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# **The Thermal Structure of the Venus Atmosphere: Intercomparison of Venus Express and Ground Based Observations of Vertical Temperature and Density Profiles**

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## Abstract

The Venus International Reference Atmosphere (VIRA) model contains tabulated values of temperature and number densities obtained by the experiments on the Venera entry probes, Pioneer Venus Orbiter and multi-probe missions in the 1980s. The instruments on the recent Venus Express orbiter mission generated a significant amount of new observational data on the vertical and horizontal structure of the Venus atmosphere from 40 km to about 180 km altitude from April 2006 to November 2014. Many ground based experiments have provided data on the upper atmosphere (90-130 km) temperature structure since the publication of VIRA in 1985. The "Thermal Structure of the Venus Atmosphere" Team was supported by the International Space Studies Institute (ISSI), Bern, Switzerland, from 2013 to 2015 in order to combine and compare the ground-based observations and the VEX observations of the thermal structure as a first step towards generating an updated VIRA model. Results of this comparison are presented in five latitude bins and three local time bins by assuming hemispheric symmetry. The intercomparison of the ground-based and VEX results provides for the first time a consistent picture of the temperature and density structure in the 40 km - 180 km altitude range. The Venus Express observations have considerably increased our knowledge of the Venus atmospheric thermal structure above ~40 km and provided new information above 100 km. There are, however, still observational gaps in latitude and local time above certain regions. Considerable variability in the temperatures and densities is seen above 100 km but certain features appear to be systematically present, such as a succession of warm and cool layers. Preliminary modeling studies support the existence of such layers in agreement with a global scale circulation. The intercomparison focuses on average profiles but some VEX experiments provide sufficient global coverage to identify solar thermal tidal components.

The differences between the VEX temperature profiles and the VIRA below 0.1 mbar/95 km are small. There is, however, a clear discrepancy at high latitudes in the 10-30 mbar (70-80 km) range. The VEX observations will also allow the improvement of the empirical models (VTS3 by Hedin et al., 1983 and VIRA by Keating et al., 1985) above 0.03 mbar/100 km, in particular the 100-150 km region where a sufficient observational coverage was previously missing. The next steps in order to define the updated VIRA temperature structure up to 150 km altitude are (1) define the grid on which this database may be provided, (2) fill what is possible with the results of the data intercomparison, and (3) fill the observational gaps. An interpolation between the datasets may be performed by using available General Circulation Models as guidelines.

An improved spatial coverage of observations is still necessary at all altitudes, in latitude-longitude and at all local solar times for a complete description of the atmospheric thermal structure, in

particular on the dayside above 100 km. New in-situ observations in the atmosphere below 40 km are missing, an altitude region that cannot be accessed by occultation experiments. All these questions need to be addressed by future missions.

## 1. Introduction

A systematic global coverage of the temperatures in the Venus atmosphere was obtained by the VORTEX experiment (also called Orbiter Infrared Radiometer, Taylor et al. 1980; Schofield and Taylor, 1983) on board Pioneer Venus Orbiter (PVO) using a six channel filter radiometer. Density and temperature profiles retrieved from PVO radio occultations provided latitudinal variations of the structure in the 40-75 km altitude range. Deep atmospheric temperature profiles were obtained from the Venera probes (6 to 14) and by the four Pioneer Venus probes (named Large, Day, Night, and North) in 1982. These were the basic observations that led to the development of the thermal structure model compiled for the Venus International Reference Atmosphere (VIRA) published through the efforts of Pioneer Venus and Venera scientists (Kliore et al., 1985; Seiff et al., 1985). In the interim, limited compilations were prepared by Seiff (1983) and Moroz (1981). An empirical model of the Venus thermosphere (VTS3) was also developed based on the available data by Hedin et al. (1983). Many results, including ground based results that were developed just prior to the development of the VIRA model could not be included in the model. The VIRA profiles from the low atmosphere were compiled from Venera measurements and Pioneer Venus probes profiles. The latter were extrapolated adiabatically by the hydrostatic law from 12 km to the surface assuming a composition of 96.5% CO<sub>2</sub> and 3.5% N<sub>2</sub> (the Pioneer Venus probes suffered from an electrical failure when the probes were at 12 km above the surface).

The VIRA thermal structure model was found to be very useful by the Venus scientific community for further investigations of the planetary atmosphere. Many new observations of the thermal structure of Venus have been obtained since its publication: from Venera-15 Fourier spectrometry (1983), from Venera 15, Venera 16 and Magellan radio occultations in 1992, by the Venus Express orbiter since April 2006-till late 2014 from five independent experiments and by numerous ground based observations. A temperature profile of the low atmosphere was measured *in situ* with high vertical resolution by the VEGA-2 Lander in 1984. These new observations provided spatial and temporal overlap, extended the knowledge of the temperature structure downwards to the surface, revealed temporal and spatial variations. It is now possible to compare these results in order to understand the differences and reconcile them by looking at the experimental approaches, their inherent limitations and potential errors. Such intercomparison is a pre-requisite step for developing a new VIRA thermal structure model.

Zasova and Moroz (1992) and Moroz and Zasova (1997) reviewed the datasets that were collected between the publication of VIRA in 1985 and the publication of the respective papers. It was suggested to update the thermal structure model in view of the new datasets, particularly from the VeGa 2 lander, the two VeGa balloons and the detailed thermal structure of the mesosphere from

the Venera-15 Fourier spectrometer. Potential other sources for the improvement of VIRA were also addressed, using later radio occultation results from Pioneer Venus, Magellan, Galileo NIMS experiments and ground-based observations.

This study presents the results of an intercomparison of data from the Venus neutral atmosphere obtained after the publication of VIRA by an international team sponsored by the International Space Science Institute (ISSI), Bern, Switzerland from July 2013 to February 2015. Thermal structure observations that were obtained after the publication of VIRA in 1985 (Table 1) and prior to the arrival of Venus Express (VEX) in April 2006 include the following:

- Extended mission radio occultation profiles from Pioneer Venus Orbiter (Kliore, 1985)
- Radio occultation profiles from Venera-15, 16 (Yakovlev et al., 1991)
- VeGa 1 and VeGa 2 balloon data (Sagdeev et al., 1986)
- VeGa 2 Lander data (Linkin et al., 1987; Zasova et al., 2006)
- Retrievals of thermal profiles from Venera 15 Fourier spectrometer data (Schaefer et al., 1990, Zasova et al., 2006, 2007, Haus et al., 2013)
- Galileo NIMS fly-by observations (Roos-Serote et al., 1995)
- Magellan Orbiter radio occultation profiles (Jenkins and Hinson 1994)

Earth-based thermal structure observations obtained since Venus Express commenced operations include:

- Thermospheric ground based temperature structure profiles (Clancy et al., 2008; 2012a, Rengel et al., 2008a, 2008b, Sonnabend et al. 2008; 2010)
- 2012 Venus transit observations, deriving the temperature from the sunlight refraction in the mesosphere (Tanga et al., 2012; Pere et al., 2016)

The list of ground based measurements included in this study is certainly not exhaustive, but representative. The principal idea was to include those data sets which are accessible in digital form as much as possible to facilitate the detailed comparison.

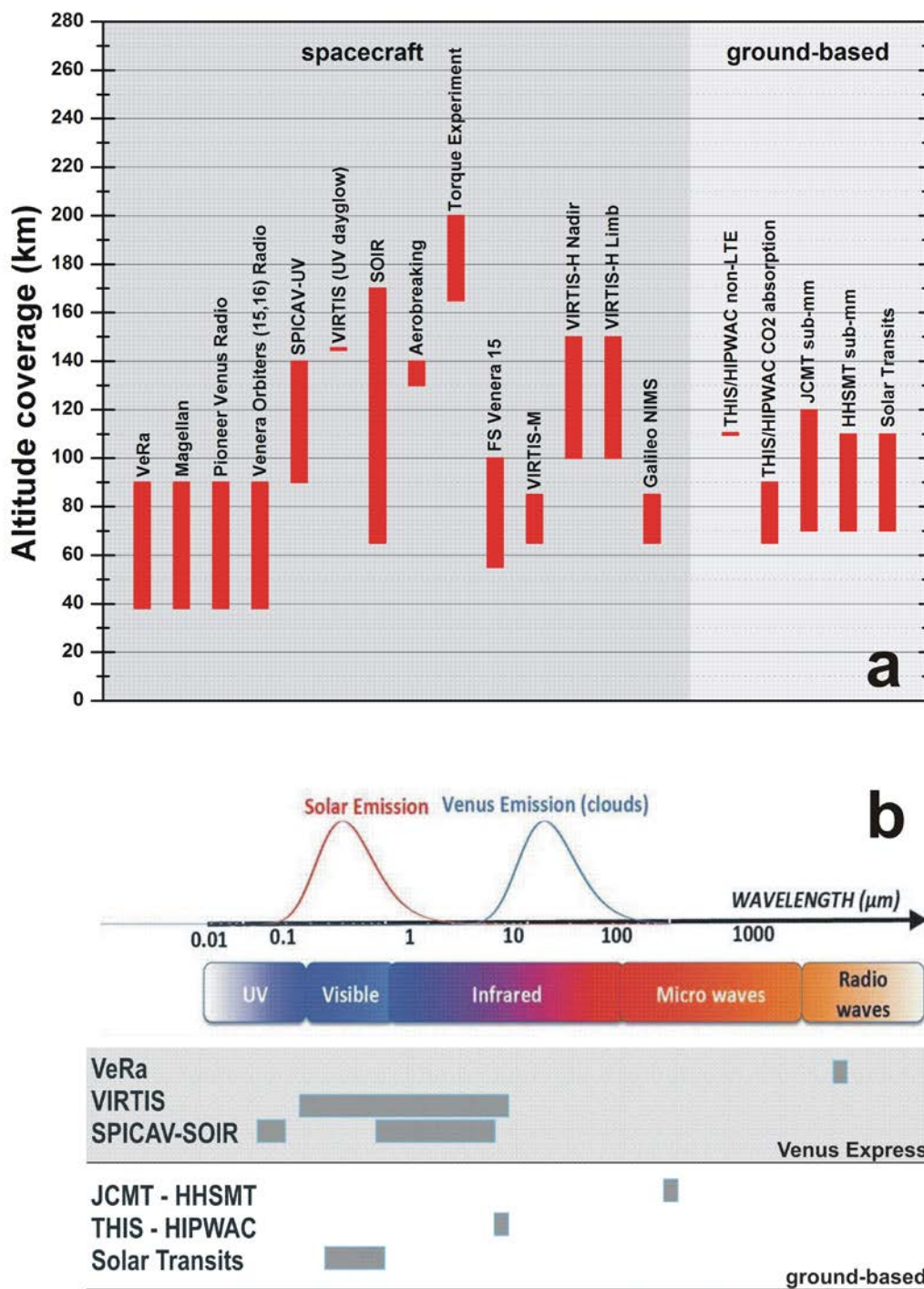
Five experiments operated from the Venus Express orbiter that yield atmospheric profiles of neutral number density and temperature versus altitude or pressure and in-situ atmospheric mass density from drag or aerobraking experiments:

- *Solar Occultation in the Infra-Red (SOIR)*: The solar occultation method retrieves vertical profiles of carbon dioxide abundance and atmospheric temperature from CO<sub>2</sub> number density as well as molecular rotational temperatures from CO<sub>2</sub> spectral structure (Bertaux et al., 2007) at the morning and evening terminators at occulted latitudes

- *Spectroscopy for Investigation of Characteristics of the Atmosphere of Venus (SPICAV)*: Stellar occultations allow the determination of vertical profiles of CO<sub>2</sub> abundances and derive the temperature from the CO<sub>2</sub> number density (Bertaux et al., 2007)
- *Venus Express Radio Science (VeRa)*: Radio occultations allow the derivation of vertical profiles of temperature, pressure and total neutral number density between 40 km and 100 km altitude (Häusler et al., 2006; 2007)
- *Visible and Infra-Red Thermal Imaging Spectrometer (VIRTIS)*: VIRTIS observations provide thermal maps at medium spectral resolution and profiles from nadir and limb locations at high spectral resolution (Piccioni et al., 2007)
- *Venus Express Atmospheric Drag Experiment (VEXADE)*: VEXADE retrieved atmospheric mass density (i) between 130-140 km from accelerometer readings during aerobreaking (Müller-Wodarg et al., 2006; 2016), (ii) from the torques acting on the solar panels by the atmospheric drag between 166 km to 186 km at high planetary latitudes (Persson, 2015), and (iii) from Precise Orbit Determination (POD) when the orbiter was between 166 – 186 km altitude (Rosenblatt et al., 2012).

The VEX and ground-based data sets which were considered in this study are described, as well as the experimental approaches, a discussion of the comparison and recommendations regarding future observations are given.

The altitude ranges of the post-VIRA experiments conducted at Venus and the spectral ranges of the experiments considered in this study are shown schematically in Figure 1.



**Figure 1:** Panel (a): vertical coverage of the post-VIRA atmospheric structure experiments. Panel (b): spectral ranges of the experiments considered in this study

**Table 1: Observations of the vertical structure of the Venus atmosphere by spacecraft and ground based experiments (data from experiments in light lines are not used in the intercomparison)**

Instrument / Experiment	Method	Years covered	Nominal vertical coverage [km]	Pressure range [mbar]	Vertical resolution [km]	Temperature uncertainties [K]	Latitudinal coverage	Local time coverage	Horizontal resolution	Sensitive to clouds?	section
<b>Active remote sensing observations from spacecraft</b>											
VeRa/VEX	radio occultation	2006-2014	38 - 100	3000 - 0.03	0.5 - 1	0.1 - 1	both hemispheres	night & day side	slant paths, 400 km	NO	2.3.1.2
Magellan	radio occultation	1992	38 - 100	3000 - 0.03	0.5 - 1	0.1 - 1	both hemispheres	night & day side	slant paths, 400km	NO	n/a
Pioneer Venus	radio occultation	1978-1991	38 - 100	3000 - 0.03	0.5 - 1	0.1 - 1	both hemispheres	night & day side	slant paths, 400km	NO	2.3.1
Venera 15, 16 orbiters	radio occultation	1983	38 - 100	3000 - 0.03	0.5 - 1	1 - 10	both hemispheres	night & day side	slant paths, 400km	NO	2.3.1.1
SPICAV-UV/VEX	stellar occultation	2006-2014	90 - 140	$10^{-1} - 10^{-7}$	0.5 - 7	<25% 1 - 20 K Altitude dependent	both hemispheres	night side	slant paths, 400km	NO	2.3.3
SOIR/VEX	solar occultation	2006-2014	70 - 170	$100 - 10^{-8}$	0.3 - 5 (lat. dep.)	1 - 20	both hemispheres	terminator	slant paths, 400km	NO	2.3.2
VExADE-AER/VEX	Aerobraking	2014	130 - 140	$10^{-5} - 10^{-6}$	0.3	23	70°N - 80°N	morning terminator	10 km	NO	2.1.2.1
VExADE-TRQ/VEX	Spacecraft torque measurement	2008-2013	165 - 200	$10^{-7} - 10^{-9}$	1.0	~ 30 K	70°N - 90°N	78 - 98° Solar Zenith Angle	Slant paths	NO	2.1.2.2
VExADE-POD/VEX	Precise Orbit Determination	2008-2013	175-185	$10^{-8}$	n/a	n/a	70°N - 90°N	terminators	n/a	NO	2.1.2.2
<b>Passive Remote Sensing (IR/Microwave) from spacecraft</b>											
FS VENERA-15	15 $\mu\text{m}$ CO <sub>2</sub> temperat.-aerosol ,	1983	55-100 km	300 - 0.03	3 - 5 (scale)	2 - 5 (altitude)	mostly Northern	4 -10 AM 4 -10 PM	60 km at pericenter	YES, self-consistent retrieved	2.1.1

	7-30 $\mu\text{m}$				height)	dependent)				from each spectrum	
VIRTIS-M/VEX	4.3 $\mu\text{m}$ CO <sub>2</sub> band	2006-2008	65 - 85	100 - 0.1	8	< 5	mostly southern hemisphere	night side	60 km at pericenter	YES	2.2.3.1
VIRTIS-H/VEX Nadir	4.3 $\mu\text{m}$ CO band	2006-2012	65 - 80	100 - 4	8	< 5	both hemispheres	night side	50km footprint (individual retrievals)	YES	2.2.3.2
VIRTIS-H/VEX Limb	4.7 $\mu\text{m}$ CO <sub>2</sub> band	2006-2012	100 - 150	0.03 - 10 <sup>-7</sup>	15 – 25 (altitude dependent)	> 30	North hemisphere	day side	115x38 km	likely	2.2.3.2
Galileo NIMS	4.3 $\mu\text{m}$ CO <sub>2</sub> band	1990	65 - 85	100 - 0.1	8	< 5	South hemisphere	night side	n/a	NO	2.2.2
<b>Ground-based observations</b>											
THIS / HIPWAC	CO <sub>2</sub> non-LTE emission	1990-1991 2007-2014	110 km	2·10 <sup>-3</sup>	+/- 10	10	both hemispheres	day side	0.9" to 1.6"	NO	3.2
THIS / HIPWAC	CO <sub>2</sub> absorption	2012	65 - 90	100 - 0.8	10	< 10	both hemispheres	night side	0.9" to 1.6"	NO	3.2
JCMT sub/mm line absorption	CO absorption	2001-2015	75 – 120	20 - 10 <sup>-4</sup>	4 10 above 100 km altitude	7	both hemispheres	mapping PM/AM night side Dayside PM/AM average	13.5" to 14.5" 4000 km sub-earth footprint	NO	3.1.1
HHSMT sub/mm line absorption	CO absorption	2007	75 - 110	20 - 0.002	4 10 above 100 km altitude	< 15	both hemispheres	PM night & day side averages	13.5" to 14.5" 13.5" to 14.5" 10000 km sub-earth footprint	NO	3.1.2
Various space and ground-based telescopes	Photometry (imaging of Venus transits)	2004 2012	70 - 110	100 – 0.002	5	10 - 20	All, simultaneous	terminator	Slant path, 400 km	YES	3.4

## 2. Spacecraft Observation Methods

### 2.1 Direct (in-situ) measurements

#### 2.1.1 Entry Probes/Landers and Balloons

Atmospheric *in-situ* measurements after the observations by the Venera 13 and Venera 14 entry probes in 1982 were made by the two VeGa balloons (Sagdeev et al., 1986) and by the VeGa landers in 1985 (Linkin et al., 1987). Each of the two VeGa spacecraft consisted of a carrier spacecraft with a Venus lander descending to the surface and a balloon that was deployed from a separate entry capsule at an altitude of about 50 km. VeGa-1 entered the Venus atmosphere on 11 June 1985, VeGa-2 followed four days later. The VeGa 1 lander communications failed and no data could be transferred. The two carrier spacecraft went on to rendezvous with comet Halley in 1986.

##### 2.1.1.1 VeGa balloons

The VeGa 1 balloon entered the atmosphere at 8.1°N latitude, 176.9°E longitude, and the VeGa 2 balloon at 7.45°S latitude and 179.8°E longitude. All measurements on the VeGa balloons were performed successfully during their journey through the middle clouds at an altitude of about 54.5 km (Linkin et al., 1986, Sagdeev et al., 1986). The two balloons observed a near constant temperature difference of about 6.5 K when carried westward by the ambient winds at average speeds of 69 m/s and 66 m/s. The VeGa-1 balloon moved almost exactly along the 8° North latitude and travelled nearly 8,500 km in the darkness of the Venus night before crossing the morning terminator. The trajectory of VeGa-2 was shifted by about 500 km southward and floated at a mean altitude of 53.6 km (535 mbar) and experienced temperatures ranging from 308 K to 316 K. The communication with the balloons was lost when the batteries drained after 40 hours of operations. The values of pressure and temperature along the trajectories of the balloons are given in Table 2.

**Table 2: VeGa-1 and VeGa-2 balloons**

	Start of operation					Temperature and pressure at the balloon altitude			
	Date	time (hours UT)	latitude	longitude	LT (hours)	pressure (mbar)		temperature (K)	
						(*)	(**)	(*)	(**)
VeGa-1	11 June 1984	02:06	8° N	77°	00:18	540	630	308	322
VeGa-2	15 June 1984	02:06	7.5° S	180°	01:00	535	900	302	338

\* at maximum floating altitude (54 km)

\*\* at minimum floating altitude (VeGa-1: 53 km, VeGa-2: 50 km)

The pressure dependence of the temperature is close to adiabatic at the floating altitude of both balloons, with a temperature difference of a few Kelvins. This was interpreted as an indication of the existence of sufficiently extended non-mixing atmospheric masses (Linkin et al., 1986). Each balloon during the flight was inside its own region of this type.

#### 2.1.1.2 VeGa-1 and -2 Landers

The VeGa-1 (7.2°N, 177.8°E entry location into the atmosphere) and VeGa-2 (8.5°S, 164.5° entry location) landers were designed like the earlier Venera landers and carried well calibrated and redundant temperature sensors. The VeGa-1 lander experienced a strong updraft during its descent, well before reaching the surface, causing the control electronics to believe that it had landed and thus some instruments were deployed prematurely and consequently not all the planned measurements were successfully acquired.

The VeGa-2 lander remains the only probe of all Venus landers which observed the atmospheric temperature all the way from 64 km down to the surface. VeGa-2 landed at 6.45°S latitude and 181.08° longitude, which implies a drift toward the equator during its descent through the atmosphere.

#### 2.1.2 Mass Density Measurements from the Venus Express Atmospheric Drag Experiment

Thermosphere densities were measured in-situ by the Venus Express Atmospheric Drag Experiment (VExADE), which itself consisted of 3 separate experiments, aerobraking (VExADE-AER), Precise Orbit

Determination (VExADE-POD) and torque measurements (VExADE-TRQ). Both the POD and TRQ experiments were carried out during the main science phase of Venus Express (2008-2013) during campaigns when pericentre altitude ranged from 165-190 km, while aerobraking (VExADE-AER) was performed in June/July 2014, at the end of the nominal science mission when the pericentre altitude was lowered to 130 km. For pericentre altitudes of 165-190 km the atmospheric drag experienced by the spacecraft is strong enough to affect its orbit and be measured by radio tracking techniques (POD), though too weak to be detected by the onboard accelerometers. A series of Atmospheric Drag Experiments (ADE) was performed by lowering the pericentre to altitudes between 165 km and 190 km. By tracking the spacecraft at high resolution with the Deep Space Network and subsequently modeling the spacecraft orbits, the integrated deceleration experienced during each pericentre pass was derived, and thereby a single density value for the pericentre location of each spacecraft pass (Rosenblatt et al., 2012). This provided the first *in-situ* measurements of thermospheric mass density at high latitudes (75°N – 90°N) and at low solar activity, finding mean densities to be around 60% of those predicted for the same latitudes by the VTS3 model (Rosenblatt et al. 2012). The spacecraft torque measurements (TRQ) were carried out during the same campaigns but consisted in analyzing the response of the spacecraft's Inertial Mass Unit (IMU) to the torque experienced by asymmetric orientation of the two solar panels relative to the ram direction. Thereby, the torque experiment obtained vertical density profiles from 165-190 km, similar to what the aerobraking experiment obtained for lower altitudes (130-140 km), while the POD experiment gave a single density value at the pericentre altitude during every orbit of the POD campaign.

