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# Periodic signals from the Circinus region: two new cataclysmic variables and the ultraluminous X-ray source candidate GC X–1

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## ABSTRACT

The examination of two 2010 *Chandra* ACIS exposures of the Circinus galaxy resulted in the discovery of two pulsators: CXO J141430.1–651621 and CXOU J141332.9–651756. We also detected 26-ks pulsations in CG X–1, consistently with previous measures. For  $\sim 40$  other sources, we obtained limits on periodic modulations. In CXO J141430.1–651621, which is  $\sim 2$  arcmin outside the Circinus galaxy, we detected signals at  $6120 \pm 1$  s and  $64.2 \pm 0.5$  ks. In the longest observation, the source showed a flux of  $\approx 1.1 \times 10^{-13}$  erg cm<sup>-2</sup> s<sup>-1</sup> (absorbed, 0.5–10 keV) and the spectrum could be described by a power-law with photon index  $\Gamma \simeq 1.4$ . From archival observations, we found that the luminosity is variable by  $\approx 50$  per cent on time-scales of weeks–years. The two periodicities pin down CXO J141430.1–651621 as a cataclysmic variable of the intermediate polar subtype. The period of CXOU J141332.9–651756 is  $6378 \pm 3$  s. It is located inside the Circinus galaxy, but the low absorption indicates a Galactic foreground object. The flux was  $\approx 5 \times 10^{-14}$  erg cm<sup>-2</sup> s<sup>-1</sup> in the *Chandra* observations and showed  $\approx 50$  per cent variations on weekly/yearly scales; the spectrum is well fit by a power law with  $\Gamma \simeq 0.9$ . These characteristics and the large modulation suggest that CXOU J141332.9–651756 is a magnetic cataclysmic variable, probably a polar. For CG X–1, we show that if the source is in the Circinus galaxy, its properties are consistent with a Wolf–Rayet plus black hole binary. We consider the implications of this for ultraluminous X-ray sources and the prospects of Advanced LIGO and Virgo. In particular, from the current sample of WR–BH systems we estimate an upper limit to the detection rate of stellar BH–BH mergers of  $\sim 16$  yr<sup>-1</sup>.

**Key words:** galaxies: individual: Circinus – novae, cataclysmic variables – X-rays: binaries – X-rays: individual: CXOU J141332.9–651756 – X-rays: individual: CXO J141430.1–651621 – X-rays: individual: CG X–1 (CXOU J141312.3–652013)

## 1 INTRODUCTION

The *Chandra* ACIS Timing Survey at Brera And Rome astronomical observatories project (CATS @ BAR; Israel et al., in preparation) is a Fourier-transform-based systematic search for new pulsating sources in the *Chandra* Advanced CCD Imaging Spectrometer (ACIS; Garmire et al. 2003) public archive. As of 2015 April 30, 10,282 ACIS Timed Exposure observations have been examined and  $\sim 457,000$  sources were detected. Data taken with gratings or in Continuous-Clocking mode were not considered. The  $\sim 93,600$  light curves of sources with more than 150 photons were searched

for coherent signals with an algorithm based on that of Israel & Stella (1996). The limit of 150 counts is related to the intrinsic ability of the Fourier transform to detect a signal with 100 per cent modulation at a minimum confidence level of  $3.5\sigma$  in  $10^5$ – $10^6$  trials. CATS @ BAR has so far discovered 43 new certain X-ray pulsators; see Esposito et al. (2013b,c,a) for the first results and Esposito et al. (2014, 2015) for our analogous *Swift* project.

In this paper, we report on the CATS @ BAR results for the galaxy ESO 97–G13 (the ‘Circinus galaxy’, hereafter CG; Freeman et al. 1977) and its surroundings in the Circinus constellation. CG is a nearby Seyfert II active galaxy that lies close to the plane of our own Galaxy (J2000 Galactic coordinates:  $l = 311^\circ.3$ ,  $b = -03^\circ.8$ ; distance  $d = 4.2$  Mpc; Tully et al. 2009). The CG con-

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tains hydrogen-rich star forming regions in the inner spiral arms and, due to its closeness, offers a good opportunity to study its population of X-ray sources (Bauer et al. 2001; Sambruna et al. 2001), which includes several ultraluminous X-ray sources (ULXs, Winter et al. 2006; Mapelli et al. 2010b; Walton et al. 2013; see Mushotzky 2004; Fabbiano 2006; Zampieri & Roberts 2009; Feng & Soria 2011 for reviews on ULXs).

The CG was observed several times with *Chandra*, but it was in two long Timed Exposure observations carried out in late 2010 that CATS @ BAR pinpointed two new X-ray pulsators in the Circinus region: CXO J141430 and the uncatalogued CXOU J141332. The pipeline detected also CG X–1 (CXOU J141312.3–652013), whose emission is modulated at  $\sim 7$  h. CG X–1 is not a new pulsator, but the long-standing debate about its nature (Smith & Wilson 2001; Bauer et al. 2001; Weisskopf et al. 2004) prompted us to include it in our study.

The plan of the paper is as follows. In Section 2 we give details on the X-ray observations used in our study. The rest of the paper is divided into two main parts. The first one focuses on the new pulsators and comprises Sections 3 to 7. In Section 3 we describe the timing analysis that led to the discovery of the new pulsators and allowed us also to set upper limits on the pulsations for dozens of other X-ray sources. The detailed study of CXO J141430 is presented in Section 4, that of CXOU J141332 in Section 5. To study these sources, we used also data from *XMM-Newton* and optical observations taken with the VLT Survey Telescope (VST). The optical observations and their analysis are described in Section 6. The nature of CXO J141430 and CXOU J141332 is discussed in Section 7.

The second part of the paper is dedicated to CG X–1. In Section 8 we recall the main facts about this source. The results from the analysis of the 2010 *Chandra* data, which were not used before to study CG X–1, are summarised in Section 9. In Section 10, we propose that CG X–1 might be a Wolf–Rayet plus black hole (WR–BH) binary system, and consider the implications of this possibility for ULXs and for the prospects of detection of gravitational radiation from BH–BH mergers. A summary with conclusive remarks follows in Section 11.

## 2 X-RAY OBSERVATIONS

All the observations used in this work are summarised in Table 1. The most important observations are the one in which CATS @ BAR detected the new pulsators and its companion, 12823 and 12824, marked with a hash mark in Table 1. They were carried out in a week in 2010 December to study the central region of the CG (Mingo et al. 2012). Their combined exposure is  $\sim 190$  ks. In both cases, three ACIS-S and two ACIS-I CCDs were used in full frame mode, ensuring a wide coverage over the Circinus region. The *Chandra* data were processed and analysed with the *Chandra* Interactive Analysis of Observations software package (CIAO, version 4.7; Fruscione et al. 2006) and the calibration files in CALDB version 4.6.7. The Circinus field as imaged with *Chandra* in observation 12823 is shown in Fig. 1. In the data sets 12823 and 12824, CG X–1 and CXOU J141332 were positioned in the back-illuminated CCD 7 (S3). The photons from these sources were accumulated within a circle with radius 1.5 arcsec and an ellipse with semi-axes 3.5 and 3 arcsec, respectively. CXO J141430 fell on the front-illuminated CCD 3 (I3) and the source counts were extracted from an ellipse with semi-axes of 15 and 14 arcsec. The choice of regions of different size is due to the point-spread function at the

off-axis angles of the sources. For each source, the background was estimated locally, using source-free regions as close as possible to the target. The Solar system barycentre correction to the photon arrival times was applied with AXBARY. The spectra, the redistribution matrices, and the ancillary response files were created using SPEXTRACT.

The second most useful observation for our study is that performed with *XMM-Newton* in 2001 August with a duration of  $\sim 100$  ks (obs. ID 0111240101; Molendi, Bianchi, & Matt 2003; it is marked by a star in Table 1). We used the data collected with the European Photon Imaging Camera (EPIC), which consists of two MOS (Turner et al. 2001) and one pn (Strüder et al. 2001) CCD detectors. The raw data were reprocessed using the *XMM-Newton* Science Analysis Software (SAS, version 14.0) and the calibration files in the CCF release of 2015 March. The observation suffered of intense soft-proton flares. The intervals of flaring background were located by intensity filters (see e.g. De Luca & Molendi 2004) and excluded from the analysis. This reduced the net exposure time by  $\sim 30$  per cent in the pn back-illuminated CCDs and  $\sim 10$  per cent in the MOS front-illuminated CCDs. The source photons were extracted from circles with radius of 25 arcsec for CXO J141430 and 15 arcsec for CXOU J141332 (these radii were essentially imposed by CCD gaps and/or the presence of neighbouring sources) and the backgrounds from regions in the same chip as the sources. CXOU J141332 was in the unread part of the central CCD of the MOS 2 operated in a partial window mode, so only pn and MOS 1 data exist for it. Photon arrival times were converted to the Solar system barycentre using the SAS task BARYCEN. The ancillary response files and the spectral redistribution matrices for the spectral analysis were generated with ARFGEN and RMFGEN, respectively. Due to the low number of photons, we combined for each source the spectra from the available EPIC cameras and averaged the response files using EPICSPECCOMBINE.

We made use of other *Chandra* and *XMM-Newton* observations, which were reduced and analysed in a similar way; they provided only detections and flux estimates, or upper limits for the two new pulsators. Six data sets were collected with *Chandra* from 2000 to 2009 with various instrumental setups and durations from  $\sim 1$  to 25 ks, and one with *XMM-Newton* in 2014 with exposure of  $\sim 30$  ks. Apart from these observations, listed in Table 1, no other *Chandra* pointing of the CG covered the positions of CXO J141430 or CXOU J141332, while in a  $\sim 60$  ks observations performed with *XMM-Newton* in 2013 (obs. ID 0701981001), both sources fell either in gaps or at the edge of CCDs, or outside the field of view of the instruments.

## 3 CATS @ BAR TIMING ANALYSIS

The CIAO WAVEDETECT routine detected 156 sources in the ACIS field of view of observation 12823; they are marked by circles in Fig. 1. The 44 sources with more than 150 photons (those marked by magenta and red circles in Fig. 1) were searched for periodic signals. The CATS @ BAR search algorithm is based on a fast Fourier transform and takes into account also the possible presence of additional non-Poissonian noise components in the Leahy-normalised (Leahy et al. 1983) power spectra (see Israel & Stella 1996 for more details). Correspondingly, the CATS @ BAR signal threshold in a power spectrum, which takes into account the number of independent Fourier frequencies, is a function also of the local underlying noise. For the Circinus data, the maximum frequency of the search ( $\sim 0.16$  Hz) is dictated by the sampling time of 3.14 s, while the

**Table 1.** Summary of the X-ray observations used in this work. The hash marks and the star indicate the *Chandra* and *XMM-Newton* observations from which most of the information was obtained.

Satellite	Instrument	Obs. ID	Date	Exp. (ks)	Mode <sup>a</sup>
<i>Chandra</i>	ACIS-23678	355	2000 Jan 16	1.3	TE FAINT (3.24 s)
<i>Chandra</i>	ACIS-235678	356	2000 Mar 14	24.7	TE FAINT (3.24 s)
<i>Chandra</i>	ACIS-235678	2454	2001 May 02	4.4	TE FAINT (3.24 s)
<i>XMM</i> *	pn / MOS 1 / MOS 2	0111240101	2001 Aug 6–7	100.6 / 104.1 / 104.0	FF (73.4 ms) / FF (2.6 s) / LW (0.9/2.7 s)
<i>Chandra</i>	ACIS-456789	10873	2009 Mar 01	18.1	TE HETG VFAINT (2.04 s)
<i>Chandra</i>	ACIS-456789	10850	2009 Mar 03	13.8	TE HETG VFAINT (2.04 s)
<i>Chandra</i>	ACIS-456789	10872	2009 Mar 04	16.5	TE HETG VFAINT (2.04 s)
<i>Chandra</i> #	ACIS-23678	12823	2010 Dec 17–19	154.4	TE VFAINT (3.14 s)
<i>Chandra</i> #	ACIS-23678	12824	2010 Dec 24	39.4	TE VFAINT (3.14 s)
<i>XMM</i>	pn / MOS 1 / MOS 2	0656580601	2014 Mar 01	32.7 / 37.8 / 37.9	FF (73.4 ms) / FF (2.6 s) / LW (2.7 s)

<sup>a</sup> TE: Timed Exposure, HETG: High Energy Transmission Grating, VFAINT: Very Faint telemetry format, FF: Full Frame, LW: Large Window; the readout time is given in parentheses, for the central and peripheral CCDs in the case of the MOS 2 in LW.

longest period to which the observation is realistically sensitive (because of its duration) is  $\approx 80$  ks; 32 768 frequencies were searched. The search resulted in the detection of five sources with significant signals in their power spectra.

In two cases the power peaks were coincident with the frequencies of known spurious signals due to the spacecraft dithering pattern. The CATS @ BAR pipeline automatically performs check for these artificial signals by means of the DITHER\_REGION CIAO task.<sup>1</sup> Furthermore, every candidate signal is crosschecked with the CATS @ BAR database of recurring signals of instrumental origin, and repeating or dubious signals are carefully inspected and rejected. These two sources are the supernova remnant candidate CG X–2 (Bauer et al. 2001), labelled 7–2 in Fig. 1, and Circinus XMM2, which is classified as an ULX (Winter et al. 2006), labelled 7–17. In a third object, CG X–1, the detected  $\sim 27$  ks-period modulation was already known (Bauer et al. 2001; object 7–3 in Fig. 1). This source is discussed in detail in Sections 8 to 10. The remaining two sources, 3–1 = CXO J141430 (Section 4) and the uncatalogued 7–20 = CXOU J141332 (Section 5), are genuine new X-ray pulsators, as it was confirmed also by the other data sets.

For all the other sources with more than 150 events, a  $3\sigma$  upper limit to the pulsed fraction of any sinusoidal signal was calculated (throughout the paper, we will give  $3\sigma$  upper limits on non-detections and limits at 90 per cent confidence level on poorly constrained quantities; all uncertainties will be given at  $1\sigma$  confidence level). The pulsed fraction was defined as the semi-amplitude of the sinusoidal modulation divided by the mean count rate. As expected, around 100–200 photons the upper limits start crossing the 100 per cent threshold, above which no meaningful information related to any coherent signal can be inferred. For many other sources, the upper limits are not constraining. For future reference, all the results are summarised in Fig. 2.

## 4 THE 1.7 / 17.8 h PULSATOR: CXO J141430

### 4.1 Timing analysis

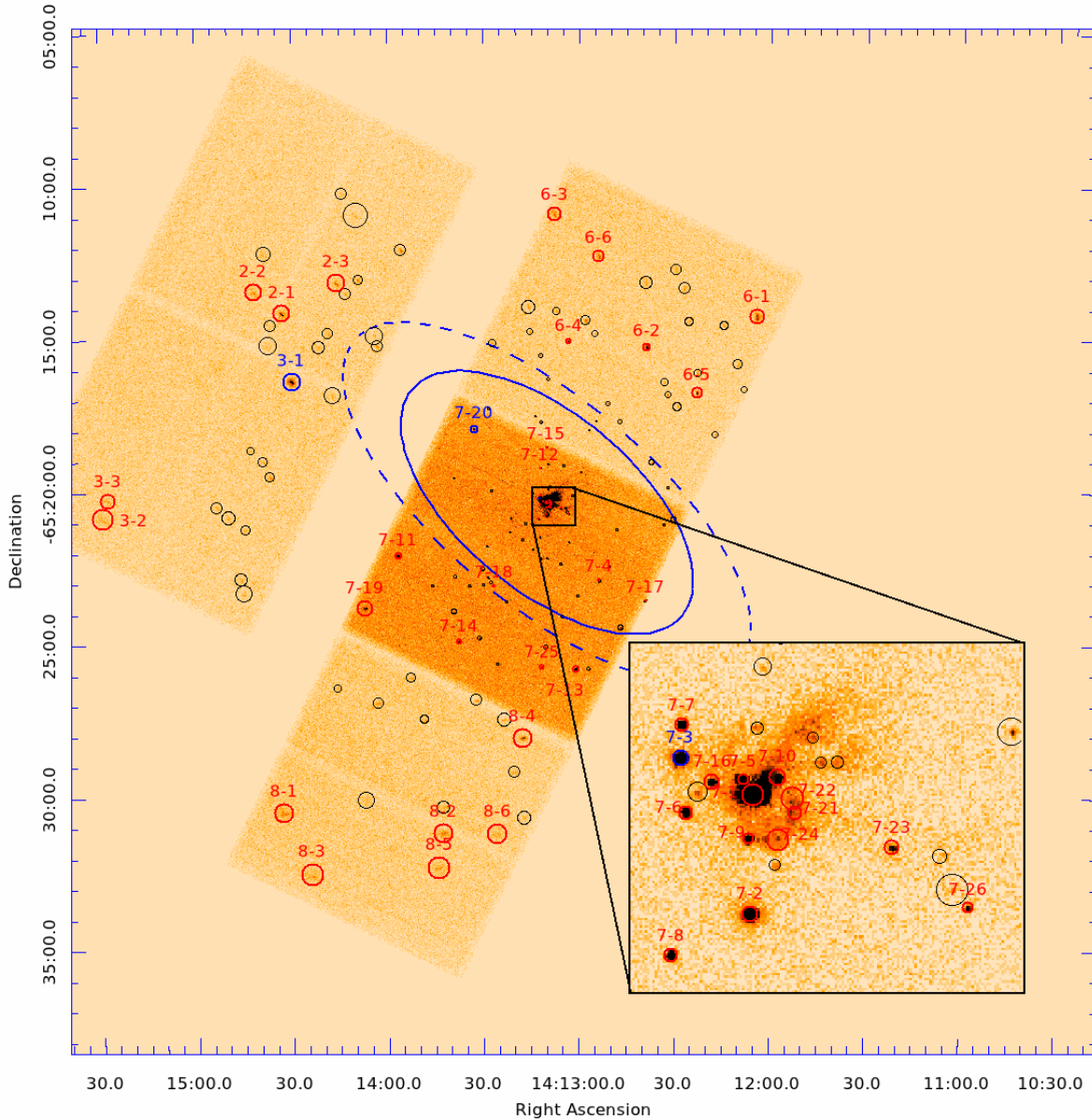
CXO J141430 is the brightest of the two new CATS @ BAR pulsators. It shows two distinct periodic signals: a  $\sim 100$  per cent modulation at about 6.1 ks and another large-amplitude signal at

about 64 ks. The power spectrum is shown in Fig. 3. When the 32 768 frequencies analysed and the number of sources for which the search was carried out (44) are taken into account, both signals were detected at a confidence level larger than  $10\sigma$ .

To hone the estimates of the periods, we made use of both the *Chandra* 12823 and 12824 pointings, where about 820 and 250 net counts, respectively, were collected. For the short signal, we used a phase-fitting technique and found  $P_1 = 6.120 \pm 2$  s. For the 64 ks period, the number of sampled cycles, approximately three, is too small for the phase fitting. We therefore binned the light curve to 6 120 s, so to avoid beat-frequency signals produced by the shorter periodicity, and fit a sinusoidal function to it. The fit has a  $\chi^2$  of 54 for 30 degrees of freedom (dof) and we derived the period  $P_2 = 64.2 \pm 0.5$  ks. The 0.5–10 keV background-subtracted light curves folded on our best periods are shown in Fig. 3. We measured the following pulsed fractions:  $100 \pm 4$  per cent ( $P_1 = 6.1$  ks; this value is to be regarded as a lower limit) and  $70 \pm 4$  per cent ( $P_2 = 64.2$  ks). Within the statistical uncertainties, the shape and the pulsed fraction of both signals are energy-independent. In the soft ( $< 2$  keV) and hard ( $> 2$  keV) bands we measured pulsed fractions of  $96 \pm 5$  per cent and  $106 \pm 5$  per cent for the 6.1 ks period, and  $69 \pm 5$  per cent and  $64 \pm 7$  per cent for the 64.2 ks period.

The 6.1-ks signal is significantly detected also in the 2001 *XMM-Newton*/EPIC data ( $\sim 700$  net counts between the three EPIC cameras), while the observation is too short for the 64.2-ks period (Fig. 3). We measured the period  $P_1 = 6.04 \pm 0.04$  ks and a pulsed fraction of  $88 \pm 12$  per cent. CXO J141430 is detected with  $\sim 170$  net counts in the 2014 *XMM-Newton* pointing (pn plus MOS 2, in the MOS 1 the source fell in one of the failed CCDs). The short-period pulsations are clear also in that data set, but the low count statistics hampers a precise estimate of the period. Finally, CXO J141430 was in the field of view of *Chandra* also in the observations 355 (2000 January, 1.3-ks) and 356 (2000 March 25-ks; Table 1). In observation 355, CXO J141430 is detected with a dozen of photons only, and a signal-to-noise ratio  $\text{SNR} \gtrsim 3$ : the short duration and the very small number of photons preclude any analysis of the periodic signals. In the data set 356, the source is detected with about 90 photons ( $\text{SNR} > 9$ ). The 6.1-ks signal can be clearly observed but, similarly than in the 2014 *XMM-Newton* observations, the uncertainty on the period is very large.

<sup>1</sup> See [http://cxc.harvard.edu/ciao/ahelp/dither\\_region.html](http://cxc.harvard.edu/ciao/ahelp/dither_region.html).

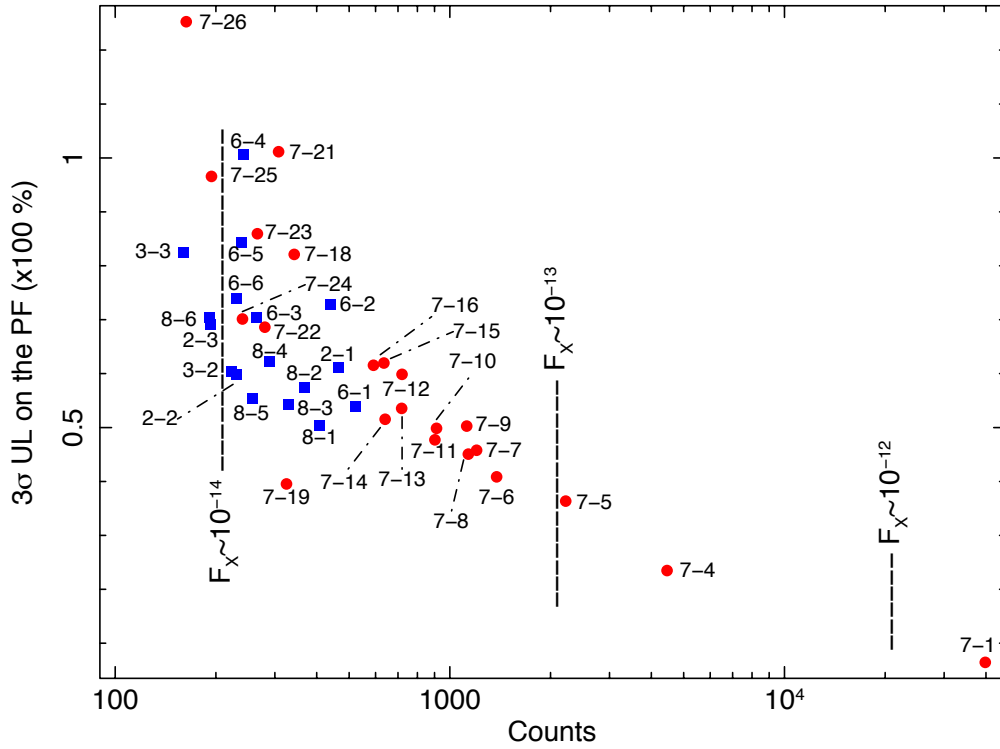


**Figure 1.** *Chandra*/ACIS image of the Circinus field. North is top and east is left; each CCD subtends an  $8'.4 \times 8'.4$  square on the sky and the zoomed area is  $1'.2 \times 1'$ . Circles mark the detected sources. Red circles indicate sources with  $>150$  counts for which we performed a timing analysis, blue circles the sources from which periodic signals were detected. The labelled sources (with the ACIS CCD number and an identification number) are those for which either signals were detected (source 7–3 is CG X–1, 7–20 is CXOU J141332, and 3–1 is CXO J141430) or an upper limit on the pulsed fraction could be placed (see Fig. 2). The solid and dashed ellipses indicate the size of the CG from 90 per cent total  $B$  light and total infrared (2MASS) magnitude, respectively (from the NASA/IPAC Extragalactic Database, NED, see <http://ned.ipac.caltech.edu/>). Source 7–1 is the CG’s active galactic nucleus (AGN), 7–2 is the young supernova remnant candidate in the CG (CG X–2; Bauer et al. 2001); other notable sources are the ultraluminous X-ray sources 7–5 = Circinus ULX3, 7–16 = Circinus ULX4, 7–4 = Circinus XMM1 = ULX5, 7–17 = Circinus XMM2 (while Circinus XMM3 is undetected; we used the nomenclature of Winter et al. 2006; Gladstone et al. 2013; Walton et al. 2013).

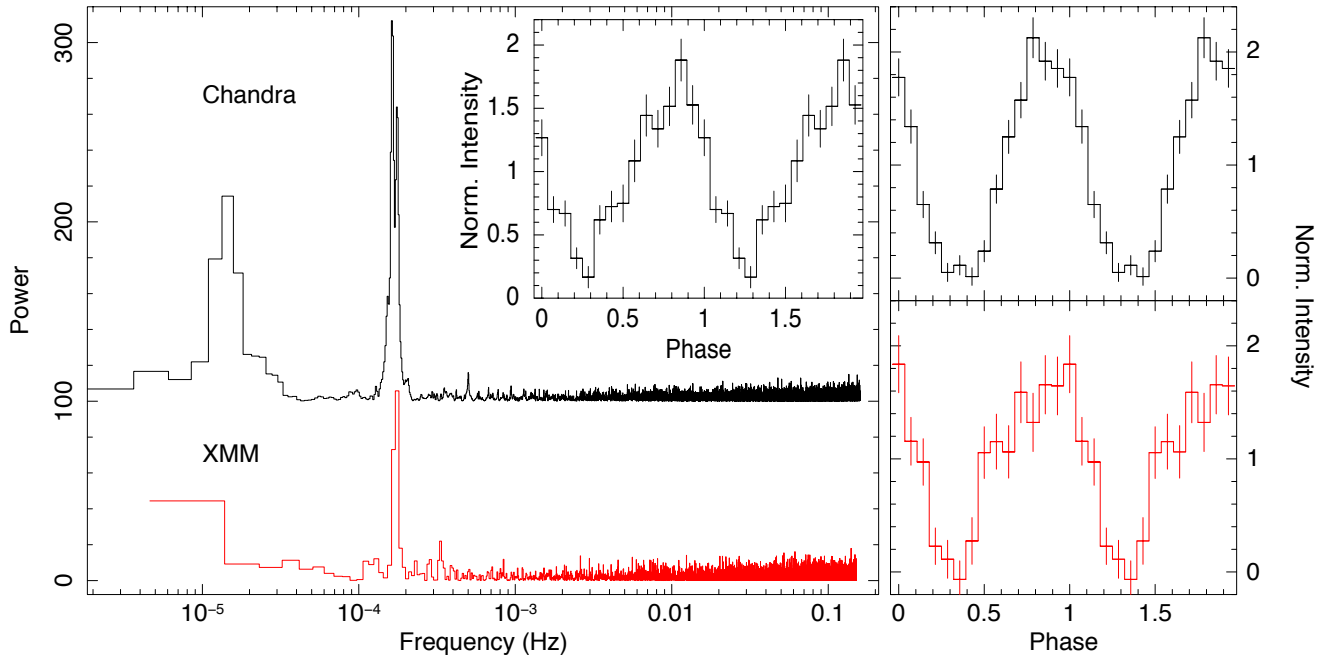
## 4.2 Spectral analysis

For the spectral analysis, we started from the long *Chandra* observation 12823. The fits were performed between 0.6 and 6 keV because of the very low signal of CXO J141430 outside this range. We fit to the data a power law model, a blackbody, and an optically-thin thermal bremsstrahlung. The blackbody model yielded a re-

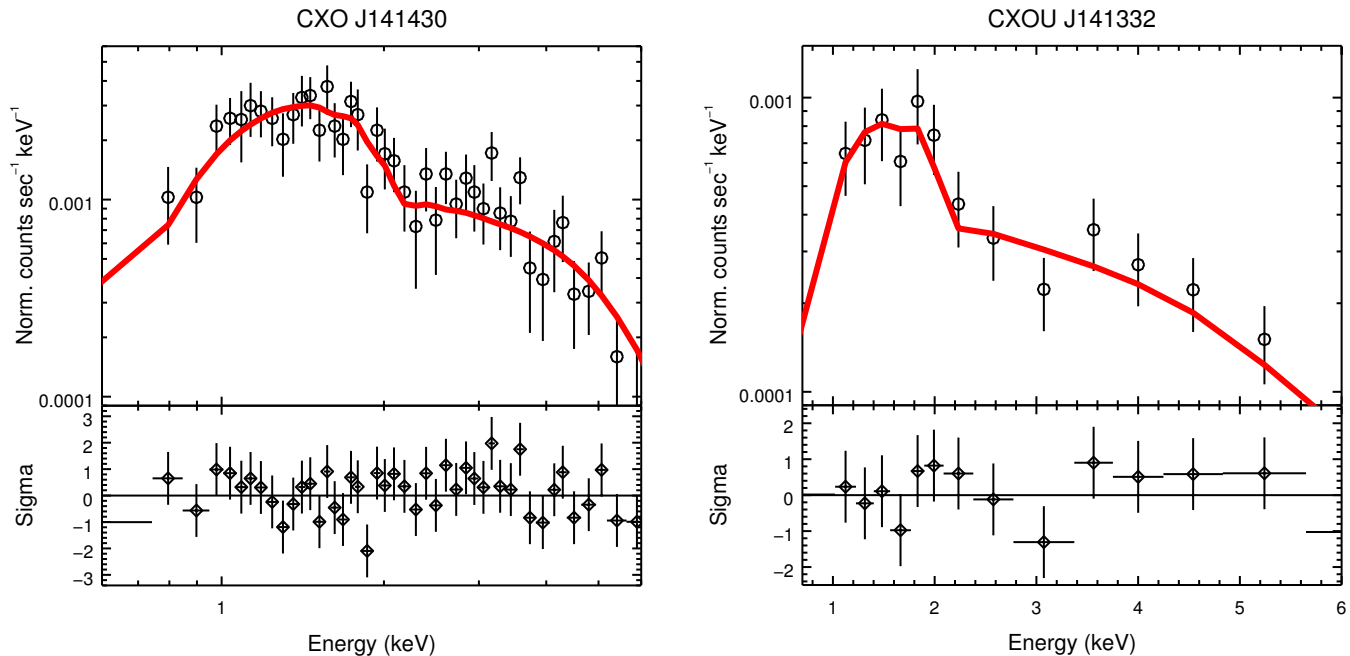
duced  $\chi^2_\nu = 1.37$  for 40 dof with clearly structured residuals; the derived temperature is  $kT = 0.81 \pm 0.03$  keV, while for the absorption there is only an upper limit of  $N_{\text{H}} < 0.8 \times 10^{22}$   $\text{cm}^{-2}$  at 90 per cent confidence. The observed flux was  $F_{\text{X}} = 6.8^{+0.5}_{-0.4} \times 10^{-14}$   $\text{erg cm}^{-2} \text{s}^{-1}$  (0.5–10 keV). The power law and the bremsstrahlung gave somewhat better fits,  $\chi^2_\nu = 1.27$  and 1.24, respectively, and better residuals. The parameters of the



**Figure 2.** Upper limits on the pulsed fraction of the Circinus sources. Red circles indicate sources detected in the back-illuminated ACIS-7 (S3) CCD, blue squares sources from front-illuminated CCDs. The sources are identified by the same labels as in Fig 1. Dashed vertical lines mark approximate fluxes for the S3 CCD sources. Source 7-1 is the CG’s AGN; other notable sources are: 7-5 = Circinus ULX3, 7-16 = Circinus ULX4, and 7-4 = Circinus XMM1 = ULX5.



**Figure 3.** Left panel: 0.5–10 keV power spectra of CXO J141430 obtained from *Chandra*/ACIS (top, black line, observations 12823 and 12824 combined) and *XMM-Newton*/EPIC (bottom, red line) data. For displaying purposes, the *Chandra* data have been shifted in power by +100. The corresponding *Chandra* (top) and *XMM-Newton* (bottom) light curves folded to the best periods are shown in the panels on the right for  $P_1 = 6.1$  ks. In the inset is shown the *Chandra* light curve folded to the longer second period  $P_2 = 64.2$  ks.



**Figure 4.** *Chandra*/ACIS spectra and best-fitting power-law models (red solid line) for CXO J141430 and CXOU J141332 (as indicated on each panel) from observation 12823. Bottom panels: residuals in units of standard deviations.

power law fit (Fig. 4, left) are  $N_{\text{H}} = 0.36^{+0.14}_{-0.13} \times 10^{22} \text{ cm}^{-2}$ ,  $\Gamma = 1.51^{+0.16}_{-0.15}$ , and  $F_{\text{X}} = (1.03 \pm 0.09) \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$ . For the bremsstrahlung,  $N_{\text{H}} = 0.31^{+0.11}_{-0.10} \times 10^{22} \text{ cm}^{-2}$ ,  $kT = 13^{+12}_{-5} \text{ keV}$ , and  $F_{\text{X}} = 0.97^{+0.09}_{-0.10} \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$ .

In the second *Chandra* observation, the flux was  $\approx 30$  per cent higher. The absorption was only poorly constrained ( $< 0.6 \times 10^{22} \text{ cm}^{-2}$  at 90 per cent confidence for the power law and  $< 0.5 \times 10^{22} \text{ cm}^{-2}$  for the bremsstrahlung), while  $kT$  and  $\Gamma$  were consistent with those measured in observation 12823. We thus decided to fit the two spectra simultaneously, with the normalisations free to vary and the other parameters tied up between the data sets. The results are summarised in Table 2.

The 2001 *XMM-Newton* data flatly reject the blackbody model, with  $\chi^2_{\nu} = 2.33$  for 28 dof. The power-law provides a good fit to the data, with an observed flux similar to that of the first *Chandra* observation. The bremsstrahlung fit was equally good. Its temperature, however, could not be constrained, as it always pegged to the highest allowed value, showing that in the *XMM-Newton* data its curvature is indistinguishable from that of a power law. For this reason, in Table 2 we give only the parameters derived from the power law fit. The quality of the spectrum from the 2014 *XMM-Newton* data is too poor for a spectral analysis. We thus used the models in Table 2 to estimate the flux of CXO J141430, obtaining  $F_{\text{X}} = (4.8 \pm 0.6) \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$  for the power-law model and  $(3.4 \pm 0.5) \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$  for the bremsstrahlung. Similarly, for the 2000 *Chandra* observation 356 ( $\sim 90$  counts in a front-illuminated ACIS-S CCDs) we evaluated a flux  $F_{\text{X}} = (1.1 \pm 0.2) \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$  for the power-law model and  $(8.0 \pm 0.9) \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$  for the bremsstrahlung. In the short pointing 355, where only a dozen of photons were detected in one of the ACIS-I CCDs, we converted the count rate into a flux with PIMMS<sup>2</sup> and found  $F_{\text{X}} = (2.0 \pm 0.7) \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$

for the power-law model and  $(1.5 \pm 0.5) \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$  for the bremsstrahlung.

## 5 THE 1.8 h PULSATOR: CXOU J141332

### 5.1 Timing analysis

CXOU J141332 displays a periodicity at roughly 6.4 ks (Fig. 5). This signal was detected at about  $3.5\sigma$  confidence level in observation 12823 ( $\sim 270$  counts) and is present also in the shorter pointing 12824 ( $\sim 80$  counts), though at a lower confidence level. By the phase-fitting analysis, we derived the value  $P = 6377 \pm 4 \text{ s}$ . In the corresponding 0.5–10 keV background-subtracted folded profile (Fig. 5), we measured a pulsed fraction of  $56 \pm 8$  per cent. The modulation of CXOU J141332 changes in shape and pulsed fraction as a function of energy: the asymmetry in the profile becomes more accentuated, while the pulsed fraction decreases as the energy increases:  $78 \pm 7$  per cent in the soft range ( $< 2 \text{ keV}$ ) and  $40 \pm 8$  per cent in the hard range ( $> 2 \text{ keV}$ ). Of all the other observations in Table 1, CXOU J141332 was detected only in the 2001 *XMM-Newton* data, where the harvest was  $\sim 440$  source photons between the pn and MOS 1. We measured a period of  $6.4 \pm 0.1 \text{ ks}$  and the pulsed fraction was  $59 \pm 10$  per cent.

### 5.2 Spectral analysis

Given the paucity of photons (all in the 1–6 keV range), we performed the spectral analysis by fitting the same spectral models adopted for CXO J141430 to the two *Chandra* data sets 12823 and 12824 simultaneously, with the normalisations free to vary and the other parameters tied up. The blackbody model provides

<sup>2</sup> We used the web version of PIMMS (Portable, Interactive Multi-Mission

Simulator) available at <http://heasarc.gsfc.nasa.gov/cgi-bin/Tools/w3pimms/w3pimms.pl>.

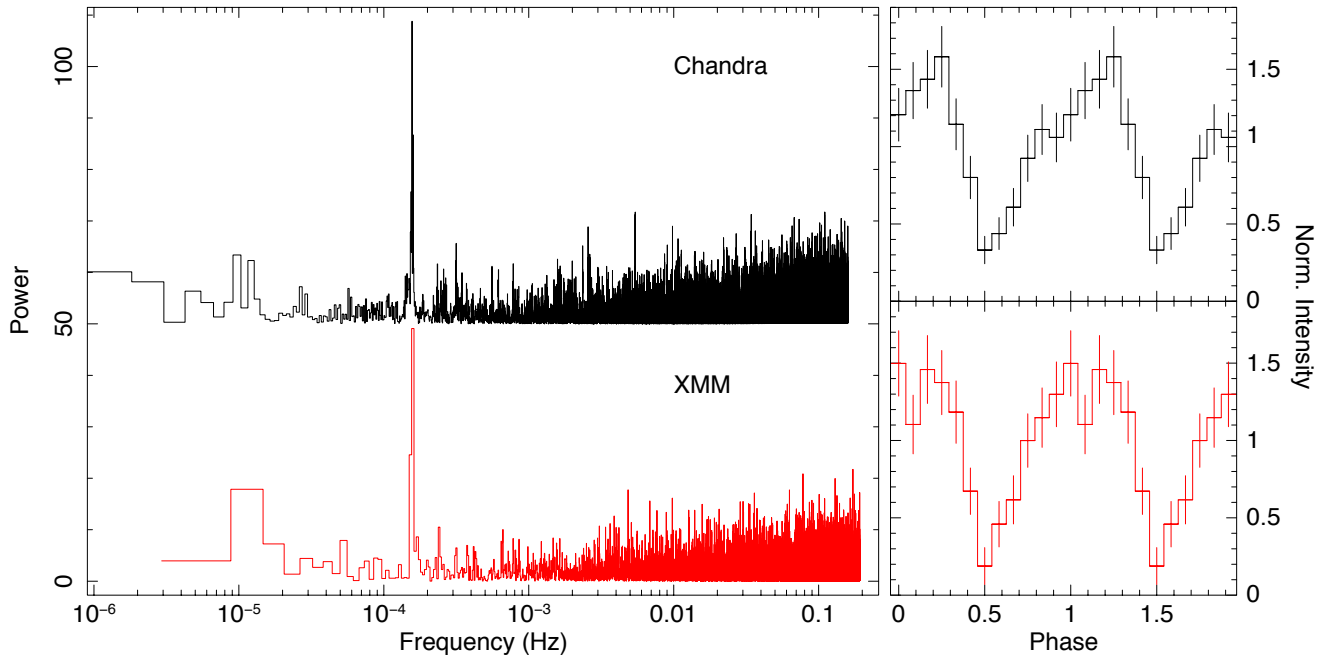
**Table 2.** Spectral results of CXO J141430. Errors are at a  $1\sigma$  confidence level for a single parameter of interest.

Model	Data	$N_{\text{H}}^a$ ( $10^{22} \text{ cm}^{-2}$ )	$\Gamma$	$kT$ (keV)	Flux <sup>b</sup> ( $10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$ )	Unabsorbed flux <sup>b</sup> ( $10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$ )	$\chi^2_{\nu}$ (dof)
PHABS(POWERLAW)	<i>Chandra</i> /12823	$0.30^{+0.13}_{-0.11}$	$1.43^{+0.14}_{-0.13}$	–	$1.07 \pm 0.08$	$1.21 \pm 0.07$	1.12 (52)
	<i>Chandra</i> /12824				$1.40^{+0.15}_{-0.14}$	$1.58^{+0.13}_{-0.12}$	
PHABS(BREMSSTRAHLUNG)	<i>Chandra</i> /12823	$0.27 \pm 0.10$	–	$18^{+24}_{-7}$	$1.01^{+0.10}_{-0.09}$	$1.13 \pm 0.08$	1.10 (52)
	<i>Chandra</i> /12824				$1.32^{+0.16}_{-0.15}$	$1.47 \pm 0.14$	
PHABS(POWERLAW)	<i>XMM</i> /0111240101	$<0.3^c$	$1.10 \pm 0.15$	–	$0.87 \pm 0.08$	$0.89^{+0.07}_{-0.06}$	1.02 (28)

<sup>a</sup> The abundances used are those of Wilms, Allen, & McCray (2000);  $N_{\text{H}}$  values  $\approx 30$  per cent lower are derived with those by Anders & Grevesse (1989). The photoelectric absorption cross-sections are from Balucinska-Church & McCammon (1992).

<sup>b</sup> In the 0.5–10 keV energy range.

<sup>c</sup> Upper limit at 90 per cent confidence level.


**Figure 5.** Same as Figure 3 but for CXOU J141332. For displaying purposes, the *Chandra* data have been shifted in power by +50.

**Table 3.** Spectral results of CXOU J141332. Errors are at a  $1\sigma$  confidence level for a single parameter of interest.

Model	Data	$N_{\text{H}}^a$ ( $10^{22} \text{ cm}^{-2}$ )	$\Gamma$	$kT$ (keV)	Flux <sup>b</sup> ( $10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$ )	Unabsorbed flux <sup>b</sup> ( $10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$ )	$\chi^2_{\nu}$ (dof)
PHABS(POWERLAW)	<i>Chandra</i> /12823	$<0.9^c$	$0.8^{+0.3}_{-0.2}$	–	$4.3^{+0.8}_{-0.6}$	$4.5^{+0.6}_{-0.5}$	1.02 (13)
	<i>Chandra</i> /12824				$6.3^{+1.4}_{-1.3}$	$6.5^{+1.2}_{-1.0}$	
PHABS(POWERLAW)	<i>XMM</i> /0111240101	$0.13^{+0.09}_{-0.07}$	$0.93 \pm 0.13$	–	$5.4 \pm 0.5$	$5.6^{+0.5}_{-0.4}$	0.82 (16)

<sup>a</sup> The abundances used are those of Wilms, Allen, & McCray (2000);  $N_{\text{H}}$  values  $\approx 30$  per cent lower are derived with those by Anders & Grevesse (1989). The photoelectric absorption cross-sections are from Balucinska-Church & McCammon (1992).

<sup>b</sup> In the 0.5–10 keV energy range.

<sup>c</sup> Upper limit at 90 per cent confidence level.

a poor fit to the data, with a  $\chi^2_\nu = 1.68$  for 13 dof, while for the bremsstrahlung ( $\chi^2_\nu = 1.31$ ) the temperature could not be constrained. For this reasons, in Table 3 we give only the parameters obtained from the power law fit (see Fig. 4, right), which gave a good fit with a rather hard power law with photon index  $\Gamma \approx 0.8$  and a best-fitting absorption  $N_{\text{H}} = 0.2 \times 10^{22} \text{ cm}^{-2}$  (with a 90 per cent upper limit of  $0.9 \times 10^{22} \text{ cm}^{-2}$ ). The observed flux was  $F_{\text{X}} \approx 4 \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$  during the first observation, and  $\approx 6 \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$  in the second; when one considers the large uncertainties, the flux increase is however only marginally significant.

The 2001 *XMM-Newton* observation offers a better spectrum, covering with a few more photons the band 0.4–8 keV. The power-law fit ( $\chi^2_\nu = 1.02$  for 16 dof) yielded photon index and flux similar to those derived with *Chandra* (see Table 3) and made it possible to constrain better the absorbing column. This was measured at  $N_{\text{H}} = 0.13^{+0.09}_{-0.07} \times 10^{22} \text{ cm}^{-2}$ . In the 2014 *XMM-Newton* observation, CXOU J141332 was not detected. The  $3\sigma$  upper limit on its observed flux, derived with PIMMS assuming the *XMM-Newton* power law model in Table 3, was  $2 \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$  in the 0.5–10 keV band in the MOS data (in the pn the position of the source occurred in proximity to streaks of out-of-time events due to the nucleus of the CG).

The position of CXOU J141332 was imaged also in the *Chandra* observations 2454, 355, 356, 10850, 10872, and 10873 (see Table 1); the source was never detected and for each data set we derived in a like manner the following upper limits (for the grating observations, we considered only the zero-order data):  $1.1 \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$  (obs. 355, ACIS-I),  $10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$  (obs. 356, ACIS-S),  $7 \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$  (obs. 2454, ACIS-S),  $5 \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$  (obs. 10873, ACIS-S),  $8 \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$  (obs. 10850, ACIS-S), and  $6 \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$  (obs. 10872, ACIS-S).

## 6 ASTROMETRY AND OPTICAL OBSERVATIONS OF CXO J141430 AND CXOU J141332

### 6.1 X-ray astrometry

In order to improve the absolute astrometry of the *Chandra* data to search for optical counterparts to CXO J141430 and CXOU J141332, we cross-correlated the X-ray source list obtained using WAVEDETECT with sources in the Two-Micron All-Sky Survey (2MASS; Skrutskie et al. 2006) catalog, which has an astrometric accuracy better than  $0''.2$ . We found 17 2MASS point sources coincident within  $0''.4$  from an X-ray source and used them to register the *Chandra* images on the accurate 2MASS reference frame by fitting a transformation matrix which includes a rotation, scale factor, and translation. We note that the *Chandra*–2MASS superposition did not require a significant transformation: the corrections are of the same order of the residuals ( $<0''.15$ ). The resulting positions (J2000.0) of the new CATS @ BAR pulsators are RA =  $14^{\text{h}}14^{\text{m}}30^{\text{s}}.1 (\pm 0''.45)$  and Decl. =  $-65^{\circ}16'23''.3 (\pm 0''.30)$  for CXO J141430, and RA =  $14^{\text{h}}13^{\text{m}}32^{\text{s}}.9 (\pm 0''.30)$  and Decl. =  $-65^{\circ}17'56''.5 (\pm 0''.25)$  for CXOU J141332, where the  $1\sigma$  uncertainties combine the *Chandra* localisation accuracy, the residuals of the *Chandra*–2MASS frame superposition, and the 2MASS absolute astrometric accuracy.

**Table 4.** VST/OmegaCAM images of regions around CXO J141430 and CXOU J141332.

MJD	Exp. (s)	Filter	Archive name
56442.149	25	$r'$	OMEGA.2013-05-30T03:34:12.715
56442.156	25	$r'$	OMEGA.2013-05-30T03:44:47.517
56460.067	25	$r'$	OMEGA.2013-06-17T01:07:16.243
56460.054	25	$r'$	OMEGA.2013-06-17T01:18:02.485
56487.025	25	$r'$	OMEGA.2013-07-14T00:36:46.111
56487.014	40	$r'$	OMEGA.2013-07-14T00:20:12.159
56488.033	25	$r'$	OMEGA.2013-07-15T00:47:30.554

### 6.2 VST data

Optical images of the regions around CXO J141430 and CXOU J141332 were retrieved from the ESO Science Archive Facility. The observations were originally obtained with the 2.6-m VST located at Paranal Observatory using the OmegaCAM instrument (Kuijken 2011), as part of the VST Photometric  $H\alpha$  Survey of the Southern Galactic Plane and Bulge (VPHAS+; Drew et al. 2014). Table 4 lists the details of the images examined.

We aligned and stacked the images using the Image Reduction and Analysis Facility (IRAF) software and the IMALIGN and IMCOMBINE packages following standard procedures. A world coordinate system was then applied using the IMWCS utility<sup>3</sup> and the Third US Naval Observatory CCD Astrograph Catalog (UCAC3; Zacharias et al. 2010). The mean error in the world coordinate system is  $0''.35$  ( $3\sigma$ ) using 22 local stars. We performed aperture photometry on the nearby fields using the IRAF package APPHOT and estimate a limiting magnitude depth of  $r' \approx 22.5$  mag. No sources are observed within  $2''$  of the coordinates of CXO J141430 and CXOU J141332.

## 7 DISCUSSION I: THE NATURE OF CXO J141430 AND CXOU J141332

With a new X-ray pulsator discovered by CATS @ BAR every  $\sim 2,200$  light curves, the detection of two previously unknown pulsating sources in the Circinus data set had a formal probability of about 0.04 per cent. While this does not qualify as a statistical anomaly, the 2010 *Chandra* observations have certainly been bountiful for the CATS @ BAR project.

### 7.1 CXO J141430

In the case of CXO J141430, which is displaced by  $\sim 2$  arcmin from the extreme edge of CG (Fig. 1), the two periodicities nail it down as an intermediate polar (IP) with spin period of  $P_{\text{spin}} = 6.1$  ks (1.7 h) and orbital period  $P_{\text{orb}} = 64.2$  ks (17.8 h). IPs, also known as DQ Herculis stars, and polars (AM Herculis stars) are the two main subclasses of magnetic cataclysmic variable stars (CVs). CVs are close binaries hosting a white dwarf (WD) accreting from a late-type Roche-lobe filling companion, either a main-sequence or a sub-giant star. IPs are characterised by asynchronous rotation ( $P_{\text{spin}} < P_{\text{orb}}$ ), while polars are phase-locked ( $P_{\text{spin}} \simeq P_{\text{orb}}$ ) and generally display strong circular polarisation (whence the name) at optical and near infrared wavelengths (see Patterson 1994; Warner

<sup>3</sup> See <http://tdc-www.harvard.edu/wcstools/>.

2003; Smith 2006 for reviews). Although it is still matter of debate, these differences are generally interpreted as due to a magnetic field in IPs which is weaker than that typically measured for polars ( $B \gtrsim 10^7$  G).

The orbital period of CXO J141430 locates the system above the 2–3 h so-called ‘orbital period gap’, where most of the IPs are found. Also the spin-to-orbit period ratio of  $\sim 0.095$  is typical of an IP (it is generally in the range 0.25–0.01, with most systems around 0.1).<sup>4</sup> In IPs, the accreted material generally passes through a disc and is then channelled onto the magnetic polar regions of the WD (at variance with polars, where the magnetic field inhibits the formation of the disc). There, a shock develops and the hot gas cools while it settles onto the WD surface emitting X-rays via thermal bremsstrahlung and cyclotron radiation (Aizu 1973). Because of their strong magnetic field, in polars the cooling takes place mainly via cyclotron, whereas IPs are expected to show bremsstrahlung-dominated emission. While the relatively poor statistical quality of the available spectra precluded a good characterisation of the X-ray emission of CXO J141430, the results of the spectral analysis are consistent with this picture (Table 2). Also the large-amplitude modulation at the orbital period is rather common in IPs (e.g. Parker, Norton, & Mukai 2005), and the  $\sim 100$  per cent pulsed fraction hints a high-inclination system. We finally note that the X-ray luminosity of CXO J141430 in the deep *Chandra* and *XMM-Newton* observations was  $L_X \approx 2 \times 10^{31} d_1^2 \text{ erg s}^{-1}$  (with a  $\approx 50$  per cent variability, see Section 4), where  $d_1$  is the distance in units of 1 kpc. Typical values for IPs,  $10^{32}$ – $10^{34} \text{ erg s}^{-1}$  (Sazonov et al. 2006), suggest that CXO J141430 is either on the lower side of the luminosity distribution or the distance to the source is substantially larger than  $\sim 1$  kpc.

For an IP with orbital period of 17.8 h, a K5V star would be a likely companion (e.g. Smith & Dhillon 1998). Using a value of  $N_H = 0.3 \times 10^{22} \text{ cm}^{-2}$  derived from our model fit to the X-ray spectra, and a conversion of  $N_H/A_V$  of  $1.79 \times 10^{21} \text{ cm}^{-2} \text{ mag}^{-1}$  (Predehl & Schmitt 1995), we obtain  $A_V = 1.7$  mag. Assuming an absolute magnitude  $M_{r'} \approx 7.1$  mag (Bilir et al. 2008),<sup>5</sup> the limiting magnitude  $m_{r'} \approx 22.5$  suggests  $d > 5$  kpc.

## 7.2 CXOU J141332

The nature of the fainter CXOU J141332 is less obvious. The source is located inside the 90 per cent total  $B$  light contour of the CG. So, the first question that needs to be addressed is whether it is a Galactic or an extragalactic source. Based on the cumulative Galactic X-ray source density versus flux distribution ( $\log N$ – $\log S$ ) from the *Chandra* Multi-wavelength Plane survey (ChaM-Plane; van den Berg et al. 2012), we estimated the probability of a foreground Galactic object of that flux within this area to be  $\approx 11$  per cent. Moreover, its absorbing column measured with *XMM-Newton* ( $N_H = 1.3_{-0.7}^{+0.9} \times 10^{21} \text{ cm}^{-2}$ ; Table 3) is much lower than the total Galactic value of  $\sim 6 \times 10^{21} \text{ cm}^{-2}$  (Dickey & Lockman 1990; Kalberla et al. 2005). CXOU J141332 is therefore most likely a Galactic source.

The hard X-ray spectrum, a power law with photon index  $\Gamma \sim 0.8$ – $0.9$ , and the modulation at 6.4 ks (1.8 h) point to a binary system consisting of a compact star accreting from a low-mass

companion, where the 1.8 h period likely traces the orbital motion. The period is in fact too long to be the spin of a typical neutron star (NS; with very few possible exceptions; Mattana et al. 2006; Esposito et al. 2011). The low flux and—chiefly—the smooth and  $\sim 60$  per cent pulsed-fraction modulation favour a magnetic CV nature also for CXOU J141332 (the period is also too short for the orbit of a standard high-mass X-ray binary). Indeed, CVs are the most abundant population of Galactic compact interacting binaries, and also the most frequent new pulsating sources in the CATS @ BAR sample (Israel et al., in preparation). In particular, since no second periodicity was detected (and lacking any information about optical polarisation), the source could be a polar. Polars are generally found at short orbital periods, most of them below the 2–3 h orbital gap, and CXOU J141332 would lie in the peak of their period distribution (e.g. Ritter & Kolb 2003). The profile of the folded light curve (Fig. 5) and its variability as function of energy may indicate a two-pole system. The luminosity during the long 2010 *Chandra* observations was  $L_X \approx (5\text{--}8) \times 10^{30} d_1^2 \text{ erg s}^{-1}$ , with the upper limits from the other observations implying a variability of  $\approx 50$  per cent or larger. For distances of the order of a few kpc, this is in good agreement with typical values for polars ( $L_X < 10^{32} \text{ erg s}^{-1}$ ; Sazonov et al. 2006).

For a polar with orbital period of 1.8 h, the companion is likely a M5V star, which has an absolute magnitude of  $M_{r'} \approx 12.5$  mag (Bochanski et al. 2011). The value of  $N_H = 0.13 \times 10^{22} \text{ cm}^{-2}$  derived from our model fit to the X-ray spectra implies an  $A_V = 0.7$  mag. The limiting magnitude  $m_{r'} \approx 22.5$  suggests  $d \gtrsim 0.7$  kpc. In the IP hypothesis, assuming for CXOU J141332 a K5V as for CXO J141430, the non-detection would indicate a distance larger than  $\sim 8$  kpc.

## 8 THE CONTROVERSIAL SOURCE CG X–1

CG X–1 (CXOU J141312.3–652013), about 15 arcsec northeast of the Circinus’ nucleus, had been known for long to be a bright and variable (possibly periodic) X-ray source. Using high-quality light curves collected with *Chandra*, Bauer et al. (2001) discovered a strong modulation at a period of  $\sim 27$  ks. The measured X-ray flux was  $9 \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$  (0.5–10 keV), and deep *HST* observations did not detect any optical counterpart to CG X–1, with a limit  $m_{F606W} > 25.3$ . They observed that the source might be either a black hole (BH) binary in the CG radiating at  $\sim 4 \times 10^{39} \text{ erg s}^{-1}$  (and hence qualifying as an ULX) or a Galactic cataclysmic variable of the polar type with a particularly long period (in both cases, the 27-ks modulation would reflect the orbital period of the system).

In the polar hypothesis, for a M2V to M6V companion star, the *HST* limit puts the source at a distance larger than 1.2 kpc, implying a luminosity of at least  $3 \times 10^{32} \text{ erg s}^{-1}$ , a rather extreme value for a polar. Bauer et al. (2001) also noticed that the association of CG X–1 with the CG is convincing: based mainly on the results from the *ASCA* Galactic plane survey (Sugizaki et al. 2001), they evaluated that the possibility of a foreground or background X-ray source is  $\lesssim 0.06$  per cent. Moreover, Smith & Wilson (2001) noticed that the absorption toward CG X–1 ( $N_H > 10^{22} \text{ cm}^{-2}$ ) is much larger than the total Galactic column ( $N_H \sim 6 \times 10^{21} \text{ cm}^{-2}$ ), further supporting the association with the CG. Overall, Bauer et al. (2001), Smith & Wilson (2001), and Bianchi et al. (2002) favoured a very bright extragalactic BH binary, harbouring a BH possibly in excess of  $50 M_\odot$ .

Despite recognising the robustness of the association, Weis-

<sup>4</sup> See for example the catalog available at the Intermediate Polar Home Page, <http://asd.gsfc.nasa.gov/Koji.Mukai/iphome/catalog/members.html>.

<sup>5</sup> Uncertainty in the assumed absolute magnitudes of the companion star may be as large as 2 mag. The same holds for CXOU J141332, see below.

skopf et al. (2004) were more open toward the possibility of a foreground polar. They observed that while the period and the luminosity would be somewhat atypical, CG X–1 would be neither the longest-period, nor the brightest known polar. On the other hand, they argued that if CG X–1 belonged to the CG, because of the short orbital period it should be a BH low-mass X-ray binary (LMXB) with a  $\lesssim 1 M_{\odot}$  companion. In this case, the huge X-ray luminosity of the BH would drive the star out of thermal equilibrium and evaporate it within  $\sim 10^3$  yr. They regard as very unlikely the possibility that a system this short-lived could be observed.

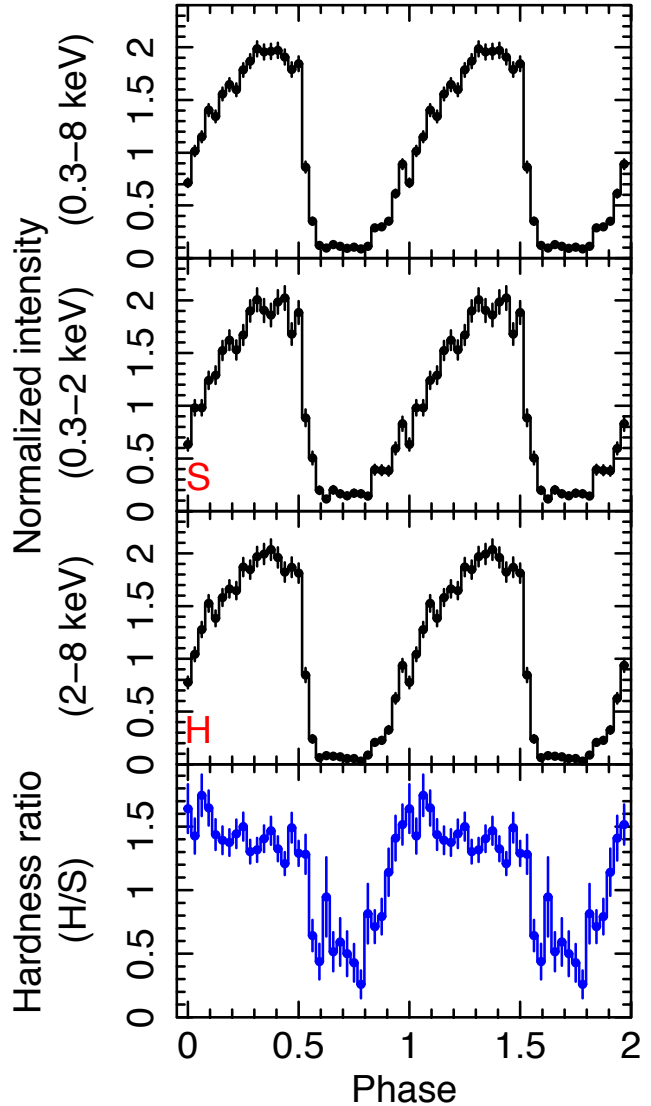
Weisskopf et al. (2004) revised the period of CG X–1 at  $26.25 \pm 0.15$  ks and noticed a possible optical counterpart with  $m_{F606W} = 23.5$ . However, they did not estimate the significance of the association or of the source, and gave instead a limiting magnitude of 24.3. Ptak et al. (2006) confirmed the limit by Bauer et al. (2001), while a recent work by Gladstone et al. (2013) proposed a counterpart with  $m_V = 24 \pm 6$  (presumably, the same excess/source detected by Weisskopf et al. 2004). In literature, CG X–1 is generally considered to be an ULX in the CG (e.g. Bianchi et al. 2002; Swartz et al. 2004; Liu & Mirabel 2005; Ptak et al. 2006; Berghea et al. 2008; Gladstone et al. 2013).

## 9 THE NEW CHANDRA DATA OF CG X–1: ANALYSIS AND RESULTS

Since the CG has been observed many times in X-rays, in particular with *Chandra*, a wealth of data exist for CG X–1. Detailed studies of CG X–1 with *ROSAT*, *BeppoSAX*, *XMM-Newton*, and *Chandra* were presented in the aforementioned works by Smith & Wilson (2001), Bauer et al. (2001), Bianchi et al. (2002), and Weisskopf et al. (2004) (but see also Matt et al. 1996; Guainazzi et al. 1999; Sambruna et al. 2001; Massaro et al. 2006; Bauer et al. 2008; Yang et al. 2009; Shu et al. 2011; Walton et al. 2013; Arévalo et al. 2014). A systematic analysis of all the available data is beyond the scope of this paper, and here we will only present results from the analysis of observations 12823/4, which represent the deepest and highest-quality data set available for this source and, to our knowledge, have not been used so far to study CG X–1.

The coordinates we derived for CG X–1 are RA =  $14^{\text{h}} 13^{\text{m}} 12^{\text{s}}.24 (\pm 0^{\text{s}}.25)$  and Decl. =  $-65^{\circ} 20' 13''.82 (\pm 0''.25)$ . Using recent results from the ChaMPPlane survey (van den Berg et al. 2012), we confirm the low probability of a foreground or background object ( $\lesssim 0.06$  per cent) estimated by Bauer et al. (2001). We actually believe that the probability of an interloper is substantially lower (approximately two times smaller), considering that the flux reached by the source in subsequent observations is  $\sim 5$  times higher (Weisskopf et al. 2004) and the fact that a background AGN can be excluded by the phenomenology of CG X–1.

For the spectral and timing analysis, we extracted the source counts within a 1.5-arcsec radius, while for the background we used an annulus with radii of 3 and 5 arcsec (see Fig. 1 and Section 2 for more details). The 0.3–8 keV source net count rate was  $(6.90 \pm 0.07) \times 10^{-2}$  counts  $\text{s}^{-1}$  in observation 12823, and  $(6.1 \pm 0.1) \times 10^{-2}$  counts  $\text{s}^{-1}$  in observation 12824. These rates are high enough to cause pileup in the ACIS detector (as we checked with a pileup map). For the spectral analysis, the pileup was dealt with by using the pileup model by Davis (2001). This procedure involves some uncertainty, because the pileup fraction in CG X–1 is strongly dependent on the orbital phase. However none of our results crucially depends on the exact value of the parameters derived from the spectral fitting.



**Figure 6.** Background-subtracted folded profile of CG X–1 (observation 12823/4) in different energy bands. The hardness ratio between the hard and soft bands is also plotted at the bottom.

About six 26-ks cycles were recorded in obs. 12823 and two in obs. 12824. There is a moderate pulse-to-pulse variability, both in shape and in the flux at maximum ( $\approx 30$  per cent). Also, the average count rate was  $\approx 15$  per cent lower during the second observation. To measure the period, we fit a sinusoidal function to the light curve. We obtained  $P_{\text{CG X-1}} = 26.1 \pm 0.1$  ks. The corresponding folded profile in different energy bands is shown in Fig. 6. The profile is asymmetric and the modulations is large (but the count rate is non-zero also at minimum). The pulsed fraction, defined as  $^6 (M - m)/(M + m)$ , where  $M$  is the maximum count rate and  $m$  the minimum, is  $91.6 \pm 1.5$  per cent in the whole band,  $89.2 \pm 2.5$  per cent in the 0.3–2 keV band, and  $97.1 \pm 1.2$  per cent in the 2–8 keV band. A conspicuous spectral softening around minimum is evident from the ratio of the hard and soft counts along the cycle (Fig. 6).

<sup>6</sup> Here we adopted a different definition of the pulsed fraction than used before, because the value inferred from a sinusoidal fit would misrepresent the amplitude of the modulation of the ‘sawtooth’ profile of CG X–1.

For the spectral analysis we fit three simple models to the data: a power law, a multicolour disc (MCD; Mitsuda et al. 1984; Makishima et al. 2000), and an optically-thin thermal bremsstrahlung, all corrected for the interstellar absorption. While a multicolour disk with  $kT \simeq 1.3\text{--}1.4$  keV gives the lowest  $\chi^2$ , all models provide an acceptable fit to the data (Table 5). All fits confirm that  $N_{\text{H}}$  towards CG X–1 is substantially larger than the total Galactic absorbing column in that direction. Weisskopf et al. (2004) reported a possible feature, probably a blend of Fe lines, in the 2001 *XMM-Newton* spectrum of CG X–1. No line is required to fit the *Chandra* data (see Fig. 7 for the longest observation); the  $3\sigma$  upper limit on the equivalent width of any line with central energy between 6 and 7 keV is 0.18 keV in observation 12823. This limit is formally compatible with the equivalent width of  $0.23 \pm 0.06$  keV derived by Weisskopf et al. (2004). However, since there is no trace of such feature in the ACIS data, it is possible that, as noticed also by Weisskopf et al. (2004), the feature observed with *XMM-Newton* is due to residual contamination from the emission of the nuclear region of the CG, which has very strong lines at 6.4 and 7 keV.

We also extracted pulse-resolved spectra from the soft (phase between 0.55 and 0.9 in Fig. 6) and hard (all other phases) parts of the hardness ratio. The data were prepared with DMTCALC, and the spectra of the two data sets from the same phase bins were combined using COMBINE\_SPECTRA, which also averaged the response matrices. Although the few counts in the soft spectrum ( $\sim 700$  versus more than 12 200 in the hard spectrum, after combining the two observations) preclude a detailed comparison, we found that the spectral variation can be described equally well by either a decrease in the absorption during the softening ( $N_{\text{H}} = (1.00 \pm 0.03) \times 10^{22} \text{ cm}^{-2}$  in the hard phase and  $(0.24 \pm 0.07) \times 10^{22} \text{ cm}^{-2}$  in the soft phase for the MCD model;  $(1.45 \pm 0.04) \times 10^{22} \text{ cm}^{-2}$  in the hard phase and  $(0.67 \pm 0.08) \times 10^{22} \text{ cm}^{-2}$  in the soft phase for the power-law model;  $(1.29 \pm 0.03) \times 10^{22} \text{ cm}^{-2}$  in the hard phase and  $(0.51 \pm 0.07) \times 10^{22} \text{ cm}^{-2}$  in the soft phase for the bremsstrahlung model) or a change in the pivotal parameter of the model ( $kT = 1.82 \pm 0.05$  keV in the hard phase and  $0.95 \pm 0.06$  keV in the soft phase for the MCD;  $\Gamma = 1.70 \pm 0.03$  in the hard phase and  $2.54 \pm 0.11$  in the soft phase for the power law;  $kT = 9.4 \pm 0.8$  keV in the hard phase and  $2.4 \pm 0.3$  keV in the soft phase for the bremsstrahlung).

## 10 DISCUSSION II: IS CG X–1 A WOLF–RAYET/BLACK-HOLE BINARY IN THE CIRCINUS GALAXY?

While we cannot completely exclude the chance superposition of a Galactic polar in the direction of the inner part of the CG, we regard the association of the CG X–1 to the CG as rather compelling. In the following, we will discuss the nature of the source in this framework.

Weisskopf et al. (2004) disfavoured the possibility of a BH LMXB because such a system would be rather short lived, and thus unlikely to be observed. This argument is not conclusive when dealing with an individual source (one can be lucky enough to observe a rare system!), but we too deem a LMXB as unlikely. This is because of its orbital profile (Fig. 6). In fact, most LMXBs show no or very low-amplitude modulation on their orbital period. In those which display orbital modulation, dipping and/or eclipsing systems, the morphology of the profile is very different. In dipping systems, the dips are produced by absorption of X-rays due to accreting matter located in the bulge at the outer edge of the accretion disc (e.g. Díaz

Trigo et al. 2006); the X-ray minima tend to be rather sharp and to show a harder than average emission, which is the opposite of what is observed in CG X–1 (Fig. 6, bottom panel). LMXB eclipses display sharp and abrupt ingresses and egresses, due to the small size of the X-ray-emitting regions. Even the smoothly-modulated accretion disc corona sources (systems viewed nearly edge-on where the outer edge of the dense disc modulates the X-rays from the central source scattered into the line of sight by an extended ionised corona; Mason & Cordova 1982; White & Holt 1982; Somero et al. 2012) have different profiles. Moreover, LMXBs containing BHs are usually transient X-ray sources, with outburst durations of the order of weeks to months (e.g. Remillard & McClintock 2006), while CG X–1 is variable but persistent.

On the other hand, the light curve and the folded profile of CG X–1 bear a strong resemblance to those observed in high-mass X-ray binaries (HMXBs) with a Wolf–Rayet (WR) star companion: Cyg X–3 in the Milky Way ( $P = 4.8$  h; Zdziarski et al. 2012), IC 10 X–1 in IC 10 ( $P = 35$  h; Prestwich et al. 2007), NGC 300 X–1 in NGC 300 ( $P = 33$  h; Carpano et al. 2007a), and the candidates CXOU J123030.3+413853 in NGC 4490 (CXOU J123030;  $P = 6.4$  h; Esposito et al. 2013a) and CXOU J004732.0–251722.1 in NGC 253 (with a candidate periodicity  $P \sim 14\text{--}15$  h; Maccarone et al. 2014).<sup>7</sup> The large modulation of these systems, showing a slow rise and a faster decay, is different from what observed in almost any other class of X-ray sources (see, e.g., the arguments summarised in Esposito et al. 2013a about the source in NGC 4490), and can be considered the signature of WR HMXBs.

Although its properties remain still poorly understood, the best studied source of this kind is Cyg X–3, where the orbital modulation has been ascribed to attenuation by electron scattering in the strong WR wind (Hertz et al. 1978). A further significant contribution to the scattering medium is thought to be a bulge of ionised matter formed by the collision of the stellar wind with the outer accretion disc (Zdziarski et al. 2010). Alternative explanations involve orbital modulation by absorbers in different phases (hot/ionised and cold/clumpy, triggered by the BH jet bow shock; Vilhu & Hannikainen 2013), or the presence of the accretion wake, a large scale asymmetry around the compact object (Okazaki & Russell 2014). The orbital profile of CG X–1 is even more asymmetric than in Cyg X–3, probably due to rather extreme properties of the absorbing/scattering medium or because of a higher inclination of the system. The variation of the hardness ratio along the orbit (Fig. 6), showing a clear softening during the X-ray minimum, could be due to direct X-rays almost completely blocked by dense matter (probably the innermost regions of the WR wind when the X-ray source is at the superior conjunction). The softer X-rays observed could be due to down-scattering into the line of sight of central X-rays by cold and less dense material located farther away from the WR star (similarly to what usually observed during eclipses in some HMXBs; Haberl 1991).

Weisskopf et al. (2004) mentioned for CG X–1 the possibility of a WR or, more in general, a naked He donor, but did not discuss it in detail owing to the scarcity of information of such systems. In particular, they observed that for a  $\sim 2 M_{\odot}$  companion, their argument against LMXB systems was still relevant. However, thanks to their compactness, even more massive WR stars can

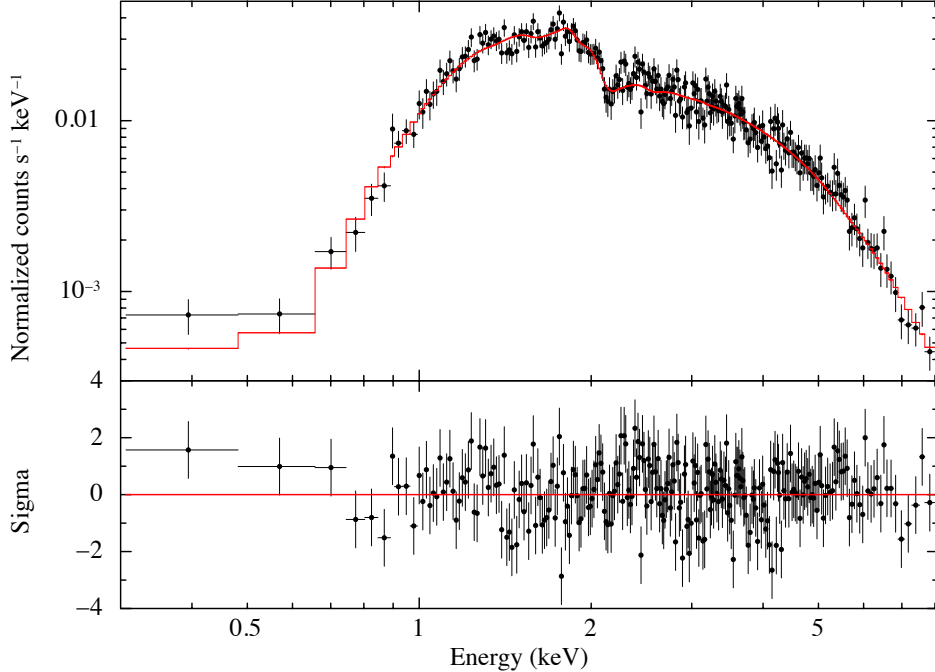
<sup>7</sup> Apart from these objects, the only other known WR HMXB is ULX–1 in M101 (Liu et al. 2013, see also Section 10.1). A period of 8.2 d was inferred from radial velocities of optical emission lines, but no X-ray light curves of good quality are available for this source.

**Table 5.** Spectral results of CG X–1. Errors are at a  $1\sigma$  confidence level for a single parameter of interest.

Model	Obs. ID	$N_{\text{H}}^a$ ( $10^{22} \text{ cm}^{-2}$ )	$\Gamma$	$kT$ (keV)	Flux <sup>b</sup> ( $10^{-12} \text{ erg cm}^{-2} \text{ s}^{-1}$ )	Unabsorbed flux <sup>b</sup> ( $10^{-12} \text{ erg cm}^{-2} \text{ s}^{-1}$ )	$\chi^2_{\nu}$ (dof)
PHABS(DISKBB)	12823	$1.06 \pm 0.04$	–	$1.36 \pm 0.07$	$0.83^{+0.13}_{-0.12}$	$1.17 \pm 0.17$	0.97 (271)
PHABS(POWERLAW)	12823	$1.38 \pm 0.05$	$1.72^{+0.06}_{-0.05}$	–	$1.23^{+0.10}_{-0.09}$	$1.84^{+0.12}_{-0.10}$	1.04 (271)
PHABS(BREMSSTRAHLUNG)	12823	$1.28 \pm 0.04$	–	$6.5^{+1.0}_{-0.8}$	$1.02^{+0.04}_{-0.05}$	$1.50^{+0.05}_{-0.06}$	1.00 (271)
PHABS(DISKBB)	12824	$1.03^{+0.08}_{-0.07}$	–	$1.33^{+0.14}_{-0.10}$	$0.68^{+0.21}_{-0.16}$	$0.85^{+0.29}_{-0.22}$	0.98 (92)
PHABS(POWERLAW)	12824	$1.38^{+0.11}_{-0.10}$	$1.77^{+0.12}_{-0.10}$	–	$1.00^{+0.15}_{-0.14}$	$1.52^{+0.18}_{-0.17}$	1.11 (92)
PHABS(BREMSSTRAHLUNG)	12824	$1.27 \pm 0.08$	–	$5.6^{+1.7}_{-1.2}$	$0.83 \pm 0.07$	$1.25 \pm 0.09$	1.05 (92)

<sup>a</sup> The abundances used are those of Wilms, Allen, & McCray (2000);  $N_{\text{H}}$  values  $\approx 30$  per cent lower are derived with those by Anders & Grevesse (1989). The photoelectric absorption cross-sections are from Balucinska-Church & McCammon (1992).

<sup>b</sup> In the 0.5–10 keV energy range.

**Figure 7.** *Chandra*/ACIS spectrum and best-fitting MCD model (red solid line) for CG X–1 from observation 12823. Bottom panel: the residuals of the fit in units of standard deviations.

comfortably fit in the orbit of a system like CG X–1 (the stellar radius of a  $20\text{-}M_{\odot}$  WR star is  $< 2 R_{\odot}$ ; e.g. Langer 1989; Schaerer & Maeder 1992). Also the other main observational properties of CG X–1 fit well in the scenario of a WR–BH HMXB in the CG.<sup>8</sup>

If one considers the reddening toward the CG (4 mag) and its distance module (28.1 mag, from NED), the limit on the optical counterpart by Bauer et al. (2001) and Ptak et al. (2006) implies  $M_V > -6.8$ . This value is compatible with a WR star, for which  $M_V$  is typically in the range from  $-2.5$  to  $-7$  (e.g., Massey 2003).

The X-ray luminosity of CG X–1 is variable by a factor of  $\approx 10$  (Bianchi et al. 2002; Weisskopf et al. 2004). The highest flux reported in literature ( $5.2 \times 10^{-12} \text{ erg cm}^{-2} \text{ s}^{-1}$  for a power-law fit or  $5 \times 10^{-12} \text{ erg cm}^{-2} \text{ s}^{-1}$  for a MCD fit, in the 0.5–8 keV band;

Weisskopf et al. 2004) would imply, for a distance of 4.2 Mpc, a 0.5–10 keV luminosity of  $L_X = (1.5\text{--}2) \times 10^{40} \text{ erg s}^{-1}$ . If the system is Eddington-limited, the lower limit on the mass of the accreting BH is  $M_{\text{BH}} \gtrsim 75 M_{\odot}$  for a He or C/O donor. For the system to shine in X-rays, the velocity of the WR star wind has to be slow enough to allow the formation of an accretion disk. This condition corresponds for CG X–1 to the requirement  $M_{\text{BH}} \gtrsim 1.5 v_w^4 v_{w,1000} \delta^2 M_{\odot}$ , where  $M_{\text{BH}}$  is the BH mass,  $v_w$  is the wind velocity in units of  $1000 \text{ km s}^{-1}$ , and  $\delta \approx 1$  is a dimensionless parameter (adapted from Carpano et al. 2007a, see also Illarionov & Sunyaev 1975). In the simplest wind-accretion case (e.g. Edgar 2004), the luminosity can be estimated as

$$L_X \approx \eta \frac{\dot{M}_w c^2 G^2 M_{\text{BH}}^2}{a^2 (v_{\text{orb}}^2 + v_w^2)} \quad (1)$$

where  $\eta$  is the efficiency,  $\dot{M}_w$  is the wind mass loss rate,  $a$  is the orbital separation,  $v_{\text{orb}}$  is the orbital velocity, and  $v_w$  is the wind velocity at the BH orbit. Assuming  $\dot{M}_w = 10^{-5} M_{\odot} \text{ yr}^{-1}$  and  $v_w = 1000 \text{ km s}^{-1}$  for the WR star (e.g. Crowther 2007),

<sup>8</sup> Even bright ( $L_X > 10^{40} \text{ erg s}^{-1}$ ) ULXs may contain a NS (King 2009; Bachetti et al. 2014), although such systems are probably not the majority (Fragos et al. 2015). Here, we will not discuss this possibility.

$a = 5.8 \times 10^{11}$  cm (for a  $10 M_{\odot}$  companion),  $M_{\text{BH}} = 75 M_{\odot}$ , and the formation of a disc with  $\eta = 0.1$ , the corresponding luminosity is  $L_X \simeq 2 \times 10^{40}$  erg s $^{-1}$ . More in general, for  $M_{\text{BH}} > 10 M_{\odot}$  and all the other things being equal, one finds  $L_X \gtrsim 3 \times 10^{39}$  erg s $^{-1}$ . In case of Roche lobe overflow, even higher X-ray luminosity could be achieved. However we note that if CG X–1 is indeed a WR–BH binary, the WR star is probably not filling its Roche lobe (unless it is very massive; see for example the discussion of the case of Cyg X–3, where the orbital period is much shorter, in Szostek & Zdziarski 2008). An X-ray luminosity of  $\sim 2 \times 10^{40}$  erg s $^{-1}$  can be therefore accounted for. We finally notice that, although we do not regard the question as crucial, the problem of the lifetime of the system discussed by Weisskopf et al. (2004) would be significantly attenuated, since the WR phase of a massive O-type star is thought to last a few  $\times 10^5$  years (Meynet & Maeder 2005).

### 10.1 Statistics, environment, and WR–BH binaries as ULXs

All known WR HMXBs have been mentioned in the previous Section. Three have been established so far as certain WR–BH systems: IC 10 X–1, NGC 300 X–1, and M101 ULX–1. For the fourth WR–compact object binary, Cyg X–3 in our Galaxy, it is still debated whether the compact object is a BH or a NS. There are however several pieces of evidence (radio, infrared and X-ray emission properties) that point to a 2–5  $M_{\odot}$  BH, as suggested also by evolutionary models (e.g. Lommen et al. 2005; Szostek & Zdziarski 2008; Szostek et al. 2008; Shrader et al. 2010; Zdziarski et al. 2013). Apart from CG X–1, discussed in this paper, two additional WR–BH binary candidates were found in the last two years: CXOU J123030 in NGC 4490 and CXOU J004732.0–251722.1 in NGC 253.

We expect WR–BH binaries to be associated with star forming regions, since WR are young stars with massive progenitors (with zero-age main sequence mass  $\gtrsim 25 M_{\odot}$ ). Table 6 shows the star formation rate (SFR) of the host galaxies of WR–BH binaries and binary candidates. The average SFR is  $\sim 2 M_{\odot} \text{ yr}^{-1}$ , which is quite high for nearby late-type galaxies. All metallicities listed in Table 6 are sub-solar. Recent  $N$ -body and population synthesis simulations of young star clusters (Mapelli et al. 2013; Mapelli & Zampieri 2014) suggest that  $\approx 2$  per cent of all HMXBs powered by BHs in star forming regions are BH–WR binaries, independent of star cluster metallicity.

The average luminosity of CG X–1 in the *Chandra* observations presented here makes it a bona fide ULX. The maximum luminosity reached by the source, following Weisskopf et al. 2004, is  $L_X = (1.5\text{--}2) \times 10^{40}$  erg s $^{-1}$ . At such flux levels, in high counting statistics *XMM–Newton* spectra, a persistent ULX typically shows the hallmark of the ultraluminous state (Gladstone, Roberts, & Done 2009), with the presence of two thermal components, one of which centred at rather soft energies ( $\sim 0.1\text{--}0.2$  keV) and the other producing a shallow but significant rollover around 3–5 keV. Such a spectrum is usually interpreted as the imprint of a super-Eddington accretion regime (e.g. Middleton, Sutton, & Roberts 2011).

As discussed in Section 8, the X-ray spectrum of the *Chandra* observations considered here is satisfactorily fitted with a single component model. The lack of evidence of multiple components may be real, but may also be caused by the comparatively low statistics (the flux is  $\sim 5$  times lower than at maximum) and/or by the fact that *Chandra* has a lower sensitivity than *XMM–Newton* below  $\sim 0.5$  keV, where the soft component peaks. If the spectrum is

intrinsically single component, this may suggest that CG X–1 does not enter into the ultraluminous state and hence does not accrete above Eddington.

The latter possibility is consistent with the scenario discussed above in which accretion proceeds through a wind and the accretion rate does not exceed the Eddington limit. For wind accretion from a compact WR star, the very formation of a standard accretion disc is uncertain and the accretion efficiency can in general be smaller than that of a standard disc (e.g. Frank, King, & Raine 2002). A scenario of rather efficient accretion from a wind of a massive WR star onto a BH of a few tens  $M_{\odot}$  has been invoked also for M 101 ULX–1. This source is a WR–BH ULX system with a dynamical mass measurement and a cool disc X-ray spectrum at maximum (Liu et al. 2013), and it has a significantly larger orbital period ( $\sim 8$  days, see Table 6) and smaller luminosity ( $\sim 3 \times 10^{39}$  erg s $^{-1}$ ) than CG X–1. Assuming a similar scenario also for CG X–1, simple BH mass estimates based on Eddington-limited accretion from a He–WR star (as those reported above) give  $M_{\text{BH}} \gtrsim 70 M_{\odot}$  for the observed maximum luminosity. Such a massive BH would populate the high-mass tail of the distribution of BHs formed through direct collapse of a massive star in a low metallicity environment (Mapelli et al. 2009; Zampieri & Roberts 2009; Belczynski et al. 2010), scenario indeed consistent with the metallicity inferred for the CG (see Table 6). In the same hypothesis, an independent limit can be obtained from the normalisation of the disc component at maximum luminosity (*Chandra* Obs. ID 365; Weisskopf et al. 2004), giving  $M_{\text{BH}} \simeq 8[(d/4.2 \text{ Mpc})/\sqrt{\cos i}] M_{\odot}$  ( $i$  is the inclination angle of the disc, e.g. Lorenzin & Zampieri 2009).

On the other hand, the lack of a high counting statistics spectrum at maximum luminosity prevents us from reaching a robust conclusion. We thus briefly consider also the possibility that the system is accreting from a particularly massive and big WR companion via Roche-lobe overflow, for which the mass transfer rate is expected to significantly exceed the Eddington limit (e.g. Lommen et al. 2005). In this assumptions, the maximum observed luminosity of CG X–1 would place it in the populated part of the ULX luminosity distribution (e.g. Swartz et al. 2011), where the observed flux can be produced by moderately beamed, super-Eddington emission from accretion onto a BH of a few tens  $M_{\odot}$  (or even a canonical stellar-mass BH).

It is interesting to compare the properties of the two candidate WR–BH ULX systems that we have tentatively identified, CG X–1 and CXOU J123030 in NGC 4490 (Esposito et al. 2013a). While the orbital periods are similar (7.2 and 6.4 h, respectively), the X-ray luminosity is significantly different (CG X–1 being  $> 10$  times more luminous). Assuming a similar scenario of sub-Eddington accretion from a wind with the formation of a disc, the different luminosity could be mostly ascribed to the different BH mass, with CG X–1 being about  $\sim 10$  times more massive than CXOU J123030. For non-extreme inclinations, this is consistent with the BH mass inferred from the normalisation of the disc component of CXOU J123030, for which we obtained  $M_{\text{BH}} \simeq 2.8[(d/8 \text{ Mpc})/\sqrt{\cos i}] M_{\odot}$  (Esposito et al. 2013a). On the other hand, in case of non-standard super-Eddington accretion via Roche-lobe overflow, the larger luminosity of CG X–1 may be caused by the larger accretion rate and we would then be witnessing the different X-ray outcome produced by similar BHs in very different accretion environments.

**Table 6.** Properties of observed WR–BH binaries and candidates (denoted by stars), and of their host galaxies.

Host galaxy	Source	Period (h)	BH mass <sup>a</sup> (M <sub>⊙</sub> )	WR mass <sup>a</sup> (M <sub>⊙</sub> )	SFR <sup>b</sup> (M <sub>⊙</sub> yr <sup>-1</sup> )	Z <sup>b</sup> (Z <sub>⊙</sub> )	t <sub>GW</sub> <sup>c</sup> (Gyr)
IC 10	X-1	34.9	33	35	0.07	0.22	1.4
NGC 300	X-1	32.8	20	26	0.14	0.19	1.7
NGC 4490	CXOU J123030.3+413853 *	6.4	–	–	4.5	0.23	0.038
NGC 253	CXOU J004732.0–251722.1 *	14.5	–	–	4.0	0.24	0.33
Circinus	CG X-1 *	7.2	–	–	1.5	0.10	0.052
M 101	ULX-1	196.8	20	19	3.1	0.17	200
Milky Way	Cyg X-3	4.8	3	7	0.25	0.31	0.051

<sup>a</sup> For the BH and WR masses, we list only the fiducial values that we use to derive the merger rates  $R$  in equation 2. Most of these masses are very uncertain, as discussed in Prestwich et al. (2007); Silverman & Filippenko (2008); Carpano et al. (2007b,a); Crowther et al. (2010); Esposito et al. (2013a); Maccarone et al. (2014); Liu et al. (2013); Shrader et al. (2010).

<sup>b</sup> The values of SFR and metallicity of the host galaxy come from the compilation of Mapelli et al. (2010b). The metallicity  $Z$  refers to the value at  $0.7 R_{25}$ , where  $R_{25}$  is the Holmberg radius of the galaxy, for NGC 253, NGC 300, NGC 4490, M 101, and for the Milky Way, while it is the total metallicity for IC 10 and Circinus. We assume  $Z_{\odot} = 0.02$ .

<sup>c</sup> The parameter  $t_{\text{GW}}$  is the timescale for the binary to coalesce, under the assumptions discussed in the main text.

## 10.2 WR–BH binaries as precursors of BH–BH binaries

There is great uncertainty on the expected rate of BH–BH mergers in the frequency range that will be observed by Advanced LIGO and Virgo ( $\sim 10\text{--}10^4$  Hz; Abadie et al. 2010). In fact, while for NS–NS binaries the expected merger rate can be derived from the properties of the observed NS–NS binaries (e.g. Kim et al. 2003) and from the rate of short gamma-ray bursts (e.g. Coward et al. 2012; Fong et al. 2012), no evidence has been found of BH–BH systems yet. WR–BH binaries can provide us with essential clues, since they are a possible precursor of BH–BH binaries (or BH–NS binaries), provided that the system is so tight that it remains bound when the WR star evolves into a compact remnant.

Here, we use the properties of all known WR–BH binaries and candidates to infer the BH–BH merger rate in the instrumental range of Advanced LIGO and Virgo. The main source of uncertainty is represented by the fate of the WR: we do not know the natal kick and the mass of the compact remnant that will form from the evolution of the WR. In particular, the natal kick might unbind the binary or transform it into a loose system that will not merge in a Hubble time. Thus, population synthesis simulations of the evolution of each WR–BH system are necessary, in order to predict the binary fate (e.g. Belczynski et al. 2013). On the other hand, our knowledge of the current orbital parameters and of the masses of the binary members is poor for most WR–BH systems: the WR star has been detected only in Cyg X-3, IC 10 X-1, NGC 300 X-1 and M101 ULX-1, and the mass of the BH is poorly determined even in these four cases, due to the uncertainty on the inclination and because the WR winds might affect the radial velocity estimate (e.g. Maccarone et al. 2014). Given these uncertainties, population synthesis models can hardly constrain the fate of most WR–BH candidates. Thus, we adopt a much simpler approach in order to obtain an upper limit to the BH–BH merger rate from the observed WR–BH binaries and candidates. As already done by Maccarone et al. (2014), we assume that all seven WR–BH candidates in our sample will become BH–BH binaries through direct collapse of the WR star (leading to a secondary BH mass  $m_2 = 10 M_{\odot}$ ), and that the orbital properties of the binary will be substantially unchanged after the collapse of the WR star. When even the mass of the primary BH is not known (Table 6), we assume  $m_1 = 10 M_{\odot}$ .

Thus, the rate of BH–BH mergers per Mpc<sup>3</sup> ( $R$ ) can be approximately estimated as (see e.g. Mapelli et al. 2010a, 2012; Ziosi

et al. 2014):

$$R = \rho_{\text{SFR}}(z) \sum_i (t_{\text{GW},i} + t_{\text{evol},i})^{-1} (\text{SFR}_i)^{-1}, \quad (2)$$

where  $\rho_{\text{SFR}}(z)$  is the cosmic SFR density (Hopkins & Beacom 2006),  $t_{\text{GW},i}$  is the coalescence timescale of the  $i$ -th binary ( $i = 1, \dots, 7$  in our sample, see Table 6),  $t_{\text{evol},i}$  is the time elapsed from the formation of the  $i$ -th binary (as a binary of two main sequence stars) to the birth of the second BH, and  $\text{SFR}_i$  is the current SFR of the  $i$ -th galaxy. For our sample, we assume  $t_{\text{evol},i} \approx 3 \times 10^6$  yr, which is the main sequence lifetime of the most massive stars. The values of  $\text{SFR}_i$  are from Table 6. Equation 2 gives a strong upper limit for the sample, for the fact that we neglect local galaxies that do not host any WR–BH binary.

The coalescence timescale  $t_{\text{GW}}$  can be expressed as (Peters 1964):

$$t_{\text{GW}} = \frac{5}{256} \frac{c^5 a^4 (1 - e^2)^{7/2}}{G^3 m_1 m_2 (m_1 + m_2)}, \quad (3)$$

where  $c$  is the light speed,  $G$  is the gravitational constant,  $a$  is the semi-major axis and  $e$  the eccentricity of the orbit. We assume  $e = 0$  for all considered binaries. Table 6 shows the values of  $t_{\text{GW}}$  for each considered binary, under these assumptions. The value of  $t_{\text{GW}}$  for M 101 ULX-1 is much larger than the Hubble time: the impact of this system is negligible for any estimate of the Advanced LIGO and Virgo detection rate. The second longest value of  $t_{\text{GW}}$  is  $\sim 2$  Gyr (for NGC 300 X-1): this value is well below the Hubble time, and implies that we might detect at present time the merger of systems like NGC 300 X-1 that formed  $\sim 2$  Gyr ago, i.e. at redshift  $z = 0.3$ , when the SFR density was a factor of two higher than now [ $\rho_{\text{SFR}}(z = 0.3) \sim 3.6 \times 10^{-2} M_{\odot} \text{ yr}^{-1} \text{ Mpc}^{-3}$ , Hopkins & Beacom 2006]. To obtain the best favourable upper limit, we use  $\rho_{\text{SFR}}(z = 0.3) = 3.6 \times 10^{-2} M_{\odot} \text{ yr}^{-1} \text{ Mpc}^{-3}$  for all sources in equation 2, and we obtain  $R \sim 4 \times 10^{-9}$  and  $1 \times 10^{-9} \text{ yr}^{-1} \text{ Mpc}^{-3}$ , if we include or do not include Cyg X-3 among the WR–BH candidates, respectively. Since Advanced LIGO and Virgo will detect BH–BH mergers out to  $\sim 1$  Gpc, we can infer upper limits to the detection rate of BH–BH mergers of  $\sim 16$  and 5 events per year, if we include or do not include Cyg X-3, respectively. Such rates are consistent with those predicted in previous papers that study WR–BH binaries (e.g. Belczynski et al. 2013; Maccarone et al. 2014). It is worth noting that our upper limit is quite close to the

'pessimistic' estimate of the BH–BH merger rate by the Advanced LIGO–Virgo collaboration (whose pessimistic and optimistic rates are  $R = 10^{-10}$  and  $3 \times 10^{-7} \text{ yr}^{-1} \text{ Mpc}^{-3}$ , respectively; Abadie et al. 2010).

## 11 SUMMARY AND CONCLUSIONS

We reported on our timing survey of the deep *Chandra*/ACIS observations of the Circinus galaxy and its surroundings (Fig. 1). Approximately 150 X-ray sources were detected and for about 40 of them, enough photons were collected to search for periodic signals by means of a Fourier transform. We discovered two new X-ray pulsators, CXO J141430 and the uncatalogued CXOU J141332, and bumped into the already known periodic modulation of the notable and controversial source CG X–1 ( $P \simeq 7.2 \text{ h}$ ; Bauer et al. 2001). For the other sources, we set upper limits on the presence of periodic signals (Fig. 2).

CXO J141430 is more than 2 arcmin out of the border of the CG. A spin period of 1.7 h and an orbital period of 17.8 h give the source away as an intermediate polar seen at large inclination. Its X-ray spectrum can be modelled by a power-law with photon index  $\Gamma \simeq 1.4$ . The flux in the longest *Chandra* observation was  $\approx 1 \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$  and other observations imply a variability of  $\approx 50$  per cent on time-scales of weeks–years. The typical luminosity of IPs and the nondetection of the optical counterpart suggest a distance larger than  $\sim 5 \text{ kpc}$ .

Albeit at  $\sim 3.5$  arcmin from the nucleus, CXOU J141332 appears inside the CG. However, the low absorption column, which is much smaller than the total Galactic density, argues against an extragalactic source. Indeed, the probability of a foreground Galactic X-ray source is substantial ( $\approx 10$  per cent). The period of this source is 1.8 h and the spectrum is hard: it can be described by a power law with  $\Gamma \simeq 0.9$ . The observed flux was  $\approx 5 \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$ , with  $\approx 50$  per cent variations on weekly/yearly scales. We believe that also CXOU J141332 is a Galactic magnetic cataclysmic variable, most probably of the polar type. Assuming that the companion is a M5V star (or similar), the nondetection of its optical counterpart is not particularly constraining for the distance, implying only  $d \gtrsim 0.7 \text{ kpc}$ . On the other hand, if the system is within a few kpc, its luminosity is in the normal range for polars.

We deem the association of CG X–1 with the CG convincing. Prompted by the similarity of its modulation to the distinctive ones of the known WR–HMXRBs and candidates, we advanced the possibility that CG X–1 might be one of such systems. The observations of CG X–1 (in particular, the high luminosity implied and the limits on the optical counterpart) are consistent with this hypothesis. The dearth of observed WR–HMXRBs is puzzling, since they should be relatively common and very bright (Lommen et al. 2005; Linden, Valsecchi, & Kalogera 2012). Though there is no obvious bias against their detection, it is possible that some of these objects might be misclassified or unrecognised. At any rate, in the last few years, the sample of WR–HMXRBs has grown rapidly to four confirmed sources and three candidates (including CG X–1), all of them outside our Galaxy, with the exception of Cyg X–3. Besides their relevance for the population of X-ray binaries and ULXs, WR–BH systems are very important as they might be progenitors of double-BH systems. We used the information from the current sample of systems to estimate an upper limit to the detection rate of stellar BH–BH mergers with Advanced LIGO and Virgo, which turned out to be  $\sim 16 \text{ yr}^{-1}$  for a distance range of 1 Gpc.

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