



























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TOI–1278 B: SPIRou Unveils a Rare Brown Dwarf Companion in Close-in Orbit around an M Dwarf

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Abstract

We present the discovery of an $18.5 \pm 0.5 M_{\text{Jup}}$ brown dwarf (BD) companion to the M0V star TOI–1278. The system was first identified through a percent-deep transit in Transiting Exoplanet Survey Satellite photometry; further analysis showed it to be a grazing transit of a Jupiter-sized object. Radial velocity (RV) follow-up with the SPIRou near-infrared high-resolution velocimeter and spectropolarimeter in the framework of the 300-night SPIRou Legacy Survey carried out at the Canada–France–Hawaii Telescope led to the detection of a Keplerian RV signal with a semi-amplitude of $2306 \pm 10 \text{ m s}^{-1}$ in phase with the 14.5 day transit period, with a slight but nonzero eccentricity. The intermediate-mass ratio ($M_{*}/M_{\text{comp}} \sim 31$) is unique for having such a short separation ($0.095 \pm 0.001 \text{ au}$) among known M-dwarf systems. Interestingly, M-dwarf–BD systems with similar mass ratios exist with separations of tens to thousands of astronomical unit.

Key words: Brown dwarfs – Radial velocity – Transit photometry – Transits – M dwarf stars

1. Introduction

Over the past 25 yr, the search for and characterization of planets around other stars has moved from speculative to one of the most active fields in astronomy. Among planetary systems, those orbiting the coolest stars, M dwarfs, are of particular interest for a number of reasons. They provide insight into planetary formation in a regime of host-star mass that differs significantly from our own solar system, with some systems

more reminiscent in their architecture and scale to the Galilean moon system. Furthermore, as M dwarfs have radii and masses that are significantly smaller than those of our Sun (M0V to M9V dwarfs ranging, respectively, from 0.6 to $0.1 R_{\odot}$ and 0.6 to $0.1 M_{\odot}$), planets are generally easier to discover and characterize through transits than similarly sized planets around sunlike stars. This is particularly true for older, slowly rotating M dwarfs, with activity being a limiting factor for a number of

low-mass stars. Furthermore, the lower luminosity of M dwarfs (0.1 to $10^{-3.5} L_{\odot}$; Veeder 1974; Faherty et al. 2016) implies that the habitable zone (HZ) is much closer to the host star than in our solar system, implying that the orbits of planets in the HZ can be sampled on timescales of days. Given that M dwarfs make up the majority of stars among nearby stellar systems (being five times more numerous than FGK stars), it is likely that most of the HZ planets within 10 pc are found around M dwarfs (Figueira et al. 2016). This is exemplified by our closest stellar neighbor, the M5.5V Proxima Centauri, having a terrestrial planet in the HZ (Anglada-Escudé et al. 2016).

As they provide an easier path to the characterization of HZ terrestrial planets, planets around M dwarfs have been the subject of considerable efforts. Transit searches dedicated to or optimized for M dwarfs (Irwin et al. 2009) have been successful in identifying systems such as the seven Earth-sized planets around TRAPPIST-1 (Gillon et al. 2017), and the ongoing Transiting Exoplanet Survey Satellite (TESS) mission has a top-level Requirement of uncovering 1000 small ($R_p < 4 R_{\oplus}$) planets and measuring the masses of 50 such planets (Ricker et al. 2014). This requirement is not restricted to M dwarfs, but the TESS bandpass was designed to extend significantly to the red to maximize its sensitivity to planets around M dwarfs, compared to the Kepler mission, which had a bluer bandpass better matched to sunlike stars. While Earth-sized HZ planets receive significant attention, a number of statistical properties of the M-dwarf planets have emerged recently regarding more massive companions to M dwarfs.

Jupiter-mass planets are remarkably rare around M dwarfs at short separations. While 0.4%–1% of sunlike stars host a hot Jupiter (occurrence rate estimates differ depending on detection method considered; Wang et al. 2015), they are rarer around M dwarfs (Meyer et al. 2018); although, a few systems have recently been discovered (e.g., Bayliss et al. 2018). This is most unlikely to be an observational bias, as a hot Jupiter around an M dwarf would induce a $>100 \text{ m s}^{-1}$ radial velocity (RV) signal (from Kepler’s laws, combining a $1 M_{\text{Jup}}$ on a 4 day orbit around a $0.1\text{--}0.5 M_{\odot}$ host) and a readily detectable transit ($\sim 4\%$ to $\sim 100\%$ deep for M0V to M9V; from radius relations in Boyajian et al. 2012; Filippazzo et al. 2015).

Here, we report the discovery of an $18.5 M_{\text{Jup}}$ transiting brown dwarf (BD) companion on a 14.5 day orbit around the M0V star TOI-1278 (see Table 1 for basic properties of the host star), above the maximum mass expected to form through a protoplanetary disk; typically $<1 M_{\text{Jup}}$ with only a few objects between 1 and $10 M_{\text{Jup}}$ in the 90-star sample by Ansdell et al. (2016). It occupies a nearly empty part of the mass distribution of close-in companions of M dwarfs, between the Jupiter population and the massive BD companions. In Section 2 we present discovery and follow-up observations of the system. In Section 3 we assess the properties of the host star, and in Section 4 we detail the data analysis. Finally, in Section 5 we discuss the properties of the system in comparison to other planetary and BD companions, its dynamical fate, and characterization prospects.

2. Observations

We present below the TESS discovery data of TOI-1278, as well as its ground-based follow-up observations in photometry,

high-resolution imaging, low-resolution spectroscopy, and precise near-infrared (NIR) velocimetry, all required to establish the nature of the transiting object TOI-1278 B.

2.1. TESS Photometry

TOI-1278 was observed by TESS (Ricker et al. 2014) in sector 15 with CCD 4 on Camera 1 from 2019 August 15 to September 11. This target was selected in the Cool Dwarf List (Muirhead et al. 2018), a specially curated list of late-K and M dwarfs included in the TESS Input Catalog (TIC; Stassun et al. 2018, 2019) and the Candidate Target List for 2 minute cadence light-curve sampling. We used the publicly available³¹ Presearch Data Conditioning Simple Aperture Photometry (PDCSAP; Smith et al. 2012; Stumpe et al. 2012, 2014) light-curve produced by the NASA Ames Science Processing Operations Center (SPOC; Jenkins et al. 2016). The PDCSAP flux values are corrected for long-term systematic trends seen among other stars in the same sector/camera/CCD using Cotrending Basis Vectors. No long-term stellar variability on timescales shorter than the time series considered here are removed. Furthermore, these PDCSAP fluxes account for the dilution caused by nearby stars. The TIC contamination ratio defined as the flux from other sources within the aperture divided by the target star flux (in TESS-band) is 0.326395 for TOI-1278. This dilution correction is important, as any other flux contribution would lead to an underestimation of the transiting object radius by measuring a smaller transit depth. The TOI-1278 normalized PDCSAP light curve is presented in Figure 1 and clearly shows two transit events.

SPOC sector 15 data validation (DV) reports (Twicken et al. 2018; Li et al. 2019) first identified the two transits with a signal-to-noise ratio (S/N) of 16.5. Both events were reported to be V-shaped, suggestive of a high impact parameter, with an average depth of $1.03\% \pm 0.07\%$. This led to the announcement of planet candidate TOI-1278.01 with an orbital period of 14.476 days and an estimated radius of $9.3 R_{\oplus}$. Three consecutive transit detections are usually needed to confirm the previous two were produced by the same object. In this case, the two transits were found to have consistent duration, suggesting that they are caused by the same object rather than two planets within the same system or two planets orbiting different stars. We furthermore inspected the PDCSAP light curve near phase $\phi_{\text{sec}} = 0.496$, the predicted phase of secondary eclipse based on the geometry of the TOI-1278 B orbit constrained in Section 4. The eclipse signal due to reflected light is estimated at 13 ppm for a geometric albedo of 0.5. Given TESS photometric precision, no secondary eclipse is detected, and we can only infer a 3σ upper limit of 840 ppm for the eclipse depth.

2.2. Ground-based Transit Monitoring

While the position of TOI-1278 in the HR diagram (see Section 3) strongly suggests that the star is not itself an equal-luminosity eclipsing binary, the TESS detection could still be attributed to a nearby eclipsing binary (NEB) located within a few pixels. This is due to TESS’s coarse image sampling ($21''$ per pixel) and, hence, the need to confirm on-target transit events with arcsecond angular resolution ground-based observations. In the case of TOI-1278, additional transit

³¹ Mikulski Archive for Space Telescopes: <https://archive.stsci.edu/tess/>.