All three VEX drag datasets – i.e. thermospheric densities from radio tracking, from torque, and from aerobraking – show considerable and significant diurnal variability, with day-to-day mass densities often varying by over 100%. The aerobraking and torque data show significant variability even within each pass, with horizontal wavelengths on the order of 100-200 km which may be associated with gravity waves. Both the day-to-day and the intra-orbit variability are similar to phenomena which have been observed at similar pressure levels in the Martian thermosphere (e.g. Fritts et al. 2006).

The PVO, Magellan and VEx missions obtained atmospheric drag data, but at different locations and local times. While the PVO-ONMS mass spectrometer and PVO-aerobraking sampled the thermosphere at low latitudes, the in-situ data from VEx were taken at polar latitudes. The local solar time coverage is also different: PVO mass spectrometry and aerobraking covered all local times, while VEx sampled the terminators, with SZA in the range 80-100°.

### 2.1.2.1 Venus Express aerobraking (VExADE-AER)

Vertical profiles of total mass densities in the thermosphere were inferred from accelerometer measurements on Venus Express during the aerobraking campaign from 24 June to 11 July 2014 (Mueller-Wodarg et al., 2016). ESA planned this campaign in the final months of the mission in view of the risk of losing the spacecraft due to the enhanced atmospheric drag at decreasing altitude. The pericentre of the highly eccentric VEX orbit ( $e=0.84$ ) was located at  $75^{\circ}\text{N}$  at local solar times of 04:30 to 06:12 and altitudes of 130 to 134 km. Data from the on-board accelerometers could not be used at higher altitudes for the derivation of mass densities because of the insufficient sensitivity of the (engineering) instrument.

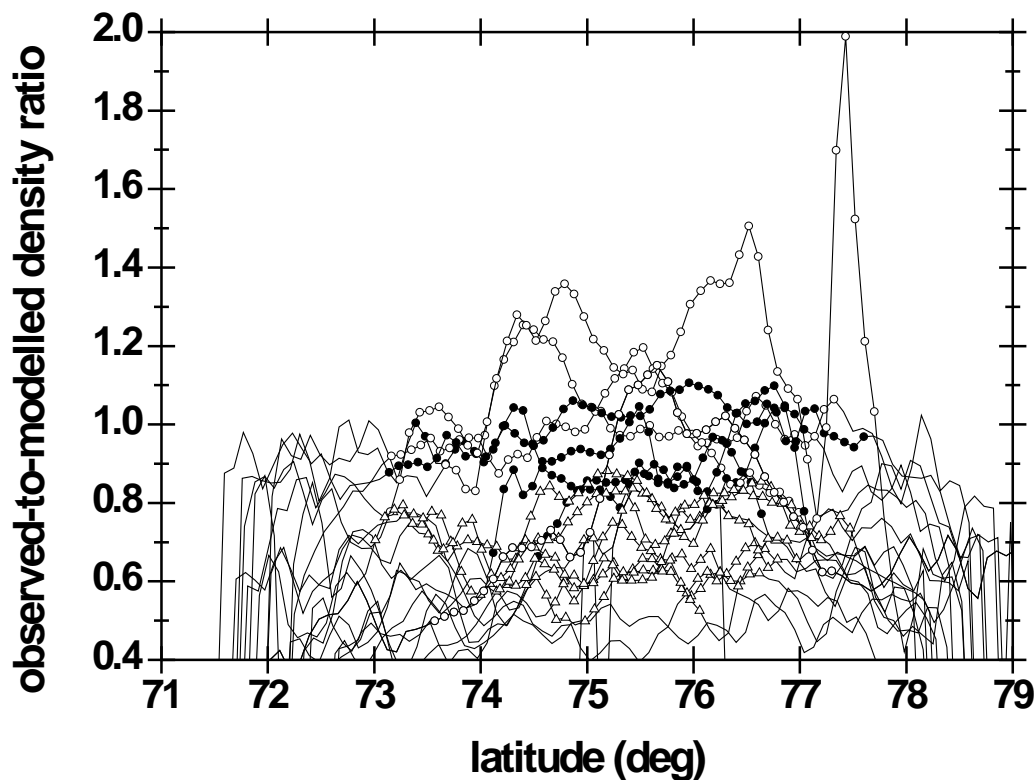
The raw accelerometer data taken at 8 samples per second around the pericentre were averaged over 2 s, and resampled at 1 Hz in the density processing. Density profiles at 1 Hz sampling extending about  $3^{\circ}$  in latitude before and after the pericenter have been obtained for each of the 18 consecutive days of the aerobraking campaign in 2014. The general method of the derivation of mass densities from accelerometer measurements is described by Bruinsma et al. (2004, 2006).

The uncertainty of the derived mass densities was computed from a systematic part caused by the uncertainty in the spacecraft aerodynamic coefficient which was estimated to be 10%, plus the measurement noise and bias of the accelerometer. This systematic uncertainty had no impact on the analysis of the relative variations within a single orbit, for example wave structures. The (formal)  $1\text{-}\sigma$  noise of the accelerometer data averaged over two seconds was found to be  $0.001\text{ m/s}^2$ . The bias of the accelerometer was estimated to  $2\cdot 10^{-4}\text{ m/s}^2$  to  $5\cdot 10^{-4}\text{ m/s}^2$  from measurements outside the sensitivity range of  $2\cdot 10^{-3}\text{ m/s}^2$  at higher altitudes. Taking this uncertainty into account, the density data can be used on average to an altitude of 139 km, which corresponds to profiles of about 80 seconds duration.

The mass densities observed by VEX are compared with an empirical model. Ratios of the observed VEX mass density with those from the VTS3 model by Hedin et al. (1983) were computed for each profile (Figure 2). Valid observed mass densities are on average about 30% smaller than densities from the VTS3 model, that means in better agreement with the Hedin model than the mass densities obtained by the Precise Orbit Determination from radio tracking (Rosenblatt et al., 2012) at higher altitudes (160 km -170 km). A high variability of the ratio of  $\sim 10\%$  is seen in form of wave-like features along the orbit. The ratio of observed densities to modelled densities is altitude-dependent, being smaller than 1 at lower altitudes (about 0.78, or 78% near 130 km altitude) and decreasing with altitude (60% near 140 km altitude). This demonstrates a systematic difference between the neutral scale heights of the

observed densities and the VTS3 model densities. These differences are most likely caused by temperature differences in the polar thermosphere and possibly uncertainties in our knowledge of the polar atmospheric composition.

Temperatures are derived from the neutral atmospheric scale heights  $H = k \cdot T / (m \cdot g)$  where  $T$  is the temperature,  $k$  is the Boltzmann constant,  $g = 8.49 \text{ m/s}^2$  is the gravity acceleration, and  $m$  is the mean molecular weight of the atmospheric species which is estimated using the VTS3 model. VTS3 predicts a mean molecular mass  $m = 34.7 - 41.8$  atomic mass units (amu) for the latitude range  $71.5^\circ\text{N}$  to  $79.0^\circ\text{N}$ , the local solar time (LST) range 04:30 h to 06:18 h and  $F10.7_{\text{mean}} = 130.7 - 134.0$  using the 10.7 cm radio flux as a proxy for the solar flux. The daily F10.7 proxy varied between 93.4 and 200.7 during the time of observations. A mean temperature of  $114 \pm 23 \text{ K}$  was derived from the observed mass density profiles. The VTS3 model temperatures are higher for the same observing conditions: 141 K - 159 K. This temperature difference is consistent with the differences in scale height mentioned above.



**Figure 2:** Ratio of observed VEXADE mass densities versus VTS3 model mass densities. Only accelerations above  $0.003 \text{ m/s}^2$  (3-sigma) are shown. Symbols mark those ranges where the acceleration is  $> 0.01 \text{ m/s}^2$ . The profiles are observations from 24 June to 11 July 2014. The ratios are computed using a 16-data point average.

### 2.1.2.2 Venus Express Torque Experiment (VExADE-TRQ)

In addition to calculating thermospheric densities from spacecraft tracking and from accelerometry, atmospheric density can be calculated by measuring aerodynamic torque exerted on the spacecraft as it travels through the thermosphere, as measured by the spacecraft's attitude control system. The total torque acting on the spacecraft also includes contributions from gravity field gradients and the solar radiation pressure: These two terms need to be modelled and subtracted from the total measured torque. The remaining torque is then caused by the atmospheric drag force acting on the solar panels. This technique was first demonstrated using the Magellan orbiter (Croom & Tolson, 1994), initially with the spacecraft in a normal flight configuration but later with its solar arrays set in asymmetrical orientations in order to create larger torque forces on the spacecraft at a given atmospheric density (the so-called "windmill" experiment, see Tolson et al., 1995). This technique was then further developed during the Venus Express mission; the observation and data reduction procedure for VEX, and its validation by comparison with using radio tracking data are described in detail by Damiani et al., 2012 and Persson, 2015. The torque experiment allowed calculation of atmospheric densities at altitudes of 165 – 200 km.

Venus Express torque measurements were performed during approximately 100 pericentre passes below 200 km altitude between 2008 and 2014 at latitudes between  $75^\circ\text{N}$  and  $90^\circ\text{N}$ . Like the density measurements from VEX aerobraking, discussed above, all density measurements from the VEx torque investigation were carried out near the terminator (SZA of  $80 - 100^\circ$ ); since, for thermal and operational reasons, pericentre-lowering for aerobraking or torque measurements was only carried out when the orbital plane was nearly perpendicular to the Sun-Venus vector.

Figure 3 shows one example of atmospheric mass densities derived by the torque method during a single pericentre passage on 18 May 2011 (Persson, 2015). The reader is reminded that Venus Express had a highly elliptical polar orbit, with a pericentre at high northern latitudes. In this particular orbit, the pericentre was at a latitude of  $84.8^\circ\text{N}$  above the dayside near the evening terminator (Local Solar Time = 16:38); the spacecraft approached pericentre travelling northwards above the dayside, crossing the terminator to the nightside 133 seconds after pericentre. It can be seen that the densities measured

after pericentre, when the spacecraft is approaching the terminator, are markedly lower than those measured before pericentre; this sharp density gradient near the terminator is consistent with previous observations such as those from Pioneer Venus Orbiter [Keating et al. 1980]. Strong oscillations in the atmospheric mass density are evident in many torque passes (Persson, 2015), with horizontal wavelength typically in the range 100 – 300 km, similar to those observed at 130-145 km altitude in aerobraking data (Müller-Wodarg et al., 2016).

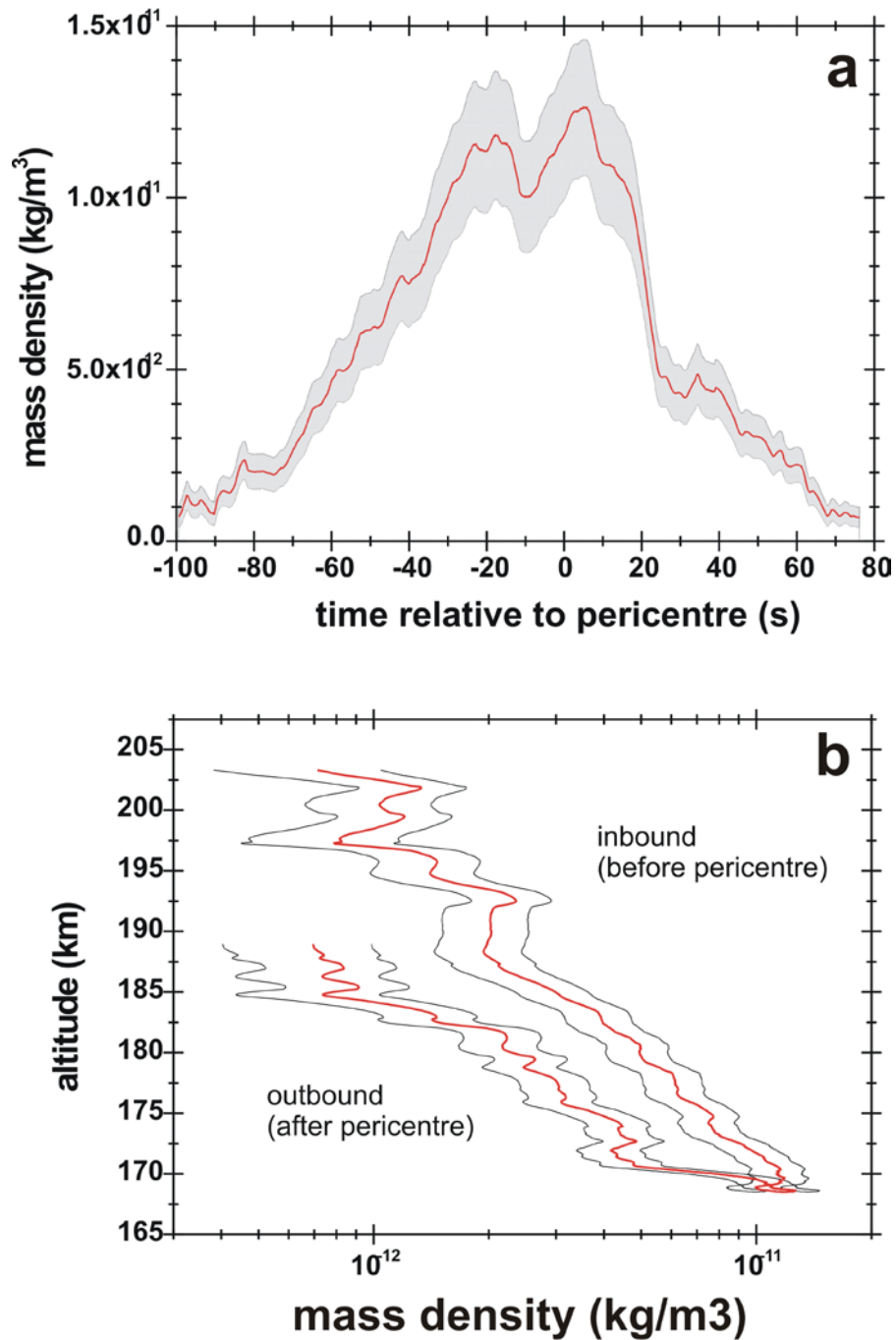
The density profile from all of the torque passes was binned by altitude; due to the strong gradient with respect to solar zenith (SZA) angle all data were also normalized to 90° SZA (for details, see Persson 2015). The resulting vertical profiles, binned separately for morning and evening terminators, are shown in Figure 4. The error bars denote the measurement error as a function of altitude; the solid lines show the +/- 1 sigma dispersion of measured densities in each altitude bin, i.e. the standard deviation of mass density variability in each altitude bin. Any differences in density between morning and evening terminators are smaller than the measurement error.

For all the VEx torque data, mass densities were found to be 40 to 45% less than those predicted by the Hedin model, as was found in results from aerobraking at 130-140 km altitude (Sec. 2.1.2.1); this again indicates a lower thermospheric mass density at polar latitudes than at low latitudes observed by Magellan and PVO missions (Keating & Hsu 1993).

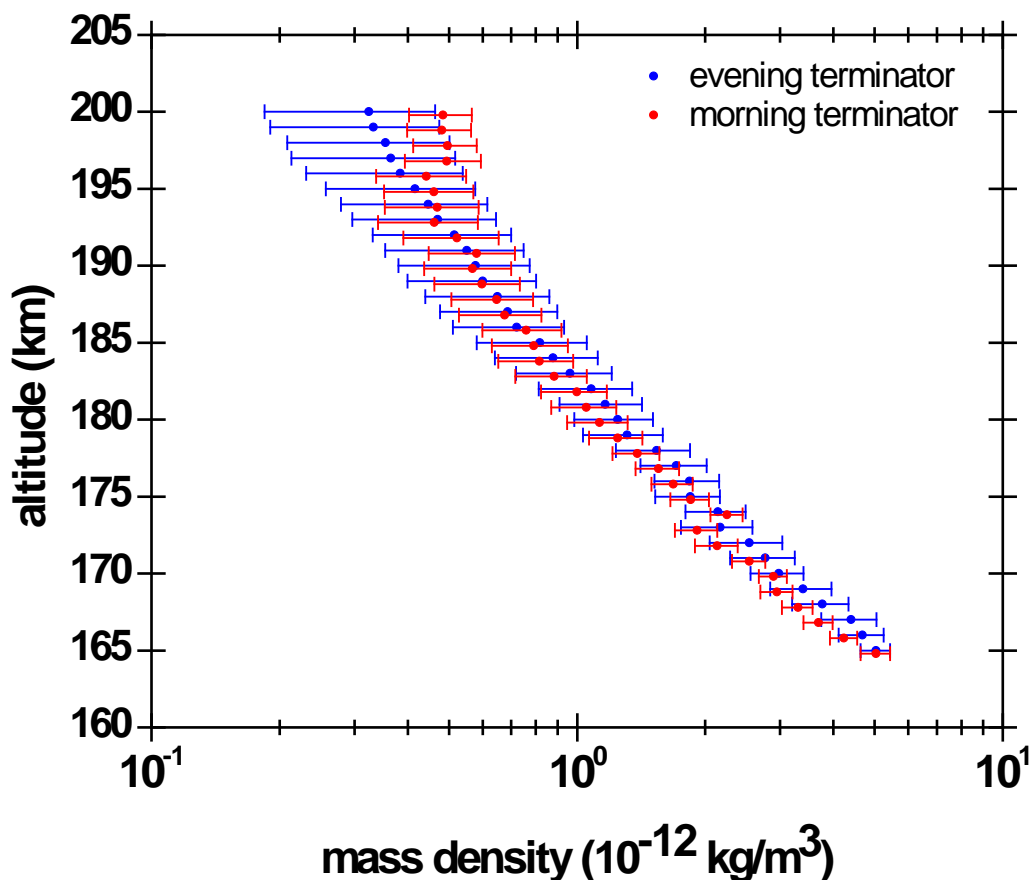
## 2.2 Passive Near-Infrared Observations

### 2.2.1 Venera-15 and Pioneer Venus

First maps of the atmospheric thermal structure were produced by the Orbiter Infrared Radiometer (OIR) on Pioneer Venus using a six channel filter radiometer (Taylor et al., 1980; Schofield and Taylor, 1983). The Fourier Spectrometer (FS-V15) on Venera 15 orbiter (Oertel et al. 1985, 1987, Moroz et al., 1986) observed emitted radiation from the Venus atmosphere in the range 250 - 1650  $\text{cm}^{-1}$  (6 – 40  $\mu\text{m}$ ) at a resolution of 4.5 or 6.5  $\text{cm}^{-1}$ . Its measurements yielded atmospheric properties above the clouds (Schaefer et al. 1987, 1990; Spänkuch et al. 1990). Spectral profiles derived from the 15  $\mu\text{m}$   $\text{CO}_2$  band (and also from both the  $\text{CO}_2$  hot 950 and 1050  $\text{cm}^{-1}$  bands and the isotopic 1260  $\text{cm}^{-1}$  band) and from spectral ranges which are free from gaseous absorptions were used to retrieve the vertical temperature and aerosol profiles from 55 km to 95 - 100 km altitude (Zasova et al., 1999, 2004, 2006, 2007).



**Figure 3:** Thermospheric neutral mass density derived from VEX torque measurements on 18 May 2011. Panel (a) derived mass density as a function of time relative to the pericentre (dashed vertical line) on 18 May 2011. Panel (b): derived mass density versus altitude. (Persson, 2015).



**Figure 4.** Mean mass densities along the morning (red) and evening (blue) terminator averaged over 1 km altitude bins. Error bars show measurement errors for individual measurements (Persson, 2015).

### 2.2.2 Galileo NIMS

The first space-based spectral maps of the Venus night side at a fairly high phase angle were produced by the NIMS experiment on the Galileo spacecraft (Carlson and Taylor, 1993; Roos-Serote et al., 1995) during its flyby at Venus in 1990. Temperature profiles were retrieved from the 4.7  $\mu\text{m}$  band (Roos-Serote et al., 1995) between 75 and 91 km altitude and latitudes between 59° S and 64° N. The temperatures were found to be about 10 K higher at 91 km and about 4 K cooler between 74 and 83 km when compared with the VIRA model, which was well within the variability of the VIRA model.

### 2.2.3 VEX VIRTIS

The Visible and InfraRed Thermal Imaging Spectrometer (VIRTIS) on Venus Express operates in two infrared modes in addition to a visible channel (200-1000 nm, M-vis channel) (Piccioni et al., 2007): a

high spectral resolution mode (H channel) and a lower spectral resolution mapping mode (M channel). Temperature profiles and spatial maps were derived by different groups using slightly different retrieval methods from the H and M channels (Grassi et al., 2008; Migliorini et al., 2011; Arnold et al., 2012, Gilli et al. 2015; Haus et al, 2014, Grassi et al., 2014; Garate-Lopez et al., 2015).

Radiation with wavelengths  $\lambda > 4 \mu\text{m}$  which is observed on the night side of Venus is driven by the thermal emission of the atmosphere. The VIRTIS-M spectral range covers in particular the strong  $\text{CO}_2 \nu_3$  band. The resolution of VIRTIS-M in this spectral region is such that for the different points on the spectral sampling grid of the instrument unit optical thickness is achieved at various altitudes between the cloud deck top (62-70 km) up to 80-85 km. This allows reconstructing the vertical temperature profile from the observed radiance at different wavelengths inside the band. The non-LTE emission by  $\text{CO}_2$  induced by direct solar radiation, however, does not allow daytime observations. Numerical experiments of simulated observations demonstrated how the random component of the retrieval error remains below 2 K in the range 3 - 50 mbar ( $\sim 81 - 68$  km altitude) and below 5 K for pressures below 1.2 mbar ( $\sim 85$  km) and above 90 mbar ( $\sim 65$  km). The main source of error in the upper atmosphere (above 85 km) is defined by the instrumental random noise and residual calibration misfits because the band is so opaque that an expected radiation level falls below the noise level. The retrieval becomes more and more difficult in the lower mesosphere (below 65 km) due to the opacity induced by the clouds. The retrieval results are sensitive to intrinsic vertical smearing related to the finite width of the weighting functions: the vertical resolution of the retrieval is roughly in the order of 7.5 km. The methods to constraint the air temperatures at different altitudes from VIRTIS-M data are described by Haus et al. (2013, 2014), Grassi et al. (2014) and references therein.

