GP simultaneously to model the activity-induced RV variations (see Section 5). The green dash line shows the 3-planet Keplerian model and the orange dash-dot line shows the GP model that accounts for activity induced RV variations. The blue solid line shows the median three-planet Keplerian + GP model. The 1-sigma credible intervals of the best fit Keplerian + GP model is indicated by the blue shaded region.

Bottom: Residuals of the RV data.

Table 3. System parameters obtained from the joint light curves and radial velocities analysis. The median values and 1-sigma uncertainty are reported.

Parameter [Unit]	Planet b	Planet c	Planet d
<i>Fitted parameters</i>			
Period P [day]	3.127463 ± 0.000005	7.493134 ± 0.000020	16.608164 ± 0.000083
Epoch T0 [BJD-2457000]	2065.564269 ± 0.000396	2068.270505 ± 0.000577	2062.017406 ± 0.001309
Planet-to-Stellar radius ratio [Rp/Rs]	0.0329 ± 0.0006	0.0377 ± 0.0007	0.0425 ± 0.0009
Impact parameter b	0.20 ± 0.12	0.75 ± 0.02	0.53 ± 0.05
Radial velocity semi-amplitude K [m s ⁻¹]	4.93 ± 0.83	5.67 ± 1.77	3.90 ± 2.54
Eccentricity e	0 (adopted)	0 (adopted)	0 (adopted)
Angle of periastron ω [°]	0 (adopted)	0 (adopted)	0 (adopted)
<i>Derived parameters</i>			
Transit duration T14 [hr]	2.06 ± 0.02	1.97 ± 0.03	3.19 ± 0.07
Transit depth [ppm]	1082 ± 37	1421 ± 55	1808 ± 78
Scaled semi-major axis a/Rs	11.73 ± 0.35	20.99 ± 0.63	35.69 ± 1.06
Orbital semi-major axis a [au]	0.0367 ± 0.0011	0.0657 ± 0.0020	0.1116 ± 0.0033
Inclination i [deg]	89.03 ± 0.61	87.97 ± 0.11	89.14 ± 0.10
Planet radius R_p [R_\oplus]	2.41 ± 0.05	2.76 ± 0.07	3.12 ± 0.08
Planet mass M_p [M_\oplus]	8.56 ± 1.54	13.20 ± 4.23	11.84 ± 7.79
Planet density ρ_p [$g\ cm^{-3}$]	3.35 ± 0.64	3.45 ± 1.14	2.14 ± 1.42
Planet surface gravity $\log g_p$	3.16 ± 0.09	3.23 ± 0.15	3.08 ± 0.30
Equilibrium dayside temperature [K]	871 ± 24	651 ± 18	499 ± 14
Stellar insolation [S_\oplus]	95.58 ± 0.07	29.81 ± 0.05	10.32 ± 0.07
TESS instrument offset σ_{TESS} [ppm]	64.0 ± 8.6		
CHEOPS instrument offset σ_{TESS} [ppm]	48.8 ± 14.7		
HARPS jitter σ_{HARPS} [m s ⁻¹]	$0.22 + / - 0.79$		
Systemic radial velocity γ_{HARPS} [m s ⁻¹]	10.73 ± 2.63		
Limb darkening parameter $u_{1,TESS}$	0.21 ± 0.18		
Limb darkening parameter $u_{2,TESS}$	0.53 ± 0.26		
Limb darkening parameter $u_{1,CHEOPS}$	0.92 ± 0.18		
Limb darkening parameter $u_{2,CHEOPS}$	-0.33 ± 0.21		
<i>GP RotationTerm parameters</i>			
GP rotation period $R_{rot,GP}$ [day]	30.63 ± 3.81		
σ_{GP}	6.62 ± 1.45		
$Q0$	0.83 ± 1.48		
dQ	1.94 ± 3.67		
f	0.70 ± 0.23		
Stellar mass M_s [M_\odot]	0.67 ± 0.06		
Stellar radius R_s [R_\odot]	0.67 ± 0.01		
Stellar density ρ_s [$g\ cm^{-3}$]	3.12 ± 0.33		

