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Shapley Supercluster Survey: Galaxy evolution from filaments to cluster cores

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ABSTRACT

We present an overview of a multiwavelength survey of the Shapley Supercluster (SSC; $z \sim 0.05$) covering a contiguous area of $260 h_{70}^{-2} \text{ Mpc}^2$ including the supercluster core. The project main aim is to quantify the influence of cluster-scale mass assembly on galaxy evolution in one of the most massive structures in the local Universe. The Shapley Supercluster Survey (ShaSS) includes nine Abell clusters (A3552, A3554, A3556, A3558, A3559, A3560, A3562, AS0724, AS0726) and two poor clusters (SC1327–312, SC1329–313) showing evidence of cluster–cluster interactions. Optical (*ugri*) and near-infrared (*K*) imaging acquired with VLT Survey Telescope and Visible and Infrared Survey Telescope for Astronomy allow us to study the galaxy population down to $m^* + 6$ at the supercluster redshift. A dedicated spectroscopic survey with AAOmega on the Anglo-Australian Telescope provides a magnitude-limited sample of supercluster members with 80 per cent completeness at $\sim m^* + 3$. We derive the galaxy density across the whole area, demonstrating that all structures within this area are embedded in a single network of clusters, groups and filaments. The stellar mass density in the core of the SSC is always higher than $9 \times 10^9 \text{ M}_\odot \text{ Mpc}^{-3}$, which is $\sim 40\times$ the cosmic stellar mass density for galaxies in the local Universe. We find a new filamentary structure ($\sim 7 \text{ Mpc}$ long in projection) connecting the SSC core to the cluster A3559, as well as previously unidentified density peaks. We perform a weak-lensing analysis of the central 1 deg^2 field of the survey obtaining for the central cluster A3558 a mass of $M_{500} = 7.63_{-3.40}^{+3.88} \times 10^{14} \text{ M}_\odot$, in agreement with X-ray based estimates.

Key words: gravitational lensing: weak – galaxies: clusters: general – galaxies: clusters: individual: A3552, A3554, A3556, A3558, A3559, A3560, A3562, AS0724, AS0726, SC1327-312, SC1329-313 – galaxies: evolution – galaxies: photometry – galaxies: stellar contents.

1 INTRODUCTION

It is well established that the properties of galaxies are correlated with the environment (e.g. Lewis et al. 2002; Goto et al. 2003;

Balogh et al. 2004; Blanton et al. 2005; Baldry et al. 2006; Pimblet et al. 2006; Haines et al. 2007; Bösch et al. 2013). At $z \simeq 0$, galaxy populations in rich clusters are dominated by ellipticals, S0s and at lower masses dwarf ellipticals (dEs) with few, if any, star-forming spirals (Dressler 1980; Dressler, Thompson & Shectman 1985). Cluster galaxies have not always been as inactive as they are at the present epoch. Butcher & Oemler (1984) showed that the fraction of

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blue (star-forming) galaxies among cluster members increases from almost zero in the local Universe to ~ 20 per cent by $z \sim 0.4$, while recent *Spitzer/Herschel* surveys have confirmed large numbers of starburst galaxies in clusters to $z \sim 1$ and beyond (e.g. Popesso et al. 2012). While some of this rapid evolution of cluster galaxies can be explained by the cosmic $\sim 10\times$ decline in star formation (SF) among field galaxies since $z \sim 1$ (e.g. Le Flocc’h et al. 2005), cluster populations have shown an accelerated evolution in SF over the last four billion years, resulting in an overall $\sim 15\times$ reduction in the total star formation rates (SFRs) per unit halo mass since $z \sim 0.3$ (Haines et al. 2013). Empirically, clusters accrete gas-rich, star-forming spirals at $z \gtrsim 0.5\text{--}1.0$ and then contribute to transforming them somehow into the passive S0s and dEs of local clusters.

Several mechanisms affecting the galaxy properties and dependent upon the environment have been proposed and investigated in detail and all of them serve to kinematically disturb spiral galaxies and/or transform their structural properties and/or deplete their reservoirs of gas, and so quench SF. These physical processes include gravitational and tidal interactions amongst galaxies (Toomre & Toomre 1972; Moore et al. 1996), between galaxies and the cluster gravitational field (Byrd & Valtonen 1990), galaxy mergers (Barnes & Hernquist 1991), group-cluster collisions (Bekki 2001), ram-pressure (Gunn & Gott 1972) and viscous stripping (Nulsen 1982), evaporation (Cowie & Songalia 1977) and ‘starvation’ (Larson, Tinsley & Caldwell 1980). Since these mechanisms are characterized by different time-scales and efficiencies which depend, in turn, on the properties of both the galaxies themselves (e.g. their stellar masses, morphologies) and their environment, they can affect the galaxy properties in different characteristic ways (Boselli & Gavazzi 2006; Haines et al. 2007). Examples of this are the different effects of tidal and hydrodynamical interactions. The ram pressure exerted by the hot and dense intracluster medium (ICM) can effectively remove the cold gas supply, truncating the gas disc and quenching SF, in massive cluster galaxies passing through the centre of rich clusters in about one crossing time ($\sim 10^9$ yr), while the interstellar gas of a massive galaxy may never be completely stripped if the galaxy moves on a tangential orbit or is member of a poor cluster. Repeated high-velocity encounters (harassment) of cluster galaxies can destroy the fragile disc of dwarf galaxies, but not significantly affect the structure of a giant galaxy.

The development of large spectroscopic surveys such as the Sloan Digital Sky Survey (SDSS), plus the availability of panoramic far-ultraviolet–far-infrared (FUV–FIR) data from the *GALEX* and *Spitzer* space telescopes have allowed the impact of environment on SF to be quantified in unprecedented detail. In these studies, however, the environment and its characteristics are usually treated as ‘static parameters’. While the galaxies move across the cluster, field and filament environments (experiencing different ram pressures, encountering other galaxies and possibly being involved in merging), the cluster potential well, the galaxy density as well as the ICM are usually not considered ‘time-dependent’. Of course, at first pass this approach is an unavoidable reduction on complexity. On the other hand in a hierarchical universe with the assembly of the structures, the galaxies evolve and move, tending towards denser regions with time, while the environments change too, thus what we actually observe is *galaxy evolution in an evolving environment*.

The most massive structures in the local Universe are superclusters, which are still collapsing with galaxy clusters and groups frequently interacting and merging, and where a significant number of galaxies are encountering dense environments for the first time. The relative dynamical immaturity of superclusters and the presence of infalling dark matter haloes make them ideal laboratories

to test the predictions of hierarchical mass assembly models, and in particular on galaxy evolution. Superclusters are not so rare systems in the Universe (see Sheth & Diaferio 2011). The observation of superclusters is often considered a challenge to the hierarchical structure formation paradigm since such extreme dense structures, but also voids, are not reproduced by the N -body simulations. Yaryura, Baugh & Angulo (2011) pointed out that the reason of this discrepancy can be due to the method used to assess the probability of finding such events in the distribution of cold matter. They proposed a new technique to analyse an ensemble of N -body simulations in a volume equal to that of the two-degree Field Galaxy Redshift Survey (2dFGRS; Colless et al. 2001) where the probability to find peculiar structures (overdense and underdense) was thus estimated to be ~ 2 per cent and not null.

Superclusters allow in principle both to study dynamical processes such as cluster–cluster collisions and group–cluster mergers and to sample different environments from cluster cores to filaments and fields. Furthermore, within a dynamically active and locally dense structure the probability to observe evidence of environmental effects on galaxy evolution is dramatically enhanced making these systems a sort of *magnifying glass* to identify the different physical mechanisms which transform the properties of galaxies. In order to address all these aspects, a careful selection of the target is fundamental. The ideal structure should map different environments with evidence of cluster–cluster interactions. Finally, in order to study in detail the galaxy properties a resolution $\lesssim 1$ kpc is required in a wide range of galaxy mass down to the dwarf regime where such galaxies are not quenched by internal processes, but are more susceptible to environmental transformations. With all this in mind, we have undertaken a study of the Shapley Supercluster which is the largest conglomeration of Abell clusters in the local Universe.

The Shapley Supercluster Survey (ShaSS) will map a 23 deg^2 region ($\sim 260 \text{ Mpc}^2$) of the Shapley Supercluster at $z = 0.048$, containing filaments and embedded galaxy groups which form a dynamically bound network connecting nine Abell and two poor clusters, in order to identify the primary locations (groups, filaments, clusters) and mechanisms for the transformation of spirals into the S0s and dEs.

Among the observational studies of superclusters, the STAGES project (Gray et al. 2009) addressed in particular the study of galaxy evolution in the Abell 901(a,b) supercluster at redshift $z \sim 0.165$ while the ORELSE survey (Lubin et al. 2009, but see also Mei et al. 2012) searched for structures on scales greater than 10 Mpc at higher redshifts ($0.6 < z < 1.3$) with the aim to investigate the properties of member galaxies. The comparison between the STAGES, ORELSE and other similar studies and ShaSS will be then unavoidable and useful to trace the evolution with redshift, although we notice that the area (in Mpc^2) of ShaSS is a factor 10 that of the STAGES survey and was chosen to map the filaments connecting the Abell clusters.

The optical survey VST-ACCESS in four bands collected at the European Southern Observatory (ESO) VLT Survey Telescope (VST) represents the core of this multiband project. The infrared coverage with Wide-field Infrared Survey Explorer (WISE), the dedicated spectroscopic survey with AAT/AAOmega and the Shapley-VISTA survey, together with other proprietary data, provide the fundamental data set to achieve the scientific goals of ShaSS. In this paper, we will give an overview of the project and present the first results. In Section 2, the motivations and main scientific objectives of the project are discussed. The target is described in Section 3. The characteristics and strategy of the survey

are described in Section 4. Details of the data reduction and analysis of the data quality are given in Sections 5 and 6, respectively. The first results concerning the characterization of the environment are presented in Section 7, where we derive the galaxy density across the whole ShaSS region and the underlying dark matter distribution for the central 1 deg^2 field which allows us to estimate the mass of the galaxy cluster A3558. In Section 8, we present the summary and conclusions.

Throughout the paper, we adopt a cosmology with $\Omega_M = 0.3$, $\Omega_\Lambda = 0.7$ and $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$. According to this cosmology, 1 arcsec corresponds to 0.941 kpc at $z = 0.048$ and the distance modulus is 36.66. The magnitudes are given in the AB photometric system.

2 SURVEY MOTIVATIONS AND OBJECTIVES

2.1 The role of large-scale mass assembly on galaxy evolution

In Λ cold dark matter cosmological models, structure formation occurs hierarchically, such that the most massive haloes corresponding to galaxy clusters form latest, doubling their masses on average since $z \sim 0.5$ (Boylan-Kolchin et al. 2009; Gao et al. 2012), and are also the most dynamically immature. Preferentially located at the nodes of the complex filamentary web, clusters are continually accreting dark matter haloes containing individual galaxies or galaxy groups. On average, 50 per cent of galaxies in local clusters have been accreted since $z \sim 0.4$ (Berrier et al. 2009), of which 40 per cent are within groups (McGee et al. 2009).

The process of central galaxies becoming satellites, as their host haloes are accreted into more massive haloes, has been shown to strongly affect their evolution, with satellite galaxies more likely to be quenched than central galaxies of the same stellar mass (e.g. Wetzel, Tinker & Conroy 2012). SF is suppressed within galaxy groups, with the fraction of star-forming galaxies declining steadily with increasing group mass (at fixed stellar mass) and proximity to the group centre (Weinmann et al. 2006; Wetzel et al. 2012; Woo et al. 2013). Galaxies which are quenched within such groups and later accreted into clusters are described as having been ‘pre-processed’.

Numerical simulations and theoretical studies predict the effect of cluster–cluster mergers on galaxy properties. Bekki, Owers & Couch (2010) claimed that merging of galaxy clusters may induce SF due to the increase of the external pressure of the ICM compressing the cold gas in the cluster galaxies. This results in a population of starbursts or post-starburst galaxies having the same age and so dating the merging event. The spatial distribution of these star-forming galaxies is expected to differ from the overall distribution of the other cluster galaxies. Similarly, for comparable group and cluster halo densities gravitational shocking as the group enters the cluster could temporarily increase its mass and pull the group members into a denser, more compact system inducing galaxy mergers (Moss 2006). Simulations also showed that starbursts induced by galaxy mergers can be amplified by a factor of ~ 2 (Martig & Bournaud 2008) if they take place in the tidal field of cluster/group of galaxies.

Owen et al. (2005) studied the excess population of radio galaxies in the cluster A2125 at $z = 0.25$, mostly located in groups outside the cluster core and with radio luminosities indicating the SF as the main mechanism responsible of the radio emission. The authors related this observational evidence with the ongoing major cluster–cluster merger and explained the increased SF as due to the variation of the tidal field experienced by the member galaxies during a cluster–cluster interaction, probably being close to the core

passage. Studying a system of merging clusters, Johnston-Hollit et al. (2008) observed an excess of star-forming galaxies aligned along the bridge of galaxies connecting A3158 to the A3125/A3128 complex, suggestive of merger-induced SF. In addition, they found that the fraction of radio-loud sources is lower with respect to a global cluster environment and similar to that measured by Venturi et al. (2000) in the cluster A3558 which is considered as an example of a late merger stage. They suggested that in the cluster cores radio emission from active galactic nuclei (AGNs) is suppressed in the late stages of cluster mergers. The suppressed radio emission observed in both bright and faint cluster galaxies has been also associated to cluster merger by Mauduit & Mamon (2007). They explained their findings with the disruption of cool cores, i.e. the gas supply to the central AGN for the bright galaxies, while the low radio loudness and fainter cluster galaxies have been affected by the enhanced ram pressure when crossing the shock front between the merging clusters. Further observational support of the important role that cluster mergers play in triggering the evolution of cluster galaxies is provided by Owers et al. (2012) studying the merging cluster A2744. They identified three rare ