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Signs of environmental effects on star-forming galaxies in the Spiderweb protocluster at $z = 2.16$

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ABSTRACT

We use multi-object near-infrared spectroscopy with VLT/KMOS to investigate the role of the environment in the evolution of the ionized gas properties of narrow-band-selected $H\alpha$ emitters (HAEs) in the Spiderweb protocluster at $z = 2.16$. Based on rest-frame optical emission lines, $H\alpha$ and $[N\text{II}]\lambda 6584$, we confirm the cluster membership of 39 of our targets (i.e. 93 per cent success rate), and measure their star formation rates (SFR), gas-phase oxygen abundances, and effective radius. We parametrize the environment where our targets reside using local and global density indicators based on previous samples of spectroscopic and narrow-band cluster members. We find that star-forming galaxies embedded in the Spiderweb protocluster display SFRs compatible with those of the main sequence and morphologies comparable to those of late-type galaxies at $z = 2.2$ in the field. We also report a mild gas-phase metallicity enhancement (0.06 ± 0.03 dex) at intermediate stellar masses. Furthermore, we identify two UVJ-selected quiescent galaxies with residual $H\alpha$ -based star formation and find signs of extreme dust obscuration in a small sample of starburst submillimetre galaxies based on their FIR and $H\alpha$ emission. Interestingly, the spatial distribution of these objects differs from the rest of HAEs, avoiding the protocluster core. Finally, we explore the gas fraction–gas metallicity diagram for seven galaxies with molecular gas masses measured by ATCA using CO(1–0). In the context of the gas-regulator model, our objects are consistent with relatively low mass-loading factors, suggesting lower outflow activity than field samples at the cosmic noon and thus, hinting at the onset of environmental effects in this massive protocluster.

Key words: galaxies: abundances – galaxies: clusters: individual: PKS 1138-262 – galaxies: evolution – galaxies: high-redshift – galaxies: star formation.

1 INTRODUCTION

Local galaxy clusters host a very distinct distribution of galaxy populations with respect to the field. While the latter is dominated by galaxies that are blue in colour, and have relatively high star formation rates (SFRs) and disc-like morphologies, the former mainly host redder objects with low levels of star formation and triaxial shapes (e.g. Dressler 1980; Balogh et al. 1998; Hogg et al. 2003). These discrepancies in terms of galaxy populations strongly correlate with the cluster total mass and with the distance to their centres, suggesting that the impact of the environment is controlled by both the accretion history of the clusters (Dekel & Birnboim 2006) and the spatial and

temporal trajectory that galaxies follow in them (e.g. Muzzin et al. 2014; Haines et al. 2015; Rhee et al. 2017).

This qualitative description has been extensively studied from a quantitative point of view over the last decades (Barsanti et al. 2021) crystallizing into several correlations that expose the gradual effect of the environment over the galaxies' physical properties. Among them, we find the morphology–density relation (Dressler 1980), the colour–density relation (Hogg et al. 2003), and the star formation–density relation (Balogh et al. 1998; Lewis et al. 2002), which describe the discrepancies between the cluster and field outlined above. Furthermore, signs of enhanced metal enrichment have also been found in the densest regions of clusters, specially for relatively low-mass galaxies (e.g. Petropoulou et al. 2011; Petropoulou, Vilchez & Iglesias-Páramo 2012; Paulino-Afonso et al. 2018; Ciocan et al. 2020). The simultaneous manifestation of all these correlations

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demonstrates the influence of the environment in the accelerated transformation of galaxy properties. However, the exact mechanism behind these changes is still a matter of debate since we need to find cluster-specific processes that account for the dynamical transformation of cluster galaxies (Mortlock et al. 2013; Swinbank et al. 2017) as well as for the evolution of their interstellar medium (ISM, Peng & Maiolino 2014; Maiolino & Mannucci 2019).

Several possibilities have been proposed to tackle this conundrum, with mergers and repeated tidal interactions (i.e. harassment, Moore et al. 1996) likely driving the dynamical transformation while ram-pressure stripping (RPS, Gunn & Gott 1972, see also Jaffé et al. 2015) and strangulation (Peng, Maiolino & Cochrane 2015) being responsible for the ISM evolution. In this scenario, cluster galaxies are first detached from the cosmic web once they enter into the cluster halo, restricting the amount of cold gas inflows they get from the cosmic web (Dekel et al. 2009). Thus, cluster galaxies start depleting their gas reservoir in the outskirts of the cluster while they keep enriching their ISM through successive generations of stars. On top of this, the outside pressure of the intracluster medium (ICM) would push back the outflowing gas from feedback processes, forcing the galaxy to recycle the already processed gas before being fully quenched due to RPS and gas exhaustion in the cluster core (Wetzel et al. 2013). This description of galaxy evolution in massive clusters holds until $z \sim 1$.

At earlier epochs, an increasing fraction of overdense regions is still in the process of being fully assembled, with most of them commonly referred to as protoclusters beyond $z = 2$ (see Overzier 2016 for a review). Interestingly, the progenitors of the quiescent galaxies that populate the inner cores of local clusters formed the bulk of their stellar budget more than $\gtrsim 10$ Gyr ago, at the same time that protoclusters were being assembled. Thus, the so-called ‘cosmic-noon’ (i.e. $z \approx 1.5\text{--}3$, Madau & Dickinson 2014) is a crucial epoch to understand the assembly of the main components of massive clusters as well as to explore the onset of the first environmental effects and their influence over the early evolution of protocluster galaxies. Over the last years, several attempts have been made to examine the star formation activity and the gas-phase metallicity of protocluster galaxies compared to the field. The star-forming population is dominant in these structures and even starburst galaxies are relatively common within them (Dannerbauer et al. 2014; Popesso et al. 2015; Casey et al. 2017). However, larger samples display a big diversity of behaviours regarding star formation. Several studies have shown that star formation is enhanced relative to the field in dense environments at high redshift (Alberts et al. 2014; Shimakawa et al. 2018a; Monson et al. 2021; Lemaux et al. 2022), suggesting protocluster galaxies may have undergone accelerated mass assembly compared to their field counterparts. However, there are also some protoclusters where no such differences are seen as compared to the field (e.g. Cucciati et al. 2014; Toshikawa et al. 2014; Sattari et al. 2021; Shi et al. 2021). The discrepancies between works may arise from the diverse galaxy populations studied, the lack of statistics due to small sample size, and the different evolutionary stages in which these protoclusters are observed (Overzier et al. 2008; Toshikawa et al. 2014).

In addition, the early metal enrichment of protocluster galaxies is still a matter of debate. Some authors found evidence of enhanced gas-phase metallicity in low-mass cluster and protocluster galaxies at $z = 1.5\text{--}3$ (e.g. Kulas et al. 2013; Shimakawa et al. 2015; Maier et al. 2019), which can be explained by the shut down of pristine gas inflows in the cluster environment due to the early onset of the ICM. However, recent works have reported various levels of metallicity deficiency compared to field galaxies at the same redshift (Valentino

et al. 2015; Chartab et al. 2021; Sattari et al. 2021). To explain these results, it has been proposed that the dark matter haloes of young and not-yet-massive ($\log M_*/M_\odot < 13.5$) protoclusters at high z would still be on a phase of powerful cold stream accretion (Dekel & Birnboim 2006). This extra supply of gas on protocluster galaxies would dilute their current metallicities with respect to the general field. Conversely, some other studies have not observed significant environmental dependence of the mass–metallicity relation (MZR) during this epoch (e.g. Kacprzak et al. 2015; Tran et al. 2015; Namiki et al. 2019). While the total mass of the overdensity seems to play a key role in explaining these metallicity discrepancies between protoclusters in terms of their accretion mode (cold versus hot, Dekel et al. 2009), there are still other potential biases that should be taken into account, such as the presence and relevance of the active galactic nucleus (AGN) fraction within these forming structures (see Macuga et al. 2019 and Monson et al. 2021 for some examples) and the different selection criteria for the parent galaxy samples under scrutiny. Furthermore, the molecular gas properties of protocluster at $z > 2$, which are key to understanding the gas feeding and consumption processes that fuel star formation, remain largely unexplored except in a few cases and over small sample sizes (Dannerbauer et al. 2017; Wang et al. 2018; Tadaki et al. 2019; Zavala et al. 2019; Champagne et al. 2021; Aoyama et al. 2022). Only very recently there has been some attempts to map the CO emission of galaxies in protoclusters (e.g. Jin et al. 2021) providing a new window to investigate galaxy evolution in these forming large-scale structures. Thus, in order to shed light into the early stages of galaxy formation and evolution in protoclusters at the cosmic noon we must trace the star formation activity, metal enrichment, and gas reservoir of the individual objects that belong to these large-scale structures in formation.

Among the many assembling clusters detected during the last years, PKS1138–262 at $z = 2.16$ (hereafter the Spiderweb protocluster) stands as one of the most massive ($M_{200} > 2 \times 10^{14} M_\odot$, Shimakawa et al. 2014) and the best-studied systems beyond $z = 2$, both in terms of its galaxy populations and large-scale structure. This protocluster was first discovered by Kurk et al. (2000) using narrow- and broad-band photometry to identify an overdensity of Lyman- α emitters (LAEs) around the radio galaxy MRC1138–262 or Spiderweb galaxy (Roettgering et al. 1994; Pentericci et al. 1997). These initial reports about the presence of an overdensity of LAEs at $z = 2.16$ were later followed up by Pentericci et al. (2000) who spectroscopically confirmed the cluster membership of 15 LAEs. From that moment, the Spiderweb protocluster has been subject to exhaustive spectrophotometric campaigns to unveil the properties and distribution of its galaxy populations. This includes the characterization of the central radio galaxy (Pentericci et al. 2000; Carilli et al. 2002, 2022; Miley et al. 2006; Hatch et al. 2008, 2009; Emonts et al. 2016, 2018; De Breuck et al. 2022), the location of several X-ray emitters (Pentericci et al. 2002; Croft et al. 2005; Tozzi et al. 2022), the observation of an emerging red sequence within the cluster core (Kurk et al. 2004; Kodama et al. 2007; Zirm et al. 2008; Tanaka et al. 2010, 2013), the discovery of a network of starburst submillimetre galaxies (SMGs, Dannerbauer et al. 2014, 2017), the location of a rich population of H α emitters (HAEs) which represent the bulk of the known cluster members up to date (Hatch et al. 2011; Kuiper et al. 2011; Koyama et al. 2013; Shimakawa et al. 2014, 2015, 2018b), and the use of (sub-)millimetre observations to trace the dust content and gas reservoirs of several protocluster members (Emonts et al. 2018; Tadaki et al. 2019) as well as for the mapping of the protocluster large-scale structure in CO(1–0) using ATCA (Jin et al. 2021).

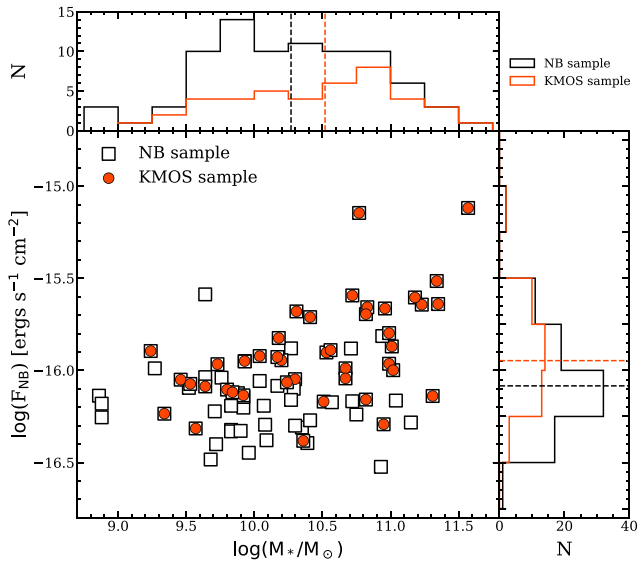


Figure 1. Stellar-mass and narrow-band flux properties of both our spectroscopic targets and their parent sample of narrow-band emitters from Koyama et al. (2013). Main diagram: The red circles represent the KMOS spectroscopic sample. The empty squares display the parent sample. Side panels: distribution (histograms) and median values (dashed lines) of the KMOS sample and its parent sample. The colours follow the same scheme applied in the main diagram.

In this work, we investigate the environmental imprints of galaxy evolution in the Spiderweb protocluster focusing on the star formation, gas-phase metallicity, stellar-disc size, and molecular gas properties of the protocluster members. This manuscript is structured as follows: Section 2 describes our new KMOS spectroscopic observations in the Spiderweb protocluster and the wealth of archival data available within this field. Section 3 outlines the methods used to analyse the physical properties of our targets and the environmental parameters measured to them. Sections 4 and 5 present our main results and the discussion of their physical interpretation in the context of galaxy evolution, respectively. Finally, Section 6 outlines the major conclusions of this study. Throughout this article, we assume a Chabrier (2003) initial mass function (IMF), and adopt a flat cosmology with $\Omega_{\Lambda} = 0.7$, $\Omega_{\text{m}} = 0.3$, and $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$. All magnitudes quoted in this paper are in the AB system (Oke & Gunn 1983).

2 OBSERVATIONS

In this section, we describe the main characteristics of the data sets we used in our analysis. In particular, we combine a new set of near-infrared (NIR) spectroscopic observation with previous multiwavelength photometry covering a significant fraction of the Spiderweb protocluster field.

2.1 KMOS spectroscopy

We carried out multi-object integral field spectroscopy observations of a sample of 42 narrow-band-selected HAEs in the Spiderweb protocluster to obtain $\text{H}\alpha$ and $[\text{N II}]\lambda 6584$ emission-line fluxes and study the galaxies’ star formation activity and gas-phase metallicities as a function of several environmental indicators. In Fig. 1, we put in context our spectroscopic targets (hereafter KMOS sample) with their parent sample of narrow-band (NB) emitters from Koyama

et al. (2013). The KMOS sample was selected to have narrow-band fluxes $F_{\text{NB}} \gtrsim 4 \times 10^{-17} \text{ erg s}^{-1} \text{ cm}^{-2}$ and a similar stellar mass range than its parent sample. The distribution and median values (dashed lines) of both the parent and the KMOS sample show that the latter is slightly biased towards massive galaxies, and that at a fixed stellar mass it includes objects with relatively high narrow-band fluxes. In addition, our targets are selected to trace both the inner core and some of the filaments around the large-scale structure of the protocluster simultaneously (see Fig. 2). As a result of the narrow-band selection criteria, our targets span a very narrow range in redshift space ($z = 2.159 \pm 0.008$) though they encompass a broad range of cluster-centric distances (up to $R = 5R_{200}$).

The observations were carried out using the K -band Multi-Object Spectrograph (KMOS, Sharples et al. 2013) installed at the Very Large Telescope (VLT) in Cerro Paranal, Chile (Program ID: 095.A-0500(A), PI: Y. Koyama). Our program employed two different arm configurations targeting 42 HAEs identified by Koyama et al. (2013). The observations were executed in service mode between 2015 April and June under average seeing conditions of 0.6 arcsec, airmass value of 1.3, and no moonlight contamination (i.e. dark time). The observations were carried out in the K band (19 300–24 600 Å), which captures $\text{H}\alpha$ and $[\text{N II}]\lambda 6584$ at $z = 2.16$ with a nominal spectral resolving power of $R \sim 4200$ around its central wavelength, which translates to $\sigma_{\text{ins}} \approx 35 \text{ km s}^{-1}$ in velocity space. Each KMOS arm carries a small IFU ($2.8 \text{ arcsec} \times 2.8 \text{ arcsec}$) with spatial resolution of 0.2 arcsec per spaxel. For each configuration, one arm was used to take sky exposures while two additional arms were fixed to monitor stars. In summary, 21 arms were used for science targets per pointing. The observing time of each configuration was divided into $\sim 1 \text{ h}$ observing blocks (OB) made of five on-source subexposures of 450 s each plus overheads. The total on-source time per configuration is 3 h. The spectroscopic data reduction was carried out using the ESO-Reflex workflow pipeline (Davies et al. 2013). The main reduction steps were bias subtraction, flat-field normalization, sky subtraction, wavelength and flux calibration, and frame stacking. The resulting data cubes conserve the spatial and spectral properties previously outlined in this paragraph.

2.2 Archival data

In addition to our spectroscopic campaign, we used abundant optical to NIR complementary archival imaging data in this field. This includes the Subaru/Suprime-Cam B and z' band, and the Subaru/MOIRCS J and K_s band from a previous MAHALO-Subaru project publication (Koyama et al. 2013). The NIR campaign carried out by VLT/HAWK-I to obtain Y, H, K_s deep imaging (PI: A. Kurk, program IDs 088.A-0754, 091.A-0106, 094.A-0104, see Dannerbauer et al. 2017 and Shimakawa et al. 2018b). Furthermore, we use the Post-BCD (PBCD) products from the Spitzer data archive library to obtain IRAC broad-band imaging at 3.6 and 4.5 μm (PI: D. Stern, campaign IDs 736 and 793, see Seymour et al. 2007). Finally, we use the reduced *Hubble Space Telescope* (HST) ACS/WFC data in filters $F475W$ and $F814W$ from the Hubble Legacy Archive (PI: H. Ford, proposal ID 10327, see Miley et al. 2006). The exposure times and seeing conditions of the retrieved coadded mosaic images are shown in Table 1. Finally, we also use the ATCA CO(1–0) map presented in Jin et al. (2021), which belongs to the COALAS project (large program ID: C3181, PI: H. Dannerbauer). The coordinates, redshifts, and general properties of our final protocluster galaxy sample are summarized in the Appendix.

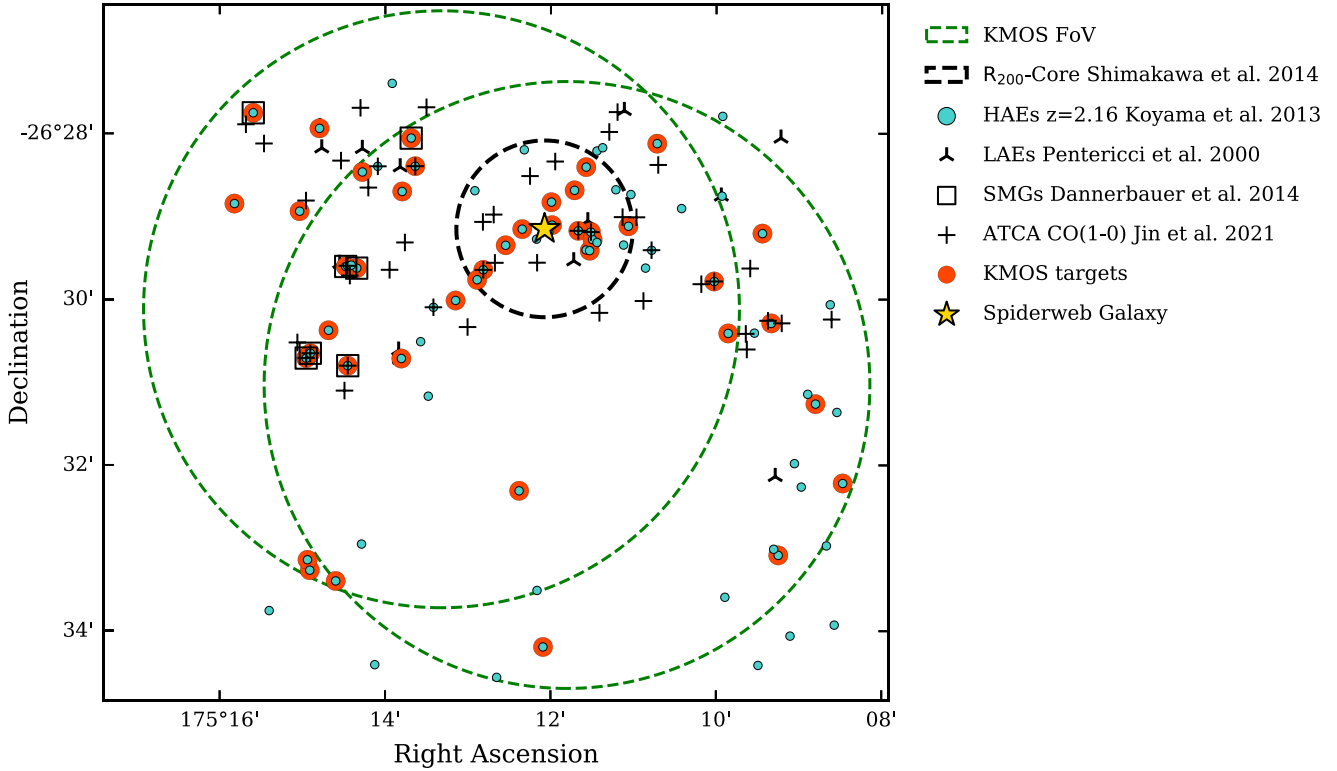


Figure 2. Distribution of spectroscopic and narrow-band-selected cluster members in the field of the Spiderweb protocluster. The red circles show the sample of KMOS-HAEs studied in this work. The blue circles depict the parent sample of narrow-band-detected HAEs from Koyama et al. (2013). The empty squares highlight a small overlapping sample of SMGs from Dannerbauer et al. (2014). The three-pointed stars display the position of LAEs from Pentericci et al. (2000). The black crosses show the positions of ATCA CO(1–0) emitters detected by Jin et al. (2021) in this field. The yellow star shows the position of the radio galaxy MRC 1138–262 (Roettgering et al. 1994). The black dashed circle encloses R_{200} according to the mass estimation performed by Shimakawa et al. (2014) for the protocluster core. Finally, the two dashed green circles show the field of view of the two KMOS configurations used to observe our targets.

Table 1. Properties of the optical to NIR photometric bands used in this work. Ellipsis are used to avoid repetitions in the instrument and reference columns.

Instrument	Filter	Exp. time (s)	FWHM (arcsec)	Pixel-size (arcsec)	Reference
<i>HST</i> /ACS	<i>F475W</i>	20 670	0.11	0.05	Miley et al. (2006)
...	<i>F814W</i>	23 004	0.11	0.05	...
Suprime-Cam/Subaru	B	6300	1.15	0.20	Shimakawa et al. (2018b)
...	z'	4500	0.70	0.20	Koyama et al. (2013)
MOIRCS/Subaru	J	9060	0.69	0.12	...
...	K_s	3300	0.63	0.12	...
...	NB2071	11 160	0.63	0.12	...
HAWKI/VLT	Y	26 880	0.37	0.11	Dannerbauer et al. (2017)
...	H	14 832	0.49	0.11	...
...	K_s	9228	0.38	0.11	...
Spitzer/IRAC	3.6 μm	3000	1.80	0.40	Seymour et al. (2007)
...	4.5 μm	3000	1.80	0.40	...

3 METHODS

3.1 Emission-line fitting

In order to confirm the cluster membership of our targets and to measure their star formation activities and metallicities, we need to carry out a line-fitting procedure. To carry out this task, we developed a self-written code in PYTHON with the Astropy library (Astropy Collaboration 2013, 2018). In the next paragraphs we describe the

most important steps that we followed during the emission-line detection and fitting.

First, we visually inspect each data cube with the software QFITSVIEW, which is a publicly available IFS visual tool developed by Thomas Ott at the MPE. We search for emission lines in the wavelength range defined by the width of the narrow-band filter used to classify our targets as HAEs (see Koyama et al. 2013). At least one emission line is clearly detected in 40 out of the 42 original KMOS

targets. In these cases, we integrate the spectral axis of each data cube over the wavelength limits of the detected emission line and define the brightest spaxel of the light distribution as the spatial centre of each target within the IFU. Next, we place 1.4×1.4 (i.e. 7×7 spaxels) squared apertures (red squares in Fig. 3) there and extract a 1D spectrum for each galaxy. The same procedure and aperture position is applied to the noise data cube.

After the extraction of the 1D spectra, we fit and subtract a local continuum around the $H\alpha$ line consisting of several windows free from skyline contamination that covers a few hundred angstroms in both directions around the $H\alpha$ observed wavelength. Finally, we perform a triple Gaussian fit for $H\alpha$ and the $[N\text{II}]\lambda 6584$ and $\lambda 6548$ lines allowing for small variations in the centre, amplitude, and width of the first two. However, the flux ratio $[N\text{II}]\lambda 6584/[N\text{II}]\lambda 6548$ is fixed to 1/3. Furthermore, their widths are tied and fixed to be equal or smaller than the one from $H\alpha$. The reason for this is that Balmer lines may exhibit broad components when the ionization mechanism is (at least partly) related to AGN activity or skewed profiles when inflows and outflows are present. These effects may broaden the width of the single-component fit applied to the $H\alpha$ emission line, and thus, this value is taken as the upper bound for the $[N\text{II}]$ line. Two objects (IDs 911 and 647) display extremely broad $H\alpha$ profiles ($\sigma > 700 \text{ km s}^{-1}$), indicative of AGN type 1 activity. These sources are also identified as X-ray emitters in Croft et al. (2005) and Tozzi et al. (2022), confirming their AGN nature. In these cases, we expand our method by including one additional component to trace the $H\alpha$ broad emission. One object (ID 897) display three emission lines within the inspected spectral window. These lines are consistent with $[O\text{III}]\lambda 5007$, $[O\text{III}]\lambda 4959$, and $H\beta$ emission from a background galaxy at $z = 3.166$ and thus, this object is removed from the subsequent analyses of the protocluster sample. In total, we report 39 $H\alpha$ spectroscopic detections with 12 of them being new sources unknown to previous spectroscopic studies in the Spiderweb protocluster (e.g. Pentericci et al. 2000; Kurk et al. 2004; Croft et al. 2005; Doherty et al. 2010; Kuiper et al. 2011; Tanaka et al. 2013; Shimakawa et al. 2014, 2018b; Tadaki et al. 2019; Jin et al. 2021).

In order to take into account the uncertainties of both the fitting procedure and the data cube noise, we conduct a Monte Carlo approach where the fitting method outlined above is repeated 1000 times allowing for random Gaussian variations on every single spectral data point with a maximum amplitude equal to the value found in the 1D noise spectra. After this, we take the mean value of this parameters and compute the total flux for the emission lines under scrutiny. The flux error is estimated by the standard deviation of the Gaussian fit parameters over the 1000 realizations. We consider a line to be detected if the ratio between the total flux and the flux error measured in the way described above is higher than two. This constrain apply both to $H\alpha$ and $[N\text{II}]$. In Fig. 3, we show the 2D image of the collapsed data cube around $H\alpha$ and the extracted 1D spectrum for each target.

3.2 Spectral energy distribution fitting

The field of the Spiderweb protocluster counts with extensive and deep photometry (see Table 1) comprising the rest-frame UV to NIR wavelength range for galaxies at $z = 2.16$. We seek to obtain reliable stellar masses and rest-frame magnitudes for our targets based on the observed fluxes using the spectral energy distribution (SED) fitting technique. Due to the variety of seeing conditions and instrumental pixel sizes, we follow two different approaches to extract the observed magnitudes in different bands: First, we take as a reference the MOIRCS/Subaru NB2071 image and degrade all

the other bands with $\text{PSF} < 1$ arcsec to its PSF value and pixel size. Then, we run SEXTRACTOR (Bertin & Arnouts 1996) in dual-image mode using the NB2071 image for source detection while measuring the observed magnitudes (MAG_AUTO) over the PSF and pixel size matched images. This way we ensure that we are measuring the fluxes over the same area of the targets for each band in a consistent way. Second, we carry out simple photometry (MAG_APER) over the images with the worst spatial resolution (i.e. Suprime-Cam B-band and IRAC 3.6 and $4.5 \mu\text{m}$).

After constructing our multiband photometric catalogue, we perform the SED fitting with the CODE INVESTIGATING GALAXY EMISSION (CIGALE, Boquien et al. 2019). CIGALE combines the modelling of composite stellar populations with nebular emission and dust attenuation while conserving the energy balance between the energy emitted by massive stars and its partial absorption and re-emission by dust grains. CIGALE follows a Bayesian fitting approach which deviates from simple χ^2 minimization algorithms that provide physical properties based on the best match of templates. Instead, the estimated properties are evaluated by weighting all the models depending on their statistical agreement with respect to the best fit, which has the heaviest weight. This naturally takes into account the uncertainties on the observations while also including the effect of intrinsic degeneracies between physical parameters. Finally, the physical properties and their uncertainties are estimated as the likelihood-weighted means and standard deviations.

Before running CIGALE over the photometry of our targets, we create a grid based on the stellar population synthesis models of Bruzual & Charlot (Bruzual & Charlot 2003) assuming an exponentially delayed star formation history with possible e-folding times between 1 and 8 Gyr. In addition, we constrain the possible ages of stellar populations to be younger than the age of the Universe at $z = 2.16$ (i.e. ~ 3 Gyr). Given the interacting nature of protoclusters, we also consider the possibility of a recent minor star-forming burst accounting for up to 1, 5, or 10 per cent of the mass fraction and with an age not older than 300 Myr. We assume a Chabrier IMF (Chabrier 2003) and subsolar metallicity (i.e. $Z = 0.004$). Nebular emission is also included in our grid of models with the ionization parameter (U) in the range $-2.4 < \log(U) < -2.8$, which describes typical values for star-forming galaxies at $z \sim 2$ (Cullen et al. 2016). Finally, we apply Calzetti's attenuation law (Calzetti et al. 2000) with extinction values ranging $E(B - V)_s = 0-1$ mag in steps of 0.1 mag. We estimate the rest-frame magnitudes to an average uncertainty of 0.1 mag for all bands and 0.13 dex for the stellar masses. We show the CIGALE-fitted SEDs for each protocluster member within the KMOS sample in the Appendix (see Fig. A1). The reduced χ^2 distribution of our targets encompass the range $0.2 \leq \chi^2 \leq 4.9$ and yields an average value of $\chi^2 = 1.4$, with two-thirds of our sample (i.e. 26 objects) displaying $0.5 \leq \chi^2 \leq 2$. To put our sample in context we plotted our galaxies into the rest-frame UVJ diagram (Fig. 4). This diagram splits the galaxies into two different groups, an old-age sequence of quiescent galaxies (upper left corner) and a star-forming sequence of galaxies. Most of our targets lie within the star-forming region, which is consistent with their previous HAE classification (Koyama et al. 2013).

3.3 Size determination

While space-based *HST* observations are ideal to measure the sizes of galaxies due to its high spatial resolution and lack of atmospheric effects, the *HST* mosaics ($3 \text{ arcmin} \times 6 \text{ arcmin}$) using the *F475W* and *F814W* filters cover only 31 out of the 39 KMOS spectroscopically confirmed protocluster members. Furthermore, given the redshift of

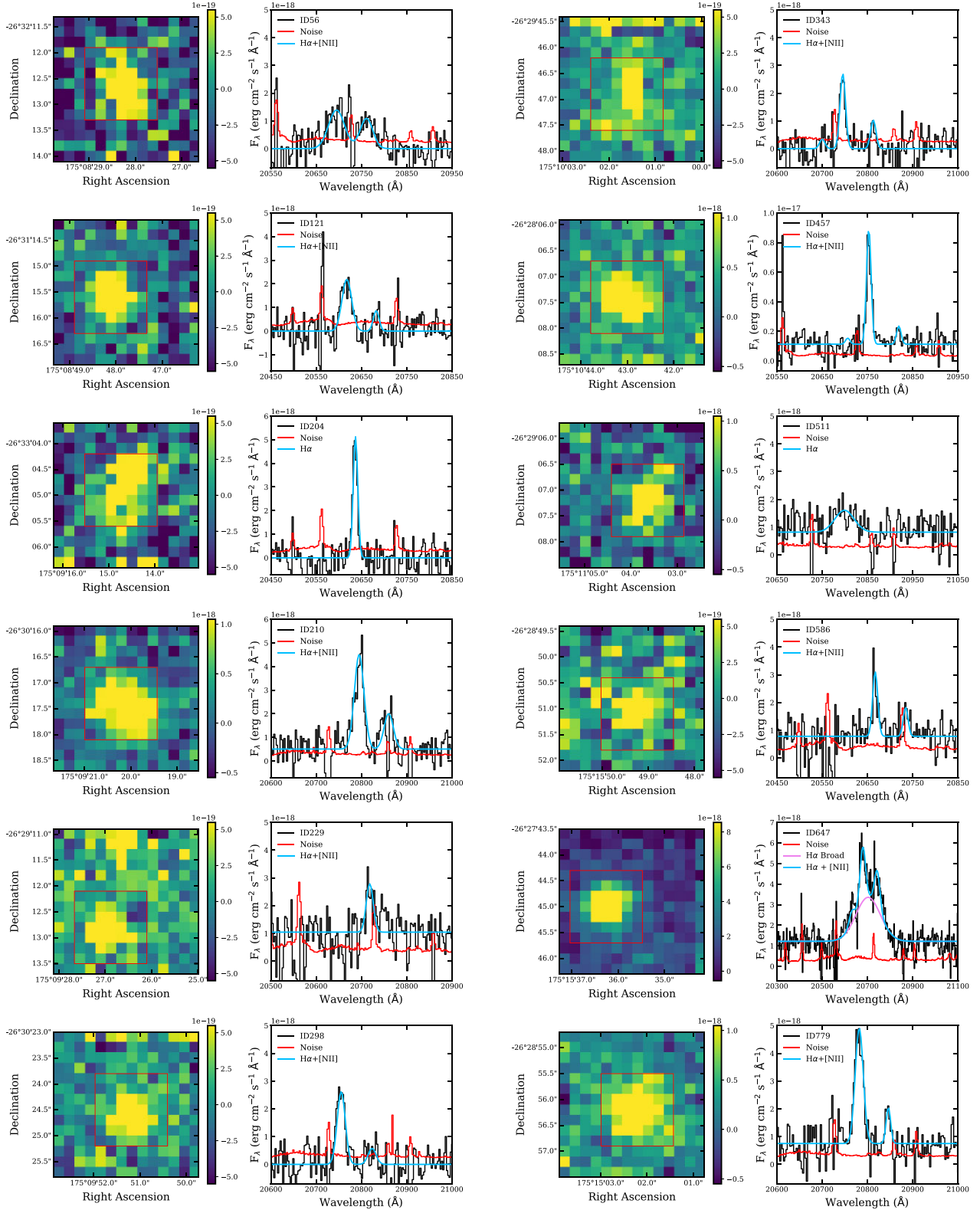


Figure 3. Left: Image of the $H\alpha$ -collapsed data cube. The red square shows the 1.4 arcsec squared aperture used to extract the 1D spectrum of each object. Right: Aperture-extracted signal spectrum (black), noise spectrum (red), and fitted spectrum (blue) to the $H\alpha$ and $[NII]$ emission lines. The used IDs follow the numbering given by Koyama et al. (2013).

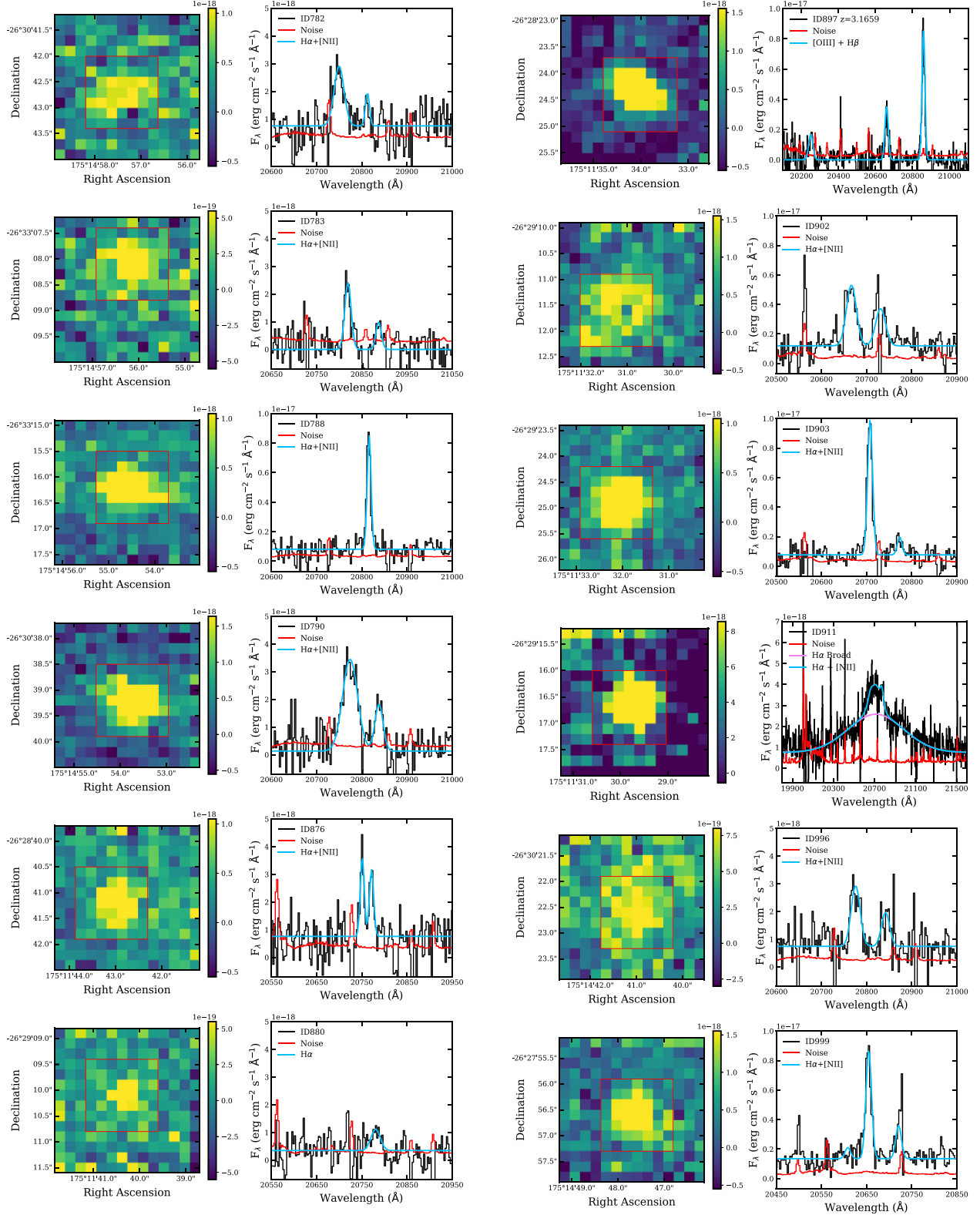


Figure 3 – continued

our targets, the *HST* imaging would only trace the rest-frame UV and FUV wavelength range, i.e. the very bursty and young star-forming regions, while being blind to the underlying older stellar populations which account for the bulk of the galaxies' stellar mass.

Thus, we carry out the size measurement of our targets using the VLT/HAWK-I K_s -band mosaic ($8.5 \text{ arcmin} \times 12 \text{ arcmin}$) whose central wavelength traces the rest-frame emission around 6800 \AA at $z = 2.16$. Unfortunately, the nearby $H\alpha$ emission line may

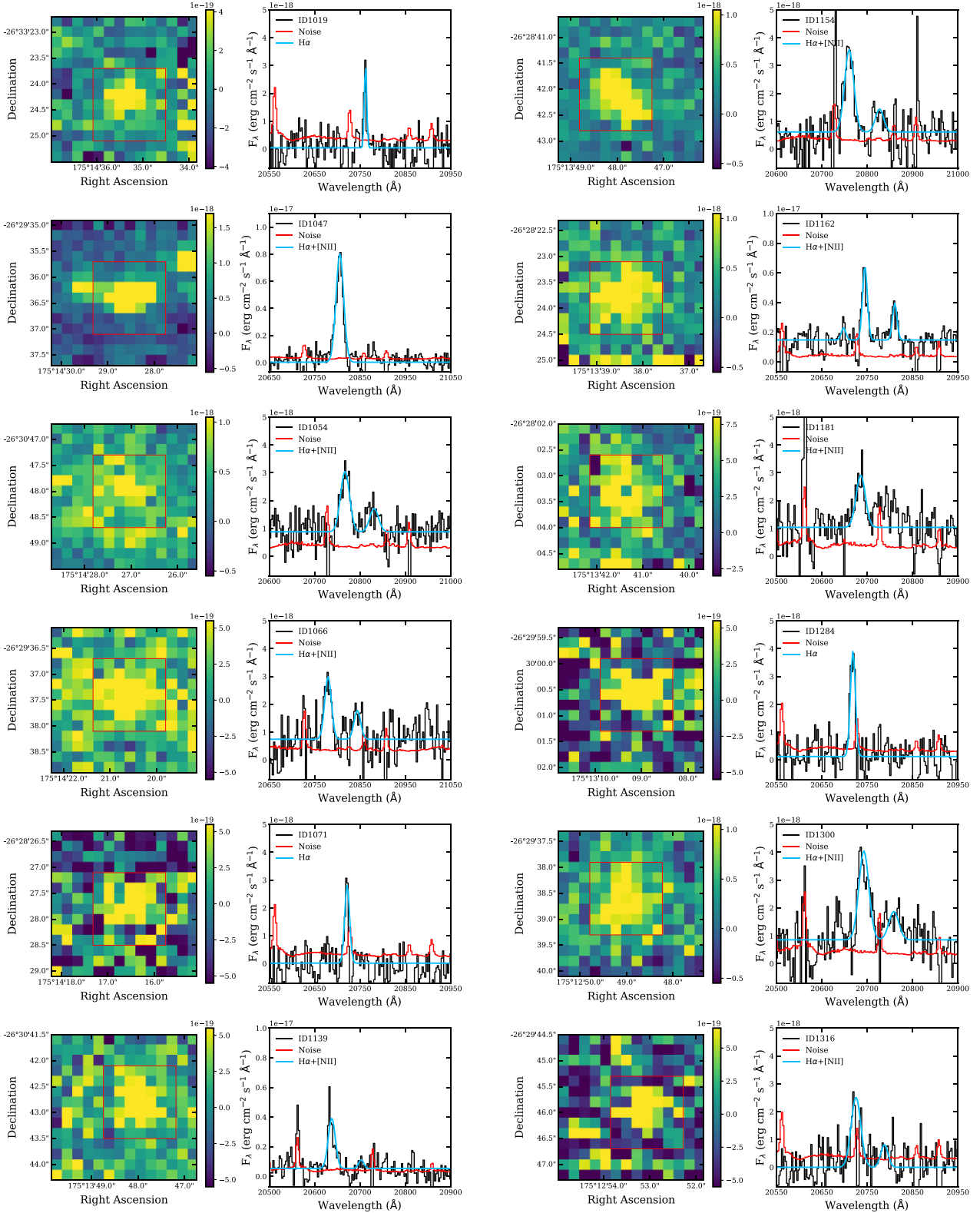


Figure 3 – continued

contaminate the continuum surface brightness profile originated from the stellar component for very active star-forming galaxies and type 1 AGNs. To solve this issue, we estimate the H α contribution to the total broad-band flux within our sample, finding that it is constrained

to $\lesssim 15$ per cent except for two cases. Therefore, we conclude that our K_s -band size measurements are not significantly affected by the H α component. Then, we measure and extract an average PSF size of ~ 0.4 arcsec using PSEX (Bertin 2011). This size is equivalent to a

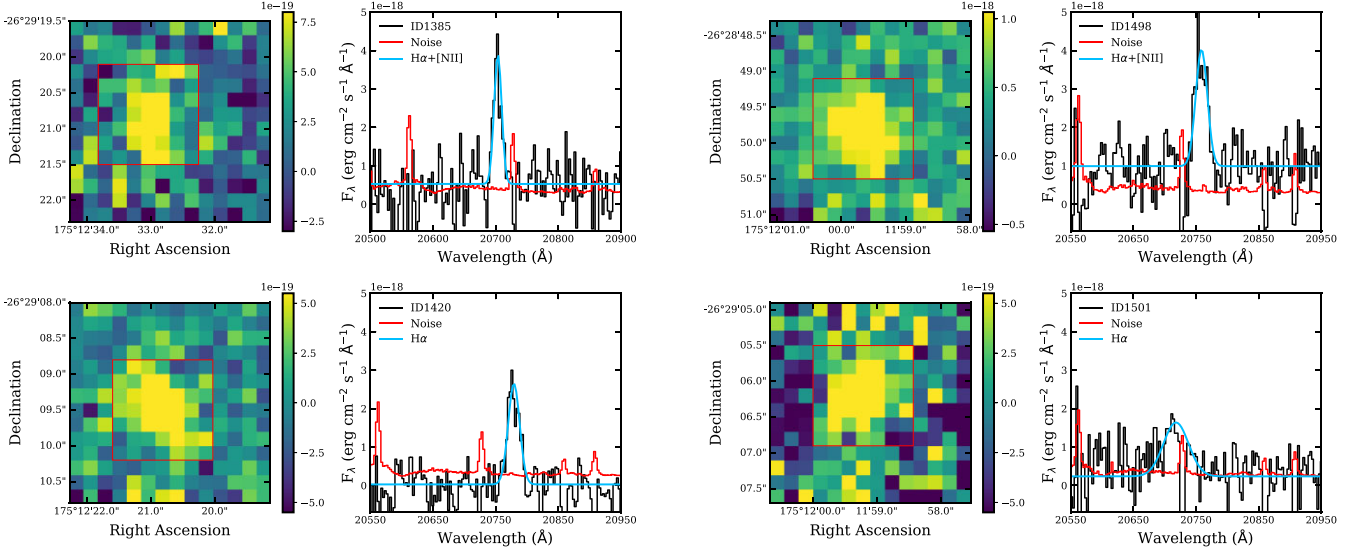


Figure 3 – continued

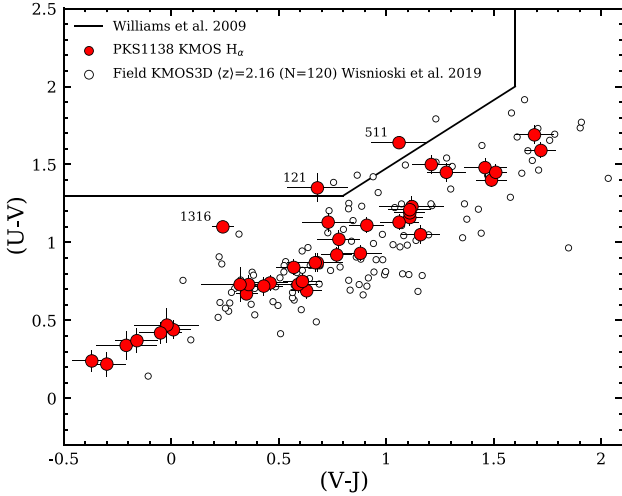


Figure 4. Rest-frame UVJ diagram. The limits separating the quiescent and star-forming regions follow the prescription given by Whitaker et al. (2013). The red circles show the colour distribution of our spectroscopically confirmed HAEs. By comparison, the white circles display the colour distribution of the KMOS3D sample of field galaxies at similar redshift (Wisnioski et al. 2019). Three objects display UVJ colours departing from the main sample and are labelled according to their HAE IDs (Koyama et al. 2013) for additional discussion in the main text.

physical diameter of 3.32 kpc assuming the cosmological parameters outlined in Section 1.

We model the surface brightness profile of our targets and measure their structural parameters (e.g. Sèrsic index, R_e , axial ratio and position angle) using GALFIT (Peng et al. 2002) over squared windows of 10 arcsec \times 10 arcsec with our targets at its centre. The models are computed following a single-component approach and allowing for free variation of all structural parameters. We fit additional components to nearby objects within the inspected window in order to remove possible light contamination on the main target. Finally, we visually check the result after subtracting the model to the original K_s -band image and determine if the inspected object shows signs of

strong residuals (e.g. bulge presence, asymmetric disc, or tidal tails). However, we find a good agreement between the modelled single-component surface brightness profiles and the K_s -band images in most cases. This argues against a secondary bulge component or visible interactions affecting the surface light profile of our targets. Finally, we discard those objects whose R_e uncertainties are greater than 50 per cent and those displaying $R_e < 1.66$ kpc (e.g. half of the PSF size) due to the difficulties to reliably resolve them. As a result, we obtained size measurements for 27 galaxies out of the 39 spectroscopically confirmed cluster members within our sample. The remaining objects did not reach enough S/N for the GALFIT models to converge into a solution or displayed too large uncertainties. The observed and modelled K_s -band surface brightness profiles for our targets, as well as their residuals after subtraction can be found in the Appendix (see Fig. A2).

3.4 Star formation activity

Our KMOS spectroscopy campaign provides us with access to the $H\alpha$ emission line for the protocluster members at $z = 2.16$. We apply the SFR calibration developed by Kennicutt (1998) modified for a Chabrier IMF. This calibration has proven to be one of the most reliable ones, both at local and high redshift (e.g. Moustakas, Kennicutt & Tremonti 2006; Wisnioski et al. 2019):

$$\text{SFR}(H\alpha) = 4.65 \times 10^{-42} L(H\alpha), \quad (1)$$

where $L(H\alpha)$ is the luminosity of the $H\alpha$ emission line. We estimate the value of this quantity by measuring the $H\alpha$ spectroscopic fluxes of our targets and assuming a Calzetti’s extinction law ($R_v = 4.05$, Calzetti et al. 2000) to account for the dust attenuation. We use the extinction A_v values obtained from the SED fitting to account for the diffuse dust attenuation in the galaxy’s continuum following $A_{\text{cont}} = 0.82A_{v,\text{SED}}$ (Wuyts et al. 2013; Wisnioski et al. 2019). However, the nebular contribution produced in the active star-forming regions of the galaxy remains unaccounted for at this stage. To tackle this problem, Wuyts et al. (2013) add an extra extinction term (A_{extra}) that can be parametrized as $A_{\text{extra}} = 0.9A_{\text{cont}} - 0.15A_{\text{cont}}^2$ and is in good agreement with the previous extinction estimates made by Calzetti et al. (2000) in the local Universe. Thus, the total extinction applied

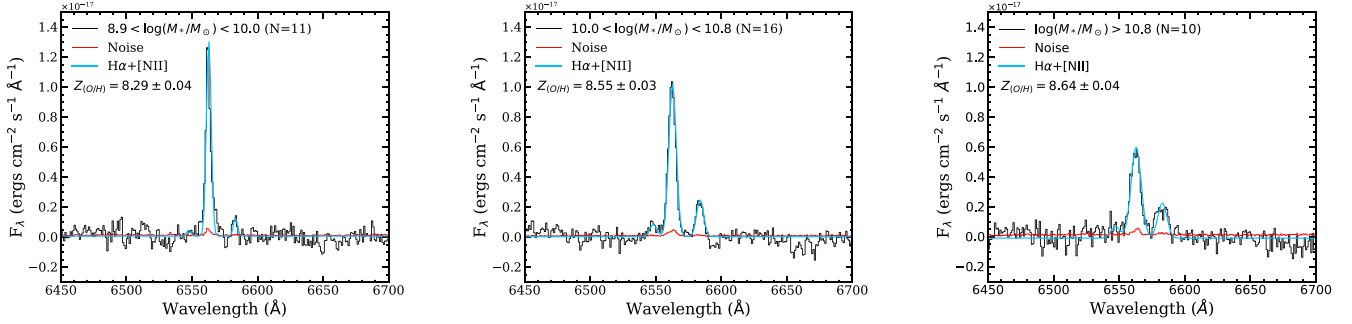


Figure 5. Stacked spectra and fit around $H\alpha$ and $[N\text{ II}]\lambda 6584$ for the three inspected stellar-mass bins.

to the measured $H\alpha$ fluxes is $A(H\alpha) = A_{\text{cont}} + A_{\text{extra}}$. This method has been recently tested by the KMOS3D team in Wisnioski et al. (2019), which we will use as our main field comparison sample in the following sections. As it was explained in Section 3.1, two of our objects (IDs 647 and 911) display extreme $H\alpha$ line widths suggesting the presence of type 1 AGNs. In these two cases, the SFR values were computed using only the $H\alpha$ narrow component in equation (1). The final SFR and reddening values of our protocluster sample can be found at the end of this work in Table A1.

3.5 Gas-phase metallicities

Over the last decades, several optical emission-line diagnostics have been developed to estimate the gas-phase metallicity of galaxies. Typically, these diagnostics rely on the flux ratio between ionized oxygen or nitrogen, and the hydrogen emission lines (see Kewley & Ellison 2008 for a review). At $z > 1$ the access to these line diagnostics shifts to the NIR and simultaneous measurements of multiple emission lines become difficult due to the stretching of the spectra. In this work, we apply the N2 calibration developed by Pettini & Pagel (2004) which involves the $[N\text{ II}]\lambda 6584/H\alpha$ ratio:

$$12 + \log(\text{O}/\text{H}) = 8.90 + 0.57 \times \text{N2}, \quad (2)$$

where N2 is equivalent to $\log(F[N\text{ II}]/F(H\alpha))$. This method relies on a local calibration, although it has been successfully applied for field galaxies up to $z \sim 3$ (Erb et al. 2006; Wuyts et al. 2016; Sanders et al. 2021). However, the accuracy of its absolute values at high redshift is still under debate (Steidel et al. 2014). Nevertheless, this work aims to investigate relative metallicity differences between the protocluster and general field populations and thus, it is not affected by calibration uncertainties as long as the chosen comparison samples are analysed following the same approach and belong to the same cosmic epoch. The detection of $[N\text{ II}]\lambda 6584$ is the limiting factor when using this method to derive gas-phase metallicities. We achieve this in 24 out of 39 confirmed HAEs. In order to investigate the average properties of our full sample, we decide to stack the individual spectra in three stellar-mass bins. We exclude two objects from our analysis due to their extreme $H\alpha$ line widths ($\sigma > 700 \text{ km s}^{-1}$), indicative of AGN activity. The stacking analysis is performed using the approach outlined in Shimakawa et al. (2015):

$$F_{\text{stack}} = \frac{\sum_i^n \frac{F_i(\lambda)}{\sigma_i(\lambda)^2}}{\sum_i^n \frac{1}{\sigma_i(\lambda)^2}}, \quad (3)$$

where $F_i(\lambda)$ is the flux density of an individual spectrum and $\sigma_i(\lambda)$ is the noise as a function of wavelength. The sums within a bin apply a median sigma clipping algorithm to reduce the contamination

by sky residuals. Then, the N2 ratio is measured for the stacked spectra by simultaneously fitting two Gaussian curves for $H\alpha$ and $[N\text{ II}]\lambda 6584$. The two objects with extremely broad $H\alpha$ components are excluded of this analysis. Finally, we estimate the average measured uncertainties to derive gas-phase abundances to be of the order of ~ 0.1 dex for the individual measurements across our stellar mass range and ≤ 0.04 dex for the stacking analysis bins. The results of our stacking process can be seen in Fig. 5.

3.6 Molecular gas masses

In this section, we will use previous works reporting radio measurements to estimate the molecular gas masses of a subsample of our HAEs. In particular, we will focus on the results recently published by Jin et al. (2021) investigating the $\text{CO}(1-0)$ luminosity function of protocluster galaxies with ATCA. In total, we share eight overlapping targets with six of them being detected in $\text{CO}(1-0)$ at more than 5σ level. The $\text{CO}(1-0)$ luminosity ($L'_{\text{CO}(1-0)}$) can be converted into a molecular gas mass estimate in a simple way assuming a conversion factor ($\alpha_{\text{CO}(1-0)}$) that traces the amount of molecular gas from optically thick virialized clouds (Dickman, Snell & Schloerb 1986, Solomon et al. 1987):

$$M_{\text{mol}} = \alpha_{\text{CO}(1-0)} \times L'_{\text{CO}(1-0)} \quad (4)$$

However, $\alpha_{\text{CO}(1-0)}$ is found to be approximately constant only for nearby galaxies with solar gas-phase metallicities ($\alpha_{\text{CO}(1-0)} = 4.36 \pm 0.90$, Bolatto, Wolfire & Leroy 2013). This conversion factor increases for low-metallicity regions due to the contraction of the CO -emitting surface relative to the area where the gas is H_2 for a fixed cloud size. Therefore, we compute our molecular gas masses assuming the metallicity-dependent conversion factor $\alpha(Z)$ outlined in Genzel et al. (2015) and later revisited by Tacconi et al. (2018):

$$\alpha(Z) = 4.36 \sqrt{0.67 \exp(0.36 \times 10^{(8.67-Z)}) \times 10^{1.27 \times (8.67-Z)}}, \quad (5)$$

where $Z = 12 + \log(\text{O}/\text{H})$ assuming the Pettini & Pagel (2004) metallicity calibration. This conversion factor is based on the geometric sum of the metallicity corrections proposed by Genzel et al. (2012) and Bolatto et al. (2013). However, these corrections strongly diverge in the low-metallicity regime. Thus, we constrain the range of application of this prescription to galaxies with $12 + \log(\text{O}/\text{H}) > 8.44$, which correspond to a maximum disagreement of 0.2 dex between methods. Only three objects within our sample display lower metallicity values but none of them are affected by this constraint as they do not count with previous $L'_{\text{CO}(1-0)}$ measurements. Given the distribution of our galaxies in the MZR (see Section 4.3), this approach allows us to constrain the molecular gas mass of most HAEs with secure metallicity measurements. Finally, we will use those

HAEs with available CO(1–0) information within our sample to investigate the relation between the molecular gas fraction, i.e. $f_{\text{gas}} = M_{\text{mol}}/(M_* + M_{\text{mol}})$, and the gas metallicity in different environments at the cosmic noon (Section 4.5).

3.7 Environment quantification

In general, cluster membership is defined as an interval in redshift space around the value given for the whole cluster structure, which usually coincides with its brightest cluster galaxy. Even though this may be sufficient to qualitatively disentangle the general field population of galaxies from objects residing in denser environments, we need a quantitative way to measure the environment in order to study its influence of the physical properties of galaxies within these dense regions.

Most studies quantify the environment of their samples in two different ways: locally as a number density of cluster members (e.g. Dressler 1980), and globally by taking into account the general properties of the cluster (M_{200} , R_{200} , and σ_{cl}) which define its phase-space (Carlberg, Yee & Ellingson 1997). In both cases, it is required a high number of confirmed cluster members to reliably map the cluster structure and local density peaks. Koyama et al. (2013) used a narrow-band MOIRCS/Subaru imaging mosaic (7 arcmin \times 8 arcmin) to identify HAEs in this field at the redshift of the Spiderweb protocluster. These HAEs form the bulk of the known galaxy populations known in this protocluster up to date (~ 90 objects) and they comprise the parent sample of our KMOS spectroscopic campaign, which has confirmed the membership of the narrow-band-selected HAEs with a ~ 93 per cent success rate (39 out of 42 targets). This is comparable to previous H α narrow-band surveys in the field at $z \sim 2$, which also find a contamination rate lower than 10 per cent (e.g. HiZELS, Sobral et al. 2013). Furthermore, this contamination rate is expected to be even lower in protoclusters as they host a high number of emitters in a relatively narrow redshift window, while for a given field of view the density of field contaminants (i.e. background or foreground emitters) remains constant. Based on this result, we assume that the majority of the remaining narrow-band-detected HAEs can be considered cluster members with very high probability. In addition, we also include spectroscopically confirmed cluster members from other studies tracing different galaxy populations: LAEs (Pentericci et al. 2000 and Croft et al. 2005), additional HAEs cluster members (Kurk et al. 2004), and CO emitters (Jin et al. 2021). Combining both the spectroscopic and the narrow-band-selected sources, we gather a cluster sample comprised of 125 independent objects. These sources are projected into a plane where distances between them are measured by their sky angular separation and transformed to the physical scale using a fixed cosmology with the scale factor (a) corresponding with the redshift of the cluster (e.g. $a = 8.288$ kpc arcsec $^{-1}$ at $z = 2.16$). Thus, it is possible to define a number surface density of objects within a given radius in the following way:

$$\Sigma_N = \frac{N}{\pi R_{N-1}^2} \quad (6)$$

In particular, we will use local densities defined by the minimum radius required to enclose three neighbouring galaxies (i.e. Σ_3). The reason behind this choice is that high local density peaks produced by small but rather compact galaxy groups may trace the places where gravitational interactions between galaxies, such as mergers and close encounters, are more frequent. Furthermore, during the early stages of cluster assembly it is common that galaxies infall towards the protocluster not only through the surrounding filamentary structure of the cosmic web but also as part of small groups of galaxies

that are accreted as a whole (Shimakawa et al. 2018a). Therefore, this approach provides us with an opportunity to study the properties of galaxies residing in local density peaks during the protocluster assembly.

However, clusters of galaxies are by definition large-scale structures and thus, local density peaks may not always correlate with the overall cluster-mass distribution. For example, some of the most characteristic clusters properties such as the presence of the ICM or the splashback radius are related to the cluster-core itself. These components are linked to hydrodynamical cluster-specific interactions such as starvation and RPS, and their effects are gradually felt by the galaxies during their infalling phase. Thus, a second environmental parameter that takes into account the radial and velocity distribution of our targets with respect to the core of the cluster (i.e. the Spiderweb galaxy) is required. Our approach to trace the global environment relies on the projected cluster-centric distance (R_{proj}) of each object and its relative line-of-sight velocity with respect to the systemic velocity of the cluster (Δv), which can be measured through each object's redshift. We follow the procedure outlined by Noble et al. (2013) who used a parameter (η) that defines caustic profiles in a phase-space diagram in the following way:

$$\eta = (R_{\text{proj}}/R_{200}) \times (|\Delta v|/\sigma), \quad (7)$$

where $|\Delta v| = |(z - z_{cl})c/(1 + z_{cl})|$ and z_{cl} is the redshift of the cluster. Attending to this parameter, Noble et al. (2013) defined three separate regions: $\eta < 0.4$ for galaxies that were accreted into the cluster core long time ago, $0.4 < \eta < 2$ for galaxies that have been recently accreted, and $\eta > 2$ for galaxies infalling into the cluster but not yet associated with its main component. Given the heterogeneous structure of a cluster in formation such as this, we decided to use η as a continuous parameter that models the environmental relation of a given galaxy to a large overdensity, without imposing an upper limit on η . Thus, we define the parameter space given by $\eta > 2$ as the outskirts of the protocluster and $\eta < 2$ as the accreted region. Our combined approach using local and global environmental parameters will allow us to investigate different environmental effects in the following sections.

4 RESULTS

In this section, we investigate the main physical properties of our spectroscopic sample and their connection with the environment. In particular, we will first compare the distribution of our targets with respect to the Spiderweb galaxy at the centre of the protocluster. Then, we will compare the physical properties (e.g. SFR, metallicity, size, and molecular gas fractions) of protocluster members with coeval field samples and use local and global environmental indicators to explore the possible effects across the protocluster structure. The measured values of the physical quantities quoted in this section for every galaxy can be found at the end of this work in Table A1.

4.1 Phase-space distribution

We present the distribution of our 39 spectroscopically confirmed protocluster members with respect to the centre of the Spiderweb protocluster in terms of cluster-centric distance (i.e. as a function of R_{200}) and velocity space (i.e. redshift) in Fig. 6. The centre of the protocluster is defined by the position of the radio galaxy MRC 1138–262 at $z = 2.156$ (also known as Spiderweb galaxy, Miley et al. 2006). Our spectroscopic sample was selected through narrow-band techniques over a large field of view (see Section 2) and thus it spreads to cluster-centric distances of up to $5R_{200}$ while

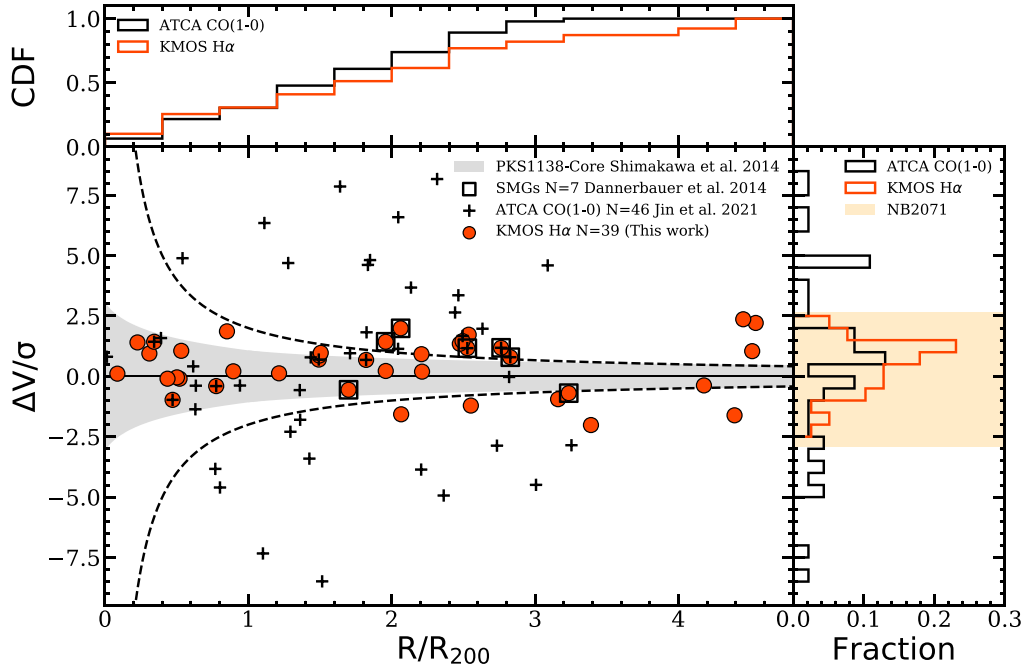


Figure 6. Phase-space diagram for PKS1138 at $z = 2.16$. Individual symbols represent the same samples as those in Fig. 2. The grey area represents the gravitationally bound region of the diagram assuming the R_{200} and M_{200} values estimated by Shimakawa et al. (2014) for the protocluster core. The dashed line shows the $\eta = 2$ contours following equation (7). The right-hand panel displays the fraction of KMOS H α emitters and CO(1–0) emitters (Jin et al. 2021) as a function of the velocity (i.e. redshift) space. The orange area in the right-hand panel mimics the width of the narrow-band filter used to identify the HAEs of the parent sample (Koyama et al. 2013). The top panel displays the cumulative distribution function (CDF) of KMOS H α emitters and CO(1–0) emitters as a function of cluster-centric radius.

encompassing a relatively narrow range of systemic velocities with respect to the Spiderweb galaxy. This is explicitly shown in the right-hand panel of Fig. 6, where we compare the narrow-band filter width (orange area) used for the target selection with the distribution of our spectroscopically confirmed protocluster members (red histogram). In contrast, the blind CO(1–0) survey carried out by Jin et al. (2021) displays a more extended distribution in redshift space, suggesting that the Spiderweb protocluster may be a multicomponent system enclosing a co-moving volume similar to the Hyperion superprotocluster at $z = 2.45$ (Cucciati et al. 2018).

Finally, we investigate the phase-space distribution of our sample with respect to two common environmental descriptors used in lower redshift clusters. First, we display the gravitationally bound region of the diagram (grey area) assuming the virialization of the protocluster core with $R_{200} = 0.53$ Mpc and $M_{200} = 1.71 \times 10^{14} M_{\odot}$ (Shimakawa et al. 2014), and the escape velocity prescription given by Jaffé et al. (2015) and Rhee et al. (2017). In addition, we show the $\eta = 2$ contours (dashed lines) marking the separation between the outskirts of the protocluster and the accreted region ($\eta < 2$) according to our definition in Section 3.7. Both approaches encompass a very similar fraction of our targets with 22 objects lying at $\eta < 2$ and 17 at $\eta > 2$. Interestingly, none of the SMGs we have in common with Dannerbauer et al. (2014) lie within R_{200} and only one source can be found at $\eta < 2$, implying that these objects are more frequent in the filamentary structure surrounding the Spiderweb protocluster and suggesting they may change their nature before reaching the inner core.

4.2 Star formation across different environments

In this section, we investigate the influence of the environment on the star-forming properties of our sample of galaxies. In Fig. 7, we

present the distribution of our galaxies over the SFR- M_{*} diagram. Field star-forming galaxies follow a tight correlation between these two parameters which is usually referred to as the ‘Main Sequence’ (Speagle et al. 2014). This correlation is represented by a solid line in Fig. 7, with its 3σ limits depicted by the grey area around it. Overall, our sample is consistent with the expectations of the main sequence, with most of our galaxies scattering within the 3σ region. Seven of our most massive targets ($\log M_{*}/M_{\odot} > 10.9$) display X-ray emission (Tozzi et al. 2022), hinting at the presence of AGNs.

Only six objects lie below the scatter of the main sequence even when taking into account their error bars: three of them (IDs 121, 511, 1316) show UVJ colours that detach them from the blue sequence, with the first two lying within the quiescent region and the latter displaying colours compatible with a post-starburst phase in Fig. 4. In addition, two of these objects (IDs 511 and 1316) lie within the inner core of the protocluster ($r < R_{200}$) and may be representative of its nascent quenched population. Another object (ID 647) is likely a type 1 AGN given its extreme H α broad component (see Fig. 3) and its X-ray-emitting nature, which is also the case of IDs 511 and 880. The remaining object (ID 1019) is classified as blue star-forming galaxies in the UVJ diagram and show no signs of AGN activity in their emission line profiles. On the other hand, two galaxies lie clearly above the main sequence scatter, though one of them has very low stellar mass ($\log M_{*}/M_{\odot} < 9.0$). Seven of our objects have previously been identified as SMGs by LABOCA 870 μm observations (Dannerbauer et al. 2014) displaying $\text{SFR}_{\text{FIR}} \sim 1000 M_{\odot}/\text{yr}$. However, their VLT/KMOS SFR(H α) range between 10 and 300 M_{\odot}/yr , making them compatible with the main sequence of star-forming galaxies except for the case of ID 647 (which likely host an AGN). This indicates the dust extinction of these

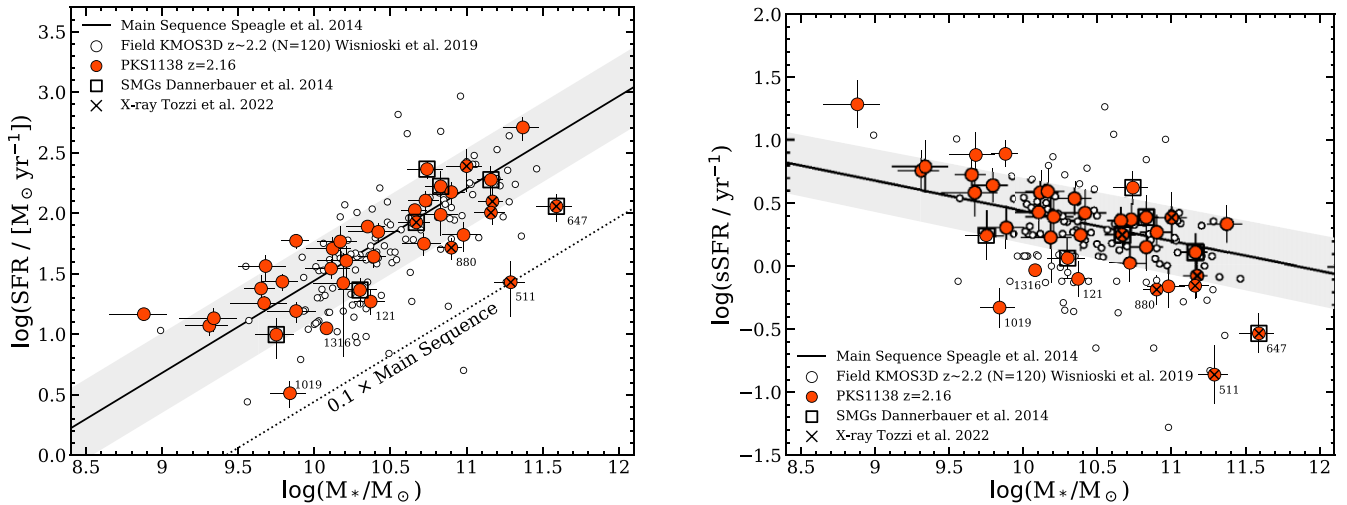


Figure 7. Left: Star-forming main sequence diagram. The red circles show the distribution of our spectroscopic sample of protocluster members. The empty circles display the field comparison sample from the KMOS3D survey (Wisnioski et al. 2019). This data set consists of 120 field galaxies at $z \sim 2.2$. The black solid line represents the star-forming main sequence parametrized for $z = 2.16$ (Speagle et al. 2014), while the grey shaded region marks the 3σ scatter around it. The dotted line depicts a sequence for the quenched or nearly quenched galaxies with one-tenth of the main sequence SFRs. Right: Star-forming main sequence using the specific star formation rate (sSFR). Symbols and colours remain the same as in the left-hand panel. Six objects display SFR values below the main sequence and have been labelled according to their HAE IDs (Koyama et al. 2013) for further discussion in the main text.

sources is highly underestimated, as was suggested by Dannerbauer et al. (2014). In general, HAEs in the Spiderweb protocluster also occupy a similar locus than the KMOS3D field sample of Wisnioski et al. (2015) at $z \sim 2.2$. These results suggest that most of our sample show no significant differences in terms of star formation compared to their field counterparts at similar redshift.

This can be seen easily by inspecting the right-hand panel of Fig. 7, where we investigate the main sequence using the specific star formation diagram (i.e. $\text{sSFR} = \text{SFR}/M_*$). The sSFR allows us to investigate the star-forming activity of a given galaxy normalized by its present stellar mass, which ease the process of identifying galaxies that are more efficient forming new generations of stars than expected per unit of stellar mass. The scatter of our sample increases with respect to the previous figure though we find similar results. We detect a group of galaxies lying below the ‘Main Sequence’ at $\log M_*/M_\odot \geq 10.0$ with sSFR values between 0.1 and 0.8 dex lower than expected, while at lower masses a few galaxies seem to display enhanced sSFR.

We also investigate the possible influence of the environment over these results by combining both local and global environmental indicators. The local density is measured as the projected surface density enclosing three neighbouring galaxies (i.e. Σ_3). This allows us to inspect the influence of local density peaks where small groups of galaxies are clustered and could be interacting in a relatively small area of the sky, although they may not necessarily lie close to the cluster-core. Fig. 8 shows the distribution of our galaxies in local density space as a function of the offsets with respect to the main sequence (i.e. $\Delta\text{MS} = \text{SFR} - \text{SFR}_{\text{MS}}$). We bin our sample into three density regimes defined by the minimum radius required to enclose three neighbouring galaxies within a circle. Following this approach, our high-density regime is defined by $R_2 < 100$ kpc, while the intermediate-density regime encompasses $200 \text{ kpc} > R_2 > 100$ kpc, and the low-density regime trace galaxies with measured $R_2 > 200$ kpc. As it can be seen in Fig. 8, the median values within each bin display a flat trend across almost three orders of magnitude in Σ_3 . A similar result is shown when applying a higher number of

neighbouring objects such as Σ_5 . These results suggest that local density peaks made of small groups of galaxies do not have a strong influence on the SFR of its galaxies, regardless of how dense they are. Intriguingly, five out of the seven X-ray emitters lie within the densest bin, while the other two reside at intermediate densities. This suggests that local density peaks where galaxy–galaxy interactions are more frequent may promote the triggering of AGN activity. On the other hand, the sample of seven SMGs spread almost evenly between the three local density regimes. A possible caveat to this approach is, however, the fact that the number surface density of galaxies in a given patch of the sky does not take into account the projection effects, i.e. the possible distance and velocities in the line of sight of the observations. This could potentially present some groups of galaxies as high-density peaks while its members lie at great distances in reality.

Finally, we use the global environmental indicator described by equation (7) and introduced by Noble et al. (2013) to check the potential influence of the protocluster’s structure over its members. This indicator relies on the combined knowledge of the cluster-centric distance of every galaxy and their redshift (or velocity) offset with respect to the systemic redshift of the cluster. We took the position and redshift of the Spiderweb radio galaxy (MRC 1138–262 at $z = 2.156$, Miley et al. 2006) as reference for our measurements. Our results are shown in the right-hand panel of Fig. 8. The median value of galaxies within the accreted region is slightly above the main sequence, while galaxies residing in the outskirts show a small SFR deficit. None the less, all median values are within the 3σ scatter of the main sequence, similarly to our previous findings using the local density environmental indicator. We will further investigate these results in Section 5.

4.3 Gas-phase metallicities

In this section, we investigate the impact of the environment over the metal enrichment of the ISM of protocluster galaxies. We follow the approach outlined in Section 3.5 to compute the gas-phase

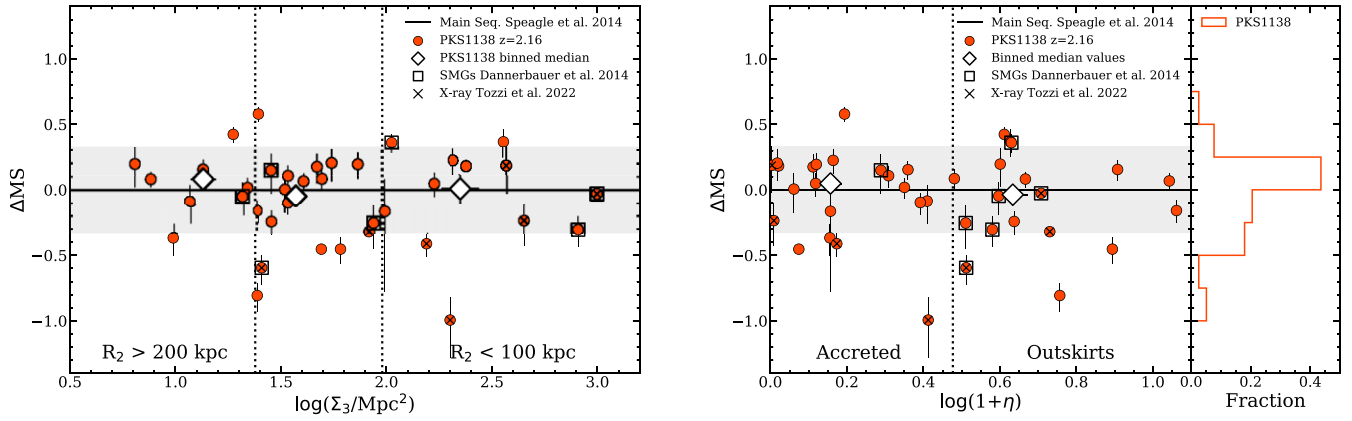


Figure 8. Left: Offsets from the main sequence of galaxies as a function of the local density. Right: Offsets from the main sequence of galaxies as a function of the general environmental parameter η . The horizontal solid line shows the position of the main sequence of star formation (Speagle et al. 2014) in both diagrams. The vertical dotted lines separate the density regimes outlined in Section 3.7. The colours and symbols follow the same scheme as in Fig. 7.

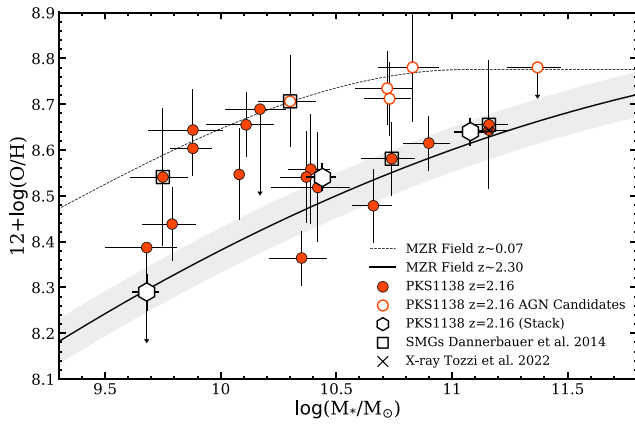


Figure 9. MZR diagram. Gas-phase metallicity measurements for individual objects in PKS1138 are shown by the filled and empty circles, with the latter being a population of AGN candidates according to their $[\text{N II}]/\text{H}\alpha$ ratios (see Section 4.3). The empty hexagons display the results of our stacking analysis in three stellar-mass bins. The solid line depicts the field MZR using a second-order polynomial fit over the combined results of the Erb et al. (2006) sample, the KMOS3D sample of Wuyts et al. (2016), and the MOSDEF sample of Sanders et al. (2021) at $z \sim 2.3$. The grey area shows the mean metallicity uncertainty for the samples used to define the field MZR at $z \sim 2.3$. The dashed line marks the position of the local MZR (Kewley & Ellison 2008) assuming the Pettini & Pagel (2004) metallicity calibration but adjusted to the Chabrier (2003) IMF.

metallicities of our targets based on the Pettini & Pagel (2004) calibration. Our results are shown in Fig. 9, where we inspect the distribution of our galaxies with respect to the field MZR at $z \sim 2.3$. We derive individual metallicity values for 24 out of 39 HAEs in our KMOS sample. Out of these 24 objects, two are excluded due to their broad $\text{H}\alpha$ profiles ($\sigma > 700 \text{ km s}^{-1}$). Another three galaxies display $\text{S/N} < 2$ in the $[\text{N II}]$ line and are shown as upper limits in Fig. 9. The remaining 15 galaxies do not achieve enough S/N in the $[\text{N II}]$ emission line to be measured, and thus they are only included as part of our stacking analysis.

Among the difficulties to study the environmental dependence of the MZR, the construction of an unbiased field comparison sample is one of the most prominent problems (Stott et al. 2013b; Namiki et al.

2019). According to the fundamental metallicity relation (Mannucci et al. 2010), it is expected that high-SFR samples are naturally biased towards lower oxygen metallicity values. Similarly, field samples that are selected using UV to optical bands (e.g. Erb et al. 2006) may overlook dusty galaxies and thus, underestimate their mean metallicity, especially at the low-mass end. In order to mitigate these effects, we build our field MZR upon the combined results of the KMOS3D survey (Wuyts et al. 2016), the MOSDEF survey (Sanders et al. 2021), and the ancillary sample (Erb et al. 2006), which share a similar SFR lower limit ($\sim 10 M_{\odot}/\text{yr}$, see Fig. 7) and reddening values with our protocluster spectroscopic sample. In Fig. 9, the field MZR (solid line) is built by fitting an error-weighted second-order polynomial of the form $y = ax^2 + bx + c$ with $y = 12 + \log(\text{O}/\text{H})$ and $x = \log(M_*/M_{\odot})$ to the mass-binned results of the three samples of $z \sim 2.3$ field galaxies. The resulting polynomial coefficients of the fitted relation (solid line) in Fig. 9 are $a = -0.03827$, $b = 1.02301$, and $c = 1.97744$, with the typical uncertainty (0.05 dex) being shown by the grey area. For reference, we also add the local MZR from Kewley & Ellison (2008) (dashed line in Fig. 9) assuming the same metallicity calibration but adjusted to the Chabrier (2003) IMF. The local relation of Kewley & Ellison (2008) saturates at $\log M_*/M_{\odot} = 11.06$. Thus, we assume a constant metallicity value beyond this limit.

Furthermore, we consider the effect of AGN contamination by splitting our sample in objects below and above $\log([\text{N II}]/\text{H}\alpha) \geq -0.35$ in the same figure (filled and empty circles, respectively). This threshold has been recently proposed by Agostino et al. (2021) as a modification of the Kauffmann et al. (2003) line to separate the purely star-forming from the composite and AGN regions in the BPT diagram (Baldwin, Phillips & Terlevich 1981) at relatively low $[\text{O III}]/\text{H}\beta$ ratios (e.g. $\log([\text{O III}]/\text{H}\beta) \leq -0.2$). This threshold can be used as a $[\text{N II}]/\text{H}\alpha$ upper limit beyond which the ionization source of most objects is either composite or dominated by the AGN contribution. Thus, we will hereafter label objects beyond this threshold as AGN candidates. Based on this criterion, we find five AGN candidates within our sample (empty circles in Fig. 9). Finally, we also consider that X-ray emission may hint the presence of AGN activity in some objects. Among the sample of 22 galaxies that entered our metallicity analysis, we only find one object with significant X-ray emission in previous works (Pentericci et al. 2002; Tozzi et al. 2022).

Most of our individual measurements display metallicity values above the field MZR, even when the AGN candidates are excluded from this analysis. The need for a significant detection of the [N II] emission line (see Section 3.1) could potentially explain this result, naturally biasing our individual measurements towards objects with high-metallicity values. To overcome this, we analyse the average properties of our sample by resorting to the stacking analysis (white hexagons in Fig. 9). Our results show that on average, our sample is compatible with the field MZR in the high-mass ($\log M_*/M_\odot > 10.8$) and low-mass ($\log M_*/M_\odot < 10.0$) ends. However, intermediate-mass galaxies display stacked gas-phase metallicities 0.06 ± 0.03 dex above the field MZR. This result is at odds with recent findings in overdense regions at $z \geq 2$ (Valentino et al. 2015; Chartab et al. 2021), which find a metallicity deficit in protocluster galaxies. However, others have reported metallicity enhancements (up to 0.2 dex) in massive protoclusters at a similar epoch (Kulas et al. 2013; Shimakawa et al. 2015) or no significant differences with respect to the field (Kacprzak et al. 2015). In Section 5, we will discuss physically motivated scenarios that may be responsible for these contradicting results in protoclusters at the cosmic noon.

One of these works (Shimakawa et al. 2015), however, is partially based on HAEs belonging to the Spiderweb protocluster. The discrepancy between our result and theirs may be related to several sample selection factors: First, their comparison sample is solely based on the Erb et al. (2006) UV-selected galaxies which tend to be biased towards lower metallicities. Furthermore, their spectroscopic campaign is based on slit spectra, which in some cases could miss a significant fraction of the emission coming from the galaxy disc. Even though the slit losses can be roughly estimated and thus corrected, this process can also be highly uncertain depending on the orientation and shape of the object with respect to the slit. In addition, the KMOS 2.8 arcsec \times 2.8 arcsec IFU allows to extract the flux from a given aperture without missing flux in the process. Finally, our spectroscopic program is based on a larger sample of sources across a wider area, improving the overall accuracy of the stacked measurements.

We also examine the relation between the metallicity enhancement measured in our sample and the environment these HAEs live in by following the two-fold approach outlined in Section 3.7. First, we consider the impact of the local environment in the left-hand panel of Fig. 10, where we show the metallicity offsets with respect to the field MZR as a function of local density indicators (Σ_3). After splitting our sample in three density bins we find a median binned metallicity enhancement of ~ 0.05 – 0.2 dex across the density regimes examined. Interestingly, there is a mild decreasing metallicity trend towards higher local density values. By analysing the median values of our sample with respect to the global environment indicator η (Fig. 10 right), we find an almost flat metallicity trend with a slightly higher median value within the accreted region of the protocluster. In summary, while the measurements coming from the highest local density peaks (i.e. Σ_3) within our sample show the lowest median metallicity values, we do not observe a clear trend from the core to the outskirts of the protocluster (in terms of η), suggesting that local and global environmental indicators trace different structures within the protocluster. Finally, we measure individual [N II]-based gas metallicities for four out of the seven SMGs within our sample. Two of them are consistent with the expectations in the field while another two show ~ 0.2 dex enhancement (Fig. 10). None the less, the three objects without [N II] detection likely have lower metallicity than the field, indicating that SMGs may share a similar distribution to the HAEs in the

MZR. We will further explore the possible causes of these results in Section 5.

4.4 Mass–size relation

We inspect the distribution of our KMOS-detected HAEs with respect to the stellar mass–size relation of field galaxies at similar redshift (van der Wel et al. 2014). The effective radius (R_e) of our targets is estimated using the VLT/HAWKI K_s -band images of this protocluster. Due to the excellent seeing conditions and depth of these images (FWHM ~ 0.4 arcsec), we are able to measure the effective radius of 27 of our KMOS HAEs. However three of these sources lie below the seeing limit (grey area). Most of our targets show surface brightness profiles compatible with the presence of extended discs (Fig. 11) in line with the predictions for late-type galaxies in the field at $z \geq 2$ (light blue band). Other works exclusively investigating HAEs find a flatter trend than van der Wel et al.’s (2014) at the high-mass end of the mass–size relation (e.g. Stott et al. 2013a; Paulino-Afonso et al. 2017). This disagreement may arise from the different selection techniques (UVJ colours and H α emission, respectively) employed to identify star-forming galaxies. Nevertheless, we find that the mean effective radius of our sample ($R_e = 3.14 \pm 0.21$) and the one from Paulino-Afonso et al. (2017) agrees within the errors ($R_{e, PA17} = 2.86 \pm 0.14$), though our sample may be biased towards higher sizes due to the seeing limitations and the higher fraction of massive galaxies.

On the other hand, five objects lie within or close to the early-type sequence in the mass–size relation, meaning that their morphology is considerably more compact than the rest of sources within our sample. Four of these objects are confirmed X-ray emitters (Tozzi et al. 2022), suggesting that they are likely hosting AGNs (IDs 511, 647, 880, and 1501). For ID 647, its K_s -band emission may be dominated by the H α broad component coming from the AGN torus, explaining its compact effective radius. In addition, IDs 511 and 880 display SFR values bellow the 1σ scatter of the main sequence in Fig. 7. These objects may be transitioning from their late-type star-forming phase towards the early-type quiescent regime. The object (ID 902) which lies in the overlapping region between the early- and late-type sequences in the mass–size diagram is also a massive AGN candidate according to its [N II]/H α ratio.

In overdense regions, environmentally driven processes such as mergers and close encounters may temporarily distort the stellar morphology of interacting galaxies or create tidal tails (e.g. Moore et al. 1996), both increasing or reducing the stellar-mass surface density depending on the geometry of the interaction and the physical properties of the galaxies involved (e.g. mass and gas reservoir, Lagos et al. 2018). On the other hand, such events may as well channel large amounts of gas towards their central regions, triggering starburst events or AGN activity in some cases, and favouring the creation of prominent bulges. In order to test these effects we inspect the stellar-mass surface density ($\mu_* = M_*/(\pi R_e^2)$) of our objects as a function of both local density and general environment indicators in Fig. 12. Our spectroscopic sample displays a rather flat trend ($\sim 10^9$ – $10^{10} M_\odot \text{ kpc}^{-2}$) across its stellar-mass range ($\log M_* = 9.5$ – 11.5) and three orders of magnitude in local density (Σ_3). This can be understood as a consequence of the tight distribution that most of our targets display along the late-type sequence in the mass–size relation (Fig. 11). Similar results are achieved after repeating this analysis by comparing the distribution of our targets as a function of the global environment parameter instead (η , right-hand panel). The small variations between the binned median stellar-mass surface values across different environmental regimes in Fig. 12 can be partly

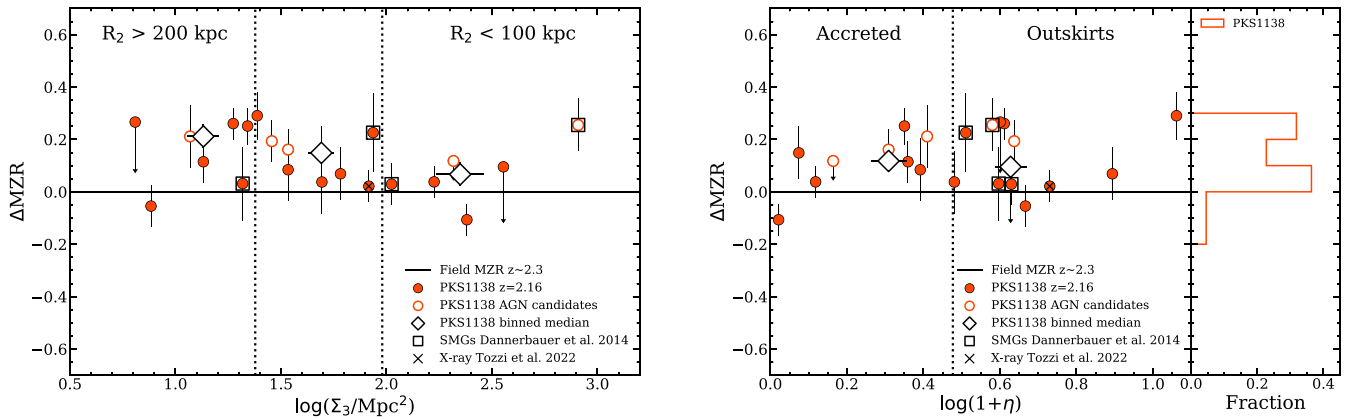


Figure 10. Left: Offsets from the field MZR as a function of local density values. Right: Offsets from the field MZR as a function of the general environmental parameter η . The horizontal solid line shows the position of the field MZR from Fig. 9. The vertical dotted lines separate the density regimes outlined in Section 3.7.

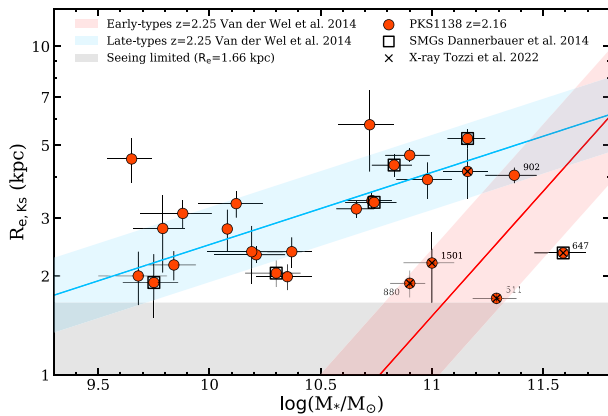


Figure 11. Mass–size relation. The red circles show the distribution of the KMOS $H\alpha$ emitters of this work. The empty squares display the sample of SMGs studied in Dannerbauer et al. (2014). The blue and red sequences, respectively, depict the position of the late- and early-type galaxies in the mass–size plane at $z = 2.25$ (van der Wel et al. 2014). The grey area displays the limiting seeing size in the physical scale. We label a few objects with their IDs for specific discussion in the main text.

explained by a combination of factors: the heterogeneous nature of our sample (including HAEs, AGN candidates, and SMGs) and the limited number of object per density bin.

4.5 Molecular gas properties

In this section, we explore the relation between the molecular gas fraction ($f_{\text{gas}} = M_{\text{mol}}/(M_* + M_{\text{mol}})$) and gas–phase metallicity of seven HAEs within our KMOS sample. Our objects were selected after cross-matching the CO(1–0) line flux catalogue published by Jin et al. (2021) with our spectroscopic sample. Out of eight overlapping targets, we obtained metallicity measurements based on the $[\text{N II}]/H\alpha$ ratio for seven of them. The molecular gas mass of every object is derived following the scaling relations published by Tacconi et al. (2018) and the procedure outlined in Section 3.6.

Bothwell et al. (2013) and Hunt et al. (2015) showed that the gas–phase metallicity gradually increases as the gas reservoir of the galaxy is consumed (i.e. towards lower gas fractions) and the SFR slowly decreases. Such combination of processes were put into

context by Peng & Maiolino (2014) in their gas regulator model, which also considers the influence of gas inflows and outflows in the evolution of the aforementioned physical properties. This model relies on four main input parameters governing the evolution of galaxies: the star formation efficiency, the gas inflow rate, the mass-loading factor ($\lambda = \text{outflows/SFR}$), and the return mass fraction. In this model, the inflowing gas is assumed to scale with the growth rate of the dark matter halo, while the outflow rate correlates with SFR and the return fraction is a constant parameter fixed for a given IMF. For a more in-depth discussion of the model we refer to Peng & Maiolino (2014).

In this work, we follow the approach outlined in Suzuki et al. (2021), who also relied on the gas regulator model to investigate the gas fraction–gas–phase metallicity relation in the field at $z \sim 3.3$. For reference, we overplot the lines of constant mass-loading factor (λ) for $\lambda = 0\text{--}2.5$ in Fig. 13. Our seven protocluster members show a mild negative trend between the gas metallicity and gas fraction. However, they are shifted towards lower mass-loading factor values compared to the field samples of Suzuki et al. (2021) at $z = 3.3$ and Seko et al. (2016) at $z = 1.5$, whose values in Fig. 13 have been recomputed to match the IMF and calibration of this work following Curti et al. (2017) transformations. This means that our objects display higher metallicity values at a fixed gas fraction. In the context of the gas regulator model and assuming that our galaxies are representative of the main sequence of star-forming galaxies at $z = 2.2$ (see Fig. 7), the shift in λ may imply that the enrichment of the ISM may be (at least partly) driven by the suppression of outflows in our targets. This effect could be the result of the influence of the protocluster environment, where a denser IGM compared to the field would exert external pressure over the gas halo of the protocluster members, preventing the outflows from star formation to leave the already enriched gas. However, the small size of our current sample and the influence of AGN candidates prevent us from discarding other possibilities at this stage. We will discuss some of these possibilities in Section 5.

5 DISCUSSION

In this work, we have searched for environmental imprints on the star formation, gas–phase metallicity, morphology, and gas fraction of a sample of 39 spectroscopically confirmed protocluster galaxies at

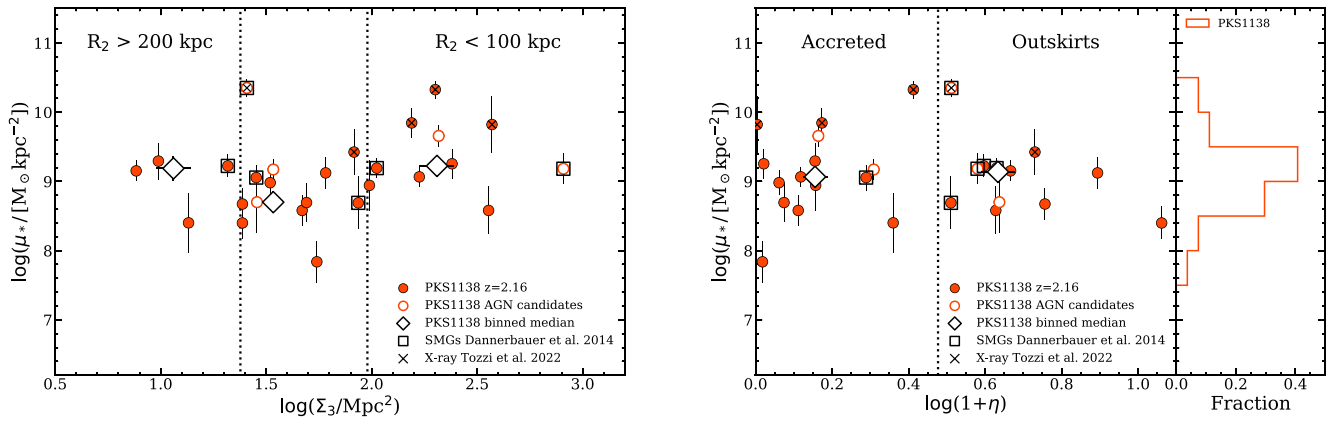


Figure 12. Left: Local density distribution of our sample as a function of mass surface density defined as $\mu_* = M_*/(\pi R_c^2)$. Right: general environment distribution (η) of our targets as a function of the mass surface density defined above. The vertical dotted lines separate the density regimes outlined in Section 3.7. The colours and symbols follow the same scheme as in Fig. 10.

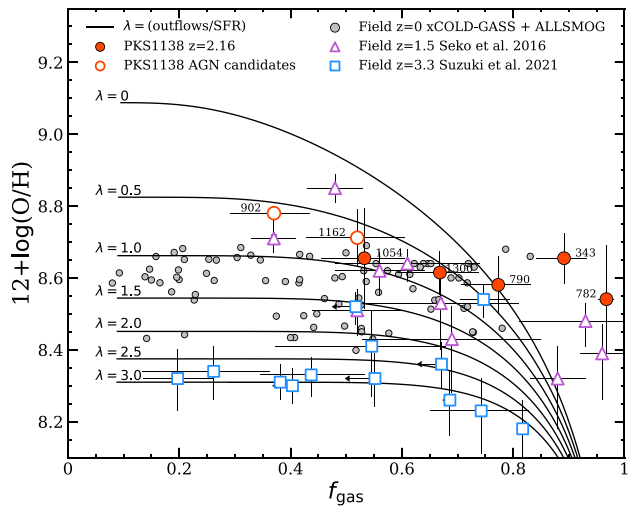


Figure 13. Molecular gas fraction versus gas phase metallicity. The solid lines show the tracks of constant mass-loading factor (outflows/SFR) assuming equilibrium in the context of the gas regulator model (Peng et al. 2010). The filled and empty circles display our subsample of seven HAEs with measured gas fractions from CO(1–0). The blue-edged squares show the distribution of field galaxies at $z = 3.3$ from Suzuki et al. (2021). The violet-edged triangles depict the field sample of Seko et al. (2016) at $z = 1.5$. The local field comparison sample is composed by galaxies from the XCOLD GASS (Saintonge et al. 2017; Catinella et al. 2018) and the ALLSMOG (Bothwell et al. 2014) surveys. We label our objects according to their HAE IDs (Koyama et al. 2013) for discussion in the main text.

$z \sim 2.2$. In the following subsections, we will discuss the implication of our results on different galaxy evolution scenarios.

5.1 Star-forming galaxies in protoclusters

Strong gravitational interactions such as mergers should be more frequent in high- z protoclusters than in local Universe clusters according to N -body and hydrodynamical simulations (Gottlöber, Klypin & Kravtsov 2001; Genel et al. 2014). This can be understood as a natural consequence of the assembling phase these structures are experiencing, with small groups infalling towards the highest density peak in the field, and the low relative velocity between the galaxies therein

(Hine et al. 2016). Therefore, it is also expected that such interactions increase the fraction of starbursting galaxies within protoclusters. However, different works on protoclusters at $z = 2–3$ have found mixed results in terms of the star formation activity of their members. While Koyama et al. (2013), Shimakawa et al. (2018b), Sattari et al. (2021), and this work find no significant differences between the protocluster and field population, others such as Shimakawa et al. (2018a) and Wang et al. (2022) find that protocluster galaxies display enhanced SFRs (specially in the low stellar mass range), and that the amplitude of such enhancement positively correlates with higher density regions within the protocluster structure.

Two considerations should be taken into account when analysing these conflicting results. First, most overdense structures that are classified as protoclusters at $2 < z < 3$ undergo a major though rapid phase of mass accretion and mass build-up until $z \sim 1.5$, when some of the most massive structures start showing a well-developed red sequence (e.g. XMMU J2235-2557 at $z = 1.4$, Rosati et al. 2009 or XLSSC 122 at $z = 1.99$ Willis et al. 2020). This implies that protoclusters belonging to the same epoch ($z = 2–3$) may be going through different evolutionary stages and thus, their galaxy population could be affected in different ways by the environment. In fact, recent studies at low redshift have shown that the dynamical properties of clusters do influence the evolution of their galaxies both in terms of star formation and AGN activity (Stroe & Sobral 2021). In that sense, the Spiderweb protocluster is one of the most massive structures at $z \sim 2$ and show signs of virialization within its inner core (< 0.5 Mpc, Shimakawa et al. 2014), while other protoclusters such as USS1558 at $z = 2.53$ (Hayashi et al. 2016) or BOSS1244 at $z = 2.2$ (Zheng et al. 2021) are formed by several less-massive clumps ($< 10^{14} M_\odot$) hosting galaxies with low proper velocities, thus favouring galaxy–galaxy interactions driven starbursts. This could explain why we find no significant differences in terms of star formation between the Spiderweb protocluster and the field population, in contrast with results in younger and less-massive structures. Second, the detection of high- z protoclusters and the analysis of their galaxy populations are carried out through a wide variety of techniques (e.g. SED, narrow-band, spectroscopic surveys, etc.) and focusing on different wavelength ranges. This could potentially bias the selection of the samples, making it difficult to establish a fair comparison between the galaxy populations of different protoclusters. Finally, the signs of SFR enhancement in protocluster galaxies are oftenly restricted to the low stellar mass

regime (e.g. at $\log M_*/M_\odot < 9.5$ in Hayashi et al. 2016), adding further uncertainties since most studies are not able to achieve a high completeness level in this mass range due to observational depth limitations, as it is the case of this work.

Recently, Tozzi et al. (2022) reported 14 protocluster members with significant X-ray emission hinting at a possible enhancement of the AGN fraction with respect to the field. Our KMOS sample contains seven of these objects. They are massive HAEs ($\log M_* \geq 10.8$) which tend to lie below the main sequence of star formation, with four of them displaying $\Delta MS \leq -0.3$ dex (Fig. 7). We find that the overlapping X-ray emitters are sparsely distributed from the outskirts to the core of the protocluster in terms of η (Fig. 8). However, five of them reside within the densest bin in terms of Σ_3 , while the other two lie at intermediate local densities. These results suggest that AGNs may preferentially be located at local density peaks but not necessarily within the innermost regions of the protocluster. The impact of the X-ray emitters can also be spotted in the mass–size relation (Fig. 11), where four out of the five objects with measured R_e display rather compact morphologies. In contrast, the majority of the HAEs that entered our analysis show sizes compatible with the presence of extended discs. We must note, however, that our R_e measurements are based on K_s band which overlaps with the $H\alpha$ emission line at $z = 2.16$. After checking the $H\alpha$ contribution to the total K_s -band flux we find that it accounts for $\lesssim 15$ per cent except for two cases. Therefore, we conclude that the size measurements are not driven by contamination from star-forming gas emission. For AGNs, a significant fraction of the K_s emission could be originated by the $H\alpha$ broad component at their centre, making them look more compact. Furthermore, three of these objects also display low SFRs (see Section 4.4) or UVJ colours similar to those of early-type galaxies, suggesting that they may be transitioning towards a quiescent state. However, it is not possible to discern between these two options with the current depth and seeing limitations of our data.

Finally, we also inspect the properties and spatial distribution of a sample of seven HAEs identified as SMGs by Dannerbauer et al. (2014). We find a factor 3–10 lower $H\alpha$ -based SFRs than the FIR-based ones, confirming previous reports of unaccounted dust obscuration in the rest-frame optical regime of this starbursting objects (Dannerbauer et al. 2014). Furthermore, these seven SMGs are found at projected cluster-centric distances larger than $1.6 R_{200}$ (Fig. 6), with six of them residing at the protocluster’s outskirts according to their global environment indicator ($\eta > 2$, Fig. 8). While the spatial distribution of our SMGs is not significantly offset with respect to the HAE population as in Zhang et al. (2022), it is clear that these objects are not yet part of the innermost regions of the Spiderweb protocluster (Fig. 2, see also Dannerbauer et al. 2014). Several authors have established links between the violent dust-obscured star-forming activity of SMGs at $z > 2$ and the surge of red ellipticals at later epochs (e.g. Swinbank et al. 2006; Michałowski, Hjorth & Watson 2010; Simpson et al. 2014; Toft et al. 2014). In fact, Smail et al. (2014) reported a similar spatial mismatch for SMGs in a cluster at $z \sim 1.6$, hypothesizing that this infalling population may constitute the progenitors of some of the fainter ellipticals dominating cluster cores by $z \sim 0$, while the pre-existing passive cluster members may have formed following the same channel at even earlier epochs. Several studies have spectroscopically identified a few members of the nascent passive population of the Spiderweb protocluster (e.g. this work, but see also Doherty et al. 2010; Tanaka et al. 2013). Assuming that these galaxies were formed through a similar SMG phase in the past, and that the duration of the starburst causing the such phase lasted a few hundred Myr (e.g. Riechers et al. 2011; Hickox et al.

2012), we estimate their formation period at $z \lesssim 2.5$. On the other hand, the current SMGs in this protocluster may deplete their gas reservoirs and start their quenching process by $z \approx 2$, concurrently to their gradual infall towards the protocluster core.

5.2 ISM enrichment at the cosmic noon

We have analysed the gas-phase metallicity properties of 22 individual protocluster members (Fig. 9) based on their $[N II]/H\alpha$ ratios. We obtained enhanced individual values (up to 0.3 dex) with respect to the field MZR at $z \sim 2.3$. However, our stacking analysis yields metallicity values only slightly above the field MZR for intermediate-mass galaxies ($\Delta MZR = 0.06 \pm 0.03$), while no significant metallicity enhancement is observed in the high- and low-mass regime. We find that there exists a population of AGN candidates according to their high $[N II]/H\alpha$ ratios (see Section 4.3). However, observation by Tozzi et al. (2022) reported X-ray emission for only one of these candidates. Even though the X-ray emission originated through AGN activity can be obscured for several reasons (Hickox & Alexander 2018), the low number of X-ray detections between the targets with high $[N II]/H\alpha$ ratios suggests that our results are not severely affected by AGN contamination and thus, we may be witnessing a population of highly metal-enriched objects.

Previous studies on the MZR evolution in protoclusters at the cosmic noon have shown mixed results. For example, Kulas et al. (2013) found a significant metallicity excess of 0.1–0.2 dex with respect to the field, in particular for galaxies in the low-mass end in a protocluster with a virial mass well above $M_{\text{vir}} \geq 10^{14} M_\odot$ (Steidel et al. 2005). A possible scenario to explain this result involves the recycling of outflowing chemically enriched gas due to the external pressure exerted by the surrounding IGM in clusters. This effect would predominantly affect galaxies in the intermediate- to low-mass regime, as the gravitational potential of the most massive systems is strong enough to retain their metals regardless of the environment (Oppenheimer & Davé 2008). Interestingly, this scenario is expected to be more efficient in massive virialized clusters, as the density of the IGM correlates with the total mass of the cluster and the accretion of pristine gas which could potentially dilute the metallicity could have stopped some time ago due to the shock heating of their main halo (Dekel & Birnboim 2006). Nevertheless, other studies in less-massive structures ($M_{\text{vir}} \sim 10^{13} M_\odot$) report a relatively constant metallicity deficit of similar magnitude across the entire galaxy stellar-mass range, which can be understood as a consequence of the cold gas accretion into these relatively young assembling protoclusters (Valentino et al. 2015; Chartab et al. 2021). Finally, several authors have reported a diversity of behaviours within their samples, with massive galaxies being less metal enriched than the low-mass ones but with absolute values close to the field MZR (Sattari et al. 2021; Wang et al. 2022), or no clear environmental dependence at all (Kacprzak et al. 2015).

Using hydrodynamical simulations, Dekel et al. (2009) proposed that protoclusters at the cosmic noon experience a gradual but rapid transformation from an early assembly phase where the cold gas accretion dominates in the halo, to a short transition period where cold streams can still penetrate an otherwise shock-heated medium, to a final stage where the cold streams are completely suppressed and the cluster achieves a virialized state. These three sequential scenarios could potentially explain the variety of metallicity results in the literature as part of an evolutionary path that involves the co-evolution of the galaxies and their host protocluster halo. In this context, the mild metallicity enhancement that we detect in the Spiderweb protocluster ($M_{200} \sim 2 \times 10^{14} M_\odot$, Shimakawa et al.

2014) suggest that the protocluster halo may have recently become shock heated to the point where inflows are no longer efficient to supply cold gas to the member galaxies.

The environmental analysis of our individual measurements in Fig. 10 reveals that the accreted region of the Spiderweb protocluster display similar median metallicity values than the outskirts in terms of the global environment indicator (η). On the other hand, a declining metallicity trend is found towards high local density peaks (Σ_3). This trend, however, may be predominately driven by massive galaxies dominating the highest Σ_3 bin ($\log M_*/M_\odot = 10.66 \pm 0.21$), while lower mass galaxies populate the intermediate regime ($\log M_*/M_\odot = 10.36 \pm 0.14$) and the least dense regions ($\log M_*/M_\odot = 10.17 \pm 0.18$). This mass segregation is not present when examining the metallicity variations as a function of η , with the accreted and outskirts regions sharing very similar stellar mass median values ($\log M_*/M_\odot = 10.39 \pm 0.15$ and 10.36 ± 0.13 , respectively). None the less, larger number statistics and deeper observations, specially for the intermediate- to low-mass star-forming population ($\log M_*/M_\odot \lesssim 10.0$), are required to determine if significant metallicity differences between environmental regimes exist within this protocluster.

5.3 The role of the molecular gas fractions

The gas regulator model (Peng & Maiolino 2014) provides us with a practical frame to investigate the relation between the possible suppression of outflows via the external pressure of the IGM and the gas-phase metallicity. As we discussed in Section 4.5, this model relies on a few input parameters, being one of them the mass-loading factor ($\lambda = \text{outflows/SFR}$). In Fig. 13, we explore the relation of this parameter with both the molecular gas fraction and the gas-phase metallicity of our targets in comparison with field samples at several cosmic epochs. It is expected that galaxies move towards higher mass-loading factors from the local to the high- z Universe (Suzuki et al. 2021). This can be explained by the different star formation histories that galaxies at different redshifts experience, with those at high z experiencing vigorous but short burst of star formation (increasing the feedback too), while the stellar-mass growth proceeds in a more gradual manner at lower redshifts. Indeed, the majority of the $z = 3.3$ field galaxy sample discussed by Suzuki et al. (2021) is consistent with $\lambda = 2.5\text{--}3$, while results at $z = 1.5$ display $\lambda = 1\text{--}2$ (Seko et al. 2016) and local Universe samples lie around $\lambda \sim 1\text{--}1.5$ (ALLSMOG from Bothwell et al. 2014; xCOLD-GASS from Saintonge et al. 2017 and Catinella et al. 2018). In contrast, a small subsample of seven HAEs embedded in the Spiderweb protocluster displays unusually low mass-loading factors for galaxies at $z \gtrsim 2$. This subsample is composed of two objects (IDs 343 and 782) with relatively low stellar mass ($\log M_*/M_\odot = 10.11$ and 9.75 , respectively) and strong MZR offsets ($\gtrsim 0.2$ dex) with respect to the field. Another three objects (IDs 790, 1054, and 1300) with high stellar masses ($10.7 < \log M_*/M_\odot < 11.2$) but small offsets ($\Delta\text{MZR} \lesssim 0.05$ dex) and two additional galaxies (IDs 902 and 1162) which share similar characteristics to the previous ones in terms of stellar mass but show $0.11 < \Delta\text{MZR} < 0.16$ are labelled as AGN candidates due to their high [N II]/H α ratios.

As we have extensively discussed in previous sections, most of our spectroscopic sample display SFRs consistent with those of the main sequence of star formation. In principle, this would suggests that the lower λ values are mostly driven by a lower outflow contribution in these protocluster galaxies, supporting the scenario where galaxies are forced to recycle their gas due to the external pressure of the IGM and thus, yielding higher metallicity values at a fixed molecular

gas fraction. Following the accretion mode evolutionary sequence depicted in Dekel et al. 2009, this scenario would also suggest that cold streams are no longer efficient on supplying pristine gas to some members of the Spiderweb protocluster, in agreement with the theoretical predictions of Dekel & Birnboim 2006 for a protocluster of $M_{\text{vir}} \geq 10^{14} M_\odot$ at $z = 2.16$.

However, we should also consider the current statistical limitations and biases of our sample. For example, these seven galaxies are drawn from those with [N II] detection within our sample, which predominately trace the metal-rich end of the protocluster distribution (see Section 4.3). This would naturally bias these objects towards higher metallicity values in Fig. 13 and thus, they may not be representative of the full population of HAEs in the Spiderweb protocluster if metallicity and gas fraction information were available for all them. In the following text, we will continue discussing our current results while keeping in mind this limitation. In particular, the two galaxies with the lowest stellar mass values are at the same time the ones with the strongest ΔMZR and highest gas fraction (IDs 343 and 782), indicating that their ISM has been significantly enriched while their gas reservoir remains nearly intact. It is tempting to think of these two objects as candidates to support environmentally driven gas recycling, specially by taking into account their low stellar mass (see Section 5.2). However, we would require significantly larger number statistics to confirm such scenario. As for the remaining five objects, their high stellar mass ($\log M_*/M_\odot > 10.7$) makes them less prone to suffer this kind of environmental effect (Oppenheimer & Davé 2008). None the less, it is still interesting to find relatively high gas fractions for such metal-enriched galaxies when the Universe was only 3 Gyr old. The rapid build-up of metals in these massive objects, if not mediated by any environmental effect, would require extreme episodes of star formation in the recent past, thus transitioning from high to low λ in a short time-scale.

6 CONCLUSIONS

In this work, we have investigated the star formation, galaxy size, gas-phase metallicity, and molecular gas fraction of a sample of HAEs within the Spiderweb protocluster at $z = 2.16$. We based our results on new multi-object NIR spectroscopy with VLT/KMOS accompanied by previously reported spectroscopic and photometric observations covering from the UV to the submillimetre wavelength regime. This wealth of data allowed us to accurately trace the environment of the protocluster using both local and global environment indicators after confirming the cluster membership of over a hundred galaxies. We have examined the physical properties of 39 HAEs as a function of environment, and discussed the implications of these results on different scenarios of environmentally driven galaxy evolution during the early stages of massive cluster assembly. In the following paragraphs, we summarize the main conclusions of this work:

(i) We analysed a sample of 42 narrow-band-selected H α emission-line candidates with VLT/KMOS in the field of the Spiderweb protocluster at $z = 2.16$. We spectroscopically confirm the protocluster membership of 39 HAEs and detect one foreground [O III] emitter at $z > 3$. Thus, we achieved a ~ 93 per cent success rate on the cluster membership confirmation based on narrow-band-selected HAEs.

(ii) We measure the star formation activity of our spectroscopically confirmed HAEs finding that most of these objects are consistent with the field main sequence of star formation at $z \sim 2.16$ (Speagle et al. 2014; Wisnioski et al. 2015). Furthermore, these

results show no correlation with several environmental indicators (e.g. local density and global environment), suggesting that the star formation of HAEs has not been strongly affected by the environment of the Spiderweb protocluster.

(iii) Our morphological analysis reported that our HAEs sample is predominantly composed of late-type galaxies with observed K_s -band effective radius consistent with those of field galaxies at $z \sim 2.2$ (van der Wel et al. 2014). Only five HAEs show compact morphologies compatible with those of early-type galaxies at this redshift. However, no environmental correlation was found when exploring the distribution of stellar-mass surface density (μ_*) of our targets across different density regimes.

(iv) We extract [N II]-based metallicities for a subsample of 22 HAEs. We find that five of them have [N II]/H α ratios consistent with those of AGN candidates according to Agostino et al. (2021). Most of the remaining protocluster members still show individual enhanced metallicities by up to ~ 0.25 dex with respect to the field MZR at similar redshift (Erb et al. 2006; Wisnioski et al. 2019; Sanders et al. 2021). However, our stacking analysis reveal just a mild metallicity enhancement at intermediate masses ($\Delta\text{MZR} = 0.06 \pm 0.03$) and no significant differences at the high- or low-mass ends.

(v) We also studied a subsample of seven SMGs in this protocluster from Dannerbauer et al. (2014). Even though their distribution in the star-forming main sequence is similar to the parent sample of HAEs when using SFR(H α), the comparison with FIR-based SFRs from Dannerbauer et al. (2014) suggests that these sources are highly dust-obscured. SMGs are also consistent with the MZR traced by HAEs of similar stellar mass, and their R_{e,K_s} are comparable to those of late-type galaxies at $z \sim 2.2$. However, SMGs are preferentially found in the outskirts (six) of this protocluster rather than in their accreted region (one), in agreement with previous works in (proto-)clusters at $1.5 < z < 2.5$ (Dannerbauer et al. 2014; Smail et al. 2014; Zhang et al. 2022).

(vi) After exploring the gas metallicity gas fraction relation, we find that the seven HAEs with available metallicities and molecular gas fractions are distributed along mass-loading factor tracks comparable to those of the local Universe but in contrast with previous studies at $z > 2$. Taking into account that these are normal star-forming galaxies, we could interpret this as a sign of suppressed outflow activity, indicative of environmental effects in the form of external pressure over the galaxies' halo by a denser IGM in the Spiderweb protocluster. However, a systematic analysis of the metal enrichment and gas fraction across the protocluster structure would be required to confirm or discard such scenario.

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DATA AVAILABILITY

The VLT/KMOS spectroscopic data used in this work are publicly available through the ESO Science Archive Facility (program ID 095.A-0500, PI: Y. Koyama). The ATCA CO(1–0) map belongs to the COALAS project (large program ID: C3181, PI: H. Dannerbauer). The photometric data used in this work are publicly available through the Subaru Mitaka Okayama Kiso Archive (SMOKA) system, the Hubble Legacy Archive (proposal ID 10327, PI: H. Ford), the NASA/IPAC Infrared Science Archive (campaign ID 1634, PI: D. Stern), and the ESO Science Archive Facility (program IDs 088.A-0754, 091.A-0106, 094.A-0104, PI: A. Kurk). Any other data set within this article will be shared on reasonable request to the corresponding author.

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¹<http://www.astropy.org>

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APPENDIX A: ADDITIONAL MATERIAL

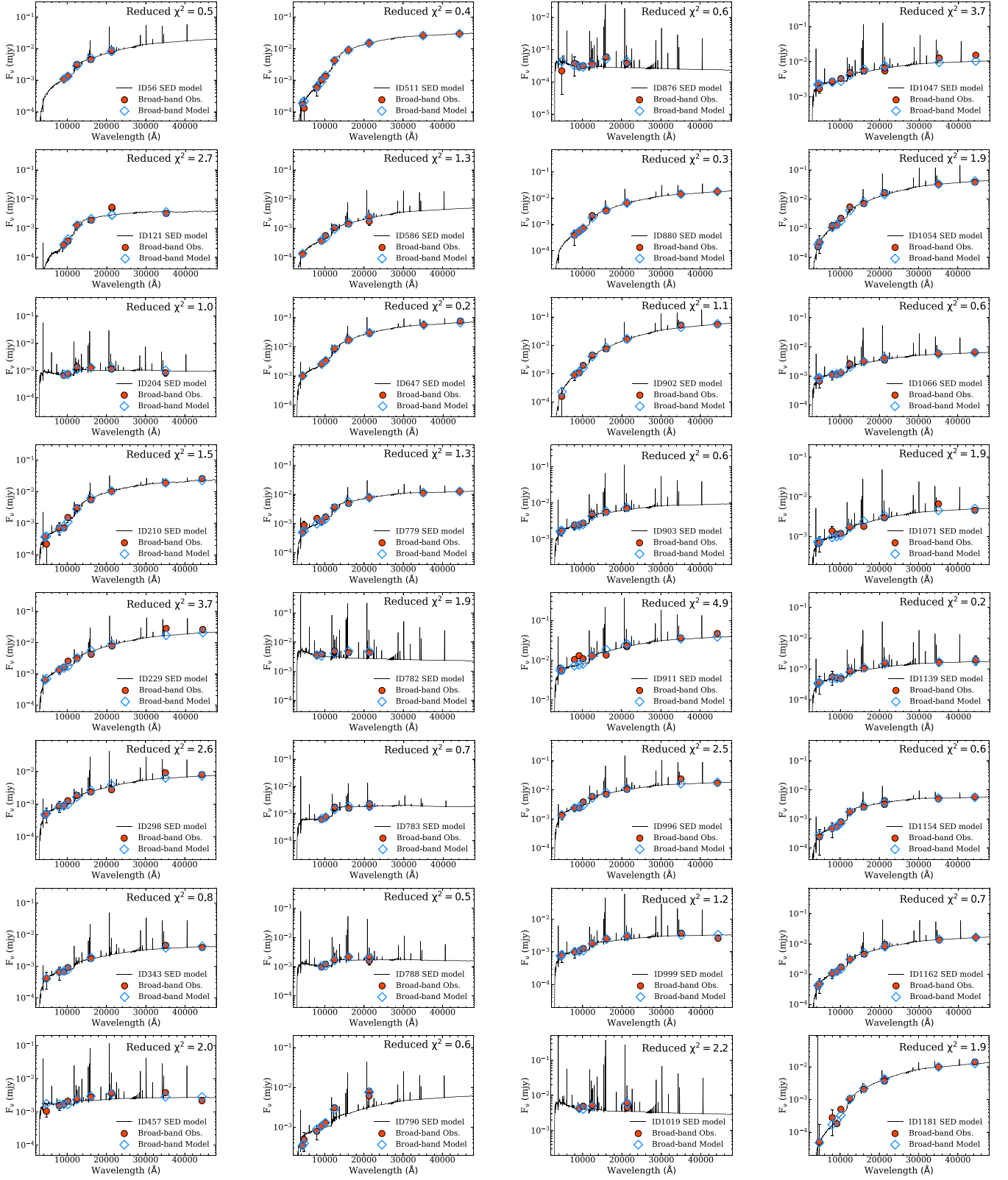


Figure A1. CIGALE SED best model and broad-band observations for every member of the protocluster sample.

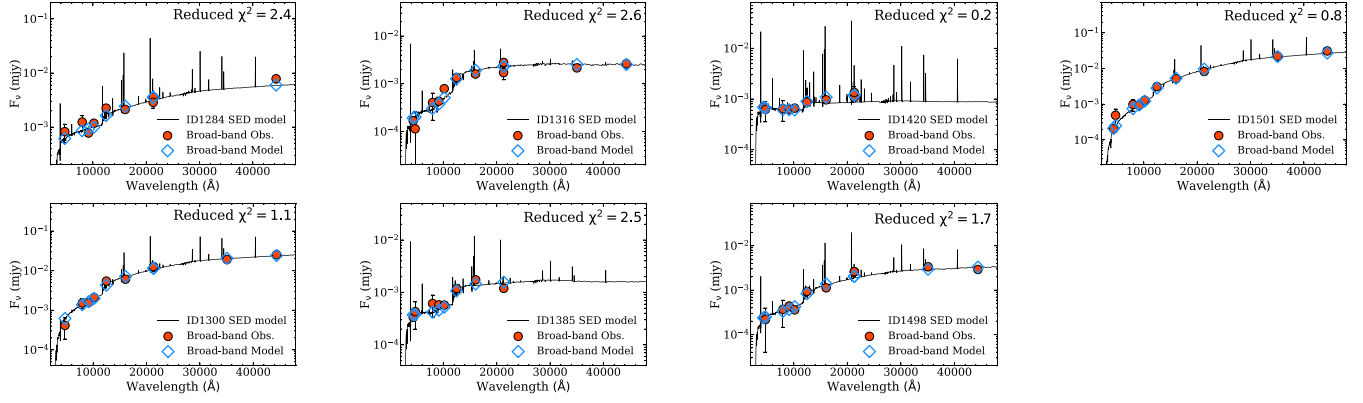
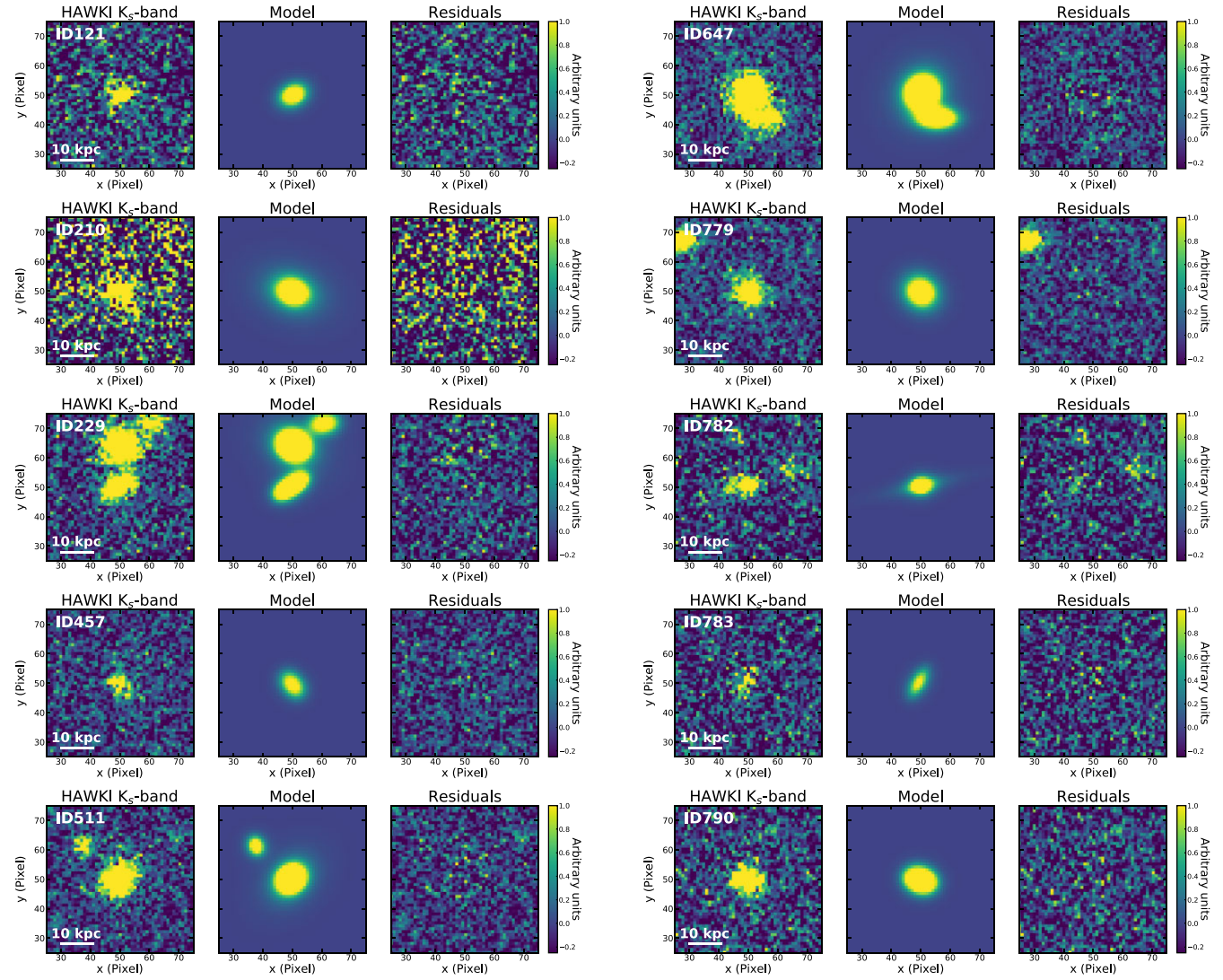


Figure A1 – continued

Figure A2. VLT/HAWKI K_s -band images, Galfit models, and residuals for the objects included in our mass–size relation (Fig. 11).

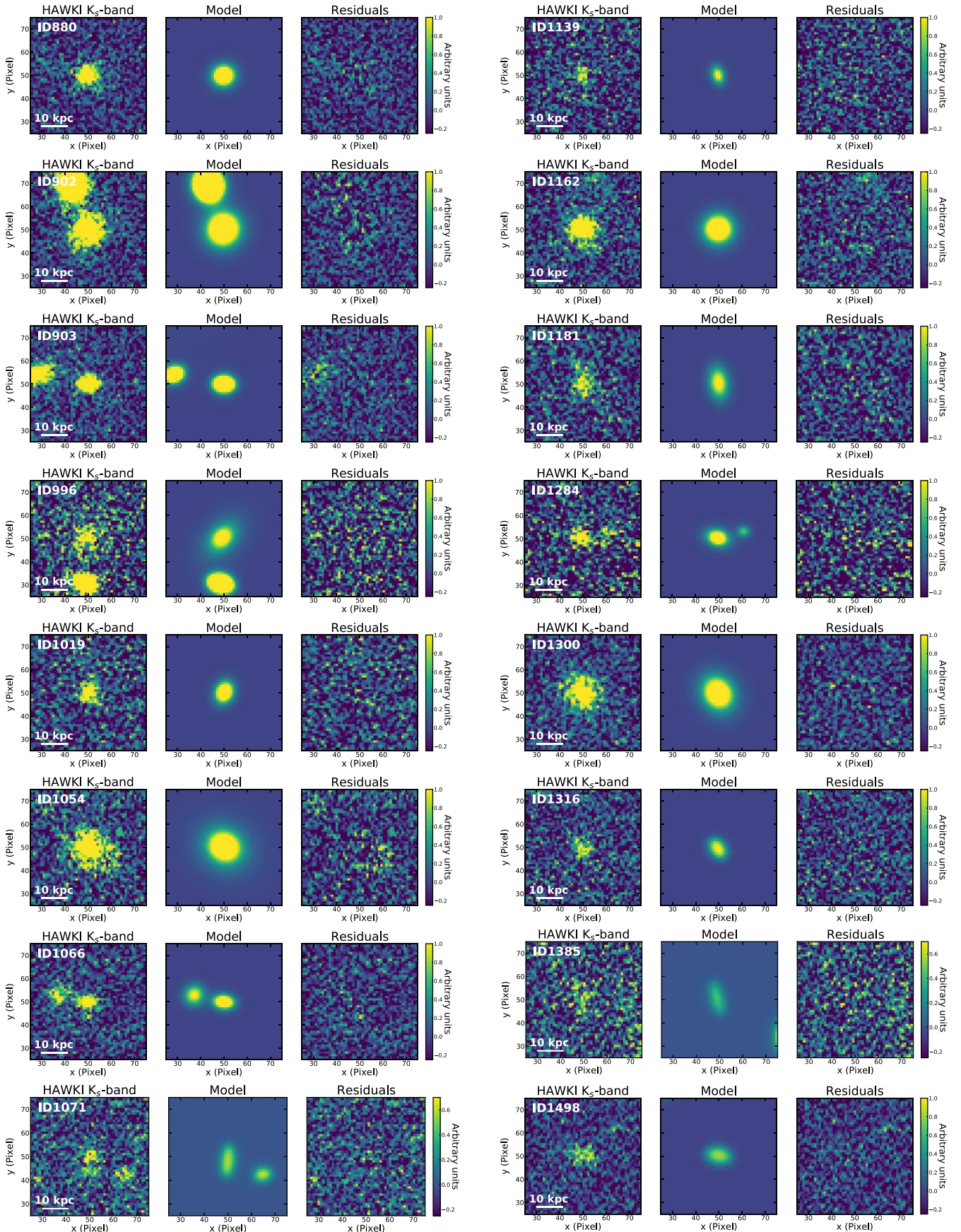


Figure A2 – continued

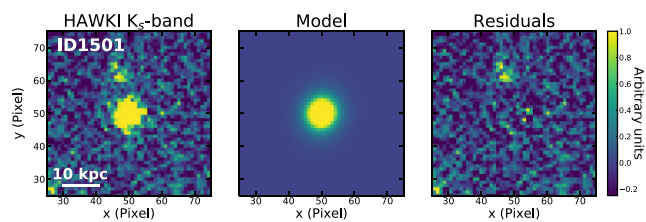


Figure A2 – *continued*

Table A1. Physical properties of the galaxies studied in this work. Coordinates (RA and Dec.) are shown in degrees. The redshifts (z) are based on the $H\alpha$ measurements of this work. The molecular masses have been computed following the metallicity-dependent calibration by Tacconi et al. (2018) and using the ATCA CO(1–0) information published by Jim et al. (2021). See Section 3.6 for more details. $E(B - V)_s$ represents the stellar contribution to the reddening assuming the Calzetti et al. (2000) extinction law (see Section 3.4). In the last two columns we highlight the presence of X-ray emission according to Tozzi et al. (2022) and the identification of sources as SMCs according to Dannebauer et al. (2014).

ID	RA (J2000)	Dec. (J2000)	z	$\log M_*/M_\odot$	SFR (M_\odot/yr)	$R_{c,Ks}$ (kpc)	$12 + \log(O/H)$	$\log M_{\text{mol}}/M_\odot$	$U - V$	$V - J$	$E(B - V)_s$	Σ_3 (Mpc^{-2})	η	X-ray	SMG
56	175.1411667	-26.5369722	2.1533	$10.83^{+0.11}_{-0.15}$	97 ± 31	-	8.78 ± 0.12	-	1.23	1.12	0.44	11.72	1.57	-	-
121	175.1466667	-26.5210278	2.1415	$10.37^{+0.09}_{-0.11}$	19 ± 4	2.4 ± 0.3	8.54 ± 0.10	-	1.35	0.68	0.12	60.56	6.84	-	-
204	175.1541250	-26.5514444	2.1444	$9.31^{+0.12}_{-0.16}$	12 ± 2	-	-	-	0.37	-0.16	0.02	13.55	7.09	-	-
210	175.1555833	-26.5048056	2.1684	$11.16^{+0.09}_{-0.11}$	101 ± 8	4.2 ± 0.7	8.65 ± 0.06	-	1.50	1.21	0.34	82.52	4.38	Yes	-
229	175.1573750	-26.4867500	2.1574	$10.98^{+0.11}_{-0.14}$	66 ± 18	3.9 ± 0.5	-	-	1.05	1.16	0.49	9.76	0.43	-	-
298	175.1642500	-26.5068333	2.1626	$10.42^{+0.14}_{-0.20}$	70 ± 13	-	8.52 ± 0.12	-	0.93	0.88	0.38	49.44	2.03	-	-
343	175.1670833	-26.4963889	2.1609	$10.11^{+0.12}_{-0.17}$	35 ± 6	-	8.66 ± 0.07	11.02 ± 0.12	0.87	0.68	0.31	21.98	1.24	-	-
457	175.1785833	-26.4686944	2.1622	$9.79^{+0.10}_{-0.13}$	27 ± 4	2.8 ± 0.7	8.44 ± 0.08	-	0.44	0.01	0.09	13.56	1.29	-	-
511	175.1843750	-26.4853056	2.1694	$11.29^{+0.09}_{-0.11}$	27 ± 13	1.7 ± 0.1	-	-	1.64	1.06	0.25	200.32	1.59	Yes	-
586	175.2637083	-26.4807778	2.1492	$10.17^{+0.11}_{-0.15}$	58 ± 19	-	<8.69	-	1.13	1.06	0.47	6.44	2.99	-	-
647	175.2599167	-26.4625278	2.1510	$11.59^{+0.10}_{-0.13}$	114 ± 29^a	2.4 ± 0.1	8.80 ± 0.07^a	-	1.45	1.28	0.42	25.63	2.25	Yes	Yes
779	175.2505833	-26.4823056	2.1665	$10.66^{+0.08}_{-0.09}$	106 ± 14	3.2 ± 0.2	8.48 ± 0.08	-	1.11	0.91	0.37	7.66	3.64	-	-
782	175.2492500	-26.5118333	2.1617	$9.75^{+0.11}_{-0.14}$	10 ± 3	1.9 ± 0.4	8.54 ± 0.15	11.22 ± 0.13	0.22	-0.30	0.00	86.45	2.24	-	Yes
783	175.2489583	-26.5523889	2.1730	$9.88^{+0.13}_{-0.19}$	16 ± 3	3.1 ± 0.3	8.64 ± 0.09	-	0.73	0.32	0.14	24.41	10.53	-	-
788	175.2485000	-26.5544722	2.1719	$9.67^{+0.15}_{-0.22}$	18 ± 3	-	-	-	0.47	-0.02	0.05	40.48	10.04	-	-
790	175.2484167	-26.5108611	2.1645	$10.74^{+0.10}_{-0.13}$	231 ± 38	3.4 ± 0.2	8.58 ± 0.08	11.27 ± 0.09	1.21	1.11	0.46	105.80	3.27	-	Yes
876	175.1952083	-26.4780833	2.1636	$8.88^{+0.15}_{-0.23}$	15 ± 2	-	-	-	0.34	-0.21	0.02	24.72	0.56	-	-
880	175.1944583	-26.4862222	2.1663	$10.90^{+0.10}_{-0.12}$	52 ± 11	1.9 ± 0.2	-	-	1.48	1.46	0.57	154.84	0.49	Yes	-
902	175.1919167	-26.4864722	2.1490	$11.37^{+0.10}_{-0.13}$	511 ± 110	4.1 ± 0.2	<8.78	11.14 ± 0.06	1.59	1.72	0.70	207.93	0.46	-	-
903	175.1921667	-26.4902222	2.1553	$10.35^{+0.11}_{-0.14}$	77 ± 7	2.0 ± 0.2	8.36 ± 0.06	-	0.74	0.46	0.25	240.56	0.05	-	-
116	175.1915833	-26.4880000	2.1557	$11.17^{+0.07}_{-0.09}$	125 ± 44^a	-	8.38 ± 0.08^a	-	0.69	0.63	0.25	450.88	0.02	Yes	-
996	175.2447500	-26.5062500	2.1657	$10.72^{+0.10}_{-0.14}$	56 ± 12	5.8 ± 1.6	8.74 ± 0.08	-	0.92	0.77	0.35	28.62	3.34	-	-

Table A1 – *continued*

ID	RA (J2000)	Dec. (J2000)	z	$\log M_*/M_\odot$	SFR (M_\odot/yr)	$R_{\text{e,Ks}}$ (kpc)	$12 + \log(\text{O}/\text{H})$	$\log M_{\text{mol}}/M_\odot$	$U - V$	$V - J$	$E(B - V)_s$	Σ_3 (Mpc^{-2})	η	X-ray	SMG
999	175.2465833	-26.4656389	2.1473	$9.88^{+0.08}_{-0.10}$	59 ± 8	-	8.61 ± 0.06	-	0.67	0.35	0.21	18.79	3.09	-	-
1019	175.2432917	-26.5566389	2.1635	$9.84^{+0.10}_{-0.13}$	3 ± 1	2.2 ± 0.2	-	-	0.24	-0.37	0.00	24.41	4.71	-	-
1047	175.2412917	-26.4934167	2.1703	$10.67^{+0.09}_{-0.12}$	84 ± 6	-	-	-	0.73	0.59	0.21	992.00	4.10	Yes	Yes
1054	175.2408750	-26.5133611	2.1644	$11.16^{+0.08}_{-0.09}$	189 ± 53	5.2 ± 0.3	8.66 ± 0.14	11.21 ± 0.09	1.45	1.51	0.63	20.82	2.95	-	Yes
1066	175.2390833	-26.4937500	2.1663	$10.30^{+0.11}_{-0.14}$	23 ± 6	2.0 ± 0.2	8.71 ± 0.10	-	0.84	0.57	0.23	810.69	2.81	-	Yes
1071	175.2379583	-26.4744167	2.1576	$10.19^{+0.14}_{-0.20}$	26 ± 20	2.4 ± 0.5	-	-	0.75	0.61	0.30	97.64	0.44	-	-
1139	175.2300833	-26.5119167	2.1447	$9.68^{+0.13}_{-0.18}$	37 ± 9	2.0 ± 0.4	<8.39	-	0.72	0.43	0.22	358.48	3.25	-	-
1154	175.2299167	-26.4783333	2.1630	$10.39^{+0.08}_{-0.10}$	44 ± 8	-	8.56 ± 0.12	-	1.13	0.73	0.24	34.18	1.47	-	-
1162	175.2272917	-26.4732500	2.1610	$10.73^{+0.09}_{-0.11}$	127 ± 24	3.4 ± 0.2	8.71 ± 0.08	10.76 ± 0.10	1.16	1.11	0.49	34.18	1.04	-	-
1181	175.2281250	-26.4676111	2.1520	$10.83^{+0.08}_{-0.10}$	166 ± 55	4.4 ± 0.3	-	-	1.69	1.69	0.63	28.39	0.95	-	Yes
1284	175.2191667	-26.5002500	2.1569	$10.21^{+0.13}_{-0.19}$	40 ± 14	3.2 ± 0.4	-	-	0.87	0.67	0.29	33.06	0.15	-	-
1300	175.2135833	-26.4940833	2.1531	$10.90^{+0.09}_{-0.11}$	149 ± 31	4.7 ± 0.2	8.62 ± 0.06	11.19 ± 0.10	1.19	1.11	0.46	168.27	0.31	-	-
1316	175.2148333	-26.4960833	2.1575	$10.08^{+0.02}_{-0.02}$	11 ± 1	2.8 ± 0.4	8.55 ± 0.10	-	1.10	0.24	0.00	49.14	0.19	-	-
1385	175.2090833	-26.4891389	2.1553	$9.65^{+0.09}_{-0.11}$	24 ± 6	4.5 ± 0.7	-	-	0.73	0.36	0.19	54.99	0.04	-	-
1420	175.2057500	-26.4858889	2.1661	$9.34^{+0.15}_{-0.23}$	14 ± 3	-	-	-	0.42	-0.05	0.06	73.15	0.32	-	-
1498	175.1998750	-26.4804722	2.1628	$10.12^{+0.12}_{-0.17}$	51 ± 13	3.3 ± 0.3	-	-	1.02	0.78	0.28	46.86	0.29	-	-
1501	175.1997500	-26.4850833	2.1568	$11.00^{+0.10}_{-0.13}$	245 ± 94	2.2 ± 0.5	-	-	1.40	1.49	0.63	371.92	0.01	Yes	-

Note. ^aThese objects display broad H α profiles ($\sigma > 700 \text{ km s}^{-1}$). SFRs and gas phase metallicities have been computed using only the narrow-line component of their emission lines.

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