### 2.2.3.1 Temperature and Cloud Parameter Retrievals from VIRTIS-M-IR Data

VIRTIS-M is a mapping spectrometer like NIMS on Galileo, capable of acquiring simultaneously spectra at each of the 256 spatially-contiguous pixels along a line of the incident image. These so-called 'cubes' are acquired with multiple exposures, scanning the line over the disk of the planet by either using the instrument pointing mirror or directly by the spacecraft motion (when closer to the orbit pericenter). The cubes provide a spectrum for each pixel of the image, each pixel covering an instantaneous field-of-view of  $250 \mu\text{rad}$ . This implies to an area of  $16.5 \times 16.5$  km on the Venus cloud deck for measurements acquired at the VEX apocenter in nadir viewing mode. VIRTIS-M operates simultaneously in the visible and near infrared spectral ranges, but only the latter being relevant for the thermal structure reported here. The infrared spectral channel covers the range 1 -  $5.1 \mu\text{m}$  with an effective spectral resolution of

12 nm. VIRTIS-M-IR data were collected between 14 April 2006 and 29 October 2008 which corresponds to about 930 Earth days or 8 Venus solar days. The local time distribution of the measurements during the mission remained however quite irregular due to operational and orbital constraints which limited the downlink capability. Most of the cube data were acquired at the apocenter (located above the South Pole) covering large areas of the southern hemisphere. A small fraction of the cubes (about 5%) are long and narrow stripes over the equatorial region and the northern hemisphere at much higher spatial resolution.

Haus et al. (2013, 2014) followed an elaborated data pre-processing pipeline that includes refinements of data calibration procedures, new approaches for an effective stray light removal (Kappel et al., 2012), and data binning into a local time (LT) and latitude (lat) grid for grid spacing of  $\Delta LT = (0.5 \pm 0.1)$  h and  $\Delta lat = (5 \pm 1)^\circ$ . New methodical approaches for self-consistent temperature profile and cloud parameter retrievals are applied where combined radiative transfer and multi-window retrieval techniques simultaneously process information from different spectral ranges of an individual spectrum. The radiative transfer model is based on DISORT (Stamnes et al., 1988). Mesospheric temperature altitude profiles (58-90 km) are determined from 4.3  $\mu\text{m}$  CO<sub>2</sub> absorption band signatures using Smith's relaxation method (Smith, 1970). Specific parts of the 4.3  $\mu\text{m}$  band wings as well as of the deep atmosphere transparency window at 2.3  $\mu\text{m}$  are utilized to derive cloud parameters (cloud top altitude, mode abundance factors, opacity). Cloud parameter retrievals are based on a four-modal initial cloud model (Haus et al., 2013) where all modes are assumed to consist of spherical H<sub>2</sub>SO<sub>4</sub> aerosols at 75 wt% solution. Wavelength-dependent microphysical parameters of each mode are calculated applying a Mie scattering algorithm (Wiscombe, 1980) and log-normal size distributions and dispersions according to Pollack et al. (1993). Refractive index data is taken from Palmer and Williams (1975) and Carlson and Anderson (2011). Quasi-monochromatic gaseous absorption cross-sections are calculated on the basis of a line-by-line procedure considering spectroscopic parameters from the Venus-HiTemp and CDS line databases (Pollack et al., 1993; Tashkun et al., 2003) in case of CO<sub>2</sub>.

Zonal averages of derived temperature profiles at mid and high latitudes are in good agreement with VIRA profiles while, however, lower temperatures are found at low latitudes. The temperature decreases with increasing latitude polewards in both hemispheres starting at 30° latitude for fixed altitudes below the cold collar (50° - 75°, 58 - 70 km) while it increases above 70 km polewards starting at 40° - 50°. The cold collar and the polar vortex regions show the strongest temperature variability with standard deviations of up to 8.5 K at 75°S and 63 km altitude. The mesospheric temperature field

depends strongly on local time. The atmosphere is essentially warmer at early night and colder at late night by about 8 K in the cold collar. The temporal temperature trend reverses at higher altitudes.

Grassi et al. (2014) averaged pixels in the spatial domain on a 4x4 pixel basis before deriving the temperature: this step allowed to substantially increase the signal-to-noise ratio at the lower limit of the band and to mitigate spatial non-uniformity in the treatment of instrumental response. The resulting averaged spectra were processed using a Bayesian retrieval method (Rodgers, 2000) in order to derive air temperatures at a fixed pressure grid, an altitude-independent CO mixing ratio and a scalar multiplier for aerosol densities and to model variations at the cloud deck altitude. The retrieval method requires an initial guess to derive the temperature profile. This initial guess is taken as the mean value of the Venera-15 FTS temperature estimates. In order to cope with the limitations of the simplified forward radiative transfer model adopted for the computations, the retrievals were limited to cases with emission angles smaller than 30°. The retrieval method considers all four aerosol modes described by Knollenberg and Hunten (1980). Final retrievals were eventually classified on the basis of latitude and local time and averaged in order to produce global maps suitable to identify phenomena such as the cooler temperature in the cold collar just after local midnight, or the warmer air at the dawn terminator at altitudes around 80 km.

### 2.2.3.2 VIRTIS-H High Spectral Resolution Observations

Migliorini et al. (2012) discuss the thermal structure resulting from the VIRTIS-H data acquired during the period May 2006 – January 2010 for a total of  $3 \times 10^4$  analyzed spectra. The thermal retrieval code applied to the VIRTIS-H data is described in Grassi et al. (2008). The Northern and Southern hemispheres were observed by VIRTIS-H at a better spatial coverage in the South because of the spacecraft orbit. Despite the low VIRTIS-H data volume in the Northern hemisphere, a comparison between the thermal behavior of the two hemispheres at all Solar Local Times at the night side of the planet is possible. A recent re-analysis of VIRTIS-H data (Grassi, personal communication) was eventually able to detect a systematic calibration offset within the 4.3  $\mu\text{m}$  CO<sub>2</sub> band. This effect induced a bias in the derived temperatures that increases with altitude. Preliminary estimates indicate that the systematic offsets reported in Migliorini et al. (2012) are caused by this effect and are less than 3 K below the 10 bar level.

The retrieval procedure is complicated by non-LTE emissions during daytime whose contributions should be properly modeled and implemented into the retrieval code. A non-interactive retrieval method that includes non-LTE forward model simulations is used to derive daytime temperature between 100 km and 150 km from VIRTIS-H CO limb emissions around 4.7  $\mu\text{m}$  (Gilli et al., 2015). The method used by Gilli

et al. (2015) is summarized by two steps: 1) minimization of data-minus-model differences and 2) a linear inversion around the solution of the first step. A selection of limb measurements (FOV smaller than 10 km) has been used for the retrieval. Measurements below 100 km and above 170 km have been excluded to avoid possible scattering effects and because of signal-to-noise limitations, respectively. Those measurements were taken between June 2006 and October 2008 with a total of about 14,000 spectra.

The vertical resolution of the profile has been estimated by the full width at half maximum of the averaging kernels at four pointing altitudes (100, 115, 130 and 140 km) which is 15-20 km in the upper mesosphere and up to 25 km in the lower thermosphere. The maximum information region for the retrieval, given by the peak of the averaging kernels, occurs at about 5 km above each tangent altitude.

The observation data were averaged in latitude/local time/altitude/SZA bins before applying the retrieval method (see details in Gilli et al. (2015)). The results show large errors ( $> 30$  K) despite the averaging. The main contribution to the error is the measurement noise (particularly large in the analyzed spectral range). The daytime thermal structure observed by VIRTIS provides a valuable piece of information to the knowledge of the upper mesosphere and lower thermosphere. There is a temperature maximum around 115 km at equatorial latitudes near the terminator which is not present at noon. This is challenging to be interpreted by the current GCMs which in contrast predict an upper mesosphere in pure radiative balance with higher temperatures at the sub solar point (Brecht and Bougher, 2012).

### 2.3 Occultation methods

Three experiments on Venus Express use the occultation method to retrieve atmospheric properties – VeRa, SOIR and SPICAV. The stellar and solar occultations are spectral measurements in the infrared and ultraviolet by SOIR and SPICAV which rely on atmospheric extinction for profiling along the limb as a function of altitude. The radio occultation method by VeRa relies on the refraction of the radio ray path defined by the index of refraction as a function of altitude. SOIR observes CO<sub>2</sub> spectral lines to obtain the CO<sub>2</sub> number density as a function of the altitude. The SOIR observations are thus conducted necessarily at the morning and evening sides of the terminator, but do occur at all latitudes because of the pericenter and apocenter of Venus Express are located above the North and South poles respectively. The SPICAV stellar occultations are performed on the night side to avoid contamination by the scattered sunlight and cover also all latitudes. The locations and local times of the VeRa radio

occultations are defined by the orientation of the Venus Express orbit plane relative to the Venus-Earth geometry.

All three methods share some common assumptions: spherically symmetric atmosphere, hydrostatic equilibrium and a known composition. The composition is assumed to be constant in a spherically homogeneous well-mixed atmosphere below the altitude of the homopause ( $< \sim 125$  km, Mahieux et al., 2015a). The CO<sub>2</sub> volume mixing ratio changes with altitude between 100 km and the homopause which affects the SOIR observations slightly. The temperature is obtained from the CO<sub>2</sub> density profile only above the homopause, i.e. the composition does not need to be assumed (Mahieux et al., 2010; Keating et al., 1985).

The occultation experiments require specific pointing directions in order to perform the measurements. This is not feasible on every orbit, therefore the temporal and spatial coverage of each experiment is not optimal. SOIR observes the density and temperature between 65 km to 170 km altitude, SPICAV between 85 km to 140 km, and Vera between 40 km to 100 km. All three experiments need initial “guess” temperature values at their respective upper boundary for the derivation of the profiles. The altitudes of the respective boundary conditions are different but the solutions converge a few kilometers below the boundary altitude. There is some overlap in the altitude coverage of the three experiments but little overlap in latitude-longitude locations or local times.

One important result from the SPICAV and SOIR occultations is that the range of the homopause altitude, estimated from the inferred CO<sub>2</sub> number densities and temperatures, is between 119 km and 138 km above the mean surface, with weak latitudinal dependences: higher altitudes are observed on the night side passed the morning side of the terminator and lower values near the evening terminator. The derived profiles are based on assumed CO<sub>2</sub> mixing ratios from earlier models below 100 km (Zasova et al., 1996) and VIRA between 100 km to 140 km which have not been explicitly validated for the encountered atmospheric conditions during the Venus Express occultation seasons.

### 2.3.1 Radio Occultations

The propagation of the radio carrier through the ionosphere and atmosphere, before and after the spacecraft disappearance behind the planetary disc as seen from the Earth, leads to a bending of the signal ray path. The bending in the dense deep Venus atmosphere is so strong that it requires a special 3-axis spacecraft antenna steering to compensate partially for this effect. Vertical profiles of refractivity versus radius are obtained using standard geometrical optics methods and Abel inversion strategies (e.g.

Fjeldbo, et al., 1971; Jenkins et al., 1994). Additional information on the derivation of atmospheric profiles is given in Tellmann et al. (2009).

The spacecraft High Gain Antenna (HGA) is pointing towards the ground station antenna on Earth. It is generally necessary to adjust the high gain antenna pointing during the occultation in order to recover as much of the altitude range as possible before the signal is lost due to atmospheric absorption or critical refraction. An accurate prediction of the radio carrier frequency not perturbed by the propagation through the atmosphere based on the ephemerides of the spacecraft, Venus and Earth and other forces acting on the spacecraft is required to separate the atmospheric frequency shift from the Doppler-shifted received sky-frequency. In the neutral atmosphere the refractivity is directly proportional to the neutral number density. The standard retrieval method assumes a constant mean atmospheric mixing ratio for the derivation of vertical number density profiles (Fjeldbo et al., 1971; Tellmann et al., 2009). These density profiles can be converted to pressure and temperature profiles assuming hydrostatic equilibrium and using the ideal gas law. This requires the implementation of an upper boundary condition for the integration of the temperature (or pressure) profiles. Usually, three different temperature boundary conditions are assumed at an altitude of 100 km (170 K, 200 K, 230 K). The dependency on the upper boundary condition strongly decreases with altitude and the three profiles merge into the same profile (Pätzold et al., 2007).

The altitude resolution is defined by the Fresnel radius of the occultation geometry which is typically in the order of 500 m. Atmospheric temperature and density profiles were derived from dual-frequency (X-band at 8.4 GHz and S-band at 2.3 GHz) radio occultations from Mariner 5 (Mariner Stanford Group, 1967), Mariner 10 (Howard et al., 1974), Venera 9 and 10 (Vasilev et al., 1980), Pioneer Venus (Kliore et al., 1979), Venera 15 and 16 (Gubenko et al., 2008). Magellan performed 20 occultations in 1992 (Jenkins et al., 1994). Atmospheric profiles from Venera 9 and 10 (Gubenko et al., 2008) were derived using a slightly different atmospheric composition (97% CO<sub>2</sub>, 3% N<sub>2</sub>) compared to the currently accepted values used to derive profiles from Magellan, PVO and Venus Express.

#### **2.3.1.1 Venera 15 and 16, Magellan**

The Venera 15 and 16 orbiters performed 42 occultations in total (Gubenko et al., 2008), mostly at polar latitudes from October 1983 to September 1984 within an altitude range between 42 km and 90 km. The frequencies used by the Venera orbiters were L-band (1 GHz) and S-band (2.3 GHz). Tabulated results are not available at present.

The Magellan orbiter performed a few occultations which were, however, not part of the baseline mission (Steffes et al., 1994; Jenkins et al., 1994). Hinson and Jenkins (1995) discussed three profiles out of about 20, covering the altitude region between 35-90 km.

#### 2.3.1.2 Venus Express VeRa

The Venus Express Radio Science Experiment (VeRa) used one-way radio signals at two coherent frequencies (X-band; 8.4 GHz and S-band; 2.3 GHz) to sound the Venus atmosphere and ionosphere during Earth occultations. The two coherent radio signals allowed separation of the classical Doppler shift from the dispersive media effects. An onboard ultra-stable oscillator (USO) provided a high quality reference frequency source for the coherent one-way downlinks. The radio signals were primarily recorded at the ESA ground station in New Norcia, Australia, but were also supported by the NASA Deep Space Network (DSN) antennas. A detailed experiment overview can be found in Häusler et al. (2006; 2007). The atmospheric profiles cover the upper troposphere and mesosphere of Venus (~40 - 100 km) at a high vertical resolution of only a few hundred meters depending on the distance between the spacecraft and the planetary limb. Atmospheric absorption and defocusing losses of the radio carriers strongly increase below 40 km. At ~32 km altitude the atmosphere becomes critically refractive, and therefore inaccessible for radio sounding.

More than 800 profiles of temperature, pressure and neutral number density were retrieved between April 2006 and January 2015 (see section 2.4). The measurements cover nearly all local times, latitudes and longitudes with a certain gap in the northern middle latitudes resulting from the geometry of the highly elliptical orbit of Venus Express.

Radio occultation studies can also be used to study the stability of the atmosphere by deriving the buoyancy or Brunt-Väisälä frequency (Hinson and Jenkins, 1995; Tellmann et al., 2009). The strong attenuation of the radio carrier strength caused by the absorption of the radio signal provides the additional opportunity to study the absorptivity distribution within the Venus cloud deck (Oschlisniok et al., 2012; Jenkins and Steffes, 1991; Steffes and Eshleman, 1982). The high vertical resolution of the profiles allows the investigation of atmospheric small scale atmospheric structures like the accurate determination of the tropopause (Kliore, 1985; Pätzold et al., 2007; Tellmann et al., 2009) or study of small-scale gravity waves (Hinson and Jenkins, 1995; Tellmann et al., 2012).

### 2.3.2 Solar Occultation InfraRed (SOIR)

The SOIR instrument is an infrared spectrometer on board the ESA Venus Express spacecraft. It uses the solar occultation technique to sound the mesosphere and the lower thermosphere of the Venus atmosphere (Nevejans et al., 2006; Mahieux et al., 2008, 2009). SOIR is sensitive to the 2.3 to 4.4  $\mu\text{m}$  wavelength range (2257 to 4430  $\text{cm}^{-1}$ ) and uses an echelle grating at very high diffraction orders (from 101 to 194) to diffract the incoming infrared sunlight. The diffraction order (called simply order hereafter) is selected using an acoustic-optical tunable filter (AOTF). The full width at half maximum (FWHM) of the AOTF transfer function has a constant value of 24  $\text{cm}^{-1}$ , while the spectral width of an order on the detector varies between 19.3 and 37.1  $\text{cm}^{-1}$ , which causes an order overlapping on the detector, which needs to be taken into account when studying the SOIR spectra. Four orders are scanned during an occultation. The detector lines along its spatial direction need to be binned on board due to telemetry limitations: two bins are downlinked to the Earth for each order, leading to 8 spectral sets in 4 wavenumber region during an occultation. Around 700 solar occultations measuring  $\text{CO}_2$  were performed during the VEX mission. All measurements always occur at the terminator, i.e. at 06:00 hours or 18:00 hours local solar time covering well all latitudes except for the  $30^\circ - 60^\circ$  North region due to the geometry of the spacecraft orbit. During an occultation, the measurements are taken at a 1 s sampling rate at successive tangent altitudes, which correspond to the minimum altitude of the light path between the Sun and the instrument slit relative to the planet surface; it is also called impact point. The vertical altitude within the atmosphere probed by SOIR varies from 65 km up to 170 km. The calculation of the tangent altitude relies on the position and orientation of the spacecraft, and weakly on the light refraction in the atmosphere which can be neglected in the sounded altitude range. The uncertainty of the tangent altitude is always lower than 200 m and is latitude dependent.

The ASIMAT algorithm was developed to process the SOIR spectra by an iterative procedure. First, the logarithm of the number density profiles in each spectral set, i.e. for one given bin and order, is fitted using the Bayesian algorithm Optimal Estimation Method (OEM, Rodgers, 2000) in a so-called onion-peeling-configuration (Mahieux et al., 2012; 2015a; Vandaele et al., 2013). More than one species is fitted in each spectral set. Only those spectra that contain spectral information are considered in the procedure: with decreasing altitude, the first spectrum in a spectral set is the one in which the spectral lines are well above the noise, the last spectrum is the one in which the atmospheric saturation starts to set in. The baseline is fit as a 2<sup>nd</sup> to 5<sup>th</sup> order polynomial. Note that the temperature is not yet fit at this stage. The OEM algorithm uses a covariance equal to 25% of the *a-priori* profile. The independent profiles for the various fitted species are combined after each global iteration by a weighted linear

moving average procedure (averaging window  $\pm 2$  scale heights) (Mahieux et al., 2012). Then, the temperature profiles are derived from the CO<sub>2</sub> number density profiles using the hydrostatic law. The number density profiles are used as *a-priori* for the next iteration which also uses the new temperature profile. The iteration is terminated when both number density and temperature profiles are within the uncertainty of the previous iteration step. The results of the inversion are the CO<sub>2</sub> number density profile and the temperature profile. The total number density and the pressure profiles are also calculated assuming a CO<sub>2</sub> volume mixing ratio from a modified Venus International Reference Atmosphere (VIRA) (Hedin et al., 1983; Zasova et al., 2006).

Large variations of the CO<sub>2</sub> number density for a given altitude level are observed by up to two orders of magnitude. The uncertainties of the CO<sub>2</sub> number densities are in the order of 10%, much lower than the observed variability. These variations seem to be day-to-day variations rather than latitude or local solar time (terminator side) variations (Mahieux et al., 2012, 2015a) which might indicate the strong influence of the atmospheric dynamics, of waves of all kinds and daily variations of the solar flux. These variations are also seen in the temperature profiles: for a given pressure level, day-to-day variations may rise up 80 K, while the uncertainty on the temperature is in the order of 10 K.

Rotational temperatures are derived from the CO<sub>2</sub> ro-vibrational spectral structure measured by the SOIR instrument (Mahieux et al., 2015b). Hence, the rotational structure in a given vibrational band is function of the so-called rotational temperature, and may be derived from the spectra if the spectral resolution is sufficient to resolve the CO<sub>2</sub> rotational spectral structure. The method developed to retrieve the rotational temperature is not as computing-time expensive as the procedure to derive both the CO<sub>2</sub> number density and temperature profiles. There are, however, drawbacks, mostly because of some instruments characteristics, such as the order overlapping, the modulation by the AOTF function and the spectral noise which is the largest error source. The general shape of the terminator temperature profiles is confirmed by using this method. The rotational temperature are in good agreement with the corresponding hydrostatic temperatures, but at larger uncertainties ranging from 20 to 100 K. No rotational non-local thermodynamical equilibrium bifurcation has been observed in the datasets.

### 2.3.3 Stellar Occultations from SPICAV

The SPICAV (Spectroscopy for the Investigation of the Characteristics of the Atmosphere of Venus) instrument performs spectroscopy at ultraviolet (110 – 320 nm) and at near infrared (650 – 1700 nm) wavelengths in the limb, nadir, stellar and solar occultation mode. A detailed description of the SPICAV

instrument and its scientific objectives can be found in Bertaux et al. (2007). The UV band spectroscopy enables in the occultation mode vertical profiling of CO<sub>2</sub>, SO<sub>2</sub>, SO, O<sub>3</sub>, aerosols and temperature profiles in the ~ 90-140 km region (Bertaux et al., 2007; Montmessin et al., 2006; 2011; Piccialli et al., 2015). The ultraviolet sensor of SPICAV has a spectral dispersion of 0.54 nm per pixel and a spectral resolution varying from 1 to 2.5 nm. The vertical resolution of a profile ranges from 0.5 to ~ 7 km depending on the occultation grazing angle.

As for the solar occultation, the stellar occultation technique relies on the computation of the atmospheric transmission obtained by dividing each spectrum affected by the presence of the atmosphere along the line of sight by the reference spectrum taken outside of the atmosphere. The reference spectrum is obtained by averaging all spectra (up to 1000) acquired above a tangential altitude of 250 km. One advantage of the stellar occultation technique is the intrinsically accurate geometric registration: the uncertainty of the inferred altitude of the tangential point relies only on the precise knowledge of the spacecraft position in its orbit and not on the precise knowledge of the spacecraft pointing attitude.

Like for solar and radio occultations, each altitude position within the profile is at a slightly different latitude and longitude due to the tangential transect of the line of sight between the star and SPICAV. The difference of the geographical locations between the start and the end of the occultation may be as much as ~2° of latitude and/or longitude. Both entry and exit occultations are possible and were recorded and processed by SPICAV. A reference altitude of 85 km was defined.

The stellar occultation retrieval starts first by separating the nitric oxide airglow emission whose signature is superimposed on that of the stellar spectrum to be followed by the derivation of a wavelength-dependent atmospheric transmission at each sounded altitude. Using the same retrieval method as in Quémerais et al. (2006) and Montmessin et al. (2006), line of sight integrated densities (slant densities) for CO<sub>2</sub>, O<sub>3</sub> and aerosols are first retrieved and then inverted to yield local density and temperature profiles by assuming hydrostatic equilibrium (see Piccialli et al. (2015) for details).

The observations cover all latitudes on the night side between 18:00 hours and 06:00 hours. The error of the SPICAV temperature retrievals vary with altitude: typical values are 1 K to 20 K in the altitude range 100 km - 130 km, and 5 K to 60 K at lower and higher altitudes (Piccialli et al., 2015).

## 2.4 VEX dataset coverages and data averaging

The majority of the data are from observations and experiments on board of Venus Express. It is necessary to consider the spatial and temporal coverage of each experiment for a data intercomparison. There is no uniformity in global and temporal coverage because of the different operational and orbital constraints. The spatial coverage from the various experiments and the data binning and averaging are presented in this section.

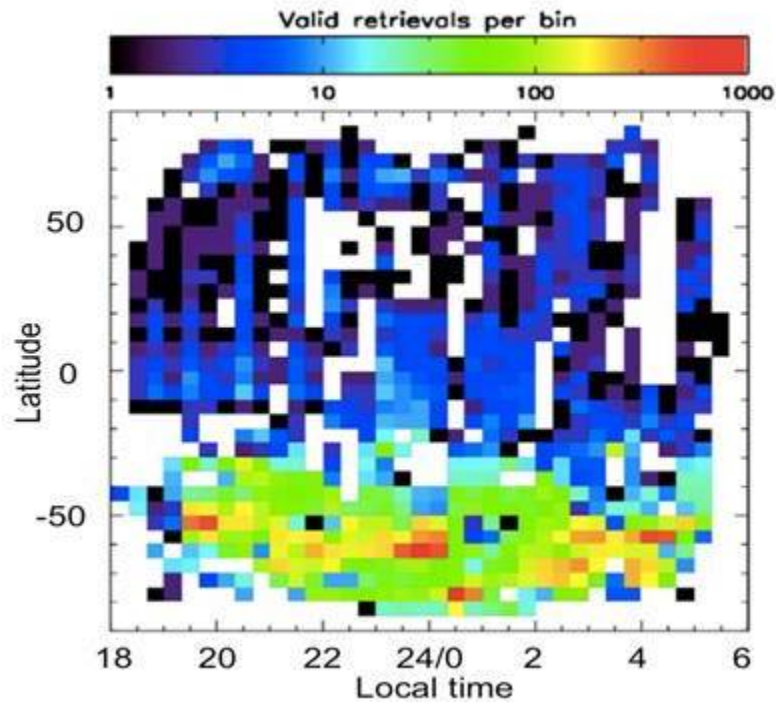
### 2.4.1 Geographical and temporal coverage of the VEX observations

#### 2.4.1.1 VIRTIS coverage

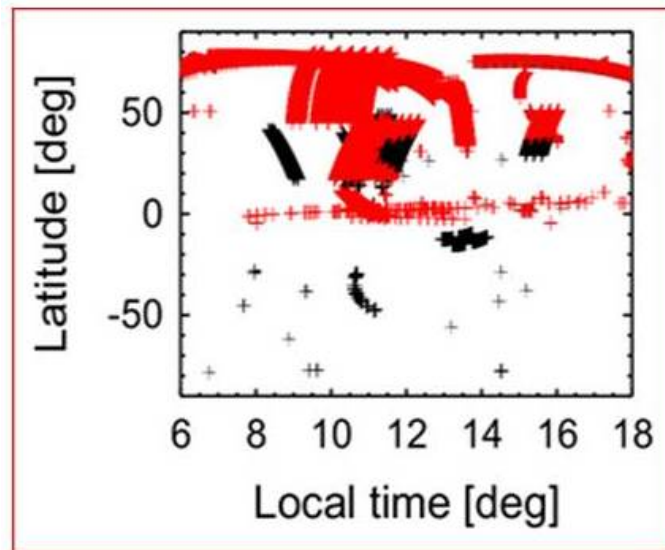
Most of VIRTIS limb data come from the Northern hemisphere because of the VEX operational strategy. The observation locations are not evenly distributed in local time and latitude as shown in Figure 5. For instance, observations between 10°N and 30°N are particularly scarce. The profiles are also not really vertical but each measurement point corresponds to a single spectrum at a given local time and latitude, with no particular geographical and vertical connection to the next data acquisition. The VIRTIS-H spectra on the night side (Figure 5) and day side (Figure 6) of Venus are analyzed separately. VIRTIS-M covers mostly the Southern hemisphere (Grassi et al., 2014).

#### 2.4.1.2 VeRa coverage

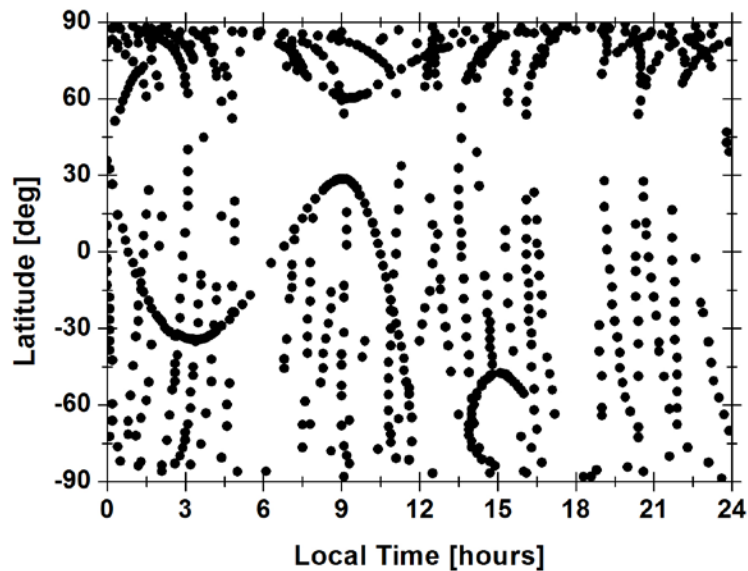
Radio occultations occur in seasons when the constellation Venus, Earth and the spacecraft orbit plane is oriented such that the spacecraft disappears behind (and reappears from) the planetary disk as seen from the Earth. It is possible to observe the ingress as well as the egress occultation because of the one-way radio link. Both ingress and egress occur at opposite hemispheres. VeRa occultations cover all latitudes and local times (Figure 7). The atmospheric profiles derived from the profile of the index of refraction are slightly slanted and cover about 4° along the meridian which means that the planetary latitude within the profile varies only slightly at constant local time.



**Figure 5:** Number of VIRTIS-H spectra (about 30,000 spectra in total) used for the night time temperature retrieval, distributed over local time and latitude (from Migliorini et al., 2012). The southern hemisphere is much better covered than the northern hemisphere. Right panel: Local-time and latitude distribution of the VIRTIS-H daytime limb observations between 100 and 170 km altitude. The red crosses represent data with a field-of-view smaller than 10 km which were used for the temperature retrieval (after Gilli et al., 2015).



**Figure 6:** Local-time and latitude distribution of the VIRTIS-H daytime limb observations between 100 and 170 km altitude. The red crosses represent data with a field-of-view smaller than 10 km which were, used for the temperature retrieval (after Gilli et al., 2015).

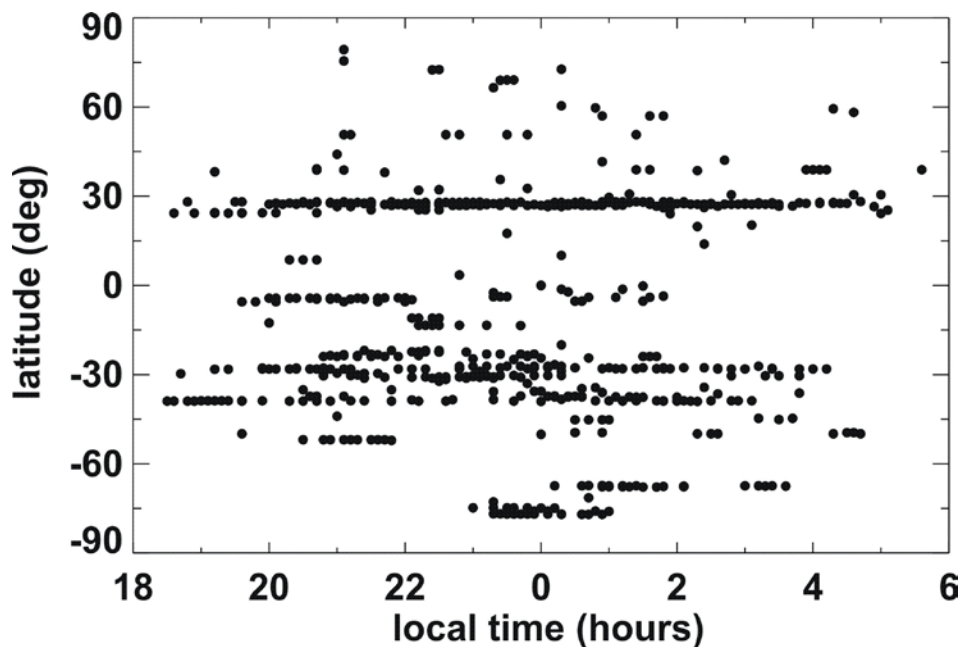


**Figure 7:** Spatial distribution of the VeRa occultation profiles as a function of local time and latitude represented by the ray pericentre at the 1 bar level (altitude  $\sim 50$  km).

### 2.4.1.3 SPICAV/UV and SOIR coverage

Although co-located in the same instrument box, stellar and solar occultations can only be observed at different times in orbit when the source of the radiation is occulted by the Venus atmosphere. Stellar occultations are performed on the nightside, to avoid contamination by stray light from the bright limb. Because of the many available UV bright stars, the stellar occultations are performed at different local times than the solar occultations. Both the solar and the stellar occultations can sample different latitudes. The vertical profiles are also slightly slanted similar to the radio occultations, each vertical location is at a slightly different latitude. Figure 8 shows the distribution of the SPICAV profile locations (latitude vs. local solar) at an altitude of 85 km . Figure 9 shows the SOIR profile locations similarly.

Figures 10 and 11 summarize the local time and latitudinal coverage from the SOIR, VeRa, VIRTIS and SPICAV experiments at specific latitude and local time bins that were used in this study.



*Figure 8: Latitudinal and local time distribution of the SPICAV stellar occultations. The latitude position and local time are represented by the ray pericentre at an altitude of ~85 km.*

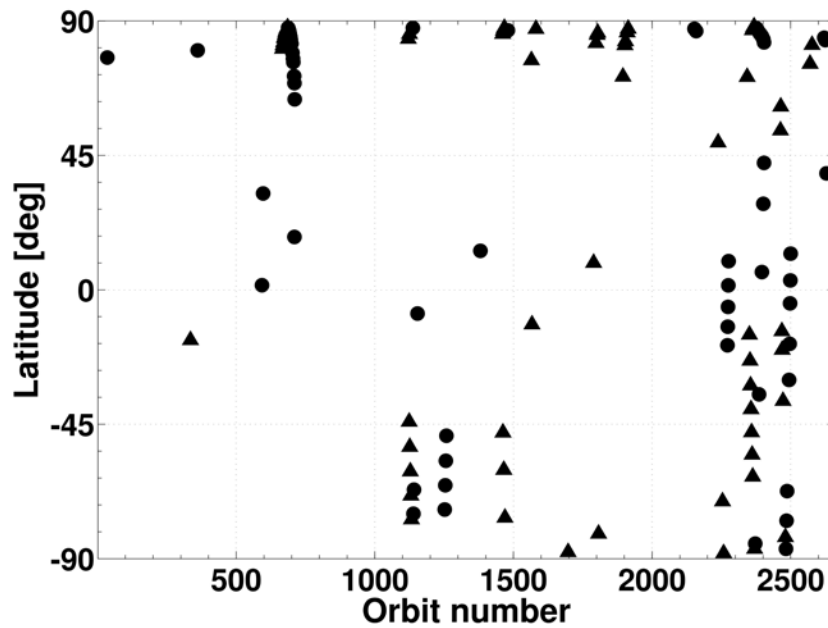


Figure 9: Latitude locations of the SOIR solar occultations when  $\text{CO}_2$  was observed at high altitudes.. Circles are for the morning terminator, triangles for the evening terminator.

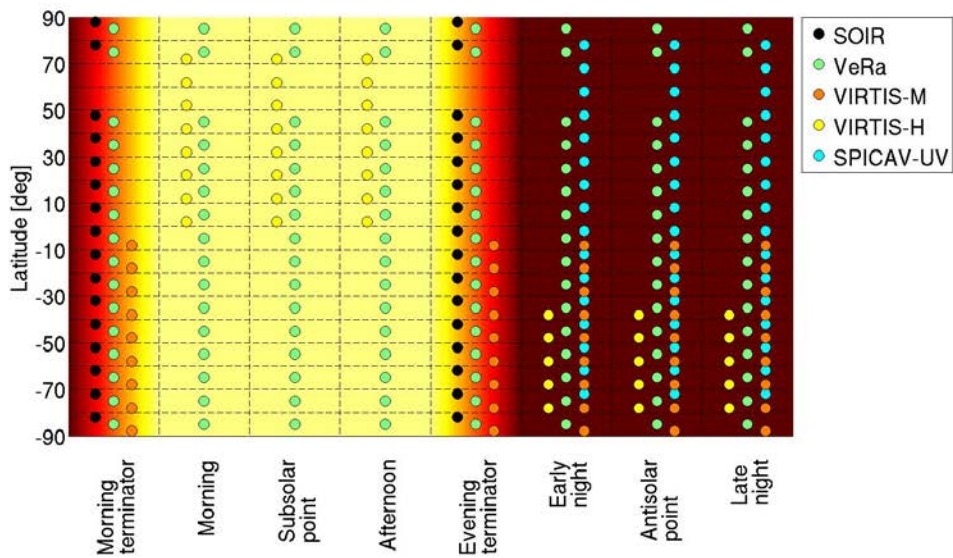


Figure 10: Local solar time versus latitude coverage for the Venus Express instruments.

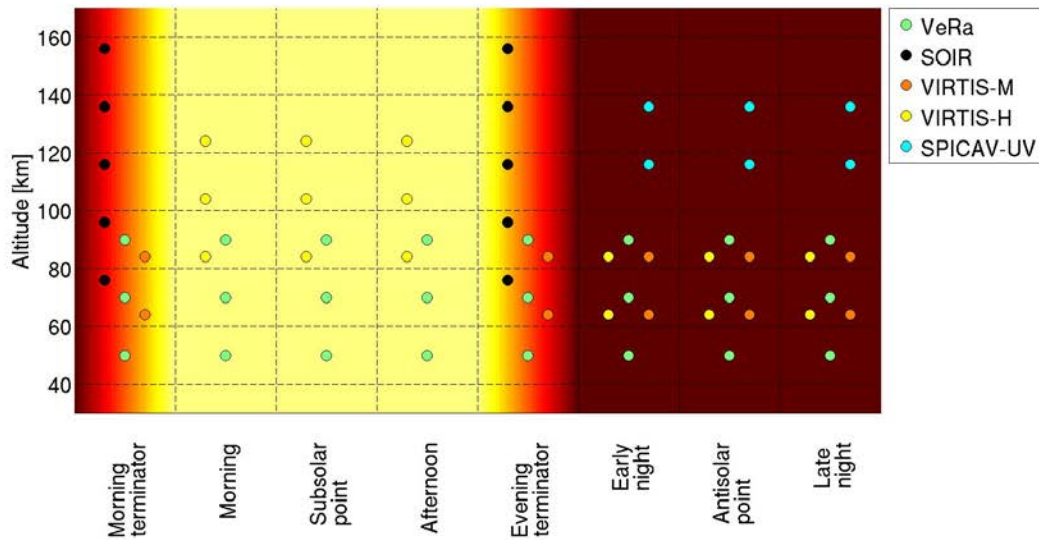


Figure 11: Local solar time versus altitude coverage for the Venus Express instruments.

#### 2.4.2 Averaging the spacecraft derived datasets

An intercomparison of the results from the different experiments is only feasible if the data are averaged over time and latitude and/or local time bins because of the different temporal and spatial sampling, coverage and respective measurement errors and uncertainties. The number of measurements from each experiment is given in Table 3 for each latitude and local time bins.

When averaging individual and independent  $N$  measurements of the same physical quantity with different uncertainties, assuming that these measurements obey a Gaussian distribution around the “true” value, the best estimate of that quantity is given by the weighted average  $\mu_w$ :

$$\mu_w = \frac{\sum_{i=1}^N w_i x_i}{\sum_{i=1}^N w_i} \quad (1)$$

where each measurement  $x_i$  at a given pressure/altitude level is multiplied by a weighting factor  $w_i$ , defined as the inverse square of the individual error  $\varepsilon_i$ . The standard deviation of a weighted sample with  $M$  nonzero weights is given by:

$$std_w = \sqrt{\frac{\sum_{i=1}^N w_i \cdot (x_i - \mu)^2}{\frac{M-1}{M} \cdot \sum_{i=1}^N w_i}} \quad (2)$$

**Table 3. Number of data sets**

	Number of data sets per latitude range (North/South)				
	0°-30°	30°-50°	50°-70°	70°-80°	80°-90°
<b>Night</b>					
VIRTIS-M (D.G.)	4545/11042	2070/185585	2049/550011	675/112416	110/4719
VIRTIS-M (R.H.)	0/1846	0/5657	0/8764	0/2974	0/429
VIRTIS-H (nadir)	307/830	114/3938	192/15312	82/4061	4/366
VeRa	32/58	18/48	40/33	46/15	97/15
SPICAV	188/154	25/119	18/33	7/26	0
JCMT	23	12	7	1	0
HHSMT	6	0	0	0	0
<b>Terminator</b>					
VIRTIS-M (D.G.) (evening)	0/18	0/7071	0/16018	0/846	0
VIRTIS-M (D.G.) (morning)	34/0	0/984	0/1200	0/21	0
VIRTIS-H (nadir) (evening)	21/19	0/34	0/13	0	0
VIRTIS-H (nadir) (morning)	4/16	11/84	3/88	0	0
VeRa	2/4	3/5	1/2	2/0	20/4
SPICAV (evening)	0	0/5	0	0	0
SOIR (evening)	9	9	11	8	33
SOIR (morning)	15	7	9	10	25
JCMT	3	0	0	0	0
HHSMT	5	0	0	0	0
<b>Day</b>					
VeRa	53/54	9/44	42/41	56/15	54/12
JCMT	2	2	0	0	0
HHSMT	4	0	0	0	0

Assuming hemispheric symmetry and combining the data from the northern and southern hemispheres, and the data from the different experiments are compared in five latitude bins – (i) 0° – 30° latitude, (ii) 30° - 50° latitude, (iii) 50° - 70° latitude, (iv) 70° -80° latitude, and (v) 80° -90° latitude. The data were grouped into three sets – Day, Night and Terminator (both morning and evening) in each latitude bin.

The ground based observations (see Section 3) have a very sparse temporal sampling and a very low spatial resolution. Those observations were therefore not included and compared as combined results from each experiment.

### 3. Ground based observations

The most extensive information regarding the atmospheric structure comes from remote or in-situ measurements of the temperature as a function of pressure (altitude). Remote observations, which provide much better spatial and temporal coverage, include spectroscopic measurements in the near-IR, submillimeter, millimeter wave and infrared atmospheric nightglow emission. In the near-IR, information about the thermal structure can be retrieved from the continuum maps and from the maps of the CO<sub>2</sub> line depths (Encrenaz et al., 2013). High spectral resolution observations at short and long wavelengths from  $\sim 1 \mu\text{m}$  to mm wavelengths enable the probing of the Venus atmospheric thermal structure from  $\sim 120 \text{ km}$  to the cloud top level (Betz et al., 1976; Clancy and Muhleman 1991; Sonnabend et al., 2010; Rengel et al., 2008a; 2008b).

#### 3.1 Sub-mm observations

Sub-mm observations of CO lines provide information about atmospheric conditions between approximately 70 and 110 km. CO is produced in this region by the photolysis of CO<sub>2</sub>. The pressure broadened rotational lines of CO provide a means to infer atmospheric properties from high resolution spectroscopy yielding a temperature profile and a line-of-sight Doppler wind velocity and the CO abundance. An optimal retrieval of temperature and CO mixing profiles requires simultaneous radiative transfer (RT) analysis of the <sup>12</sup>CO and <sup>13</sup>CO line absorption measurements, whereby a single temperature and CO mixing profile over 75–120 km altitudes is derived to provide self-consistent fits to both spectral lines (e.g. Clancy et al., 2012).

Many observations have been made by various instruments at different observatories around the world in the recent years (Clancy et al., 2012); Lellouch et al., 2008; Rengel et al., 2008a ; 2008b; Sagawa et al., 2010) : James Clark Maxwell Telescope (JCMT, Hawaii), Kitt Peak (Arizona), National Radio Astronomy Observatory (NRAO, Virginia), IRAM (Spain) and IRAM Pdb (France), Nobeyama and Heinrich Hertz Radio Telescope (HHT) on Mount Graham, Arizona.

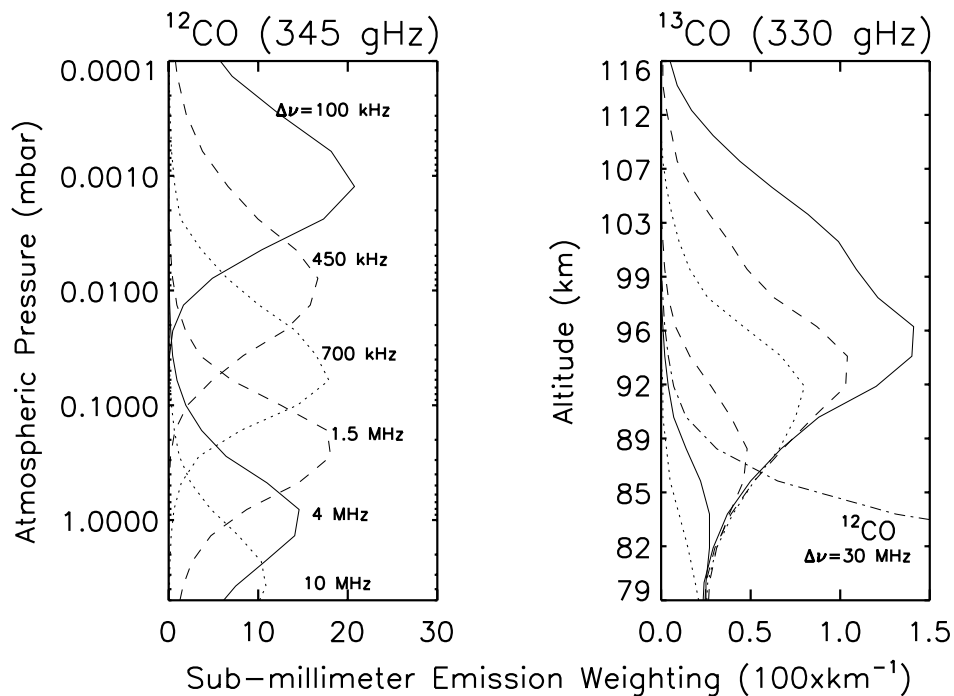
##### 3.1.1 James Clark Maxwell Telescope (JCMT)

The temperature profiles observed by the JCMT are retrieved from thermal (LTE) radiative transfer (RT) analyses of sub-millimeter optically thick (<sup>12</sup>CO, 345 GHz) and thin (<sup>13</sup>CO, 330 GHz) line absorptions formed in the mesosphere and lower thermosphere of Venus (Clancy et al., 2012). Detailed descriptions of sub-millimeter and millimeter CO line absorptions with respect to RT analysis for temperature profiles can be found in Clancy and Muhleman (1991), Lellouch et al. (1994), Rengel et al. (2008a;2008b), and

Clancy et al. (2003; 2008; 2012). In ground-based (i.e., nadir viewing) observations, pressure-broadened  $^{12}\text{CO}$  lines appear with 30-50% line center absorptions against the Venus thermal continuum, which arises from collisionally-induced  $\text{CO}_2$  opacity at altitudes of  $\sim 45 - 65$  km (e.g., Muhleman et al., 1979). Line center optical depths for  $^{12}\text{CO}$  line absorptions support thermal profiling from  $\sim 80 - 115$  km in the Venus night side atmosphere versus  $\sim 80 - 105$  km in the Venus dayside atmosphere. This day/night distinction regards the large diurnal variation in Venus CO mixing ratios above  $80 - 90$  km altitudes (see for example Clancy et al., 2003), such that the line center average optical depth for the  $345$  GHz  $^{12}\text{CO}$  transition varies from  $\tau_o \sim 12$  at the night side to  $\tau_o \sim 4$  at the day side. CO mixing ratios in the night side in lower thermosphere exhibit strong temporal and spatial variations on top of this average diurnal variation, which reflect the strong night side variation in regional down-welling, which is also dramatically exhibited by  $\text{O}_2$  singlet delta nightglow variations, e.g., Bailey et al (2008).