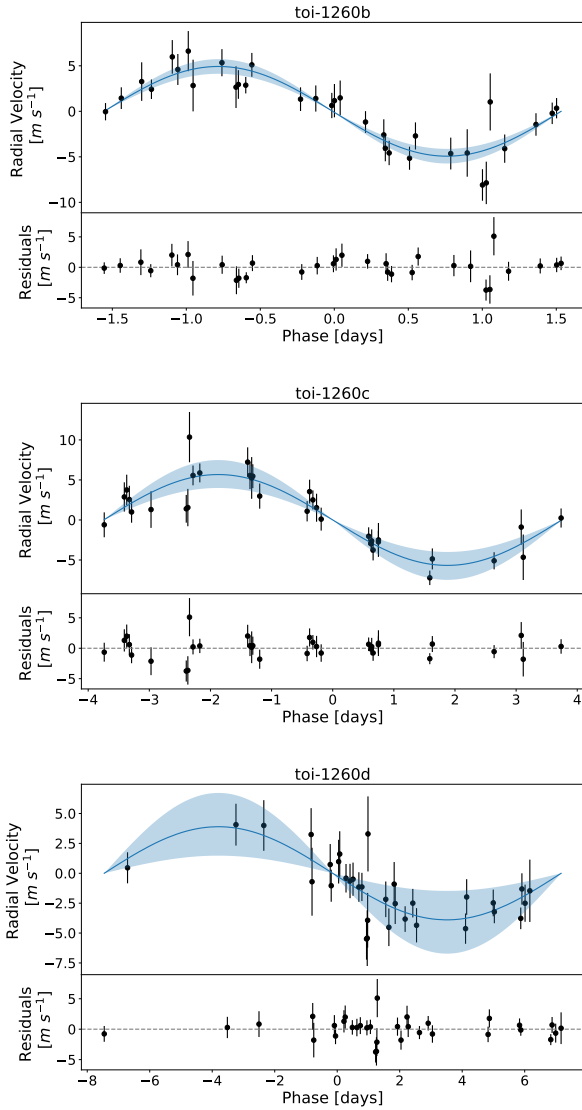


Figure 4. Phase-folded radial velocities and residuals of TOI-1260b (top), TOI-1260c (middle), TOI-1260d (bottom). The best-fit RV models are indicated by the solid blue line and the corresponding 1-sigma credible interval is shown by the blue shaded region.

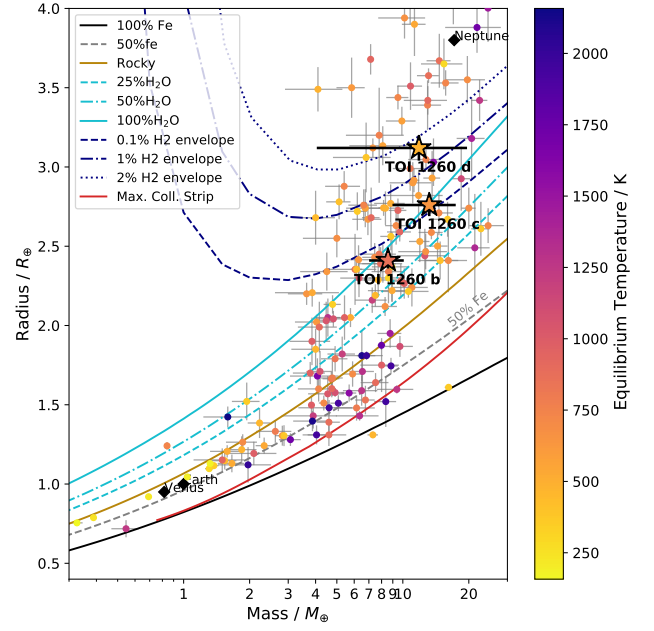


Figure 5. Mass-radius diagram showing low mass planets in the range of $0.4\text{--}30 M_{\oplus}$ which have mass and radius precision measured to better than 30% and 15%, respectively. TOI-1260b, TOI-1260c, TOI-1260d, are indicated by the star symbols. All exoplanets are colour-coded according to their the equilibrium dayside temperatures as shown in the colour bar. The different lines plotted are the theoretical mass-radius relations corresponding to the planet interior compositions (Zeng et al. 2019).

6.1 Interior composition of the planets

The TOI-1260 system has three sub-Neptune transiting exoplanets where planets b, c and d have masses of $8.56 \pm 1.54 M_{\oplus}$, $13.20 \pm 4.23 M_{\oplus}$ and $11.84 \pm 7.79 M_{\oplus}$, respectively, and their radii are $2.41 \pm 0.05 R_{\oplus}$, $2.76 \pm 0.07 R_{\oplus}$, and $3.12 \pm 0.08 R_{\oplus}$, respectively. This means that the three sub-Neptunes TOI-1260 b, c and d have bulk densities of $3.35 \pm 0.64 \text{ g cm}^{-3}$, $3.45 \pm 1.14 \text{ g cm}^{-3}$, and $2.14 \pm 1.42 \text{ g cm}^{-3}$, respectively. Figure 5 shows the distribution of known exoplanet with precise mass and radius measurements in the mass-radius diagram, alongside some theoretical mass-radius relations for different planet interior compositions. The interior of TOI-1260 b is likely to be consisted of up to 50% rocky core and a 50% H₂O layer. In the case of TOI-1260 c, the sub-Neptune planet is likely a water world or it could be composed of a water-rich core with a small fraction of H₂ atmosphere. For the outermost planet TOI-1260 d, its interior is likely to consist of a water-rich or Earth-like rocky core with up to $\sim 2\%$ of H₂ atmosphere.

The interior compositions of exoplanet correlates with the compositions of their host stars (Adibekyan et al. 2021a). This is because they were formed from accretion of the same disk material. Therefore, using physical parameters of the host star in addition to the planet’s mass and radius provides a better constrain to the planet’s interior composition. Using the values of radius, mass, and stellar properties derived in Section 5, we performed an analysis of the internal structure of the three planets in the TOI-1260 system. Our method is based on a global Bayesian model that fits the observed properties of the star (mass, radius, age, effective temperature, and the photospheric abundances [Si/Fe] and [Mg/Fe]) and planets (planet-star radius ratio, the RV semi-amplitude, and the orbital period). The hidden parameters in the Bayesian model are, for each planet, the masses of solids (everything except the H or He gas), the mass fractions of the core, mantle and water, the mass of the gas envelope, the Si/Fe and Mg/Fe mole ratios in the planetary mantle, the S/Fe mole ratio in the core, and the equilibrium temperature. All details on the methods are presented in Leleu et al. (2021).

The Bayesian analysis relies on a forward models that computes the expected planetary radius and bulk internal structure as a function of the hidden parameters. In the forward model, we assume a fully differentiated planet made of a core (composed of Fe and S), a mantle (composed of Si, Mg, Fe, and O), a pure water layer, and a H and He layer. The temperature profile is adiabatic, and the equations of state (EoS) used for these calculations are taken from Hakim et al. (2018) and Fei et al. (2016) for the core materials, from Sotin et al. (2007) for the mantle materials, and Haldemann et al. (2020) for water. The thickness of the gas envelope is determined as a function of the gas mass fraction, the equilibrium temperature, the mass and radius of the solid planet, and the age (assumed to be equal to the stellar age), using the semi-analytical model of Lopez & Fortney (2014). Importantly, the radius of the high-Z part of the planet (core, mantle and water layer) is computed independently of the thickness of the gas layer. This implies in particular that the compression effect of the gas envelope onto the core, as well as the effect of the temperature at the basis of the gas envelope are not included in the mode.

The Bayesian analysis is done assuming the following priors: the mass fractions of the planetary cores, mantles, and water layers have uniform positive priors (the mass fractions of water being limited to a maximum value of 0.5). The prior on the gas mass is uniform in log, and the bulk Si/Fe and Mg/Fe mole ratios in the planet are

assumed to be equal to the values determined for the atmosphere of the star, given above ³.

The posterior distribution of the main planetary hidden parameters are presented in Fig. 6. All planets have some fraction of gas, the mass of gas increasing for decreasing equilibrium temperatures (see Fig. 7). The fraction of water, on the other hand, is essentially unconstrained.