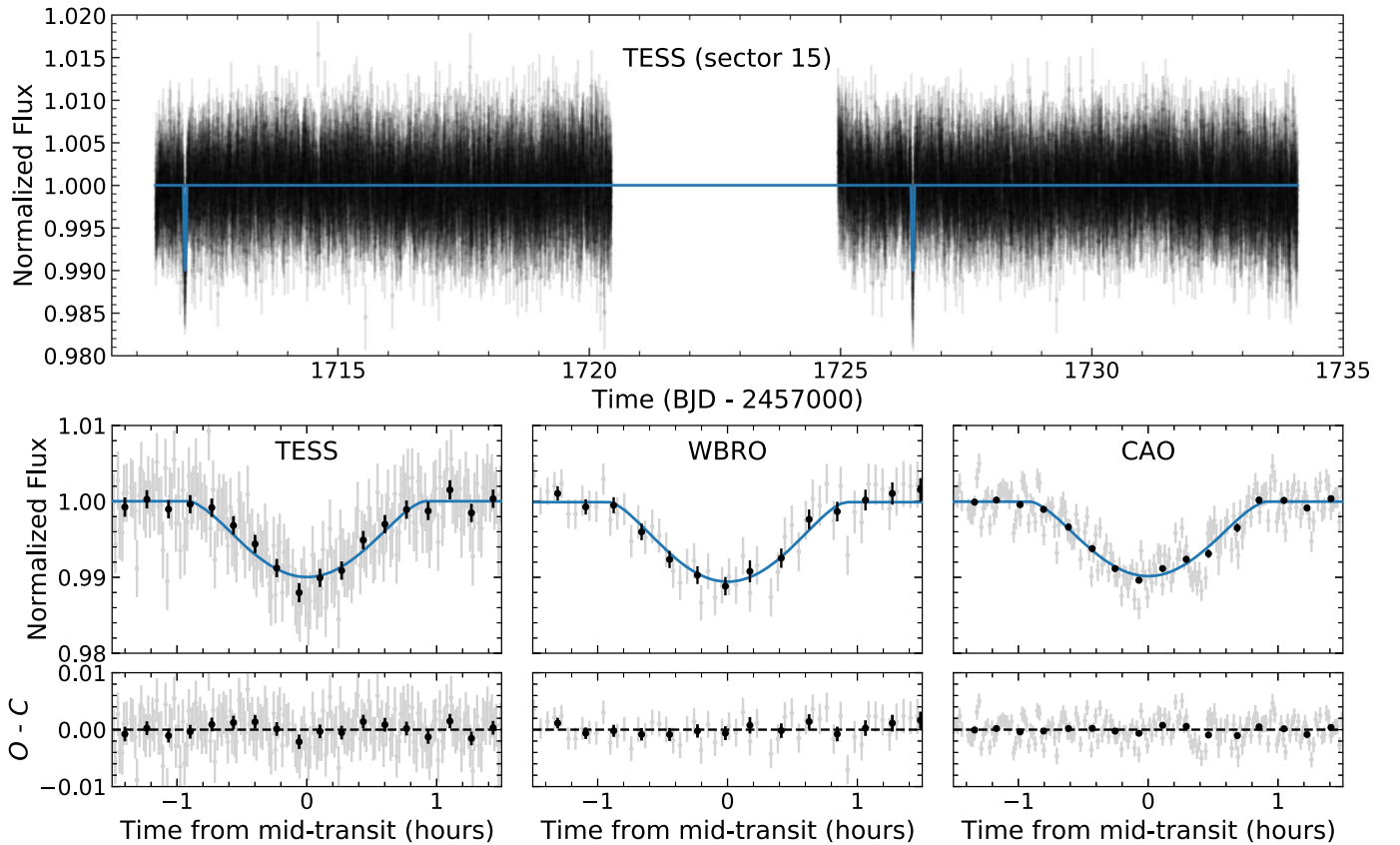


Figure 1. Top panel: TOI-1278 normalized PDCSAP light curve from sector 15. The gap in the middle of the light curve is due to data downlink when TESS is close to perigee. Lower panels: TESS, WBRO, and CAO transit light curves. The black points represent binned data (10 minutes temporal bin). The blue curves are the best-fit transit model, as discussed in Section 4. Only the baseline flux and the limb-darkening coefficients vary between the three data sets. For each instrument, the residuals (observed – calculated) are shown below the transit. The TESS data set is phase-folded and covers two transits.

observations were crucial to set stronger constraints on its light-curve determination, particularly on the orbital period P , the radius of the transiting object R_p , and the impact parameter b . Ground-based monitoring was spurred by the TESS discovery and was obtained to further constrain transit properties. We did not obtain a full out-of-transit light curve.

We scheduled two transit observations using the TESS Transit Finder (TTF), a customized version of the Tapir software package (Jensen 2013). We first observed TOI-1278 at Wild Boar Remote Observatory (WBRO) near Florence, Italy on 2019 October 27 without any filter (*clear*). A grazing, percent-deep, transit was detected on-target using a $4''.7$ aperture and was on time given the TTF prediction and uncertainty. The second observation was conducted at *M.G. Fracastoro* station (Mt. Etna, 1735 m a.s.l.) of the Catania Astrophysical Observatory (CAO), in Italy on 2019 November 25 in I_C with the 0.91 m telescope. Using a $5''.3$ aperture, a second V-shaped transit of the same depth was detected at the expected time. Both data sets were consistent with TESS in transit depth, shape, duration, and timing. Ground-based photometric follow-ups significantly increased the time baseline of our data set and further constrained the orbital fit. We have thus included these additional photometric observations, together with TESS transits, in our joint model described in Section 4. Calibration and light-curve extraction by differential photometry were accomplished using *AstroImageJ* (AIJ; Collins et al. 2017).

2.3. Ground-based Photometry

To constrain the stellar parameters of the host star, in particular the effective temperature, we obtained complementary optical photometry at the CAO, as the values reported in the TIC had relatively large uncertainties. Photometric measurements in B , V , R_C , and I_C were obtained on 2020 December 1 (see Table 2). Standard stars in the cluster NGC 7790 (Stetson 2000) were observed just after TOI-1278 to calculate the zero-points and the transformation coefficients to the Johnson-Cousins system. The errors on the B , V , R_C , and I_C magnitudes include both the photometric uncertainty coming from the photon statistics and the error of standardization based on the NGC 7790 data.

2.4. High-resolution Imaging

We searched for sources within a few arcseconds of TOI-1278 with Palomar/PHARO NIR adaptive optics imaging on 2019 November 9 in the $Br\gamma$ band. As shown in Figure 2, no companion is detected at 5σ with a contrast ratio $\Delta Br\gamma \leq 7.0$ mag for separation greater than $1''$ (76 au) and $\Delta Br\gamma \leq 6.335$ mag for separation over $0''.5$ (38 au). These contrasts thus exclude a companion brighter than $M_K \sim 11.7$, corresponding to an L4 spectral type and $L_{bol} = 10^{-4.1} L_\odot$ for field objects (using polynomial relations in Faherty et al. 2016). At an age >1 Gyr (see Section 5), this rules out stellar companions down to the upper limit of the BD regime (Phillips et al. 2020 models).

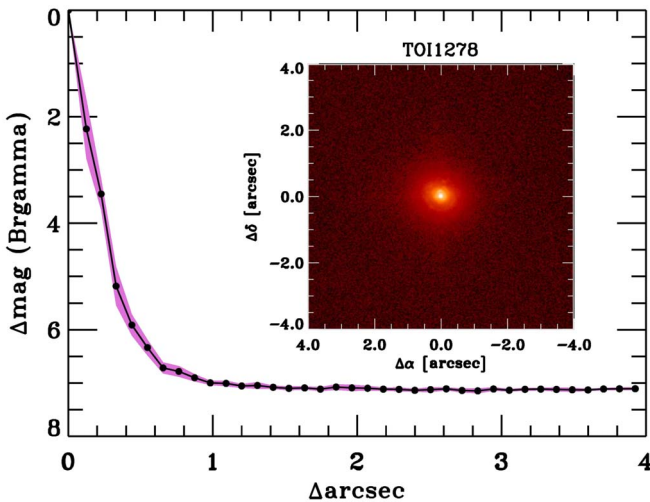


Figure 2. $\text{Br}\gamma$ 5σ contrast curve of TOI-1278 from Palomar/PHARO NIR adaptive optics imaging. No close companion is detected.

2.5. SPIRou Velocimetry

TOI-1278 was observed over 12 epochs from 2020 May 31 to November 1 using SPIRou, the NIR (0.98–2.5 μm) spectro-polarimeter installed on the Canada–France–Hawaii Telescope (Donati et al. 2018, 2020). The data set was obtained as part of the ongoing SPIRou Legacy Survey (SLS; Donati et al. 2020).

At each epoch, a four-measurement polarimetric sequence was obtained, each measurement having a 5 min integration, except for June 7, when eight measurements were obtained due to poorer observing conditions. Per-exposure S/Ns of 20–38 were obtained for the middle of the H band. These S/N values are for a spectral sampling of the SPIRou data of $\sim 2.2 \text{ km s}^{-1} \text{ pixel}^{-1}$. We rejected a handful of observations due to low S/N, resulting in 47 usable precision radial velocity (PRV) measurements at 10 different epochs. Table A1 compiles per-epoch radial velocity and associated uncertainties used in this analysis.

Observations were reduced with the standard data reduction pipeline for SPIRou (APER0, version 0.6.131; Cook et al. in prep). APER0 automatically handles all reduction steps necessary for PRV. In brief, the pipeline first corrects detector-related effects (specific to H4RG infrared arrays, i.e., the capacitive coupling between amplifiers, the $1/f$ noise, and the dark current; Artigau et al. 2018a). Then, APER0 locates the bad pixel positions, the exact positions of orders on the array, and determines the shape of the instrument slicer (Micheau et al. 2018) as well as determining the flat and blaze corrections to apply, using nightly calibration sequences. The calibrated data is then optimally extracted (Horne 1986) in both polarization channels, simultaneously (AB) and separately (A and B), as well as the simultaneous calibration channel (C). Extracted two-dimensional (E2DS, 49 orders by 4088 pixels) images are produced, and corrected for thermal emission. As the simultaneous Fabry–Perot calibration was used, leakage from the calibration channel into the science channels was also corrected. Nightly wavelength solutions are constructed using a combination of hollow-cathode and Fabry–Perot calibrations (Cersullo et al. 2019; Hobson et al. 2021). Barycentric corrections were determined within APER0 using the `barycorrpy` package (Kanodia & Wright 2018).

Telluric and night-sky emission correction is also done automatically using APER0. Night-sky emission is corrected using a principal component analysis (PCA) model of OH emission constructed from a library of high-S/N sky observations. Telluric absorption correction is done using a PCA-based approach on residuals after fitting for a basic atmospheric transmission model (TAPAS; Bertaux et al. 2014); the details of the PCA approach are described in Artigau et al. (2014), and its implementation is described in APER0 (Artigau et al. in prep).

With the telluric-corrected spectrum, we perform an RV measurement using a cross-correlation function (CCF) with a weighted mask of stellar absorption lines. The mask used must be constructed with a star that reasonably matches in temperature TOI-1278. From existing SPIRou Legacy Survey (SLS) data, we have a number of bright, nearby, M dwarfs that have been observed 100 times with an $\text{S/N} > 150$ in the H band. The best match to TOI-1278 was found to be the M0.5V Gl 846, and we use the CCF mask constructed from that spectrum through the PRV analysis.

The CCF is computed for each of the 49 SPIRou orders. Even though the spectra are telluric-corrected, parts of the domain between photometric bandpasses are unusable for PRV. We rejected all orders outside of the $YJHK$ bandpasses. The effective domain used is 980–1113 nm, 1153–1354 nm, 1462–1808 nm, and 1957–2400 nm. Per-order CCFs for usable orders were weighted by the S/N combined with their RV contents defined in Bouchy et al. (2001) and combined in a per-integration CCF. We used a cross-correlation of each observation’s CCF with the mean CCF of all observations to measure TOI-1278’s velocity. This resulted in a slightly lower dispersion of RV than fitting a Gaussian model to each individual CCF. No method that attempted to measure the velocity from the CCF profile (Gaussian fit, bisector mean position, adjustment of the mean CCF to each observation) led to a significant change in the main result: a $K \sim 2300 \text{ m s}^{-1}$ signal consistent in phase and with the period of the companion detected through the photometric monitoring (see Figure 3). The per-epoch RV uncertainty was computed by using Equations (6) and (13) in Bouchy et al. (2001) for the individual CCFs. The computed uncertainties are consistent with the point-to-point dispersion within each epoch. To prevent any biases related to underestimated RV errors, we explored the use of an additional jitter term in the RV model described in Section 4. We found that the posterior of the jitter term was consistent with zero, with a 1σ upper limit at 8 m s^{-1} , and that all other model parameters remained unaffected well within their 1σ uncertainty. We therefore conclude that the reported RV uncertainties are accurate, and we do not include the jitter term in the model presented in Section 4.

2.6. Near-infrared Spectroscopy

We obtained a low-resolution NIR (0.70–2.52 μm) spectrum of TOI-1278 on UT 2020 December 24 with the SpeX spectrograph at the NASA Infrared Telescope Facility (IRTF; Rayner et al. 2003) to improve the wavelength coverage of our spectral energy distribution and refine the estimated radius of TOI-1278. We used the $0''.8$ slit in prism mode, resulting in a resolving power $\lambda/\Delta\lambda \approx 75$.

Observing conditions were clear with a $0''.5$ atmospheric seeing. We obtained four 60 s exposures of TOI-1278 in an ABBA dither pattern along the slit at an airmass of 1.56,

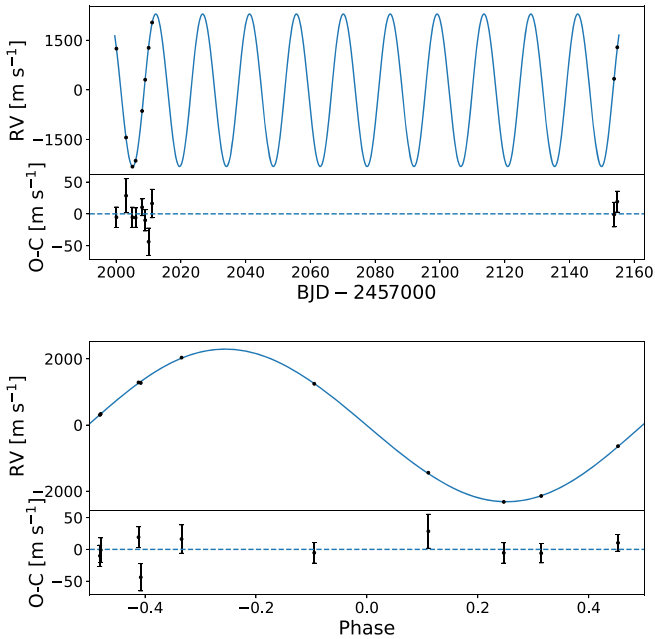


Figure 3. Top panel: RV of TOI-1278 from SPIRou monitoring with an overplot of the best-fit orbital solution. The RV measurements are binned per epoch (typically four observations per night). The median binned error bar is 16.5 m s^{-1} , and the per-epoch residual (shown below) rms is 19 m s^{-1} . The orbital fit is discussed in Section 4. Bottom panel: same as the top panel, but phase-folded using the best-fit orbital period and time of inferior conjunction.

followed by standard SpeX prism calibrations. This was immediately followed by the observation of the AOV-type spectral standard HD 203470 to correct for telluric absorption. We obtained 30 short exposures of HD 203470 in ABBA patterns along the slit at an average airmass of 1.64. We used exposure times in the range 0.5–5 s to maximize S/N while ensuring that we would obtain at least one set of exposures inside the linear range of the improved detector response of the recently upgraded SpeX detector (no sequence reached the nonlinear detector regime).

All data were reduced using `SpexTool v4.1` (Cushing et al. 2004). Two-dimensional traces were first extracted using an optimal extraction algorithm, and each resulting one-dimensional raw spectrum was visually inspected to remove bad pixel clusters that randomly occur in the two-dimensional frames during long exposures. The raw spectra are then combined in an S/N weighted average, and the raw science target is corrected for telluric absorption using the standard `xtellcor` procedure (Vacca et al. 2003).

3. Stellar Characterization

Other than for the presence of a transiting companion (see sections below), TOI-1278 is an unremarkable M dwarf at a distance of $75.46 \pm 0.07 \text{ pc}$ (Gaia Collaboration et al. 2021). In the local-neighborhood HR diagram (see Figure 4), TOI-1278 falls onto the main sequence of low-mass stars, suggestive that it is not an equal-luminosity binary nor particularly inflated due to youth. Also notable is the relatively low level of photometric activity of the star; it displayed a flat light curve, other than for the companion transits, at the part-per-thousand level over the month-long monitoring by TESS (see Figure 1). This argues against a young star that would have frequent flaring; such variability on timescales of hours to days would have been preserved by the PDCSAP photometry.

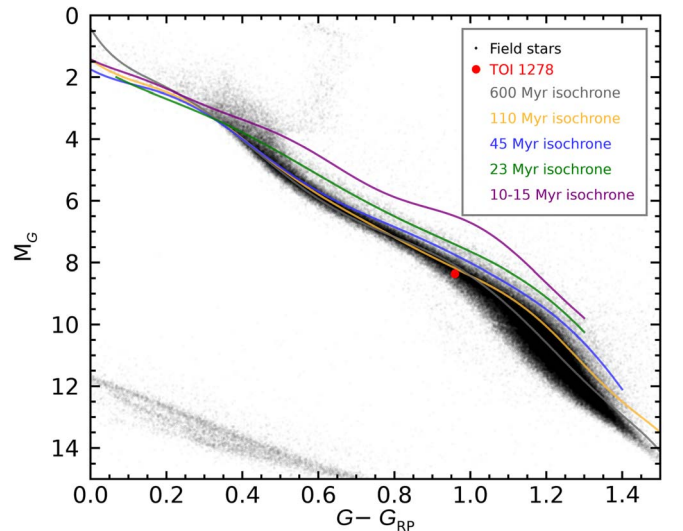


Figure 4. Position of TOI-1278 for the $<100 \text{ pc}$ volume-limited sample based on GAIA parallax and photometric measurements. TOI-1278 falls on the main sequence, and it is significantly offset from empirical tracks for young associations that are $<600 \text{ Myr}$ old. The position on the HR diagram also rules out a near-equal-luminosity binary, which would be offset from the main sequence by 0.75 mag (roughly the position of the 45 Myr isochrone). Empirical isochrones are from Gagné et al. (2018) and Gagné et al. (in prep.) and are based on the Upper Scorpius, Lower Centaurus Crux and Upper Centaurus associations (10–15 Myr), the β Pictoris and 32 Ori associations (23 Myr), TW Hydra, Columba and Carina associations (45 Myr), the Pleiades (110 Myr), and Coma Berenices cluster (600 Myr).

3.1. Spectral Energy Distribution and Radius

We used the method of Filippazzo et al. (2015) to estimate the bolometric luminosity of TOI-1278. We used the updated trigonometric parallax from Gaia eDR3 to calibrate all available broadband photometry (listed in Table 2), and our SpeX NIR spectrum was normalized to an error-weighted combination of the predicted absolute fluxes from the broadband photometric measurements that overlap with the spectrum, following Filippazzo et al. (2015). We adopted the effective temperature determined in Section 3.2 ($3799 \pm 42 \text{ K}$) and assigned a Rayleigh–Jeans and Wien tails of the appropriate blackbody temperature. Linear interpolation in logarithm space was used to construct the spectral energy distribution outside of the SpeX spectrum and between available broadband photometric measurements. Our resulting spectral energy distribution, shown in Figure 5, allowed us to measure a bolometric luminosity of $0.0614 \pm 0.0001 L_{\odot}$. This bolometric luminosity measurement is semi-empirical; one only needs to assume a temperature to set the shape of the Wien tail. Considering that it accounts for a small fraction of the overall flux, the impact of the uncertainty on the effective temperature is much smaller than the uncertainties propagated from photometry and parallax measurements. Regarding the Rayleigh–Jeans regime, the only assumption is that the two reddest photometric data points are well past the peak of the Planck function, free of strong molecular absorption and not significantly affected by a circumstellar disk. Considering that the spectral energy distribution peaks at $\sim 1 \mu\text{m}$ and that the reddest bandpasses are at 4.6 and $11.6 \mu\text{m}$, this assumption is justified. Combining this measurement with our model-dependent T_{eff} allows us to derive a semi-empirical radius measurement of $0.573 \pm 0.012 R_{\odot}$. This measurement is marginally larger (by $5\% \pm 4\%$) than predictions from the solar-metallicity models of

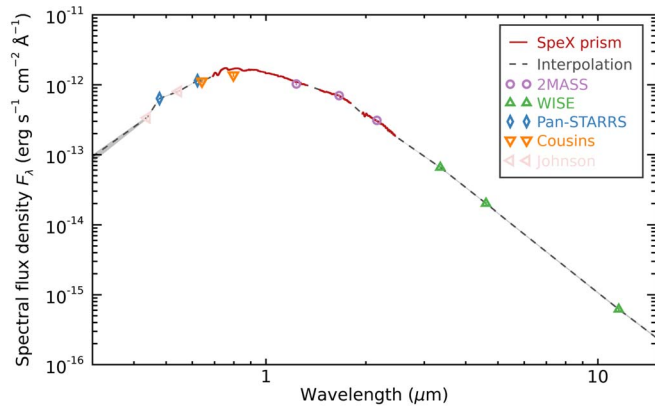


Figure 5. Spectral energy distribution of TOI-1278. Absolute fluxes estimated using photometry from various surveys are displayed as separate symbols, as indicated in the legend. Our NIR SpeX spectrum normalized with absolute fluxes from overlapping photometry is shown as a red solid line. Linear interpolation in logarithm space, and Rayleigh–Jeans and Wien tails of the appropriate effective temperature are shown as dashed lines to complete the spectral energy distribution used in our calculation of the bolometric luminosity.

Baraffe et al. (1998), and only $1\% \pm 3\%$ larger than empirical spectral type to radius relations for field stars (Pecaut & Mamajek 2013).³² This radius measurement is also consistent with interferometric measurements for nearby M dwarfs (Demory et al. 2009), which tend to be marginally inflated compared with solar-metallicity models, by about 0%–2%. This is in stark contrast with young M dwarfs, which tend to be significantly inflated (by 150%–250% at 20–40 Myr) compared with theoretical models and older M dwarfs, likely because of their faster rotation rates that drive strong magnetic fields, which in turn affect the hydrostatic balance of young M dwarfs (Malo et al. 2014). This indicates that TOI-1278 is likely not a young M dwarf, consistent with our small $v \sin i$ measurement (see Section 3.2).

Our SpeX low-resolution NIR spectrum indicates an M0 spectral type when compared with spectral type standards in the IRTF spectral library (Cushing et al. 2005; Rayner et al. 2009), consistent with our best-matching CCF mask from the M0.5V-type template Gl846. Constraining the age of TOI-1278 is key for further analysis of its properties; a young M0V would have an inflated radius and lower mass compared to gigayear-old objects. This would impact both the derived mass from RV measurements and radius estimate from transit properties. Also, uncovering a young transiting BD would be interesting to constrain evolution models.

Rare objects such as the young (93^{+61}_{-29} Myr) inflated transiting BD TOI-811 b provide complementary constraints to BD evolutionary models (Carmichael et al. 2020). This provides a complementary constraint compared to objects of similar mass and ages in young moving groups for which the radius is constrained through bolometric properties (Filippazzo et al. 2015).

Young BDs have a significant self-luminosity, and a young system would be significantly easier to detect as a double-lined spectroscopic binary. Furthermore, the TOI-1278 system could consist of an unresolved binary system with the transiting companion around one of the two components. While there are

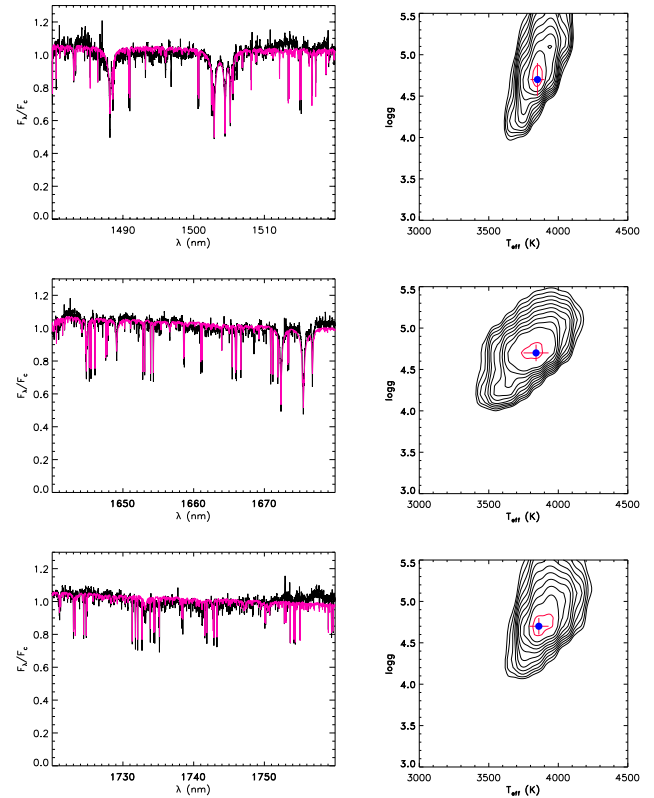


Figure 6. Left panels: continuum-normalized co-added SPIRou spectrum of TOI-1278 in three regions (black full lines) with the best-fitting synthetic spectrum over-plotted (red dotted lines). Right panels: χ^2 contour maps in the T_{eff} – $\log g$ plane. In each panel, the 1σ confidence level is denoted by the red contour. The best values and error bars on T_{eff} and $\log g$ are also indicated.

a number of well-established tools to rule out a number of scenarios for transiting companions (see Sections 2.1 and 2.4), simple color–magnitude diagram analysis can already exclude a number of configurations.

3.2. Stellar Parameters from SPIRou Spectra

The determination of the atmospheric parameters (T_{eff} and $\log g$) and projected rotational velocity ($v \sin i$) was accomplished through the code ROTFIT (e.g., Frasca et al. 2017).

We adopted a grid of synthetic BT-Settl spectra as templates (Allard et al. 2010) at $[\text{Fe}/\text{H}] = -0.5, 0.0, \text{ and } 0.3$, and effective temperature in the range 3000–4500 K (in steps of 100 K) and $\log g$ from 5.5 to 2.5 dex (in steps of 0.5 dex).

ROTFIT finds the best values of atmospheric parameters and $v \sin i$ by minimizing the χ^2 of the difference between the observed and synthetic spectra in specific spectral segments. We used the H band (eight 40 nm wide segments) as well as regions centered on strong lines within J and K (1290–1330 nm and 1960–2000 nm) for the ROTFIT analysis. These domains avoid the stronger telluric absorption bands between photometric bandpasses that may display significant correction residuals, and avoids the increased thermal noise on the red side of the K band. Some regions within Y and J are relatively clean of tellurics, but have fewer strong and sharp spectroscopic features than the regions selected for the analysis; this manifests itself in the very poor RV content of the Y and J bands for M dwarfs (e.g., Figueira et al. 2016; Artigau et al. 2018b). Figure 6 shows sample regions for the fit on the

³² See also http://www.pas.rochester.edu/~emamajek/EEM_dwarf_UBVJHK_colors_Teff.txt.

telluric-corrected co-added SPIRou spectrum. Both templates and observed spectra are normalized to the local continuum, and the templates are degraded to the SPIRou spectral resolution and resampled on the points of the target spectrum. The templates are also artificially broadened by convolution with a rotational profile whose $v \sin i$ spans a range of values to find the minimum of χ^2 . For each spectral segment, the best values of T_{eff} and $\log g$ are found by interpolation in the grid of templates, as shown in the right panels of Figure 6.

The values of $T_{\text{eff}} = 3799 \pm 42$ K, $\log g = 4.68 \pm 0.10$ dex, and $v \sin i = 1.10 \pm 0.86$ km s $^{-1}$ are the weighted means of the values from each i th spectral region, where the weights are $w_i = 1/\sigma_i^2$ (σ_i is the error of the parameter for the i th segment). The analysis was performed at three metallicity values, and the lowest χ^2 value was obtained for $[\text{Fe}/\text{H}] = 0.0$. The uncertainty of each parameter is the largest between the standard error of the weighted mean and the weighted standard deviation. The value of $v \sin i$ is significantly smaller than the spectral resolution of SPIRou (~ 4.3 km s $^{-1}$), and is sensitive to other broadening effects, such as the macro-turbulence, that are not distinguishable from rotation at low $v \sin i$ values. The $v \sin i$ should therefore be taken as a very approximate value.

3.3. Refining Metallicity Measurements

The high-resolution spectra of TOI–1278 let us determine the metallicity of the star with comprehensive examinations of individual Fe I lines in all Y , J_{2M} , H_{2M} , and K_{2M} bands. BT-Settl and Astrophysical Chemical Equilibrium Solver (ACES) models (Allard et al. 2011; Husser et al. 2013) are available at varying resolution over the SPIRou domain. The sampling of BT-Settl models is not uniform; it is typically 0.20 Å, but it can be as low as 0.05 Å, which is comparable with that of SPIRou, inside strong spectral lines. The determination of T_{eff} and $\log g$ discussed in the section above is not strongly sensitive to the resolution, but in contrast, metallicity measurements are more strongly sensitive to the data resolution. We therefore decided to use PHOENIX ACES models, available at a resolution higher than that of SPIRou, for metallicity measurements of all identified individual Fe I absorption lines. To do this, we generated multiple synthetic spectra for the fixed effective temperature and $\log g$ of 3800 K and 4.5 dex, respectively, and the metallicity range of -1.5 to 0.5 dex. The T_{eff} and $\log g$ values were adopted from ROTFIT results (see Table 2) and then rounded to the nearest 100 K and 0.5 dex.

To match the original resolution and sampling of the ACES models to that of the observed data, we degraded and convolved all of the generated synthetic data with respect to a SPIRou spectrum. Then we used the Vienna Atomic Line Database (VALD; Ryabchikova et al. 2015; Kupka et al. 2000; Piskunov et al. 1995) to identify all visible Fe I lines on the models. Finally, we performed χ^2 analyses between the data and the models on each line individually. Note that in total, 78 Fe I lines were detected and used in our analysis. The result of this chemical spectroscopy is a metallicity of $[\text{Fe}/\text{H}] = -0.01 \pm 0.28$, where the conservative uncertainty is directly from the standard deviation of all of the measurements. Hence we conclude that TOI–1278 has a solar metallicity, in agreement with the ROTFIT results.

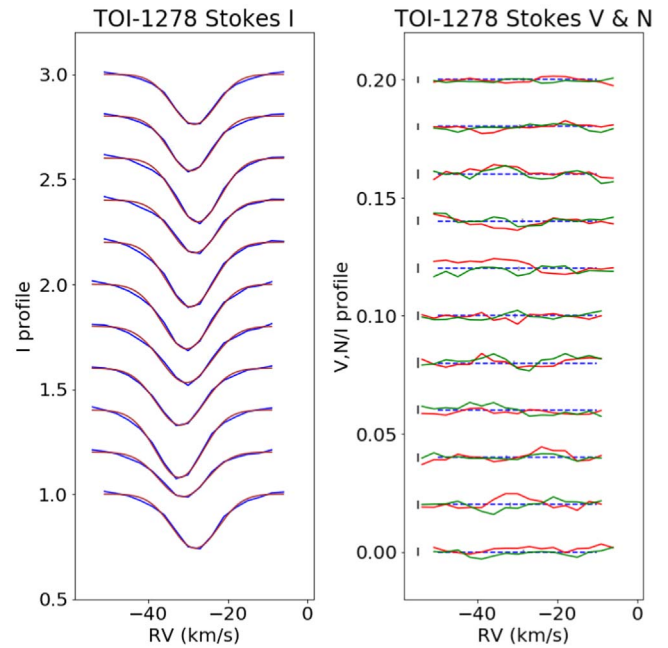


Figure 7. Stokes I (left) and Stokes V (right) of all polarization sequences obtained on TOI–1278. The error bars close to Stokes V profiles show the mean uncertainty on each data point. The green curves on the right plot show the null check profiles. Superimposed to the intensity profiles is a Gaussian fit.

3.4. Spectropolarimetry with SPIRou

Given that each visit to TOI–1278 is composed of a polarimetric sequence, it is possible to combine the individual spectra of each beam to have access to the circular polarization of the stellar surface (Donati et al. 1997, 2020). Using least-square deconvolution and a line list, a mean profile is obtained for the intensity spectrum (Stokes I profile) as well as for the polarized spectrum (Stokes V profile). The line list used for TOI–1278 is based on the VALD database for an effective temperature of 3750 K and minimum line depth of 0.15 of the continuum; it contains 500 atomic lines. The description of the data analysis method as included in the APERO pipeline is given in Martioli et al. (2020). Figure 7 shows the resulting Stokes I and Stokes V profiles for the whole data set. The noise level in the Stokes V profile corresponds to a few 10^{-4} of the intensity spectrum, with no significant detection in the available time series. Using the relationship between the mean Landé factor, the mean wavelength, and the Stokes V and I profiles as introduced in Donati et al. (1997), one can find a 1σ upper limit of 22 to 35 Gauss for the longitudinal field on these visits.