Sub-millimeter  $^{12}\text{CO}$  temperature profiling is very similar in principle to that employed by  $15 \mu\text{m}$  nadir temperature sounding for Venus with a  $\text{CO}_2$  opacity source (Taylor et al., 1980). However the temperature dependence of sub-millimeter radiation is (nearly) linear and the CO opacity source is highly variable, in altitude, LT, and latitude. The latter distinction requires that the CO mixing profile be measured simultaneously, through CO profile retrievals from RT analysis of optically thin ( $\tau_o \sim 0.1 - 0.3$ , at  $330$  GHz)  $^{13}\text{CO}$  line absorptions. The pressure-broadened line shape supports such compositional profiling up to  $\sim 105$  km. Temperature and CO contribution functions associated with sub-millimeter  $^{12}\text{CO}$  ( $345$  GHz) and  $^{13}\text{CO}$  ( $330$  GHz) profile retrieval analyses are presented in Figure 12, as reproduced from Clancy et al. (2012).

Vertical resolution for temperature profiles within the mesosphere ( $80$  km -  $100$  km) is roughly 1 scale height ( $4 - 5$  km), sufficient to resolve  $\pm 5$  K solar thermal tides with good accuracy (Clancy and Sandor, 2011). This vertical resolution degrades by a factor-of-two into the lower thermosphere ( $100 - 120$  km), due to the transition of the contributed line shape from variable pressure to (nearly) fixed thermal broadening and to decreasing vertical gradients in the CO mixing ratio. The spatial/LT resolution of temperature profiling across the Venus disk is set by the diffraction-limited telescope beam, which is  $14$  arc-seconds for  $345$  GHz JCMT observations. Hence, the large disk size of Venus when the full night side is viewed ( $\sim 60$  arc-seconds at inferior conjunction) provides  $1 - 3$  hour LT and  $20 - 40^\circ$  latitudinal resolution of night side temperature profiles up to  $\sim 70^\circ$  latitude. Day side coverage is limited from full disk (superior conjunction) to  $\sim$ half disk (elongation) resolution, such that night side (inferior conjunction) observations from JCMT are emphasized.



**Figure 12:** Representative  $^{12}\text{CO}$  (left panel) and  $^{13}\text{CO}$  (right panel) emission weighting functions are derived for different line center frequency offsets. Vertical axes indicate atmospheric pressure (left) and corresponding altitudes (right) associated with the temperature (left) and CO (right) solution profiles (reproduced from Clancy et al., 2012).

### 3.1.2 Heinrich Hertz Submillimeter Telescope (HHSMT)

The vertical thermal structure retrieved from HHSMT observations considered in this paper was derived from  $^{12}\text{CO}$  J = 2-1 at 230.54 GHz at seven positions on the Venus disc and  $^{13}\text{CO}$  J = 2-1 at 220.4 GHz at one position. Observations on the night and day sides of Venus were performed on 9 – 10 June and 14 – 15 June 2007, at  $\sim 0^\circ$  latitude. The angular diameter of Venus was 23:44" at the beginning and 25:55" at end of the campaign, the approximate FWHM beam diameters are shown in Rengel et al. 2008b. These observations are a part of a coordinated ground-based Venus observational campaign in support of the ESA Venus Express mission (Rengel et al. 2008a, 2008b). The results indicate a temperature vertical distribution and CO distribution spatially and temporally variable in the mesosphere.

The technique used to retrieve the temperature and CO profiles is described in Rengel et al. (2008a). The atmospheric model has a vertical resolution of 2 km. Normalized temperature and CO weighting functions for  $^{12}\text{CO}$  J = 2- 1, and temperature and CO weighting functions for  $^{13}\text{CO}$  J = 2-1 for several

frequency offsets around each transition can be found in Rengel et al. (2008a). Temperature and CO sensitivity are listed in Table 4.

There is evidence of changes in the thermal structure of the Venus mesosphere occurring on short time scales: small day-to-night temperature variations and short-term (Earth day to Earth day) on a time scale as short as one Earth day.

**Table 4: Vertical range where the temperature and CO density retrieved from HHSMT observations are sensitive.**

Quantity	Spectral Line	
	$^{12}\text{CO} (J = 2 - 1)$	$^{13}\text{CO} (J = 2 - 1)$
Temperature	75 - 110 km	80 - 90 km
CO distribution	70 - 100 km	80 - 95 km

### 3.2 CO<sub>2</sub> Heterodyne Observations

Heterodyne spectroscopy of CO<sub>2</sub> at mid-infrared wavelengths is a powerful tool to study temperatures and the dynamical behavior of the atmospheres of the terrestrial planets. In general, heterodyning means mixing the received signal at the telescope with a local oscillator which is usually done by a laser at IR wavelengths. The mixing yields the difference between the received frequency and the laser frequency, both typically at THz frequencies, with preserved spectral information. The down-converted signal, now at GHz, is easily amplified and analyzed with extraordinary spectral resolution. There are currently worldwide two instruments which use infrared heterodyne receivers to investigate the Venusian atmosphere. One of them is the Cologne Tunable Infrared Heterodyne Spectrometer (THIS) which was developed at the I. Physikalisches Institut, Universität zu Köln, Cologne, Germany. It operates at wavelengths between 7 and 14  $\mu\text{m}$ . The other instrument is the Heterodyne Instrument for Planetary Wind And Composition HIPWAC developed and operated by the Goddard Space Flight Center in Maryland, USA. Both receivers are transportable and can be shipped to any telescope with IR receiving capabilities. Data presented in this paper were taken at the McMath Pierce Solar Telescope, Kitt Peak, Arizona and the NASA InfraRed Telescope Facility, Mauna Kea, Hawaii. Detailed information about the instrumentation can be found in previous work on the development of the THIS receiver and in

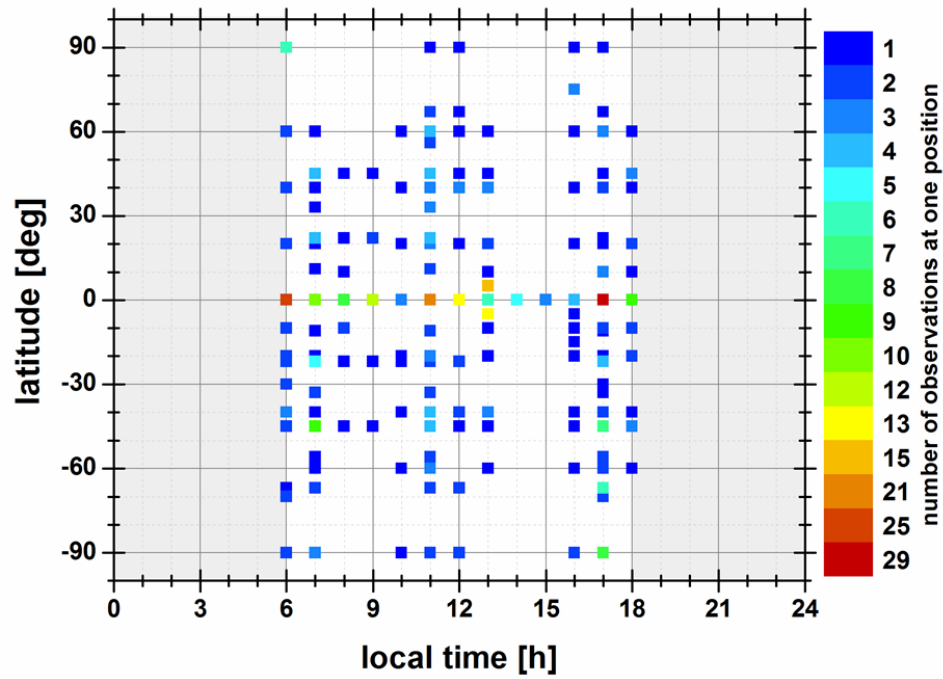
publications on the observations accomplished with HIPWAC (Sonnabend et al., 2008; Sornig et al., 2009; Kostiuk and Mumma, 1983; Kostiuk et al., 2006).

Mesospheric non-local thermodynamic equilibrium (non-LTE) emission of CO<sub>2</sub> near 10 μm is observed in the Venus atmosphere, a phenomenon first discovered in 1976 by Betz et al. (1976). The modeling of the processes which lead to the non-LTE emission (Deming et al., 1983; Roldan et al., 2000) has recently advanced (Lopez Valverde et al., 2011) and is a significant step forward on the way to a self-consistent model of the Venusian atmosphere. The kinetic temperature can be calculated from the width of the observed lines and is a good probe for the physical temperature of the emitting gas as long as stimulated emission is negligible. The high spectral resolution allows the determination of the Doppler shift of the observed CO<sub>2</sub> emission line which corresponds to the line-of-sight velocities and provides therefore a direct wind measurement. The exact altitude of the emitting region is determined by the ratio of collisions induced to the probability of spontaneous emission for the excited CO<sub>2</sub> molecules. The excitation is controlled by solar irradiation (Deming et al., 1983; Roldan et al., 2000). A recent study by Lopez-Valverde et al. (2011) finds a maximum for the non-LTE emission with a half width of 10 km at the 0.15 Pa pressure level which is equivalent to an altitude of ~110 km using a VIRA pressure-altitude profile. IR heterodyne spectroscopy offers a much higher spatial resolution in contrast to existing sub-mm observations allowing the detailed study of temperature variations as a function of latitude and local time. The observations are, however, limited to the day side and the altitude range mentioned above.

The heterodyne receivers THIS and HIPWAC have observed Venus during several campaigns in 1990/1991 and between 2007 and 2014 resulting in a comprehensive set of wind and temperature data. The data presented in this paper are temperature measurements derived between 2007 and 2014. A total of 371 individual observations were performed in 11 campaigns. An overview of the measurements and relevant observational conditions for the different campaigns is given in Table 5. The latitude-local time coverage of both instruments is shown in Figure 13. The day side of Venus is very well covered. In particular, there are a number of high quality observations at the equator.

Figure 14 illustrates the observing geometries of all targeted positions (black circles) on the planetary disk. The black circles are the size of the telescope beam relative to the diameter of the Venus apparent disk. The planetary disk is well resolved during the Venus quadrature, compared to the beam size, and allows many independent observation locations (Figure 14a, d, e, g, h, i, j, k). The beam size relative to

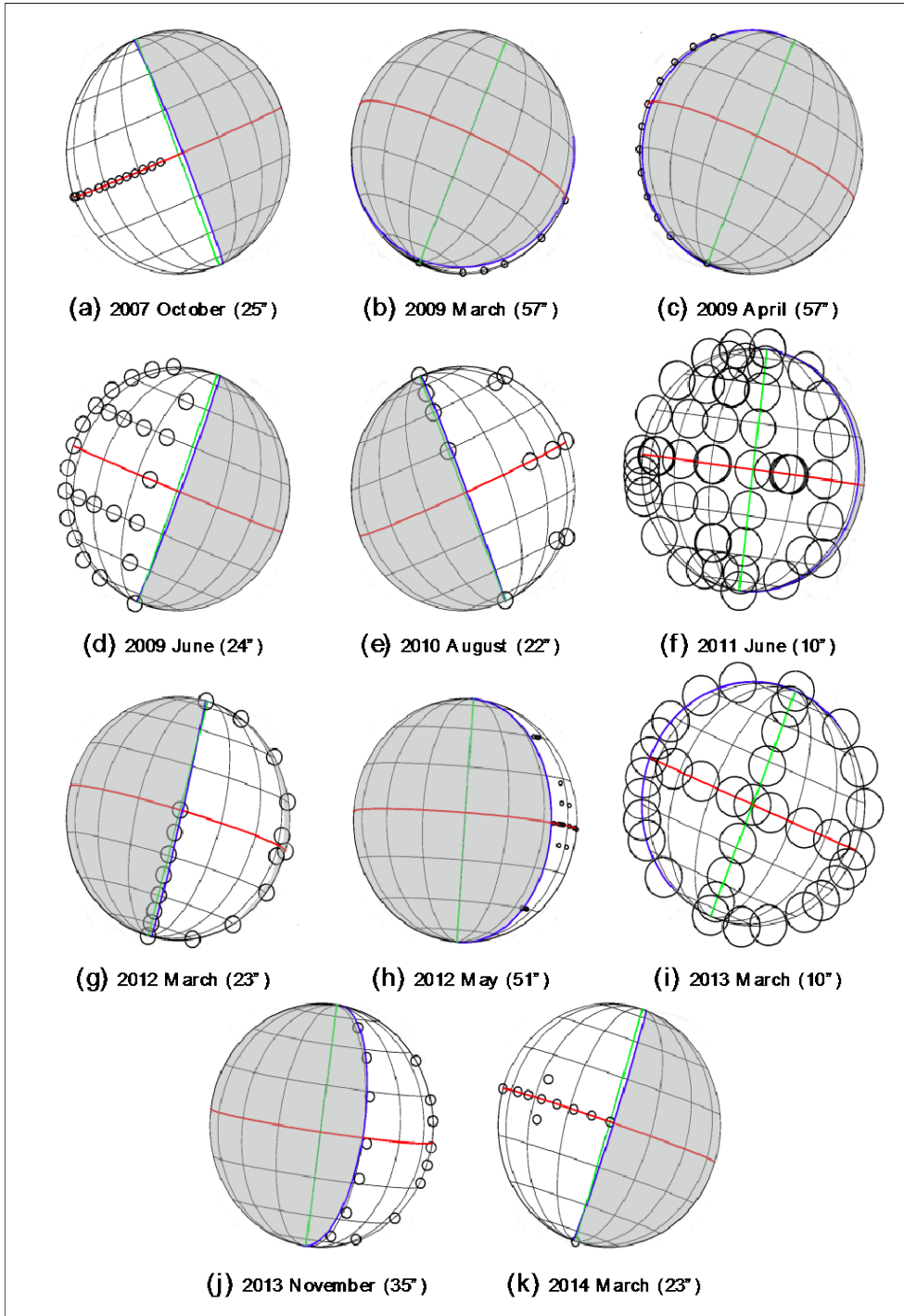
the planetary disk is, however, large and observation locations do overlap during superior solar conjunction (Figure 14f and i).



**Figure 13.** Latitude-local time coverage of both the THIS and the HIPWAC instruments. The color-code gives the number of observations.

Table 5: THIS and HIPWAC observing campaigns from 2007 to 2014.

Campaign	Date	Venus apparent diameter (arcsec)	Venus disk illumination (%)	Earth and Venus Doppler velocity (km/s)	FoV (arcsec)	number of observations
1	22.-24.10.2007	25	47	13	0.9	14
2	16.-22.03.2009	57	4	-4.8	1.6	19
3	02.-06.04.2009	57	3	3.6	1.6	32
4	02.-06.06.2009	24	50	13.9	1.6	58
5	09.-22.08.2010	22	55	-13.8	1.6	13
6	20.-25.06.2011	10	97	5.1	1.6	57
7	24.-30.03.2012	23	21	-13.3	1.6	56
8	17.-24.05.2012	51	7	7.3	0.9	23
9	11.-16.03.2013	10	100	1.4	1.6	60
10	23.11-04.12.2013	35	33	-12.3	1.6	22
11	28.-31.03.2014	23	52	13.5	0.9	17



**Figure 14:** Geometries for all Venus observing campaigns. The equator (red), the terminator (blue) and the central meridian (green) are indicated. The black circles indicate the relative size of the telescope beam to the planetary disk. The number behind the date gives the apparent diameter of Venus in arcseconds.

### 3.3 Temperatures inferred from Night Time Airglow

The  $O_2(a^1\Delta \rightarrow X^3\Sigma)$  infrared atmospheric nightglow emission is produced by three-body recombination of two O atoms in the presence of a third body. Constraints on the local temperature at the altitude of the  $O_2$  nightglow emission have been derived from the analysis of the intensity distribution within the rotational structure of the (0-0) band at 1.27  $\mu\text{m}$ . These measurements do not always have sufficient resolution to resolve individual lines. However, comparisons between observed spectra and synthetic spectra accounting for the Earth atmospheric transmission and convolved to the observed resolution of the ground-based instruments, have provided reliable estimates of the temperature. It is assumed that, in the upper mesosphere region of the emission, the rotational temperature is essentially equal to the temperature of the ambient gas.

The first measurements were made by Connes et al. (1979) who obtained high-resolution Fourier transform spectra yielding  $T = 185 \pm 15$  K. Crisp et al. (1996) also resolved rotational lines in the P and R branches using the same technique, and deduced a temperature of  $186 \pm 6$  K at  $15^\circ\text{S}$ . Without any measurements of the altitude distribution of the airglow layer, it was not possible to precisely assign these temperatures to a given altitude or pressure level. Ohtsuki et al. (2008) deduced temperatures from observations made during three different years. The average rotational temperatures from their observations were  $193 \pm 9$  K,  $182 \pm 25$  K, and  $185 \pm 20$  K. They showed cases suggesting some correlation with the regions of bright nightglow. Bailey et al. (2008a,b) derived temperatures from 181 to 196 K and also showed some relations between higher temperatures and  $O_2$  nightglow bright patches. Krasnopolsky (2010) retrieved temperatures showing a broad minimum of 171 K centered at  $4^\circ\text{S}$  increasing to 195 K at  $35^\circ\text{S}$  and 212 K at  $35^\circ\text{N}$  with an uncertainty of about 5 K. No correlation was observed between the nightglow intensity and temperature. By contrast, Bailey et al. (2008) and Ohtsuki et al. (2008) found that their measurements support the idea that compressional heating of downwelling gas heats the region of the airglow layer. They argued that dynamical effects on the nighttime thermal structure in the mesosphere-thermosphere transition region are stronger than the chemical energy released by the association of O atoms.

VIRTIS-M/VEx nightglow observations did not have sufficient spectral resolution to infer rotational temperature, however, the limb observations indicated that the peak of the  $O_2$  emission at 1.27  $\mu\text{m}$  at the limb in the northern hemisphere is located at  $96 \pm 2.7$  km (Piccioni et al., 2009). Soret et al. (2012) determined the peak altitude of the volume emission rate by an Abel inversion. These results now make it possible to assign an altitude range to the source region, which is useful to interpret the ground-based

observations. Summarizing these observations, ground-based nightglow measurements yield a mean rotational temperature of  $186 \pm 6$  K at an altitude of  $97.4 \pm 2.5$  km. These values are 15-20 K higher than temperatures listed in the VIRA model (170 K).

### 3.4 The Venus Transit on 6 June 2012 across the Solar Disk

The transit of Venus in June 2012 provided a unique case study of Venus' atmosphere transiting the Sun, while at the same time Venus Express observed the evening terminator at solar ingress and solar egress (Wilson et al. 2012). This was the first time in history that a transit of Venus occurred while a spacecraft was simultaneously in orbit around Venus. Transit observers in the past gave detailed descriptions of the telescopic aspect of Venus. In particular, during transit ingress and egress, the portion of the planet's disk outside the solar photosphere has been repeatedly perceived as outlined by a thin, bright arc ("aureole"). On June 8th, 2004, fast photometry based on electronic imaging devices allowed the rediscovery and first quantitative analysis of the phenomenon (Tanga et al., 2012). On June 5 and 6, 2012, several observers used a variety of acquisition systems to image the event – thus collecting for the first time a large amount of information on this atmospheric phenomenon. Tanga et al. (2012) had shown that the aureole photometry reflects the local density scale height at the limb and the altitude of the refracting layer. The lightcurve of each spatial resolution element of the aureole has been compared to a limb refraction model to constrain the mesospheric structure / scale height at terminator. The latitude probed by SOIR on Venus Express during orbit 2238 (+49°), at the time Venus transited the Sun as seen from Earth, provided a suitable validation to this approach (Pere et al. 2016).

The analysis of the images obtained by the *Helioseismic and Magnetic Imager* of the Solar Dynamics Observer yield temperature data at the evening terminator covering the altitude range from 70 to 110 km (Figure 15) The accuracy of the average latitudinal temperature is comparable to SOIR. The best-measured aureole signal is produced at layers at an altitude of 80 - 90 km. Table 6 reports the results obtained at 90 km. (Widemann et al., 2014 ; Tanga et al., 2016)

**Table 6: Temperatures at the 90 km altitude level from aureole observations during the 2012 Venus Transit (Widemann et al., 2014)**

Latitude	Temperature (K)	Error (K)
0° - 30°	161.0	19.5
30° - 50°	151.0	18.5
50° - 70°	154.0	11.0
70° - 80°	167.0	16.6
80° - 90°	157.2	13.7

## 4. Comparison of the Venus Express and Ground-based Temperature and Density Observations

### 4.1 Description of the Datasets

The temperature datasets available from the Venus Express VIRTIS-M, VIRTIS-H, VeRa, SOIR and SPICAV instruments were sorted in latitudinal bins, assuming a symmetry of the Northern and Southern hemisphere, as well as in local time bins. The datasets were averaged as discussed in Section 2.4.2. Figures 16 to 20 combine the temperature profiles from the VEX instruments, some Venera and Magellan profiles and the profiles from the ground-based observations as a function of vertical pressure and altitude.

The VIRTIS-M datasets were analyzed using the two different methods by Grassi et al. (2008, 2014) and by Haus et al. (2013, 2014). Dayside temperature were also derived from the VIRTIS-H non-LTE emissions (Gilli et al., 2015). These results are obtained by averaging a large number of spectra taken at same altitude/local time/latitude bins from different observations during the VEx mission. For this reason, they do not represent a real vertical profile, but an average value for each bin. The SOIR temperature profiles (Mahieux et al., 2015) were derived from observations at the evening and morning terminators.

Ground-based observations by the JCMT (Clancy et al., 2008, 2012), the HHSMT (Rengel et al., 2008a,b), HIPWAC and THIS (Sonnabend et al., 2008, 2010; Krause et al., 2014) have been binned in a similar way. The mean temperature profile observed at the evening terminator during the Venus transit (Tanga et al., 2012) is also given, as well as the average temperature deduced from O<sub>2</sub> airglow observations.

The combined temperature profiles in Figures 16 to 20 are presented in five latitude bins 0° – 30° (Figure 15) 30° – 50° (Figure 16) 50° – 70° (Figure 17), 70° – 80° (Figure 18) and 80° – 90° (Figure 19) latitude and three local times bins:

- (i) day side from 07 h LST to 17 h LST (panel a in Figures 15 to 19)
- (ii) night side from 19 h LST to 05 h LST (panel b in Figures 15 to 19)
- (iii) terminator zones: 05 h LST – 07 h LST and 17 h LST – 19 h LST (panel c in Figures 15 to 19)

The shaded color coded regions in Figures 16 to 20 mark one standard deviation for the respective experiment's averaged profile. Error bars are given for VIRTIS-H non-LTE, HIPWAC and THIS, Venus transit and O<sub>2</sub> airglow profiles which represent the total error. Error bars with respect to altitude are shown for the airglow and transit profiles.

The profiles resulting from the two different analyses are plotted for the VIRTIS-M datasets. The SOIR profiles are averaged from all available data in the latitudinal bin. Morning and evening terminator temperature profiles were computed separately for the SPICAV and SOIR datasets.

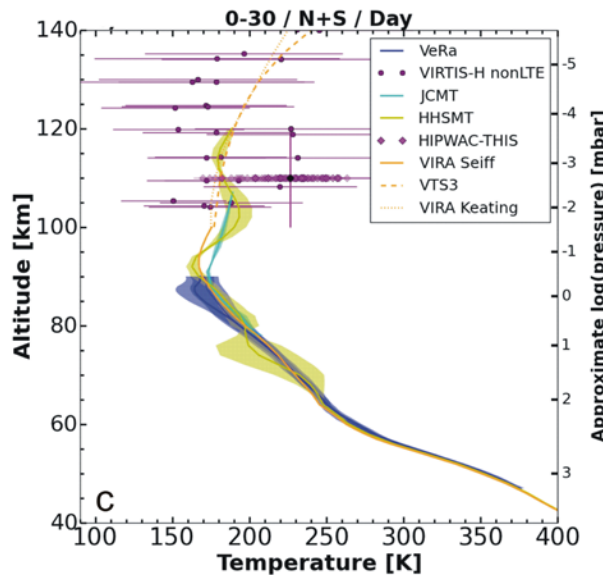
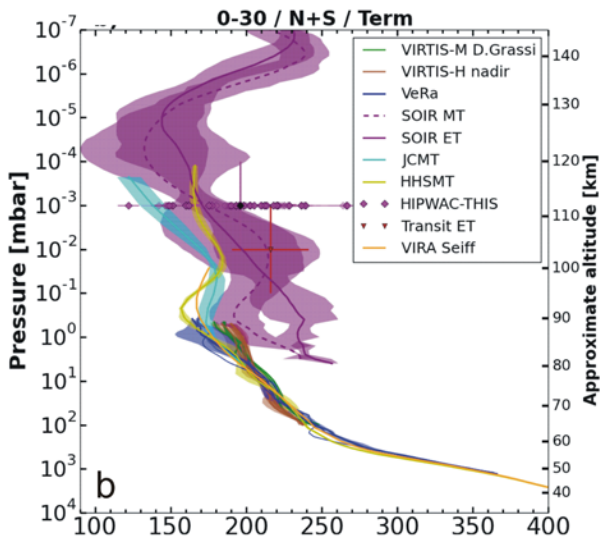
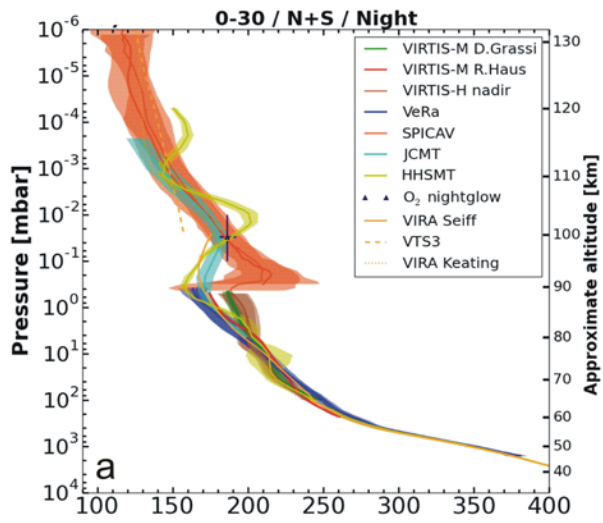
Figures 15 to 19 illustrate the huge contributions of Venus Express and the ground-based observations to the investigation of the Venus atmosphere since the publication of VIRA. The new observations are in very good agreement with those temperature values which are addressed by VIRA and provide new information about the atmospheric structure above 100 km.

## 4.2 The Troposphere and Middle Mesosphere Below 90 km Altitude

Figure 16 to 20 illustrates the overall good agreement between the datasets in each latitudinal bin for the night side and the terminator below 0.5 mbar (90 km) and within the observed variability, considering that the measurement techniques have different fields of view, are taken from ground or from space, are obtained at different times and have different spatial resolutions. The foot prints of the field of view of Venus Express instruments are all latitude dependent because of the highly elliptic orbit. This affects the temperature profiles like smoothing at increasing field of view.

The cold collar which was first detected by Pioneer Venus is identified at latitudes poleward of 50° and seen in all available datasets. Some small temperature differences appear to be present when comparing the VeRa and VIRA profiles in particular at the day side. These might be at least partially caused by differences in the spatial distribution and the sampling of the two data sets.

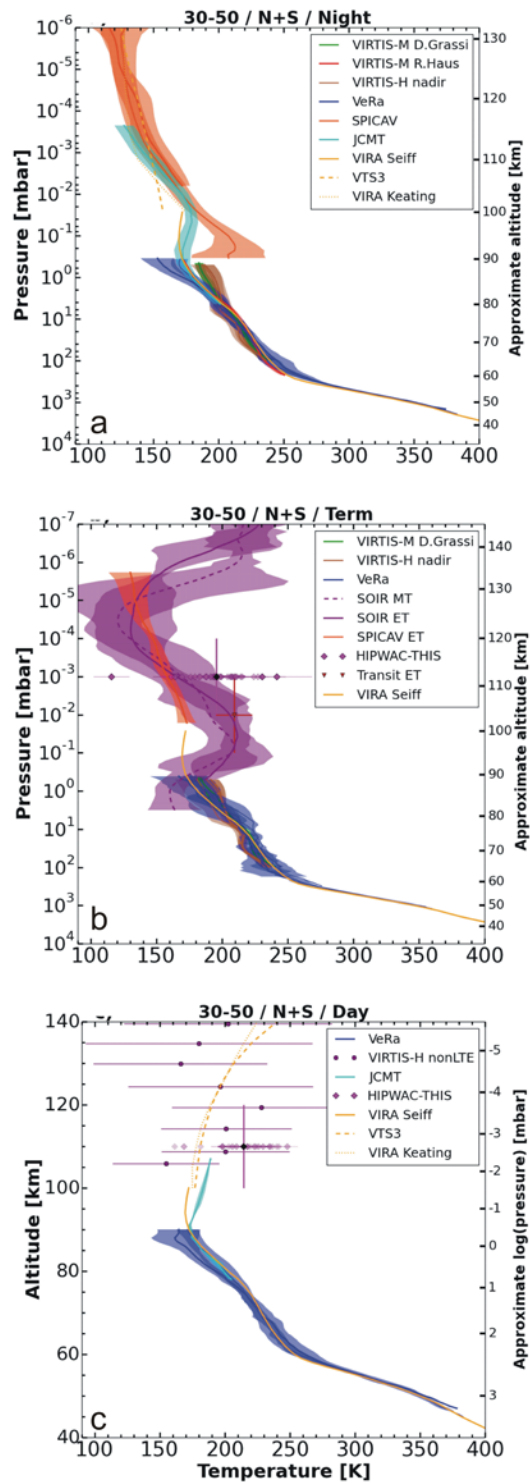
A systematic difference between the two analyses of the VIRTIS-M dataset (Haus et al., 2013; 2014; Grassi et al., 2015) is evident in particular around the 1 mbar level. There are several explanations: Discrepancies in the final VIRTIS results may be explained by differences in the retrieval methods (described in section 2.2.3.1), by the forward radiative transfer codes and/or the pre-processing procedures (required to address residual calibration issues) adopted by the two VIRTIS teams (Haus et al., 2013; 2014; Grassi et al., 2008).



2 **Figure 16:** combined temperature profiles from the Northern and Southern hemispheres between the  
3 equator and 30° latitude. Panel (a): night side, panel (b): terminators, panel (c) day side. Temperature  
4 profiles are combined from the Venus Express instruments, ground-based observations, and empirical  
5 models (VIRA; VIRA Keating; VTS3). The height above the mean planetary radius is given as pressure for  
6 the night side and terminator data (panels (a) and (b)) and in altitude for the day side observations in  
7 order to ease the comparison with VIRTIS-H data. Corresponding approximate values for  
8 altitude/pressure are also given on the right-hand side of each panel. Uncertainties (one standard  
9 deviation) are either plotted as colored areas for averaged profiles in the same bin (Venus Express  
10 datasets, JCMT, HHSMT, and Venera-15) or as error bars. The VIRTIS-H non-LTE, O<sub>2</sub> airglow and Venus  
11 transit horizontal error bars represent the total retrieval error. The vertical error bars represent the  
12 uncertainty in altitude/pressure.

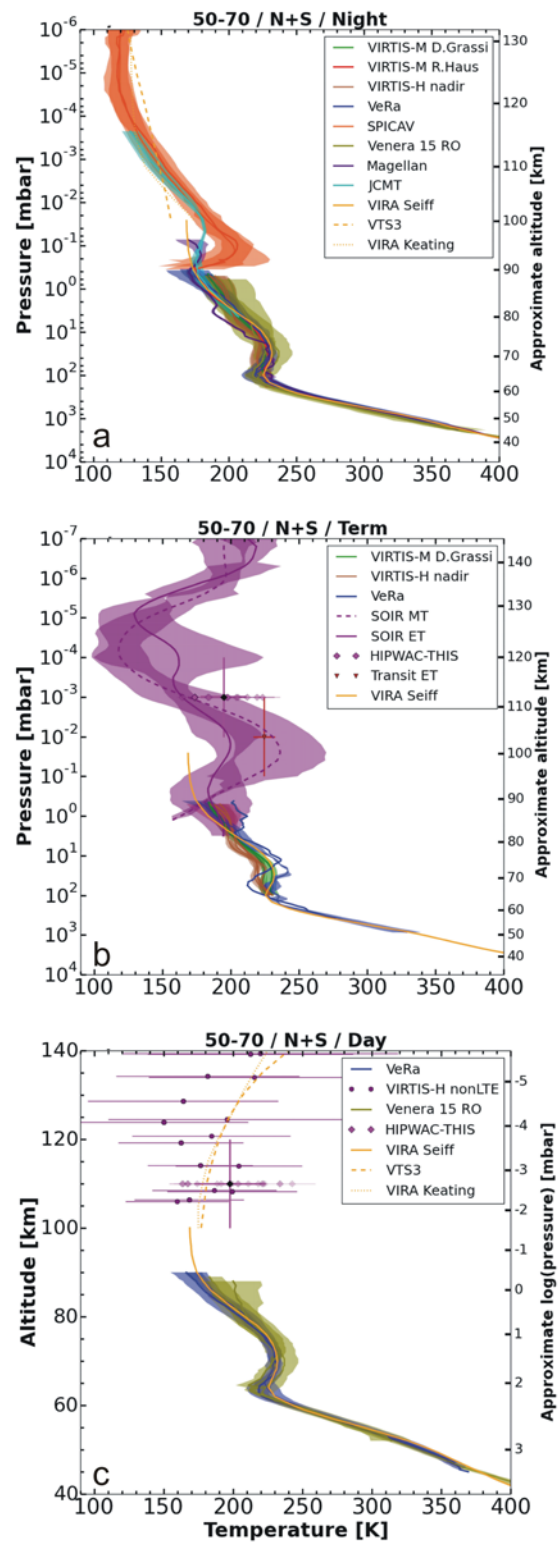
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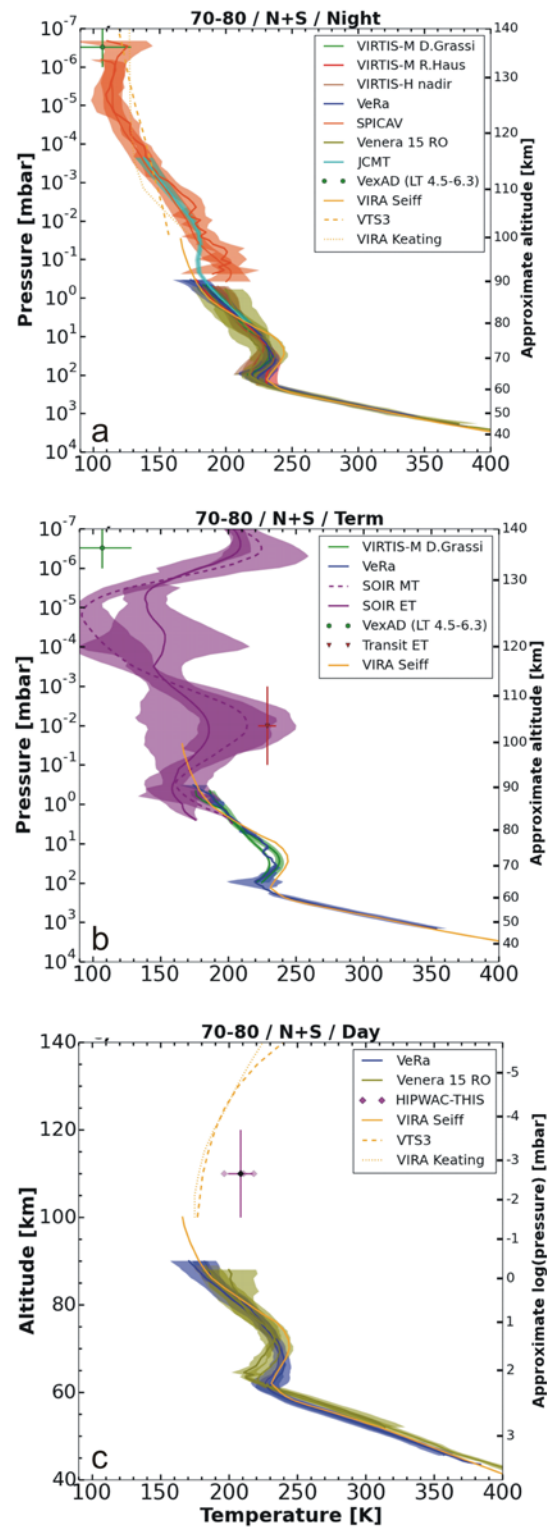
15

16 *Figure 17: same as Figure 16 but for combined temperature profiles from the Northern and Southern*  
 17 *hemispheres between 30° and 50° latitude. Panel (a): night side, panel (b): terminators, panel (c)*  
 18 *day side.*



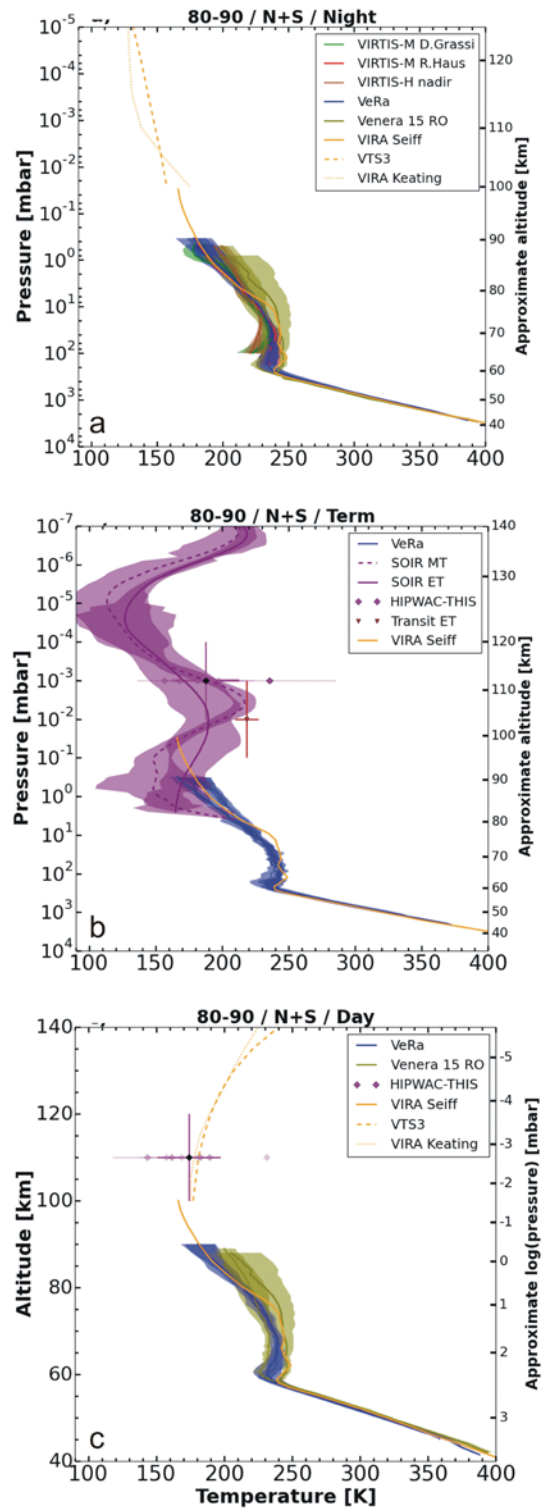
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20 **Figure 18:** same as Figure 16 but for combined temperature profiles from the Northern and Southern  
 21 hemispheres between 50° and 70° latitude. Panel (a): night side, panel (b): terminators, panel (c) day  
 22 side.



23

24 **Figure 19:** same as Figure 16 but for combined temperature profiles from the Northern and Southern  
 25 hemispheres between the 70° and 80° latitude. Panel (a): night side, panel (b): terminators, panel (c) day  
 26 side.



27

28 **Figure 20:** same as Figure 16 but for combined temperature profiles from the Northern and Southern  
 29 hemispheres between 80° and 90° latitude. Panel (a): night side, panel (b): terminators, panel (c) day  
 30 side.

### 31 **4.3 Upper Mesosphere and Thermosphere (90 km to 150 km)**

32 The combined profiles (Figures 16 to 20) show a complicated thermal structure in the 90 km to 150 km  
33 altitude range with alternating warm and cool layers rather than a gradual increase or decrease of  
34 temperature. The cold temperatures seen by SOIR of about 120 K and lower were seen also by the  
35 Pioneer Venus Orbiter drag experiments but at higher altitudes (Keating et al., 1980). A much higher  
36 variability of temperatures at each pressure level is observed by SPICAV and SOIR. The corresponding  
37 density variations are also large, up to two to three orders of magnitude. The largest temperature  
38 difference is seen in the terminator zones at all latitudes (Figures 16b, 17b, 18b, 19b, 20b) which may be  
39 caused by short-term temporal variability (all kinds of atmospheric waves) at altitudes above 100 km.  
40 The uncertainty of the SOIR temperatures peak to the extreme. The temperature inversions are seen in  
41 the terminator zones at slightly different altitudes as a function of latitude (about 100 km at  $0^\circ - 30^\circ$   
42 latitude, 95 km at  $30^\circ - 50^\circ$ , 105 km at  $70^\circ - 80^\circ$  and about 110 km at  $80^\circ - 90^\circ$ ), which may be caused by  
43 the descending circulating flow. Large variations in the vertical flow may influence the mixing of the  
44 species which in turn may affect the radiative balance. This is discussed in Section 5.

45 Temperatures derived from the VIRTIS-H (non-LTE), HHSMT, JCMT and HIPWAC-THIS experiments are on  
46 average in good agreement at the day side for latitudes lower than  $70^\circ$  (Figures 16c, 17c, 18c), but show  
47 a very large variability. Almost no data are available from the ground based experiments for latitudes  
48  $>70^\circ$ , except for a few observations by the HIPWAC-THIS experiment (Figures 19c and 20c). Compared to  
49 other observations which give averaged values the given HIPWAC-THIS data are single measurements. The  
50 variability of this single measurement is in the same range than the VEX instruments even though variability of the  
51 spatial field-of-view with the various observing runs have to be taken into account.

52 The situation is more complex above the 0.5 mbar pressure level (90 km). The SPICAV and JCMT  
53 temperatures are in good agreement above 0.03 mbar (100 km) on the night side (panel (a) of Figures  
54 16 to 20) and also agree with the average temperatures from the  $O_2$  nightglow observations at low  
55 latitudes (Figure 16a). The SPICAV profiles, however, show a maximum temperature in the 0.03 mbar to  
56 1 mbar (85 km - 100 km) range which is more pronounced and located at lower altitudes than the JCMT  
57 profiles. The HHSMT profiles are in agreement with the JCMT profiles below 1 mbar (85 km) but tend to  
58 show a higher variability than the SPICAV and the JCMT profiles above this altitude.

59 The temperature profiles from all experiments do not generally overlap temporally and spatially at the  
60 terminator zones (panel (b) of Figures 16 to 20) and the temperature variability is very high. The SOIR

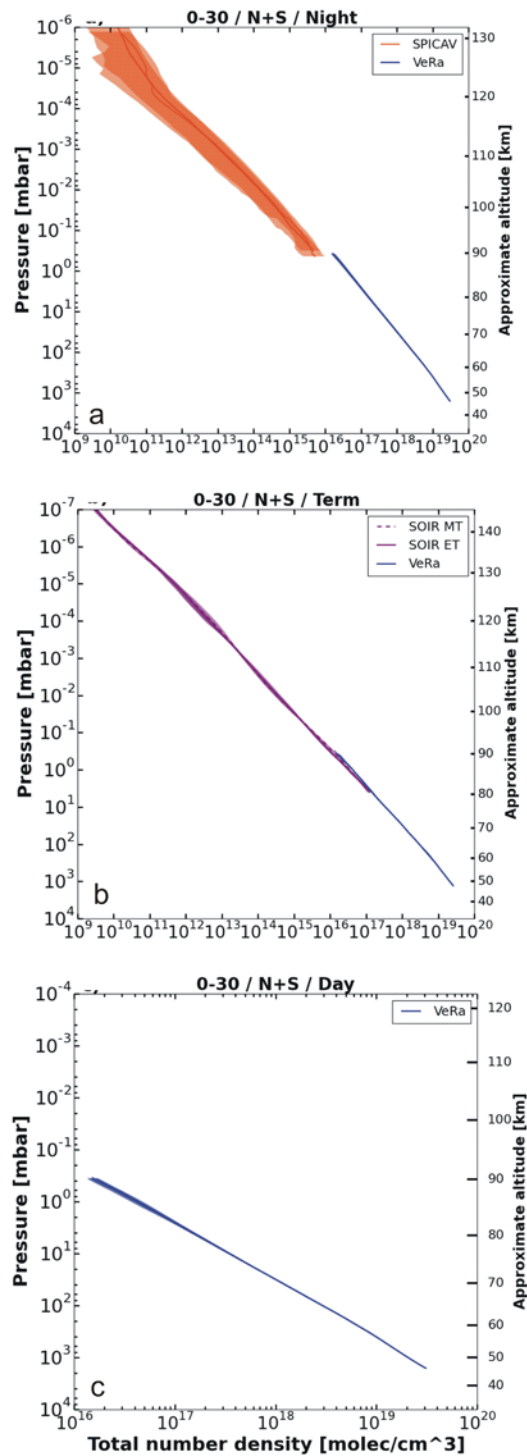
61 profiles are much warmer than the JCMT and the HHSMT profiles but agree roughly with the HIPWAC-  
62 THIS temperatures and those from the Venus 2012 transit. The SOIR, HIPWAC-THIS observations are  
63 close to the terminator ( $\leq 2$  h) addressing illuminated day side only, while the sub-mm temperature  
64 observations have a larger field-of-view. Differences between both the evening and morning  
65 terminators are also apparent.

#### 66 **4.4 Atmospheric Density**

67 The three occultation experiments on board of Venus Express, SPICAV-UV, SOIR and VeRa, return  
68 measurements of the total neutral number density. SPICAV-UV and SOIR measure directly the CO<sub>2</sub>  
69 number density from its absorption structure, and thus have to assume a CO<sub>2</sub> volume mixing ratio,  
70 which was taken from VIRA. The neutral number density profiles from SPICAV, SOIR and VEXADE were  
71 achieved only at the terminator zones and at the night side. VeRa covered the day side as well as the  
72 night side.

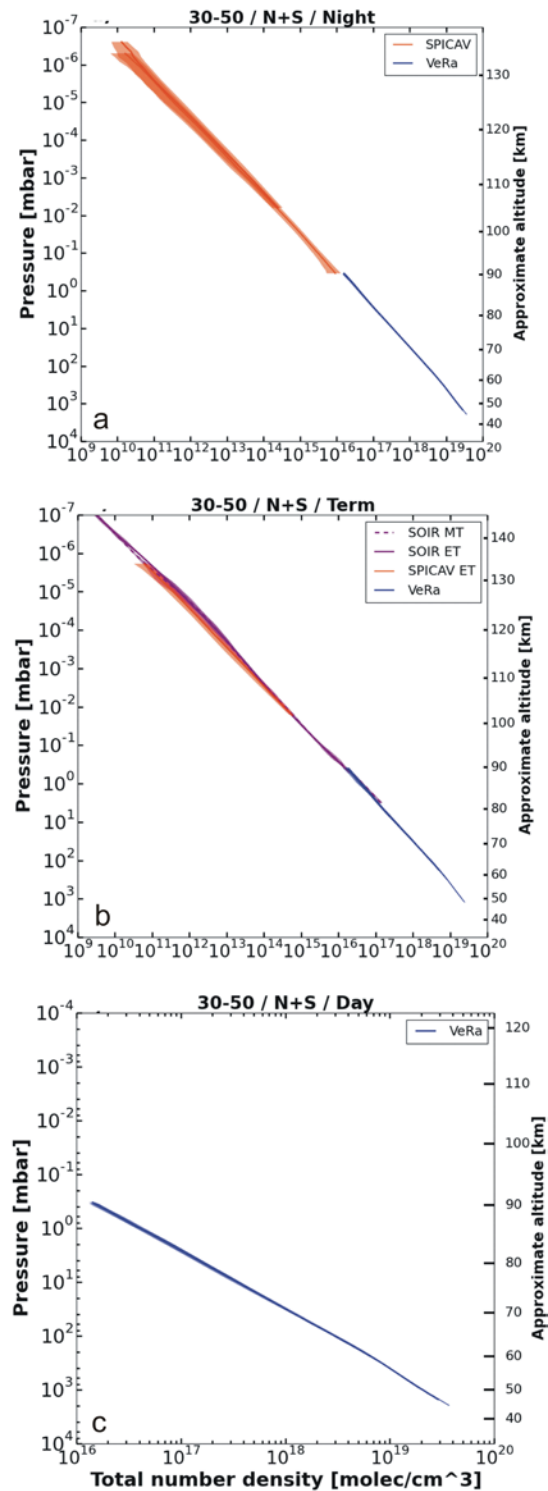
73 The SPICAV, SOIR (morning and evening) and VeRa neutral number density profiles are in very good  
74 agreement at the terminator zones from about  $10^3$  mbar to  $10^{-7}$  mbar (40 km to 150 km) for all  
75 latitude bins (Figures 21 to 25). The profiles at the night side, however, show some differences between  
76 VeRa and SPICAV at pressure levels where the profiles overlap. The uncertainties of the SPICAV profiles  
77 are significantly larger compared to VeRa. There is also a noticeable offset in the near-equatorial ( $0^\circ$  to  
78  $30^\circ$ ) latitude bin (Figure 21) and the mid-latitude ( $30^\circ$  to  $50^\circ$ ) bin (Figure 22). Similar differences are also  
79 seen in the Figures 26 to 30 (altitude versus neutral number density) which implies a change in neutral  
80 scale heights in the near-equatorial ( $0^\circ - 30^\circ$ ) latitude bin (Figure 26) and the mid-latitude ( $30^\circ - 50^\circ$ )  
81 bin (Figure 27).

82 The density values from the drag experiments at 170 km to 200 km altitude appear to be in very good  
83 agreement to an extrapolation of smoothed SOIR profiles in the high ( $70^\circ - 80^\circ$ ) and polar ( $80^\circ - 90^\circ$ )  
84 latitude bins (Figures 24a and 25a). The pressure in this altitude range is extremely low and a variability  
85 by a factor of two or more is seen in the density values from orbit to orbit. This is much lower than the  
86 variability seen in the SOIR or SPICAV profiles in other latitude bins. The reason for these differences are  
87 not yet understood. Some possible causes are discussed in Section 6.



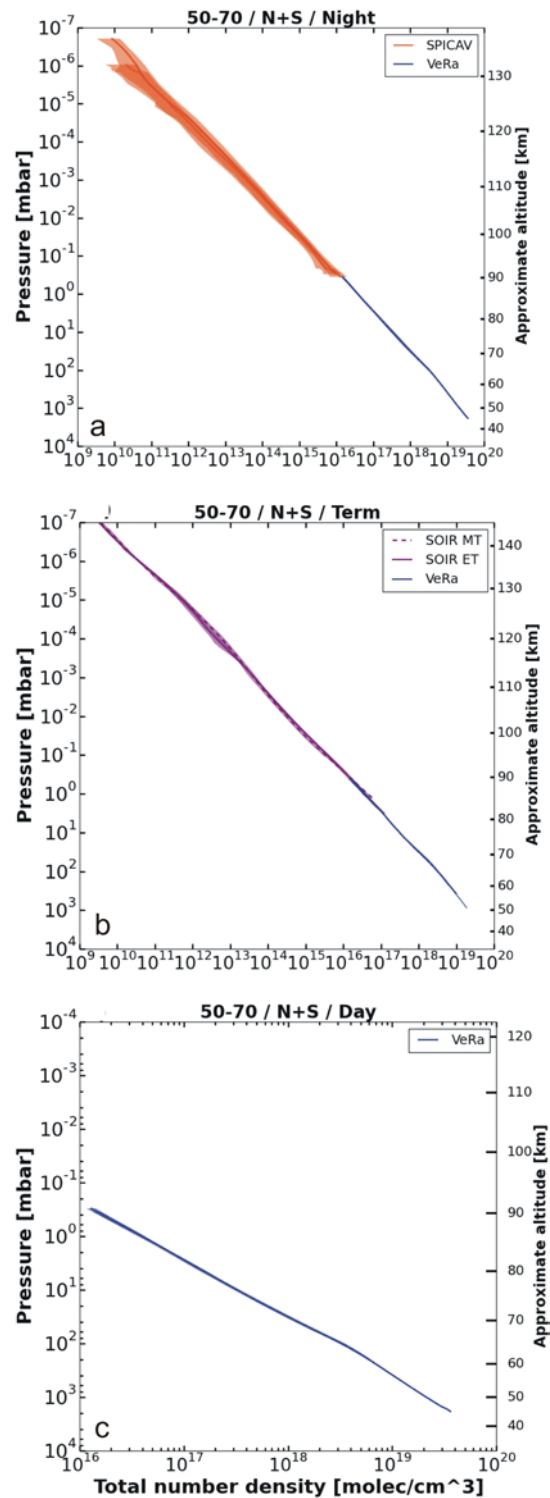
88

89 **Figure 21:** Comparison of atmospheric mean total density profiles from SOIR, SPICAV and VeRa and from  
 90 the atmospheric drag measurements as a function of pressure for the near equatorial latitude bin 0° to  
 91 30°. Panel (a): night side, panel (b): terminator zones, panel (c): day side. The colored areas mark one  
 92 standard deviation uncertainty of the average profiles for each experiment. Approximate altitudes are  
 93 shown on the right hand vertical axis.



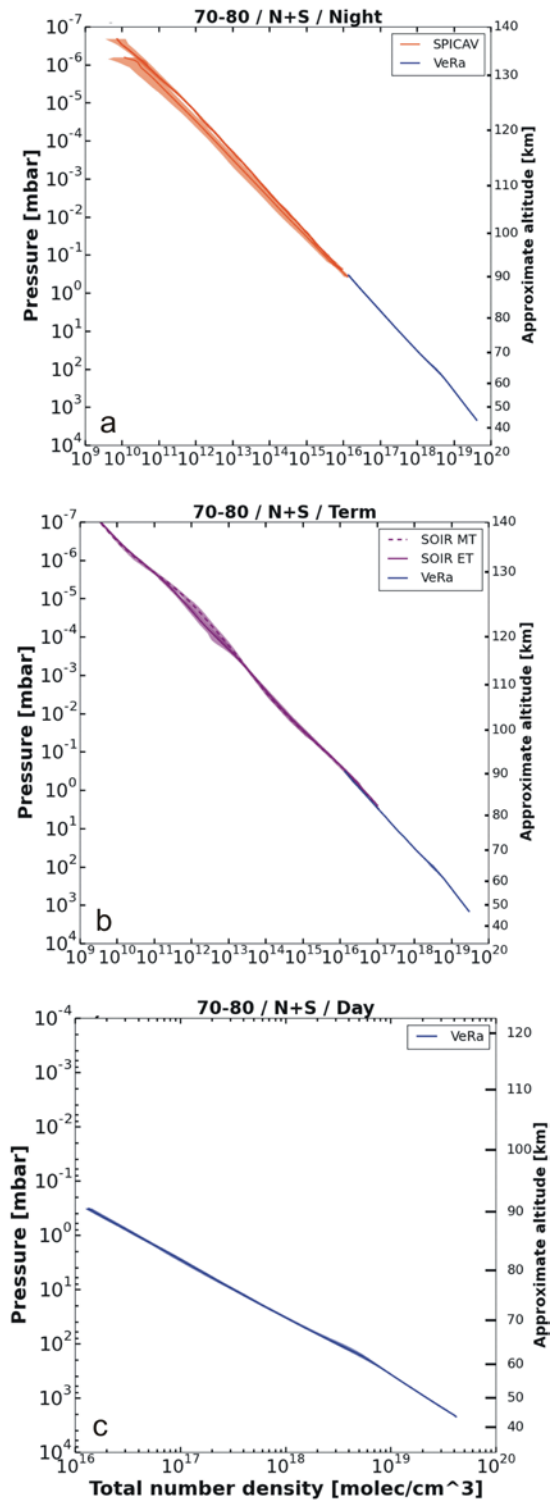
94

95 **Figure 22:** same as Figure 21 but for the latitude bin 30° to 50°. Panel (a): night side, panel (b):  
 96 terminator zones, panel (c): day side. The colored areas mark the one standard deviation uncertainty of  
 97 the average profiles for each experiment. Approximate altitudes are shown on the right hand vertical  
 98 axis.



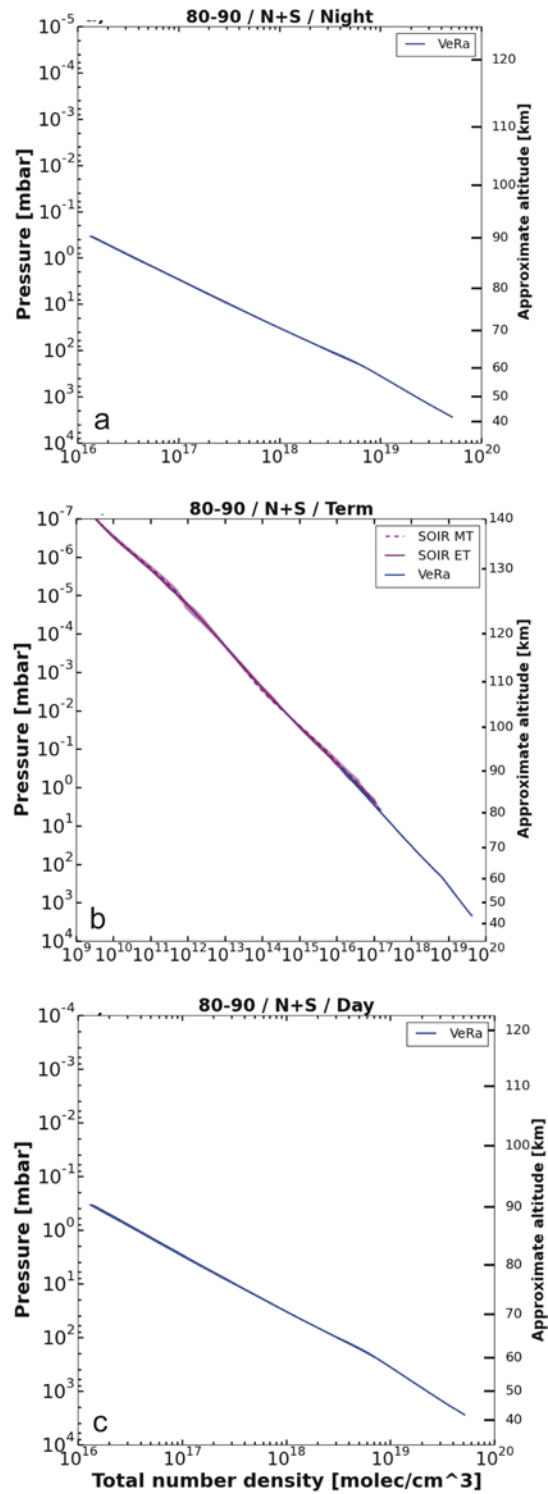
99

100 **Figure 23:** same as Figure 21 but for the latitude bin 50° to 70°. Panel (a): night side, panel (b):  
 101 terminator zones, panel (c): day side. The colored areas mark the one standard deviation uncertainty of  
 102 the average profiles for each experiment. Approximate altitudes are shown on the right hand vertical  
 103 axis.



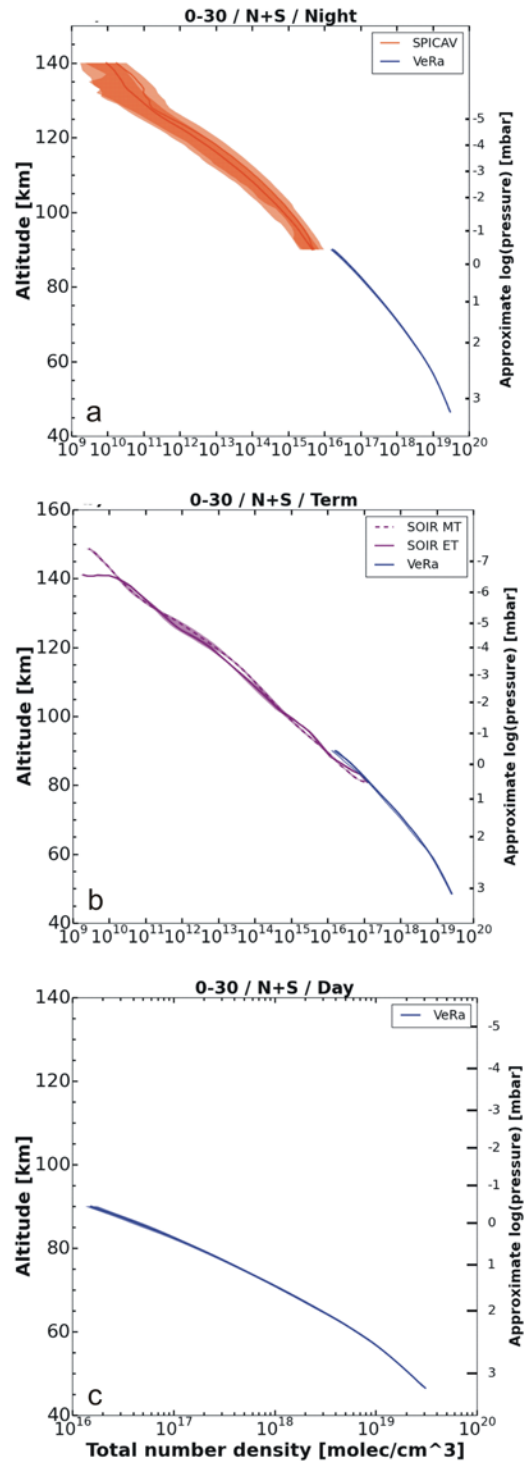
104

105 **Figure 24:** same as Figure 21 but for the latitude bin 70° to 80°. Panel (a): night side, panel (b):  
 106 terminator zones, panel (c): day side. The colored areas mark the one standard deviation uncertainty of  
 107 the average profiles for each experiment. Approximate altitudes are shown on the right hand vertical  
 108 axis.



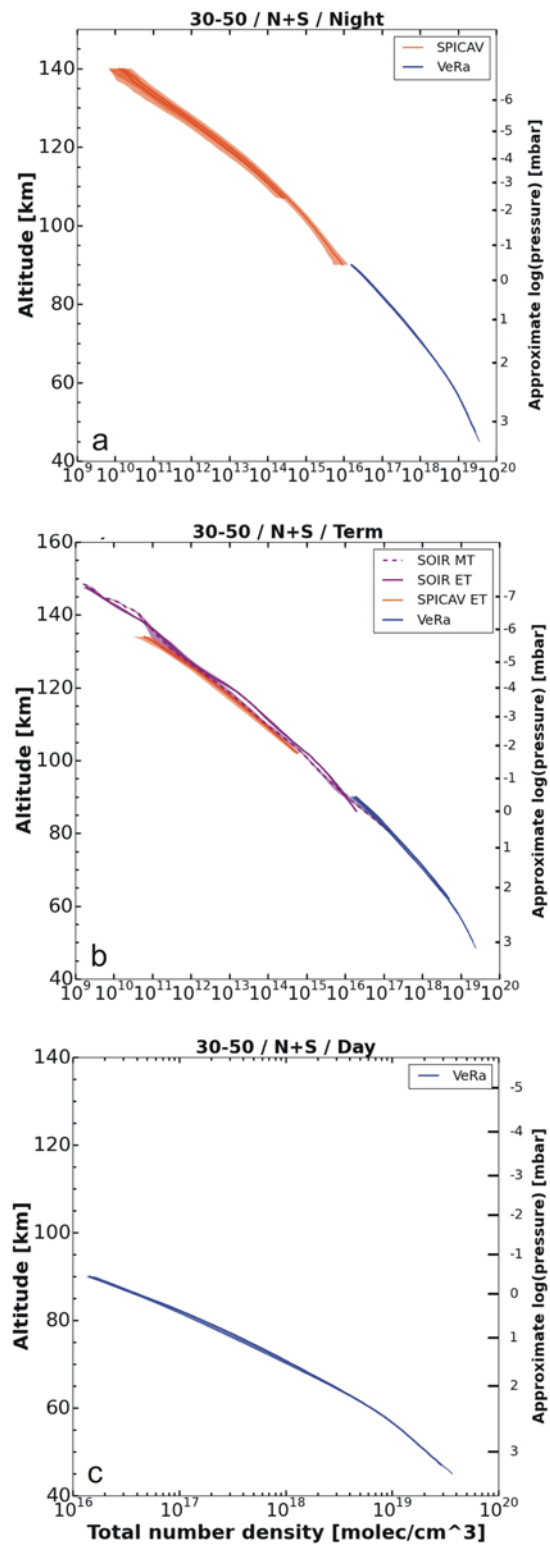
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110 **Figure 25:** same as Figure 21 but for the latitude bin 80° to 90°. Panel (a): night side, panel (b):  
 111 terminator zones, panel (c): day side. The colored areas mark the one standard deviation uncertainty of  
 112 the average profiles for each experiment. Approximate altitudes are shown on the right hand vertical  
 113 axis.



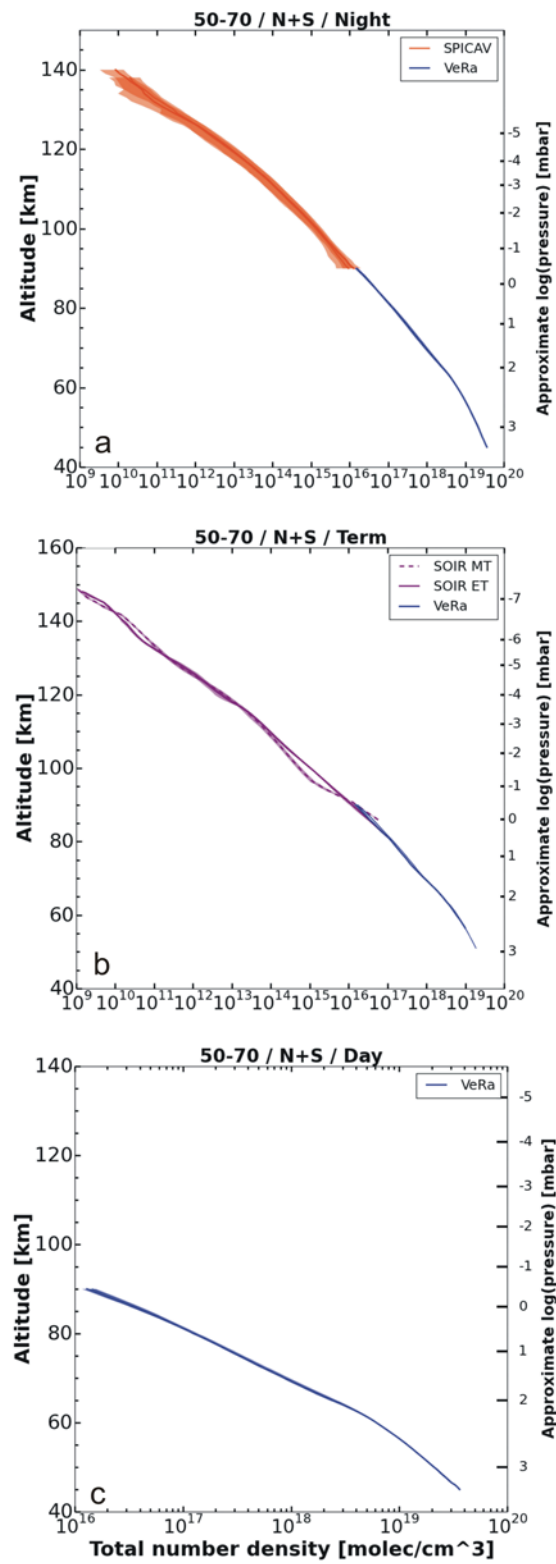
114

115 **Figure 26:** Comparison of atmospheric mean total density profiles from SOIR, SPICAV and VeRa as a  
 116 function of altitude for the near-equatorial latitudes (0° to 30°). Panel (a): night side, panel (b)  
 117 terminator zone, panel (c) night side. The colored areas mark the uncertainty of the respective average  
 118 profile as one standard deviation. Approximate pressure is shown at the right hand side vertical axis.  
 119 Night side SPICAV profiles are shown separately for the Northern and Southern hemispheres.



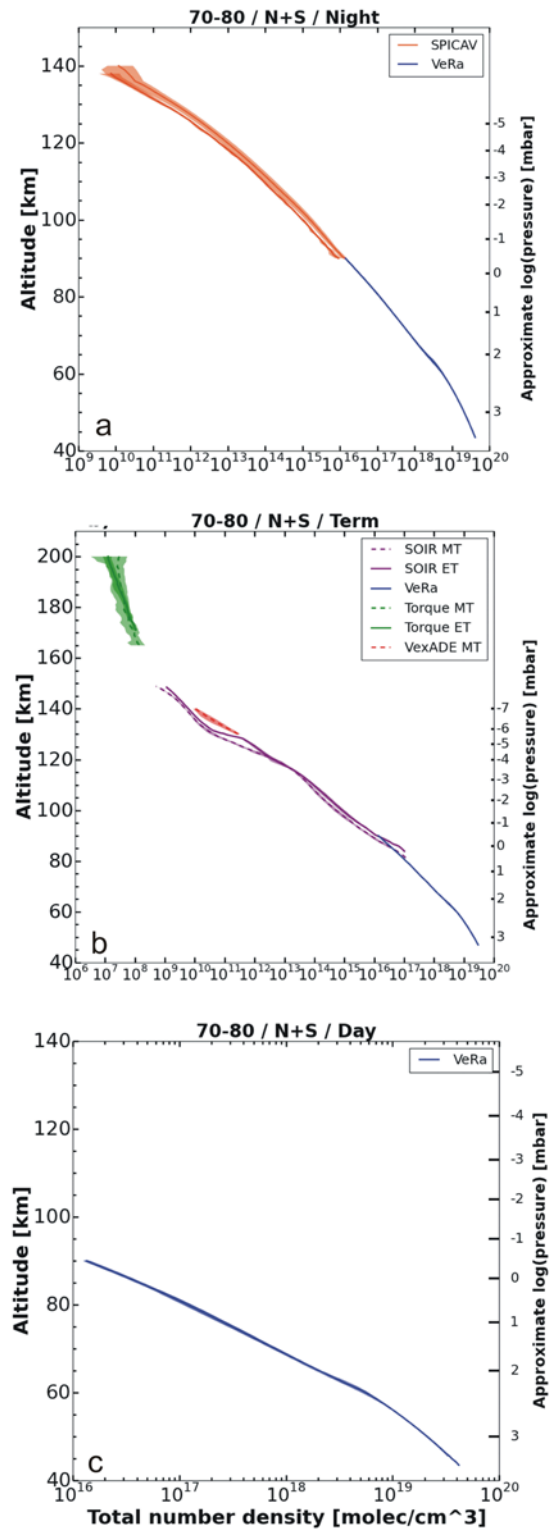
120

121 *Figure 27: same as Figure 25 but for the mid-latitudes (30° to 50°). Panel (a): night side, panel (b)*  
 122 *terminator zone, panel (c) night side.*



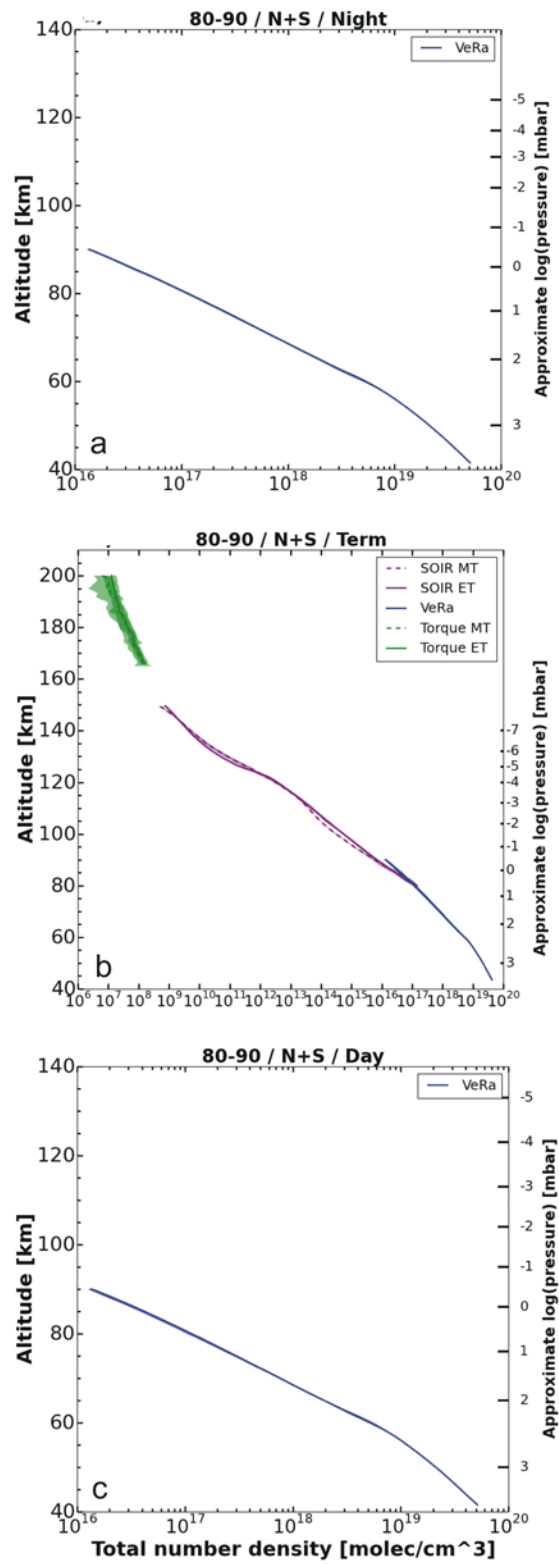
123

124 **Figure 28:** same as Figure 26 but for the mid-latitudes (50° to 70°). Panel (a): night side, panel (b)  
 125 terminator zone, panel (c) night side.



126

127 **Figure 29:** same as Figure 26 but for the high latitudes (70° to 80°). Panel (a): night side, panel (b)  
 128 terminator zone, panel (c) night side.



129

130 **Figure 30:** same as Figure 26 but for the polar latitudes (80° to 90°). Panel (a): night side, panel (b)  
 131 terminator zone, panel (c) night side.

132

## 133 **5. Comparison between Observations and Numerical Simulations**

134 It is proof of a good understanding of the thermal atmospheric structure if the observed physical and  
135 spectral properties of the atmosphere along with clouds, hazes and the insolation are reproduced by  
136 numerical simulations. Numerical models were developed based on observations by Venera, Pioneer  
137 Venus and Venus Express in order to compute the solar and thermal fluxes within the atmosphere and  
138 to determine the energy balance. The typical output products of these numerical models are  
139 temperature and neutral number density profiles (Crisp, 1986; 1989; Crisp & Titov, 1997; Bullock &  
140 Grinspoon, 2001; Eymet et al., 2009; Lee & Richardson, 2011; Lee et al., 2012; Mendonca et al., 2015).  
141 Reviews of the radiative balance of the Venus atmosphere are in Titov et al. (2007, 2013). An important  
142 *a-priori* input parameter to the simulations is the opacity distribution function and a first-guess  
143 temperature profile, usually taken from VIRA. The opacities are computed from the cloud properties  
144 and structures, the gas composition of the atmosphere, and the spectral properties of the different gas  
145 constituents.

146 It would be worthwhile to compare the output products of the various models in order to assess their  
147 capabilities. It was decided not to do so because the various models are progressively evolving. Only key  
148 aspects of the on-going modeling efforts shall be described below. The computation of atmospheric  
149 opacities is a crucial part of the modeling of the radiative transfer. Gas opacities are derived from line-  
150 by-line models which are based on spectroscopic databases such as HITRAN (Rothman et al., 2009,  
151 2013) and HITEMP (Rothman et al., 2010) and on assumptions of profiles of the atmospheric  
152 composition, in particular profiles of CO<sub>2</sub>, H<sub>2</sub>O and SO<sub>2</sub> which play an important role for the radiative  
153 transfer. The computation of the gas opacities requires assumptions on the line shapes and procedures  
154 in the line-by-line models. The continuum is highly uncertain between the dominant absorption bands,  
155 and difficult to determine experimentally (Wordsworth et al., 2010; Snels et al., 2014).

156 The cloud opacity is computed from a cloud distribution model and from assumed cloud particle  
157 properties (Knollenberg & Hunten, 1980; Zasova et al., 1999; 2007). These properties were determined  
158 from space observations mostly in equatorial regions. It is well known, however, that the cloud structure  
159 varies with latitude, with the vertical distribution (e.g. Ignatiev et al., 2009) and with the particle size  
160 (Wilson et al., 2008). The analysis of the VIRTIS-M data by Haus et al. (2013; 2014) determined the cloud  
161 structure as a function of latitude. Work is currently on-going to improve the understanding of the cloud

162 structure and their properties and characteristics ("Venus cloud structure" team supported by ISSI;  
163 Wilson et al. (2014)).

164 The properties and the distribution of the so-called unknown UV absorber within the clouds are very  
165 important for the computation of the vertical profiles of the solar flux absorption. This distribution is  
166 based on the mode-1 particles (smallest mode in the cloud particle distribution) in the upper cloud deck  
167 in many models (Crisp, 1986; Lee and Richardson, 2011; Mendonca et al., 2015), except for the most  
168 recent one (Haus et al., 2015).

169 Different techniques were used to compute the thermal cooling rates and the solar heating rates based  
170 on these opacities and the vertical structure of the atmosphere. The various radiative transfer  
171 algorithms may yield differences in the derived profiles of radiative heating and cooling rates and may  
172 therefore influence the modeled temperature structure. Large day/night variations are possible above  
173 the clouds due to short radiative time scales. Non-LTE processes and EUV heating have to be considered  
174 in order to compute correctly the thermal balance above approximately 100 km. The atmosphere is  
175 heated below 140 km by the absorption of solar radiation due to CO<sub>2</sub> near-infrared bands (2.7 μm,  
176 4.7 μm and 1-2 μm). The EUV absorption by CO<sub>2</sub>, O and a number of minor species dominates above  
177 that altitude. Thermal cooling occurs via CO<sub>2</sub> non-LTE transitions around 15 μm which competes with  
178 the heating terms together with thermal conduction (above about 150 km) to control the temperature.  
179 The modeling of these processes is quite complex, because it involves the non-LTE distribution of CO<sub>2</sub>  
180 energetic states and their associated ro-vibrational transitions. This requires models which consider the  
181 theory properly which solve simultaneously the statistical equilibrium and radiative transfer equations,  
182 very time expensive computations for the currently most advanced GCMs. Parameterizations of the  
183 15 μm-cooling and the NIR non-LTE heating based on results by Roldan et al. (2000) were already  
184 implemented into GCMs. The various authors, however, used different formulations (Brecht and  
185 Bougher; 2012; Gilli et al.; 2014). Bougher et al. (1986) used off-line simulated "look-up tables" for the  
186 solar heating rates and a parameterized scheme for the cooling which implements a line-by-line model  
187 of CO<sub>2</sub> 15 μm rates (taken from Roldan et al., 2000). Gilli et al. (2014) applied an analytical formula to  
188 reproduce the solar heating rates in those upper regions, and a complete but simplified non-LTE model  
189 for the 15 μm cooling, as it was also developed for the Mars Climate Data Base (MCD) GCM (Gonzalez-  
190 Galindo et al., 2009; 2013). It is assumed that the net absorption depends mainly on the density of the  
191 atmosphere, and to a smaller degree on the solar zenith angle, thermal structure and atomic oxygen

192 abundance. The EUV absorption is also parameterized assuming an efficiency of 20-22 %. The variation  
193 of the UV solar flux with the solar cycle is also taken into account.

194 One-dimensional radiative-convective equilibrium (RCE) procedures define the vertical temperature  
195 profile, e.g. for globally averaged conditions. Although Bullock & Grinspoon (2001) could achieve an  
196 excellent agreement with VIRA, most other RCE models could not (e.g. Lee & Richardson, 2011;  
197 Mendonca et al., 2015). The solar flux absorption, the cloud particle distribution and selected opacity  
198 parameters (Lebonnois et al.; 2015) as a function of altitude have a direct influence on the derivation of  
199 the temperature profile (e.g. Lee et al., 2012). Current GCM simulators indicate that both radiative and  
200 dynamical effects play a crucial role in the determination of the upper atmosphere thermal structure.  
201 The observed high variability of atmospheric quantities by the VEX instruments was not reflected in  
202 previous empirical models (e.g. VIRA and VTS3).

203 EUV absorption above 140 - 150 km altitude generates high temperatures above a cold layer around  
204 125-130 km where the NIR heating is weak. A local maximum produced by CO<sub>2</sub> IR bands solar  
205 absorption during daytime is advected to the terminator below 125 km (Brecht and Bougher, 2012; Gilli  
206 et al., 2014). This S-shaped structure with minima and maxima at the terminator is also observed by  
207 SOIR (Mahieux et al., 2015a). The pressure levels and magnitudes of modeled and observed  
208 temperatures, however, do not always agree. A warm region at the night side resulting from the  
209 subsidence of the day-to-night circulation air is predicted at 110 - 115 km altitude and indeed observed  
210 by SPICAV (Piccialli et al., 2015) but at lower regions.

211 Mahieux et al. (2016) attempted to reproduce the thermal structure observed by SOIR at the  
212 terminator. They developed a one-dimensional conductive radiative model which considers the heating  
213 and cooling terms of the main Venus atmospheric species CO<sub>2</sub>, N<sub>2</sub>, O, CO, H<sub>2</sub>O, HCl and SO<sub>2</sub> extending  
214 from 80 km to 180 km altitude. The modes 1 and 2 of the aerosols are considered in order to reproduce  
215 correctly the temperature profile in the mesosphere. The vertical number density profiles of the  
216 aerosols are in good agreement with the SOIR (Wilquet et al., 2009) and SPICAV-UV observations up to  
217 100 km. The aerosol profiles show constant values for both modes between 100 and 120 km which  
218 rapidly decrease at higher altitudes. The neutral number density is unfortunately lower than the  
219 detection sensitivity of both instruments, and thus cannot be confirmed by spacecraft measurements.

220 Several sources of uncertainties are investigated to improve the comparison between modeled and  
221 observed temperature profiles above the clouds: the ratio between O and CO<sub>2</sub> and its role in the cooling

222 rate, the parameterization implemented to simulate non-LTE radiative processes, potential heat sources  
223 above the cloud tops (including aerosols, upper haze layer, and unknown UV absorber), and tides,  
224 gravity waves or other sources (Zalucha et al., 2013; Gilli et al., 2016).

225 The cloud and haze particle distribution, the continuum gas opacity in the 3-7 microns spectral range for  
226 the extreme conditions in the deep atmosphere need to be sufficiently known in order to model the  
227 temperature profile. The latitudinal cloud distribution is also important to improve the understanding of  
228 the formation of the "cold collar" feature. The formulation of the non-LTE processes is still an  
229 approximation and a more accurate description is required for a GCM. Their implementation, however,  
230 considerably improved the knowledge of the energy budget in the upper atmosphere. The uncertainty  
231 of typical rate coefficients used in non-LTE simulations is still very large which is true for the uncertainty  
232 of the O-CO<sub>2</sub> collisional relaxation rate important for the cooling of the atmospheres of terrestrial  
233 planets in general.

234 The understanding of the 3-dimensional temperature structure and its variability requires a General  
235 Circulation Model which fully considers the dynamical interactions within the atmosphere. Significant  
236 progress has been achieved within the most recent Venus GCM models (Sugimoto et al., 2014a; 2014b;  
237 Ando et al., 2016; Lebonnois et al., 2016), but it is still on-going work.

## 238 **6. Discussion**

239 The Venus Express mission has considerably increased the knowledge of the Venus atmospheric  
240 structure above ~40 km and provided enough new information above 100 km to trigger new ideas for  
241 the interpretation of the observations. Three kinds of occultation experiments were performed for the  
242 first time to provide temperature profiles over a wide range of altitudes from 40 to 170 km.  
243 Considerable temperature variability is seen above 100 km. Certain features appear to be systematically  
244 present, such as a succession of warm and cool layers. Models support the existence of such layers  
245 consistent with a large scale circulation, but they are still in the process to be improved.

246 Although there is general agreement between the various experiments which observed the vertical and  
247 latitudinal temperature structure of the Venus atmosphere, the differences between individual  
248 experiments are larger than the measurement errors. The large range of temperatures found from  
249 HIPWAC-THIS is not consistent with most other observations and a temporal variability of this order  
250 appears unphysical. It is possible that the field-of-view may be inadequately accounted for or there are  
251 other analysis factors that may influence the results.

252 The processing of the occultation data is based on the common assumption of a spherically symmetric  
253 atmosphere - assuming that the atmospheric properties are the same in a given altitude region  
254 regardless of location or local time. This assumption is certainly violated at the terminators or in the  
255 presence of clouds. Vertical profiles of cloud particle sizes obtained from the Venera entry probes and  
256 balloons show significant differences when compared with the Pioneer Large Probe observations. The  
257 particle size distribution is constrained at the unit optical depth by the frequent observations of glory at  
258 four VMC wavelengths, but larger particles may exist in the deeper cloud layer. All VEx occultation  
259 experiments used the common atmospheric composition of 96.5% CO<sub>2</sub>, 3.5% N<sub>2</sub> below 90 km when  
260 starting this ISSI study. The assumed mean molecular weight dependence on altitude may be different  
261 from the actual Venus atmosphere above 90 km.

262 Radio occultation experiments experience a multi-path interference dominantly above 75° latitude at  
263 the first inversion layer at about 65 km altitude defined as the tropopause (Pätzold et al., 2007)  
264 implying an even cooler temperature inversion by an additional 15 K over an extremely short altitude  
265 range than previously thought.

266 One potential source of bias in the SPICAV stellar occultation data retrieval may be caused by the  
267 simplification of the radiative transfer complexity beyond the single scattering hypothesis. At present,  
268 only the forward scattering direction is considered by the retrieval model and the radiative transfer is  
269 approximated by a standard and simple Beer Lambert's law. The line-of-sight of the instrument becomes  
270 increasingly sensitive to contributions from scattering processes occurring in a narrow angle around the  
271 SPICAV field of view as the line of sight intersects denser and denser atmospheric layers. This has not  
272 been quantified yet but may contribute to the lesser reliability in the lowest sounded parts of the  
273 atmosphere (below 100 km).

274 One still unresolved mystery is the discrepancy seen in the neutral number densities observed by SOIR  
275 and SPICAV above ~90 km altitude and by the drag and torque experiences at much higher altitude by  
276 nearly two orders of magnitude. The radio occultation density profiles do not show such a large  
277 variation in the region of overlap (75 - 90 km). One of the key assumptions in the data reductions of the  
278 occultation experiments is the constant composition or well mixed atmosphere which is well maintained  
279 below 100 km. There are evidences, however, that the molecular composition might not be constant  
280 above 120 km. SOIR experiment results show a considerably varying homopause altitude between  
281 ~120 km and 150 km. SPICAV results see the homopause altitude between 119 and 138 km with

282 temporal and spatial variability. The homopause altitude has a clear dependence on the local solar time.  
283 It occurs at a higher altitude on the morning side than on the evening side (Piccialli et al., 2015).

284 The mean molecular weight is needed wherever the hydrostatic equation is used, and changes in its  
285 value due to change in composition at higher altitudes may become non-negligible. The assumption of  
286 hydrostatic balance may no longer be valid if large horizontal density gradients above 120 km exist as  
287 suggested by the drag experiments, SPICAV and SOIR observations and if those create very turbulent  
288 large scale motions, .

289 The Venus Express instrument data as well as the ground-based observations will definitely improve the  
290 current empirical models. Comparing VEX temperature profiles with VIRA there are small differences  
291 below 0.1 mbar (75 km), but there are large discrepancies above 0.1 mbar (75 km) at polar latitudes  
292 ( $>65^\circ$ ) (Tellmann et al., 2009). The new observations may allow the improvement of the empirical  
293 models (VTS3 and Keating et al. 1985) above 0.1 mbar (100 km). The spatial coverage is better above 90  
294 km, but still incomplete. Some experiments do provide sufficient global coverage and compare solar  
295 thermal tidal components with previous results. The focus of this inter-comparison, however, is on  
296 averaged temperature and density profiles. While the night side is very well covered by VEX instrument  
297 observations, the day side above 0.03 mbar (100 km) presents a region of ignorance with almost no  
298 observations in particular at high latitudes. GCMs must be used to predict the conditions of these high  
299 altitude day side regions where observations are lacking. The large number of experiments and  
300 investigations that have contributed to the knowledge of the thermal and density structure of the Venus  
301 atmosphere give now a consistent, coherent but still somewhat incomplete picture. Improved spatial  
302 coverage is necessary – both in latitude-longitude and at all local solar times - at all altitudes. The  
303 vertical and horizontal resolutions are different for the various experiments. Distinguishing  
304 uncertainties from variability is constrained because of the limited spatial and temporal coverage for a  
305 given instrument/experiment. It is therefore very challenging to identify temporal variability when  
306 comparing the results of various experiments.

307 Future missions must address in-situ observations by long lived aerial platforms and descent/ascent  
308 probes and landers in order to verify remote sensing observations.

309 The interpretation of SOIR and VIRTIS data requires the information on spectral line shapes. HITRAN is  
310 here the more commonly used database and there has been some improvements recently (Rothman et  
311 al., 2013).

312

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