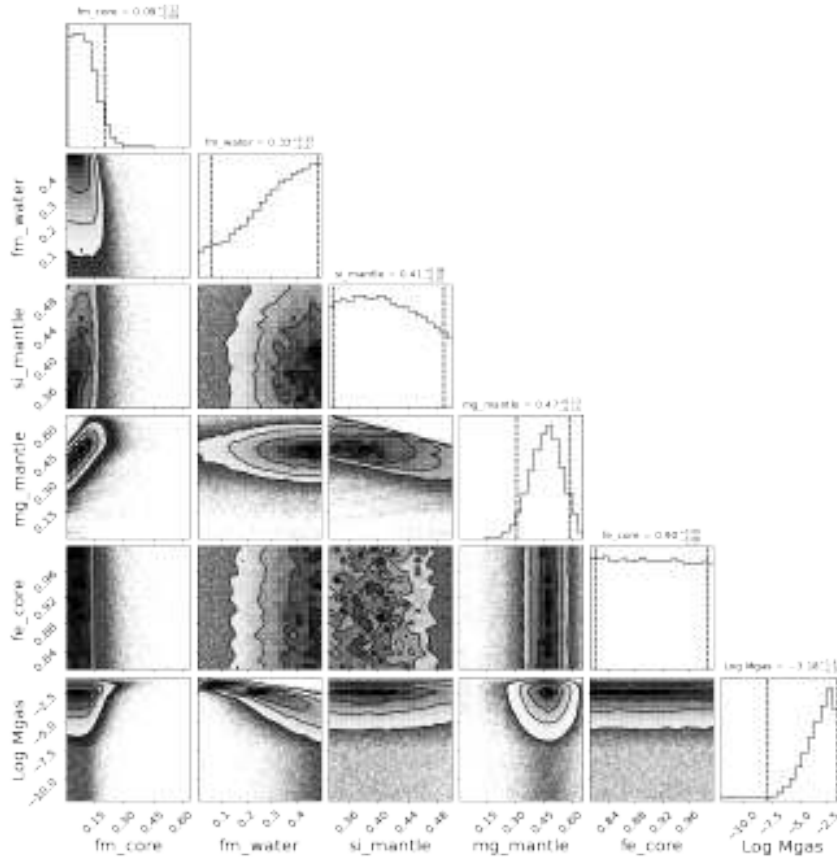
6.2 Atmospheric evolution

We considered the stellar and planetary parameters derived in our paper, as well as the present-day planetary atmospheric mass fractions presented in Section 6.1, to reconstruct the evolution of the stellar rotation rate and of the planetary atmospheres. In particular, we constrain the evolution of the stellar rotation period, which we use as proxy for the evolution of the stellar high-energy emission affecting atmospheric escape, and the predicted initial atmospheric mass fraction of the detected planets $f_{\text{atm}}^{\text{start}}$, that is the mass of the planetary atmosphere at the time of the dispersal of the protoplanetary disk, which we assume being at 5 Myr.

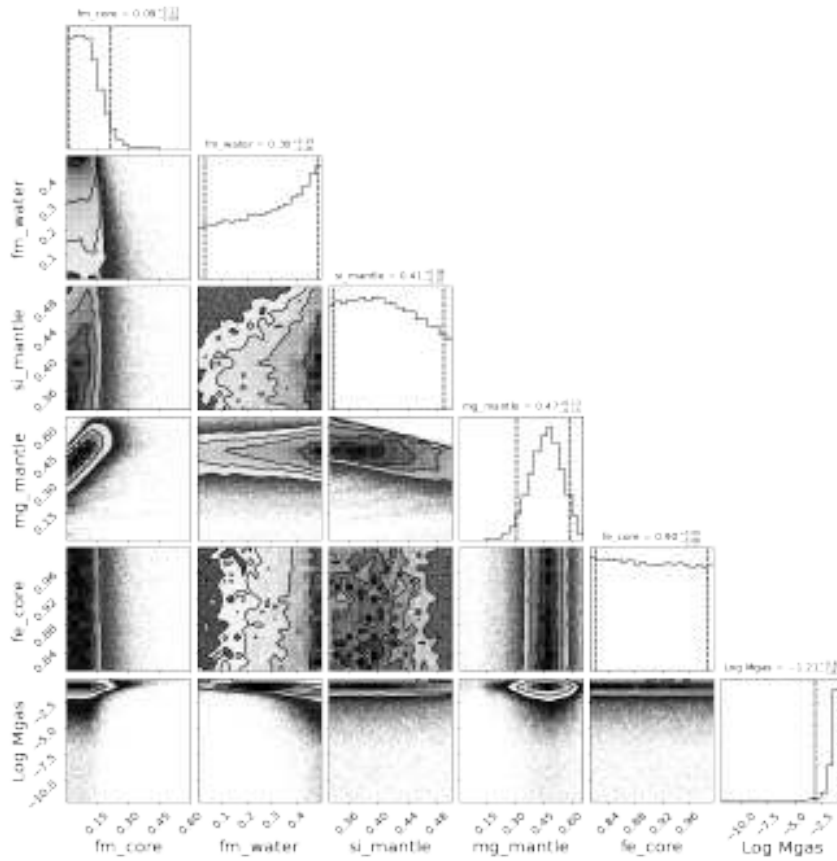
We reach these results by using the Planetary Atmospheres and Stellar RoTation RAtes (PASTA; Bonfanti et al. 2021b) code, which is an updated version of the original code presented by Kubyschkina et al. (2019c,a). In short, PASTA constrains the evolution of planetary atmospheres and of the stellar rotation rate combining a model predicting planetary atmospheric escape rates based on hydrodynamic simulations (this has the advantage over other commonly used analytical estimates to account for both XUV-driven and core-powered mass loss; Kubyschkina et al. 2018), a model of the stellar high-energy (X-ray plus extreme ultraviolet; XUV) flux evolution (Bonfanti et al. 2021b), a model relating planetary parameters and atmospheric mass (Johnstone et al. 2015b), and stellar evolutionary tracks (Choi et al. 2016). PASTA works under two main assumptions: 1) planet migration did not occur after the dispersal of the protoplanetary disk; 2) the planets hosted at some point in the past or still host a hydrogen-dominated atmosphere. PASTA returns realistic uncertainties on the free parameters (i.e. the planetary initial atmospheric mass fractions at the time of the dispersal of the protoplanetary disk, and the indexes of the power law controlling the stellar rotation period that is used as proxy for the stellar XUV emission) by implementing the atmospheric evolution algorithm in a Bayesian framework (Cubillos et al. 2017), using the system parameters with their uncertainties as input priors. All details of the algorithm can be found in Bonfanti et al. (2021b). The only difference with respect to the analysis of the systems considered by Bonfanti et al. (2021b) is that here we fit the planetary atmospheric mass fractions given in Section 6.1 instead of the planetary radii. This enables the code to be more accurate by avoiding the continuous conversion of the atmospheric mass fraction into planetary radius, given the other system parameters (see e.g. Delrez et al. 2021).

Figure 8 shows the results obtained from PASTA. As a proxy for the evolution of the stellar rotation period, Figure 8 displays the posterior distribution of the stellar rotation period at an age of 150 Myr ($P_{\text{rot},150}$), also in comparison to that of stars member of young open clusters and of comparable mass extracted from Johnstone et al. (2015a). The posterior distribution is slightly shifted towards slower rotation compared to that of the open cluster stars, indicating that the

³ It should be noted however that Adibekyan et al. (2021b) has found that despite an existing correlation between the abundances of planets and host stars, the relation is not always strictly one-to-one.



(a)



(b)

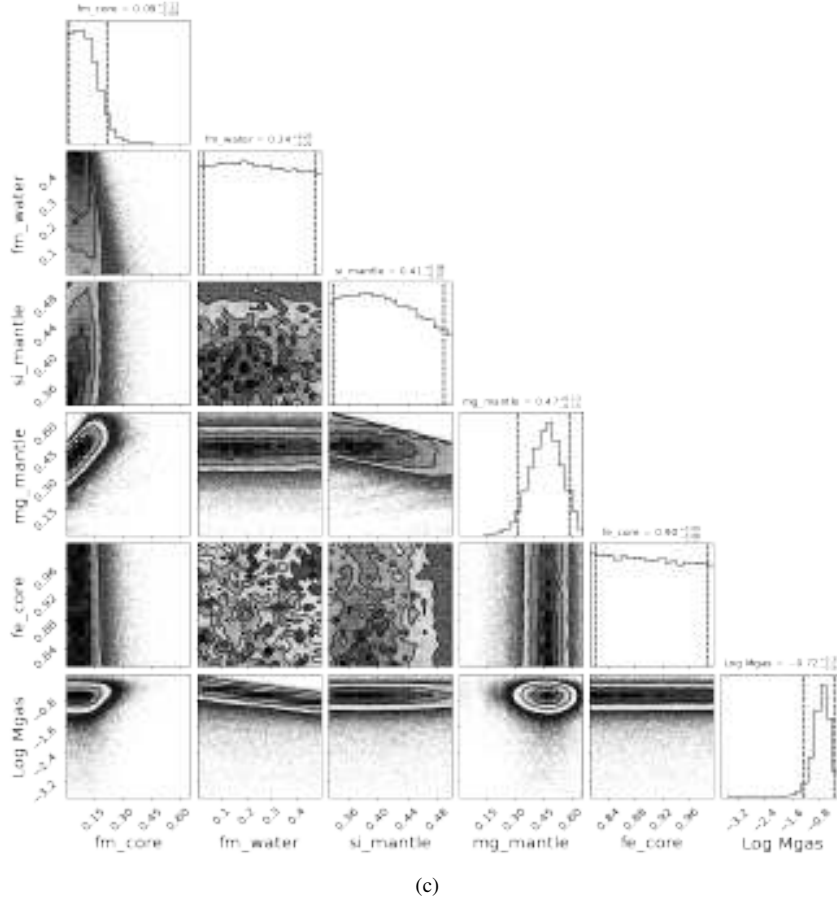


Figure 6. Corner plot showing the results on the interior composition models of (a) TOI-1260 b, (b) TOI-1260 c and (c) TOI-1260 d. The vertical dashed lines and the 'error bars' given at the top of each columns represent the 5% and 95% percentiles.

planets were likely subject to somewhat less XUV radiation than the average.

Figure 8 shows also the posterior distribution of the initial atmospheric mass fraction for planets b (in linear scale), c (in logarithmic scale), and d (in logarithmic scale) in comparison to the present-day atmospheric mass fraction (Section 6.1). The posterior distribution for planet b is flat, indicating that the planet has most likely lost (almost) entirely its primordial hydrogen-dominated envelope through escape at some point in the past, which is why PASTA is unable to constrain the initial atmospheric mass fraction. Figure 8 indicates that also planets c and d have gone through significant evolution through escape that has significantly eroded the primordial atmospheric content, which was however small in comparison to the planetary masses. Therefore, we conclude that both planets (i.e. c and d) accreted a small hydrogen envelope during the formation process compared to their masses. This may have been the result of several physical mechanisms, such as late planet formation compared to the age of the protoplanetary disk, early dispersal of the protoplanetary disk, low gas content of the disk.

As the isochronal age is loosely constrained, we performed additional evolution runs by artificially making the star much younger or older, further imposing tighter constraints on the stellar age. Despite the different evolutionary time scales, we did not find significant changes in the $f_{\text{atm}}^{\text{start}}$ of the planets. This is because (1) atmospheric mass loss is significant only during the first Myrs of evolution and (2) $f_{\text{atm,c}}^{\text{start}}$ and $f_{\text{atm,d}}^{\text{start}}$ are always found to be rather small, indicating that

the constraints given by system parameters prevent those planets to host a massive initial atmosphere regardless of the age of the system.

The current stellar XUV fluxes impinging on each planet are $F_{\text{XUV,b}} = 2.87 \cdot 10^4 \text{ erg}/(\text{cm}^2 \text{ s})$, $F_{\text{XUV,c}} = 8.97 \cdot 10^3 \text{ erg}/(\text{cm}^2 \text{ s})$, and $F_{\text{XUV,d}} = 3.10 \cdot 10^3 \text{ erg}/(\text{cm}^2 \text{ s})$. The correspondent mass-loss rate values expected for the planets right now are $\dot{M}_b = 10^{10} \text{ g/s}$, $\dot{M}_c = 1.59 \cdot 10^9 \text{ g/s}$, and $\dot{M}_d = 7.43 \cdot 10^8 \text{ g/s}$. Assuming that the stellar XUV flux does not change over time in the future, which is a reasonable assumption given the old age of the star, these values imply that in the next Gyr the planets are respectively going to lose 0.6%, 0.06%, and 0.03% of their mass. From Fig. 7 these values then imply that planet b is going to lose entirely its hydrogen-dominated envelope, while planets c and d are going to keep it. As the results of planet b are consistent with no hydrogen atmosphere at all, it is unlikely that the position of these planets in the period-radius diagram (e.g. [Fulton et al. 2017](#)) is going to change in the future.

7 CONCLUSIONS

We presented the follow-up observations of the TOI-1260 system using CHEOPS and TESS. The addition of the recent photometric dataset allow us to refine the physical parameters of the planetary system and discover a third additional transiting planet. For planets TOI-1260 b and c, we found that the radii are $2.36 \pm 0.06 R_{\oplus}$, $2.82 \pm 0.08 R_{\oplus}$, respectively, and the masses $8.52 \pm 1.45 M_{\oplus}$ and $13.29 \pm$

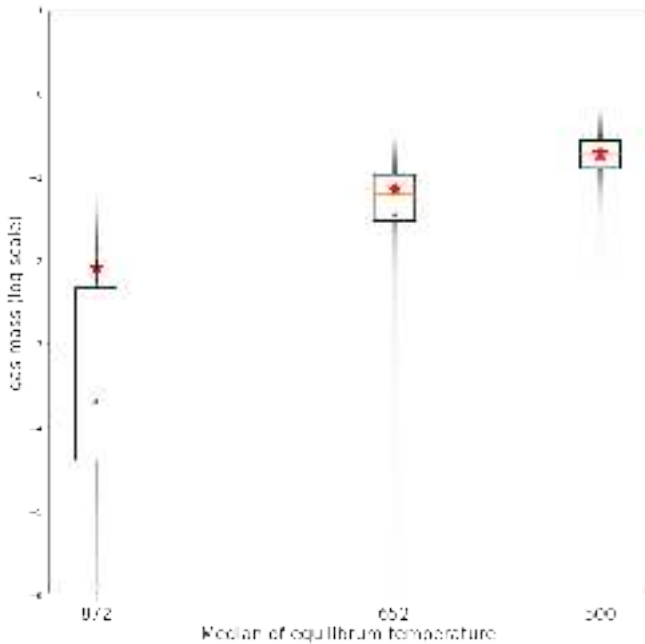


Figure 7. Gas fraction in the planets as a function of their equilibrium temperature. The box show the 25 % and 75 % percentiles, the orange line represents the median of the posterior distribution, the green triangle is the mean, and the red stars is located at the mode of the posterior distribution. Finally, the opacity of the thick vertical black line is proportional to the posterior distribution.