4. Data Analysis and Results

To constrain the orbital and physical parameters of TOI–1278 B, we performed a joint analysis of the SPIRou RV data along with the TESS and ground-based transits. We used the Python package *exoplanet* (Foreman-Mackey et al. 2020), which computes analytical transit models using *starry* (Luger et al. 2019) and solves the Kepler equation with a built-in solver. The *exoplanet* framework uses *PyMC3* (Salvatier et al. 2016) to generate Bayesian models and perform gradient-based Markov Chain Monte Carlo. Models in *PyMC3* are built with the *Theano* (Al-Rfou et al. 2016) numerical infrastructure.

With this framework, we modeled the four transits (two from TESS, one from WBRO, and one from CAO) and the full SPIRou RV time series simultaneously. We included a separate baseline flux f_0 and quadratic limb-darkening coefficients $\{u_1, u_2\}$ for each instrument. For TESS, the coefficients are from Tables 2 and 5 of Claret (2018) for Teff and log values respectively, and are held fixed to $\{0.2336, 0.3710\}$. For CAO, we used EXOFAST (Eastman et al. 2013) to interpolate tables from Claret & Bloemen (2011).³³ We assumed $[\text{Fe}/\text{H}] = 0$ (see Section 3.3) and fixed the resulting I_c bandpass coefficients to $\{0.2776, 0.3528\}$. Since the WBRO transit was observed without any filter, we chose to let $\{u_1, u_2\}$ vary following a Gaussian distribution $\mathcal{N}(0.3, 0.1)$.

We used relatively narrow uniform priors on the period P and the time of mid-transit T_0 , but their ranges are still several orders of magnitude larger than the uncertainties from the preliminary analysis by the TESS team (DV: $P = 14.4762 \pm 0.0022$ days, $T_0 = 2458711.959717 \pm 0.0015$ BJD). We also imposed Gaussian priors on the stellar radius R_s and mass M_s using values from Table 2. We used uniform priors on the orbital eccentricity e and argument of periastron ω , as well as for the systemic velocity γ . For the RV semi-amplitude K , we used a broad log-uniform prior to sample several orders of magnitude uniformly. The transit impact parameter b has a uniform prior between 0 and $1 + R_p/R_s$. Since the transit is grazing, a strong degeneracy exists between b and the planetary radius R_p . To alleviate this degeneracy and assess its importance, we compared three different priors on R_p : a broad uniform prior between 0.0 and $2.5 R_{\text{Jup}}$, a Gaussian prior based on the latest evolutionary models described in Section 5.3, and a Gaussian prior derived from the confirmed exoplanet population. To construct the latter, we follow an approach similar to that of Bayliss et al. (2018). We used the Python-based tool `masterfile`,³⁴ which relies on data from the NASA Exoplanet Archive³⁵ to obtain the population of transiting exoplanets with $M_p > 0.5 M_{\text{Jup}}$ and with an incident flux smaller than 50 times that of the Earth, yielding a sample of 21 planets. We then used the distribution of their radii to derive a Gaussian prior with a mean of $1.010 R_{\text{Jup}}$ and a standard deviation of $0.331 R_{\text{Jup}}$.

We sampled the joint posterior distribution of the 15 model parameters, $\{P, T_0, e, \omega, K, R_p, b, M_s, R_s, \gamma, f_{0,\text{TESS}}, f_{0,\text{WBRO}}, f_{0,\text{CAO}}, u_{1,\text{WBRO}}, \text{and } u_{2,\text{WBRO}}\}$, with the No-U-Turn Sampler (NUTS) from `PyMC3`. We ran four chains with 4000 tuning steps and 4000 draws. The planetary radius posterior distribution for each of the three priors investigated is shown in Figure 8. All priors yield a posterior value for R_p close to $1 R_{\text{Jup}}$, but the width of these distributions varies significantly depending on the strength of the prior. Parameters other than R_p and b were not significantly affected by this change of prior. For this reason, in the following, we only report posterior distributions on parameters that used the ‘‘Demographics’’ prior on R_p . This prior provides an insight on reasonable R_p values without completely excluding the possibility of an inflated radius due to irradiation.

For each parameter, Table 3 reports the final prior distribution used and the median posterior value, along with uncertainties corresponding to the 16th and 84th percentiles. Figure 1 shows the resulting transit model on each data set. In

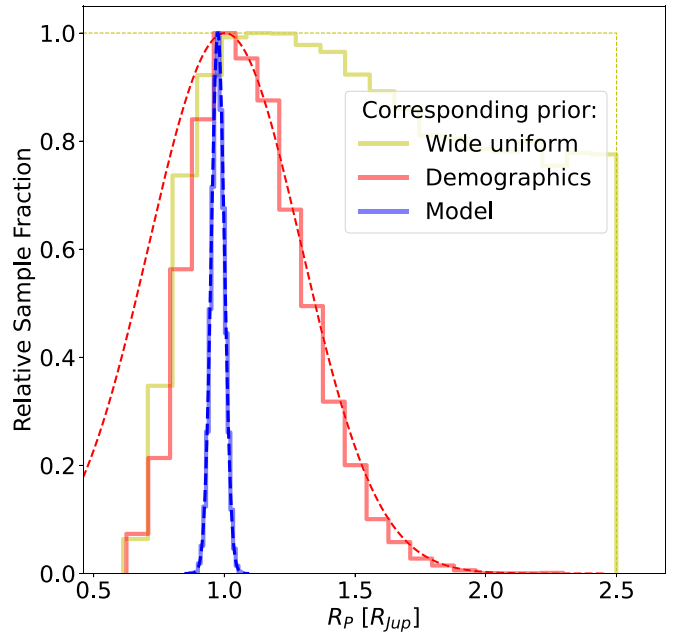


Figure 8. Planetary radius posterior distribution for three different priors. The solid lines show the histogram of the posterior distributions, and the corresponding prior probability density functions are shown as dashed lines of the same color. The prior from evolutionary models (blue), as described in Section 5.3, is $\mathcal{N}(0.975, 0.025)$ and yields $R_p = 0.976 \pm 0.025 R_{\text{Jup}}$. The prior derived from the confirmed exoplanet population (red) is $\mathcal{N}(1.010, 0.331)$ and yields $R_p = 1.09^{+0.24}_{-0.20} R_{\text{Jup}}$. The broad uniform prior $\mathcal{U}(0.0, 2.5)$ (yellow) yields $R_p = 1.56^{+0.63}_{-0.53} R_{\text{Jup}}$.

Figure 3, the full RV time series and the phase-folded RVs are both shown along with the best-fit model and the residuals. The joint posterior distribution is shown in Figure 9.

Along with a planetary mass $M_p = 18.5 \pm 0.5 M_{\text{Jup}}$, a radius of $1.09^{+0.24}_{-0.20} R_{\text{Jup}}$ yields a density $\rho_p = 18^{+14}_{-8} \text{ g cm}^{-3}$. We find a slightly eccentric orbit ($e = 0.013 \pm 0.004$). We also test a circular orbit model with $e = 0$ and $\omega = 0$ following the same modeling procedure as above. All parameters in this circular model are consistent with the values obtained with the eccentric model. However, the eccentric model is strongly favored by the Bayesian information criterion (BIC), with $\Delta\text{BIC} = 9$. Using the eccentric model, we constrain the time of secondary eclipse to be $T_{\text{sec}} = 1719.143 \pm 0.015$ (BJD—2 457 000). This corresponds to an orbital phase of 0.496 ± 0.001 . Orbital evolution and the circularization timescale are discussed in Section 5.2.

We also performed a sequential analysis where the transits were modeled first, and the resulting P and T_0 posterior distributions were used as priors in the RV Keplerian model. This approach has the disadvantage of not accounting for possible impacts of the eccentricity (which is better constrained by the RV data in this case) on the transit model. However, we note that results from both approaches are in excellent agreement, yielding similar posterior distributions with median values consistent within 1σ . This is notable as the sequential analysis does not rely on the stellar parameters as inputs.

Combining the proper motion and parallax from Gaia eDR3 with the systemic RV measurement ($\gamma = -29.334 \pm 0.007 \text{ km s}^{-1}$) derived in this section makes it possible to determine the UVW space velocities of TOI-1278 with an unprecedented precision (the Gaia eDR3 RV for TOI-1278 is $-29.6 \pm 1.1 \text{ km s}^{-1}$). However, doing so requires correcting

³³ See also <http://astrouitools.astronomy.ohio-state.edu/exofast/limbdark.shtml>.

³⁴ <https://github.com/AntoineDarveau/masterfile>

³⁵ <https://exoplanetarchive.ipac.caltech.edu>

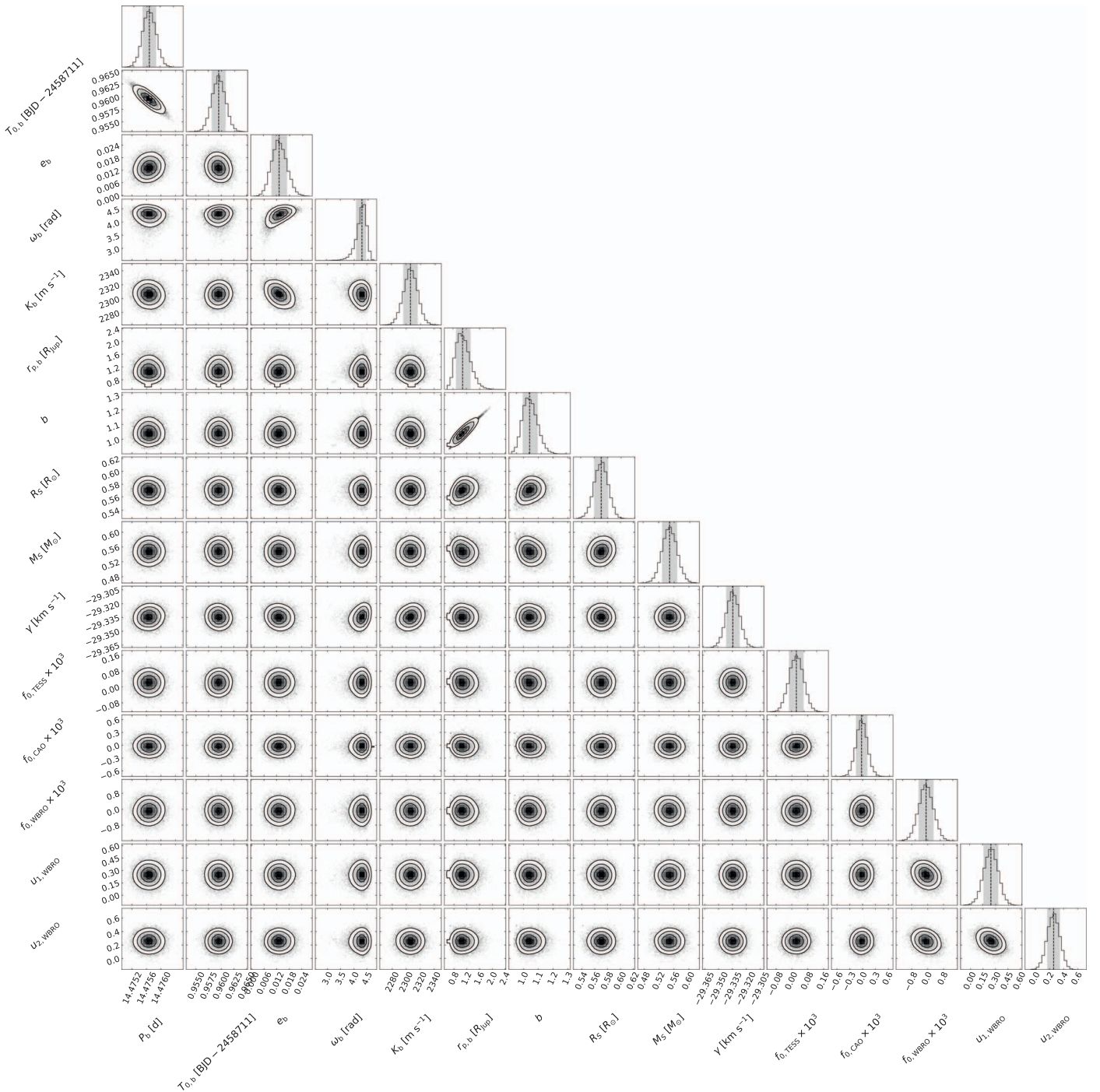


Figure 9. Posterior distribution of the joint RV and transit model parameters. The dashed lines and gray shaded areas show the median and the 1σ interval for each parameter.

two systematic effects that become important with meter-per-level instrumental precision: gravitational reddening (Einstein 1917) and convective blueshift (e.g., Meunier et al. 2017). Our mass and radius estimations allow us to calculate the expected gravitational redshift of $0.63 \pm 0.03 \text{ km s}^{-1}$. Estimating the effect of convective blueshift with a high accuracy is much harder, as empirical data on that subject is sparse, especially for M dwarfs. Löhner-Böttcher et al. (2019) estimated that this effect blueshifts the perceived RV of the Sun by about $0.3 \pm 0.2 \text{ km s}^{-1}$. Only one recent study by Baroch et al. (2020) investigated this effect for YZ CM, a binary M4 dwarf, and

found that it ranges from $+0.007$ to $+0.247 \text{ km s}^{-1}$, corresponding to a very slight redshift rather than a blueshift, which is expected for most stars. (Dai et al. 2018) investigated this effect using a larger sample from Gaia DR2, and found an empirical blueshift of $-0.2 \pm 0.2 \text{ km s}^{-1}$ for early-K dwarfs. We therefore estimate that the convective blueshift for a field-aged M0 star is likely of about $0.0 \pm 0.2 \text{ km s}^{-1}$, as a middle case between those of YZ CMi and early-K dwarfs. The total redshift to which TOI-1278 is subjected is consequently $0.6 \pm 0.2 \text{ km s}^{-1}$. We therefore estimate a heliocentric RV of $-29.9 \pm 0.2 \text{ km s}^{-1}$ for TOI-1278, and we use this value to

Table 1
Basic Data of TOI-1278

Quantity	Value
Designations	
WISEA J212154.86 + 353855.3	
Gaia EDR3 1867491607542113536	
Two Micron All Sky Survey J21215495 + 3538557	
TIC 163539739	
TOI-1278	
Astrometry	
RA	320.47842885720 ^a
DEC	+ 35.64861839399 ^a
Parallax	13.25 ± 0.01 ^a
μ_α [mas/yr]	-89.474 ± 0.010 ^a
μ_δ [mas/yr]	-46.944 ± 0.011 ^a
Galactic kinematics	
X	10.208 ± 0.009 pc ^b
Y	73.58 ± 0.07 pc ^b
Z	-13.26 ± 0.01 pc ^b
U	30.29 ± 0.04 km/s ^b
V	-31.9 ± 0.2 km/s ^b
W	16.28 ± 0.04 km/s ^b

Notes.

^a Gaia EDR3 (Gaia Collaboration et al. 2021).

^b This work, combining GAIA astrometry and SPIRou velocimetry.

calculate the space velocities listed in Table 1. Overall, the space velocities of TOI-1278 are located just within the 2σ velocity dispersion of the Galactic thin disk of Bensby et al. (2003), indicating that TOI-1278 is likely not much older than ≈ 8 Gyr.

5. Discussion

5.1. On the Paucity of Close-in BD Companions to M Dwarfs

The number of known BD companions in close-in orbits around main-sequence stars is relatively small (see full compilation in Table 4 of Mireles et al. 2020), although they are easier to find in general than planetary companions. This is expected from a formation point of view, with close-in binaries having mass ratios tending toward unity (Bate 2000). While most such companions orbit sunlike stars, TOI-1278 combines relatively rare properties; there are few close-in (< 0.1 au) companions to M dwarfs that are more massive than Saturn (see Figure 10 and Table 4 for a compilation of 0.5–78 M_{Jup} companions to M dwarfs).

An order-of-magnitude upper limit on the abundance of TOI-1278-like systems can be determined from the fact that RV surveys of the solar neighborhood have failed to uncover similar systems, even though it would have been well above their detection thresholds. The High Accuracy Radial velocity Planet Searcher survey of southern M dwarfs (Bonfils et al. 2013) surveyed 102 stars; the most massive planet in the sample (G1876 b; Delfosse et al. 1998) has a mass of 2.64 M_{Jup} , ~ 7 times lighter than TOI-1278 B. The CARMENES search for exoplanets around M dwarfs (Reiners et al. 2018) surveyed 324 nearby stars (with some overlap between samples) and did not report any such massive companions. Overall, one can safely say that the occurrence of such systems is well below 1%.

NGTS-1 b is the most massive transiting planetary companion to an M dwarf (Bayliss et al. 2018). It is ~ 20 times less

massive ($0.8 M_{\text{Jup}}$) and orbits on a shorter 2.647 day period. There are a number of comparable ($10\text{--}15 M_{\text{Jup}}$) companions to M dwarfs on distant orbits uncovered through direct imaging (e.g., GU Psc b, 2M J2126-81 b, CHXR 73 b, DH Tau b, 2M 0103-55(AB)b; Naud et al. 2014; Luhman et al. 2006; Deacon et al. 2016; Itoh et al. 2005; Delorme et al. 2013). These companions form well outside the extent of protoplanetary disks around M dwarfs, and the statistical properties in terms of orbital separation and mass suggest a different formation mechanism than for close-in planetary companions (Baron et al. 2019). The mass ratio of the TOI-1278 system (~ 31) cannot be explained through a process involving a planetary disk; the mass involved is larger than the total gas mass in planetary disks, which has a plateau around 10 M_{Jup} (Bergin & Williams 2018). As M-dwarf luminosity falls rapidly with decreasing mass ($L_\star \propto M_\star^3$), uncovering companions with small q values is challenging due to the contrast ratio involved, and there are significant biases against the discovery of low- q binaries. Detecting TOI-1278-like systems requires RV follow-ups, as gigayear-old 1 M_{Jup} to 0.1 M_\odot companions all have radii in the 0.8–1.0 R_{Jup} ; transit alone does not set a significant constraint on the nature of the companion. Conversely, the RV signal of such a companion will be huge, from 100 m s⁻¹ to many kilometers per second, implying that the companions to faint host stars can be monitored. The paucity of 10–30 M_{Jup} may also be arising from a bias in the follow-up strategy by community members; with an impact parameter slightly lower, TOI-1278 B’s transit would have a relative depth of $\sim 4\%$ and be flagged as a likely eclipsing binary. A dedicated search for similar objects within the TESS data set complemented with a follow-up of a statistically representative set of objects will be needed to assess the abundance of objects similar to TOI-1278 B. Current follow-ups, focusing on smaller-radius planets around relatively bright and nearby M dwarfs and avoiding possible eclipsing binaries, are naturally biased against the discovery of objects similar to TOI-1278 B. The true abundance of such systems therefore remains to be established.

Among transiting companions to M dwarfs, only a handful are in the BD mass regime. The LHS 6343 system is composed of a visual binary of two mid-Ms, ($0''.55$ apart; 0.358 and 0.292 M_\odot), one component orbited by a close-in BD (Johnson et al. 2011). LHS 6343 C orbits on a 12.71 day orbit around LHS 6343 A, and has a kinematically derived mass of $62.7 \pm 2.4 M_{\text{Jup}}$. It has a mass ratio of only ~ 6 with its host star. The LP 261-75 b system (Irwin et al. 2018) consists of a much tighter orbit (~ 1.8 day orbit) by a BD close to the stellar/substellar limit at 65 M_{Jup} , and a mass ratio of ~ 4.6 .

5.2. Eccentricity and Circularization Timescale

As discussed in Section 4, we find that a slightly eccentric orbit ($e = 0.013 \pm 0.004$) is preferred over a circular one. We compute the expected timescale for tidal orbital circularization of TOI-1278 B by using Equation (25) of Goldreich & Soter (1966) as rewritten by Patra et al. (2017):

$$\tau_e = \frac{e}{de/dt} = \frac{2Q_p M_p}{63\pi M_\star} \left(\frac{a}{R_p} \right)^5 P_{\text{orb}} \quad (1)$$

Q_p is the tidal dissipation parameter of the companion, depending on its internal structure. We assume here that $Q_p = 10^{3.5}$, as determined by Heller et al. (2010) for the BD. With a value of 16×10^9 yr, the timescale τ_e is comparable to

Table 2
Bulk Properties of TOI–1278

Quantity	Value
Stellar parameters	
Mass (M_{\odot})	0.55 ± 0.02^a
Radius (R_{\odot})	0.573 ± 0.012^b
$\log g$	4.68 ± 0.10^b
Luminosity (L_{\odot})	0.0614 ± 0.0001^b
T_{eff} (K)	3799 ± 42^b
$v \sin i$ (km s $^{-1}$)	1.10 ± 0.86^b
A_V (mag)	-0.01 ± 0.10^b
[Fe/H]	-0.01 ± 0.28^b
Spectral type	M0V b
Photometry	
B	15.07 ± 0.05^b
V	13.51 ± 0.03^b
R_c	12.60 ± 0.04^b
I_c	11.70 ± 0.04^b
G	12.740 ± 0.003^c
g	14.078 ± 0.007^d
r	12.898 ± 0.001^d
i	12.30 ± 0.04^d
z	12.5 ± 0.4^d
y	12.07 ± 0.05^d
J_{2M}	10.601 ± 0.025^e
H_{2M}	9.939 ± 0.021^e
$K_{s,2M}$	9.735 ± 0.012^e
$W1$	9.629 ± 0.022^f
$W2$	9.591 ± 0.020^f
$W3$	9.443 ± 0.03^f

Notes.^a Stassun et al. (2018).^b This work.^c Gaia EDR3 (Gaia Collaboration et al. 2021).^d Pan-STARRS1 (Flewelling et al. 2020).^e Two Micron All Sky Survey (Skrutskie et al. 2006).^f AllWISE (Kirkpatrick et al. 2014).

the age of the Universe, suggesting that the orbit has not had time to fully circularize. This further implies that an eccentric orbit does not require the presence of an external exciting companion to be explained. We note that τ_e scales as the fifth power of R_p , and that this radius is currently only constrained by models for M dwarfs and BDs.

5.3. Model-derived Properties and Characterization Prospects

BDs are ever cooling objects that gradually lose the heat from their initial contraction. In the absence of a luminosity measurement, only the combination of age and mass provides some constraint on the BD temperature. While the mass of TOI–1278 B is well constrained at $18.5 \pm 0.5 M_{\text{Jup}}$, its age is not. Its $\gtrsim 100$ Myr lower limit is set by the uninflated radius of the parent star, while the thin-disk kinematics sets a limit at $\lesssim 9$ Gyr (del Peloso et al. 2005). Considering a formation rate that is uniform through time in the galactic disk, it is unlikely to be at the lower end of this range when the temperature and radius rapidly decrease. Assuming a uniform prior distribution in age for the galactic disk, the corresponding 1σ confidence level interval for the age of TOI–1278 is therefore 1.4–7.6 Gyr (14th and 86th percentiles of the distribution).

Table 3
Model Parameters for the TOI–1278 System

Quantity	Prior	Posterior
P [day]	$\mathcal{U}(13.476, 15.476)$	14.47567 ± 0.00021
T_0 [BJD–2457000]	$\mathcal{U}(1711, 1713)$	1711.9595 ± 0.0013
R_p [R_{Jup}]	$\mathcal{N}(1.010, 0.331)$	$1.09^{+0.24}_{-0.20}$
b	$\mathcal{U}(0, 1 + R_p/R_S)$	$1.04^{+0.06}_{-0.05}$
M_s [M_{\odot}]	$\mathcal{N}(0.55, 0.02)$	0.54 ± 0.02
R_s [R_{\odot}]	$\mathcal{N}(0.573, 0.012)$	0.57 ± 0.01
e	$\mathcal{U}(0, 1)$	0.013 ± 0.004
ω [rad]	$\mathcal{U}(0, 2\pi)$	$4.29^{+0.15}_{-0.21}$
K [m s $^{-1}$]	$\text{LogU}(10^{-4}, 10^4)$	2306 ± 10
ϕ_{sec}	—	0.496 ± 0.001
M_p [M_{Jup}]	—	18.5 ± 0.5
i [deg]	—	88.3 ± 0.1
a [au]	—	0.095 ± 0.001
ρ_p [g cm $^{-3}$]	—	18^{+14}_{-8}
γ [k m s $^{-1}$]	$\mathcal{U}(-35, -25)$	-29.334 ± 0.007
$f_{0,\text{TESS}} \times 10^3$	$\mathcal{N}(0, 10000)$	0.02 ± 0.04
$f_{0,\text{WBRO}} \times 10^3$	$\mathcal{N}(0, 10000)$	-0.08 ± 0.46
$f_{0,\text{CAO}} \times 10^3$	$\mathcal{N}(0, 10000)$	-0.02 ± 0.12
$u_{1,\text{WBRO}}$	$\mathcal{N}(0.3, 0.1)$	0.25 ± 0.09
$u_{2,\text{WBRO}}$	$\mathcal{N}(0.3, 0.1)$	0.26 ± 0.09

We used the latest evolutionary models for cool T and Y dwarfs (Phillips et al. 2020) to determine the range of properties expected for TOI–1278 B. Figure 11 shows the radius, temperature and luminosity evolution with age for the range of masses plausible for TOI–1278. Constraints on the age of the system translate into constraints on these three parameters. These constraints on the age of the system point toward a cool-down temperature ranging from 600 K to 380 K (1σ confidence), but as high as 1360 K for the lower end of the age range; these values correspond to a bulk luminosity ranging from $10^{-4.4}$ to $10^{-6.8} L_{\odot}$, to be compared to the host-star luminosity of $5.9 \times 10^{-2} L_{\odot}$. Expressed in terms of BD spectral sequence, TOI–1278 B is therefore expected to have a spectral type of L7–L8 if we assume an age of 100 Myr, T8–T9 for an age of 1.5 Gyr or Y1 for an age of 7.6 Gyr (Faherty et al. 2016; Kirkpatrick et al. 2019).

Assuming the 1σ confidence interval for the age, the radius of TOI–1278 B is remarkably well constrained at $0.975 \pm 0.025 R_{\text{Jup}}$; degeneracy pressure dominates the radius evolution of objects in this age and mass range, implying that the object’s radius is (nearly) constant through time. This provides an important prior for transit light-curve fitting; the $b \sim 1$ impact parameter implies that the light curve alone sets little constraint on the planetary radius, with a strong degeneracy between b and R_p as seen from the diagonal in the corresponding panel of Figure 9. A much larger planet (up to $R_p \sim 2 R_{\text{Jup}}$) with an impact parameter of $b \sim 1.2$ would be allowed from transit light-curve analysis alone, but is rejected by evolutionary models.

One important output of evolutionary models is the per-band flux, which can be combined with the host-star photometry (see Table 5) to derive a flux ratio. While the system is far too close to be directly resolved with instrumentation in the foreseeable future (~ 1 mas at elongation), it is more likely to be detected through cross-correlation techniques or eclipse photometry or spectroscopy. In the most optimistic case, the contrast ratio

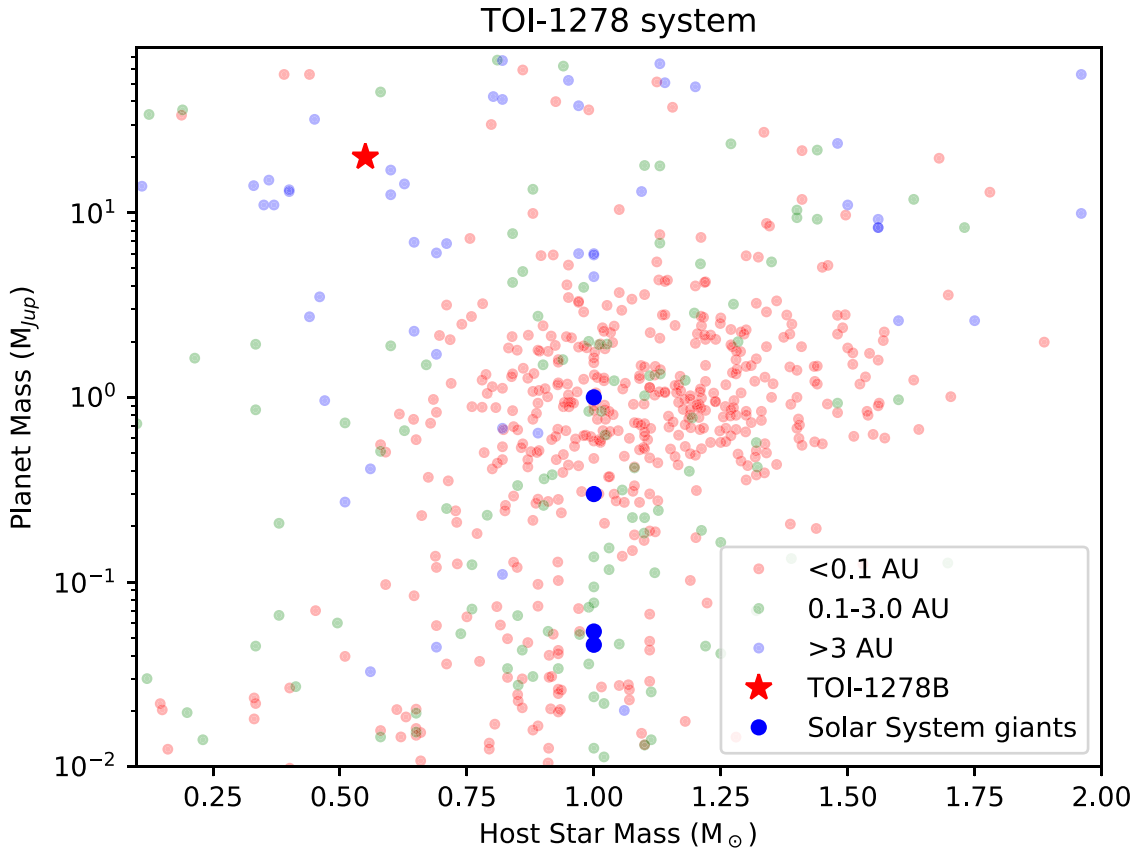


Figure 10. Substellar companion mass as a function of stellar host mass. Close-in companions (red dots) around $<0.5 M_{\odot}$ stars tend to be lighter than $0.1 M_{\text{Jup}}$. There are also a handful of massive ($>40 M_{\text{Jup}}$) BDs in close-in orbits to M dwarfs. (Data from <http://exoplanet.eu>).

Table 4
Known $0.5\text{--}78 M_{\text{Jup}}$ Companions to M Dwarfs with $P < 100$ Days

Substellar Companion	Mass (M_{Jup})	SpT	Period (days)	References	Comment
TOI-1899b	0.66 ± 0.07	M0	29.02	Cañas et al. (2020)	Lowest-mass star known to host a transiting warm Jupiter
NGTS-1b	0.81 ± 0.07	$M0.5 \pm 0.5$	2.65	Bayliss et al. (2018)	Grazing transit, kinematically old host.
GJ 876c	0.760 ± 0.002	M4	30.13	Marcy et al. (2001); Trifonov et al. (2018)	Non-transiting object, absolute mass estimated from dynamics interaction by Correia et al. (2010)
GJ 876b	2.394 ± 0.003	M4	61.08	Delfosse et al. (1998); Trifonov et al. (2018)	Non-transiting object
TOI-519 b	<14	$M3.5^{+1.0}_{-0.5}$	1.26	Parviainen et al. (2021)	Upper limit on mass from light-curve analysis. $1.07 R_{\text{Jup}}$ consistent with an object as light as $0.5 M_{\text{Jup}}$.
TOI-1278 B	18.4 ± 0.5	M0	14.48	This work	
HD 41004Bb	18.4 ± 0.22	M2	1.32	Santos et al. (2002)	Orbiting the M component of a K + M stellar binary. Only $M \sin i$ is known
NLTT 41135B	33.7 ± 2.8	$M5.1 \pm 0.5$	2.89	Irwin et al. (2010)	Transiting one of the components of a $2''3$ binary.
AD 3116B	54.2 ± 4.3	M3.9	1.98	Gillen et al. (2017)	High-probability Praesepe member.
RIK 72 B	59.2 ± 6.8	M2.5	≈ 97.8	David et al. (2019)	Member of Upper Scorpius (5-7 Myr); highly inflated.
TOI-263b	61.6 ± 4.0	M3.5	0.56	Palle et al. (2021)	Orbit synchronized with the rotation of the star
LHS6343C	62.1 ± 1.2	M ^b	12.71	Johnson et al. (2011); Montet et al. (2015)	Transiting LHS 6343A, itself member of a $0''.55$ binary.
LP 261-75B	67.6 ± 2.1	M4.5	1.88	Irwin et al. (2018)	10% deep transit at low impact parameter.
NGTS-7Ab	$75^{+3.0}_{-13.7}$	M3	0.68	Jackman et al. (2019)	Young (55 Myr), active, tidally locked in a decaying orbit (remaining lifetime of 5–10 Myr). Eclipsing one of the components of a $1''.13$ binary.

Notes.

^a The authors explore two scenarios regarding the distance of the NGTS-7 system, leading to significantly different companion masses. The GAIA EDR3 parallax (6.33 ± 0.07 mas) for the host star is consistent with scenario *i* and rules out the lighter mass of scenario *ii*.

^b Spectral subtype not specified, but 3431 ± 21 K temperature (Montet et al. 2015) for host star corresponds to $M2.5 \pm 0.5$ (Boyajian et al. 2012).

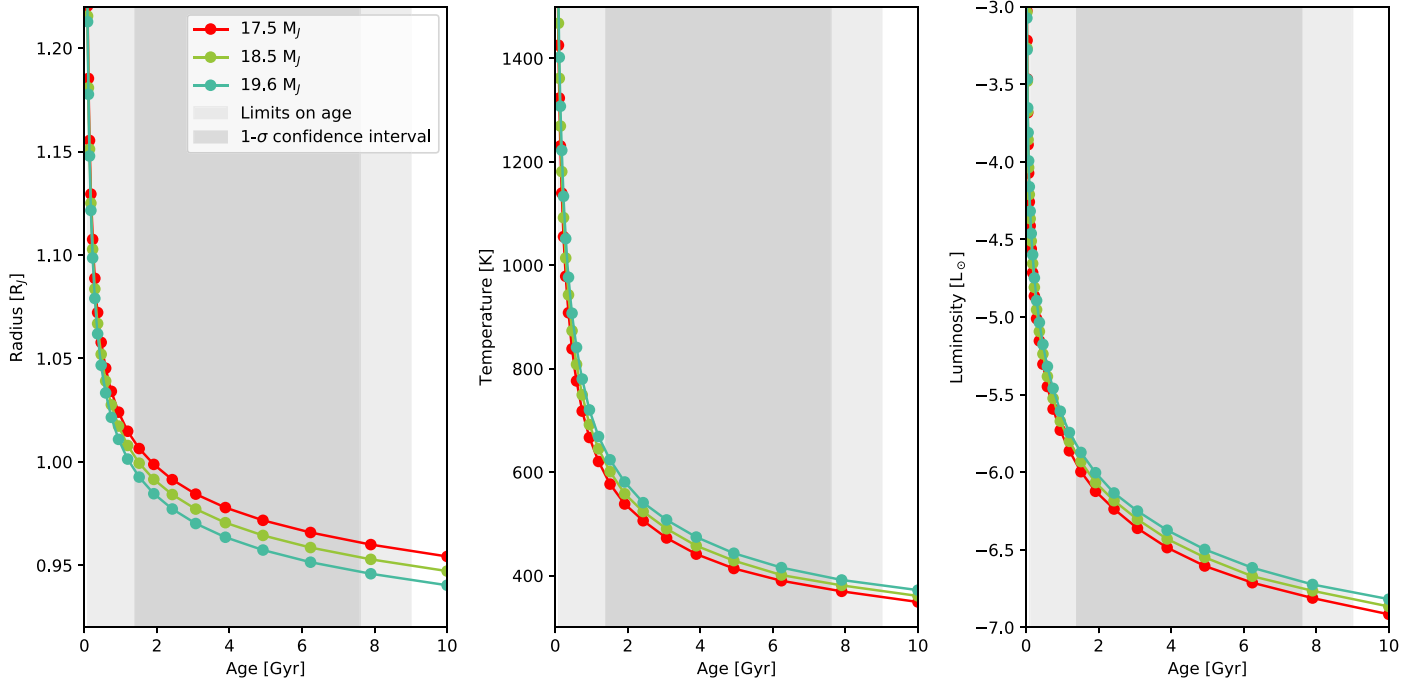


Figure 11. Radius, temperature, and luminosity of TOI–1278 B as a function of age and mass. The derived values are only mildly dependent on the mass, as it is well constrained from the PRV measurements. Unless the system is at the lowest allowed age value, somewhat unlikely but not impossible, the radius, temperature, and luminosity are constrained to within $\pm 2.5\%$, ± 110 K, and ± 0.4 dex, respectively.

Table 5

Model Properties of TOI–1278 B at Different Ages for a Mass of $18.5 M_{\text{Jup}}$

	0.1 Gyr	1.5 Gyr	7.9 Gyr
Temperature	1362 K	601 K	381 K
Luminosity	$-4.37 L_{\odot}$	$-5.93 L_{\odot}$	$-6.76 L_{\odot}$
Gravity	4.55	4.69	4.74
Radius	$0.118 R_{\odot}$	$0.100 R_{\odot}$	$0.095 R_{\odot}$
ΔJ_{MKO}	6.878	10.992	15.285
ΔH_{MKO}	7.330	12.551	16.747
ΔK_{MKO}	7.640	13.113	18.720
$\Delta W1$	7.520	11.768	14.791
$\Delta W2$	6.028	8.275	9.953
$\Delta W3$	5.134	7.620	9.157

Note. These values are interpolated from Phillips et al.’s (2020) BD models.

would be at the part-per-thousand level in the K band and, most likely, significantly more challenging (See Table 5). Overall, the prospects for eclipse spectroscopy of TOI–1278 B are poor.

One can estimate the scale height of the atmosphere of TOI–1278 B and the corresponding transit-spectroscopy signal. As its atmosphere is most likely hydrogen-dominated, one can safely assume a mean-molecular weight of 2 AMU. Assuming a temperature of 380–1300 K, which would correspond to the plausible age range, the scale height of the atmosphere would be 5–17 km, making it an impractical target for transit spectroscopy.

The expected Rossiter-McLaughlin (RM) effect (Rossiter 1924; McLaughlin 1924) for this system can be estimated. Assuming a $v \sin i = 1.1 \pm 0.9 \text{ km s}^{-1}$ and a transit depth of $\sim 1\%$, the maximum possible signal for the effect would be at the $7 \pm 6 \text{ m s}^{-1}$ level. This is below the sensitivity of our SPIRou observations and would be challenging on such a faint

target with current facilities. Considering that the transit is grazing, the signal is likely to be much smaller if the spin-orbit angle is close to 0° . A handful of hot Jupiters in polar orbits are known (e.g., Addison et al. 2013), but these orbit stars hotter than the Sun. Stars such as TOI–1278 have thicker convective layers that dampen orbital obliquity, suggesting that it is unlikely that this system displays a strong inclination and that its RM effect can be measured considering its transit configuration (Winn et al. 2010).

We note that TOI–1278 B is expected to be among the astrometric detections of substellar companions in Gaia. The semimajor axis of 0.1 au and 75 pc distance implies that the orbit subtends a 1.3 mas angle and that the astrometric orbital motion of the host star is $50 \mu\text{as}$. At $G \sim 12.7$, the per-epoch astrometric accuracy is at the $34.2 \mu\text{as}$ level (Perryman et al. 2014). The presence of the companion most likely explains the `astrometric_excess_noise_sig` = 9.09 value in the EDR3 catalog. This dimensionless parameter measures the significance of the excess in noise relative to the astrometric uncertainties when fitting a simple parallax models. For good fits of single stars, one would expect half of the targets to have a value of zero as the residuals to the fit will be (slightly) smaller than expected from uncertainties. The remaining targets will have values displaying a cumulative distribution corresponding to a Gaussian with a width of unity. Sources with a value above 2 are considered as likely having a significant noise excess in their parallax solutions. While Gaia does not have the proper time-sampling to probe short periods such as that of TOI–1278, the known orbital period from RV measurements could guide a determination of the actual astrometric orbit.

We dedicate this paper to the memory of France Allard, who until her untimely passing did so much to advance the modeling of BD atmospheres.

Based on observations obtained at the Canada–France–Hawaii Telescope, which is operated from the summit of Maunakea by the National Research Council of Canada, the Institut National des Sciences de l’Univers of the Centre National de la Recherche Scientifique of France, and the University of Hawaii. The observations at the Canada–France–Hawaii Telescope were performed with care and respect from the summit of Maunakea, which is a significant cultural and historic site.

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Facility: CFHT(SPIRou).





















Software: astropy (Price-Whelan et al. 2018).

Appendix Table of SPIRou RV Measurements

Table A1
SPIRou RV Measurements Binned per Epoch

Time BJD	RV [m s ⁻¹]	σ_{RV} [m s ⁻¹]
2459000.1000	-28111.4	17.2
2459003.0779	-30822.7	28.5
2459005.0466	-31675.7	15.6
2459006.0240	-31491.6	14.9
2459008.0282	-29992.7	13.3
2459008.9816	-29051.9	18.3
2459010.0464	-28069.4	21.4
2459011.1114	-27299.0	21.5
2459153.7576	-29025.5	17.7
2459154.7451	-28075.8	15.9

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