$3.94 M_{\oplus}$. The newly discovered TOI-1260-d has bulk properties $3.01 \pm 0.09 R_{\oplus}$ and $11.8 \pm 7.5 M_{\oplus}$.

The detailed characterization of the planetary parameters allows us to derive constraints of their internal composition and evolution that we related to the formation processes in the system and its future evolution.

The TOI-1260 system presents an exciting opportunity for comparative exoplanetology using JWST transmission spectroscopy. Moses et al. (2013) predicted that sub-Neptune sized exoplanets such as those in the TOI-1260 system can harbour a large diversity of atmospheric compositions. Multi-planet systems such as TOI-1260 give us the opportunity to test whether such diversity can exist within different sub-Neptunes in the same system. All three of the planets in the TOI-1260 system appear to be favourable for atmospheric categorisation with JWST, with transmission spectroscopy metrics (TSMs Kempton et al. 2018) of 43.6, 36.1 and 40.4 for planets b, c, and d, respectively. Figure 9 shows how this compares to similar multi-planet systems as a function of planetary radius and semi-major axis. In addition, due to its high northern declination TOI-1260 is particularly favourable for JWST visibility, with observations possible for 196 days each year (Bourque et al. 2021).

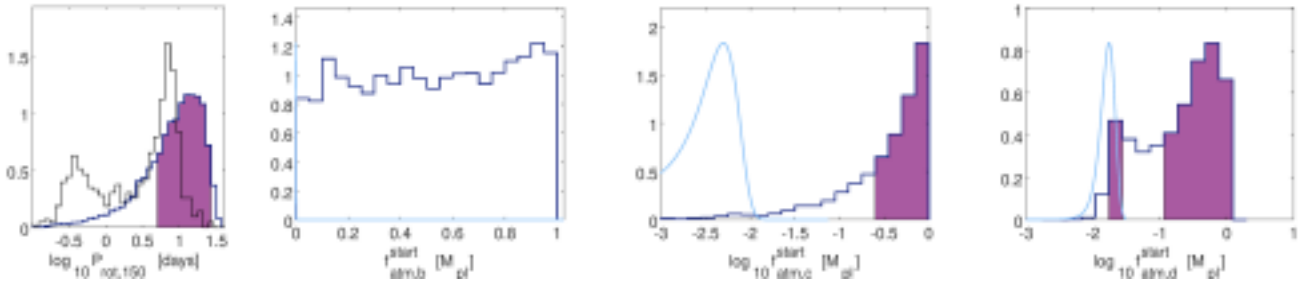


Figure 8. Posterior probability density functions (PDFs) of the stellar rotation period when TOI-1260 was 150 Myr old ($P_{\text{rot},150}$) and of the initial atmospheric mass fraction ($f_{\text{atm}}^{\text{start}}$) of the hosted exoplanets. The purple areas show the 68%-HPD (highest posterior density) interval. *Leftmost panel.* $P_{\text{rot},150}$ PDF (dark blue histogram) to be compared with the rotation period distribution of stars of comparable masses that belong to coeval open clusters (black histogram; data taken from [Johnstone et al. 2015a](#)). *Other panels.* Atmospheric mass fractions PDFs of planet b (linear scale) and of planet c and d (log scale). The light blue curve is the present-day atmospheric content, as inferred from our internal structure analysis. See text for details.

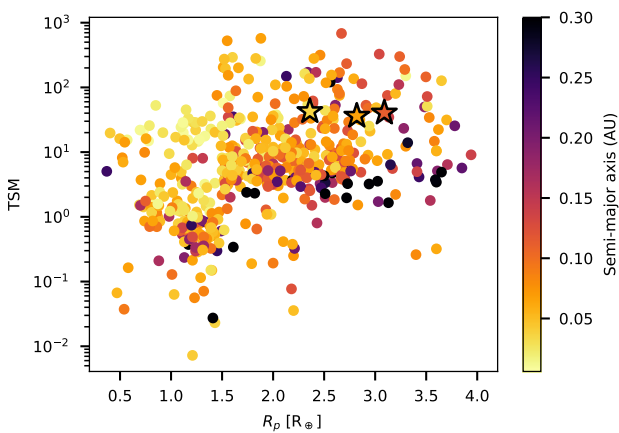


Figure 9. Unscaled transmission spectroscopy metrics (TSM) of all multi-planet systems with host stars of K-type and later as a function of planetary radius, with orbital separation visible in the colour scale. The three planets in the TOI-1260 system are marked with stars.

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KGI is the ESA CHEOPS Project Scientist and is responsible for the ESA CHEOPS Guest Observers Programme. She does not participate in, or contribute to, the definition of the Guaranteed Time Programme of the CHEOPS mission through which observations described in this paper have been taken, nor to any aspect of target selection for the programme.

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DATA AVAILABILITY

TESS data are publicly available in the Space Telescope Science Institute (STScI) at <https://mast.stsci.edu>. CHEOPS data generated and analysed in this article will be made available in the CHEOPS mission archive (https://cheops.unige.ch/archive_browser/). No new HARPS RV data were generated for this work.

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APPENDIX A: EXTRA MATERIAL

We present in Figure A1 the posterior distribution of the fitted parameters.

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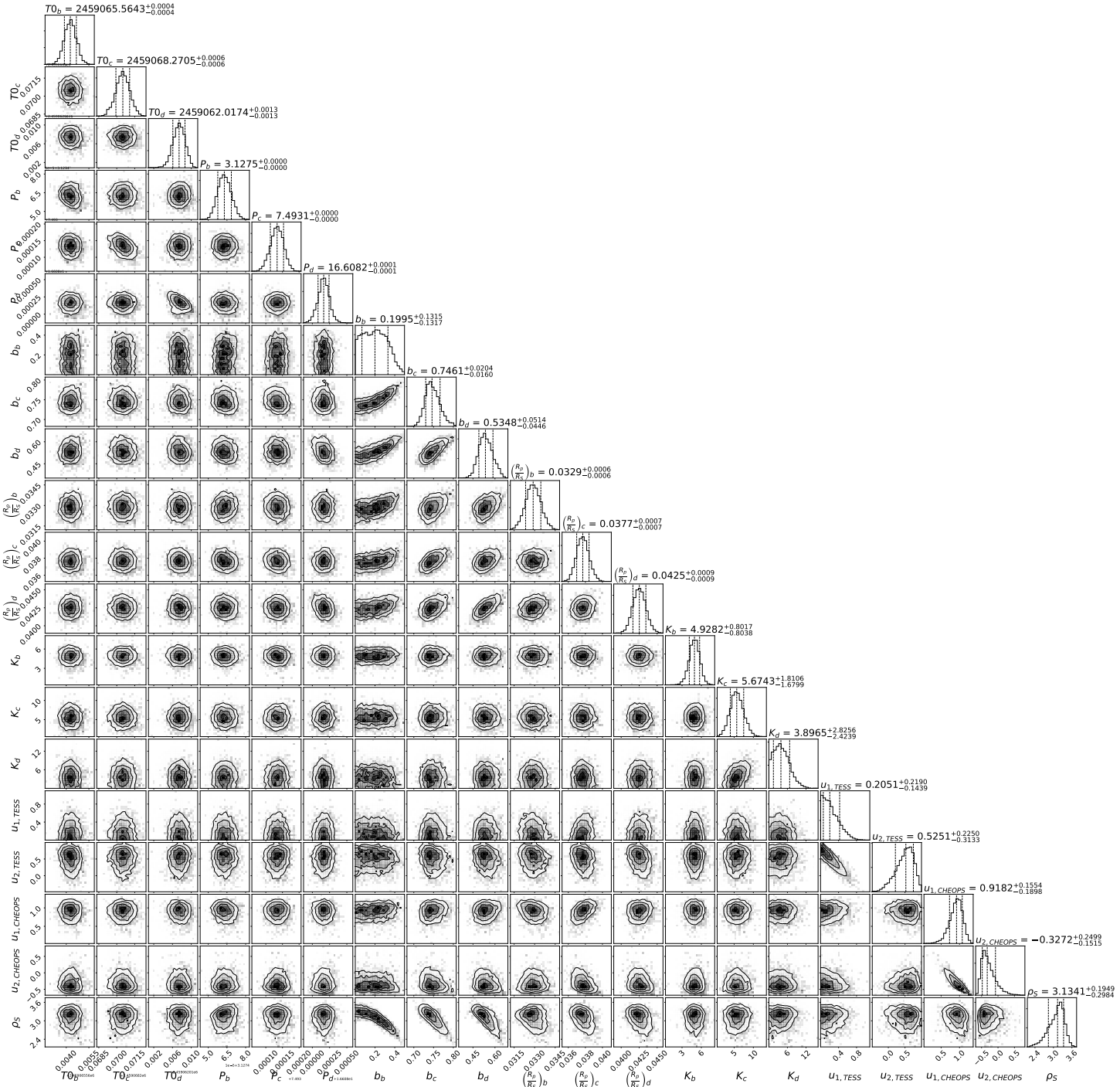


Figure A1. Corner plot showing the posterior distribution of the fitted parameters.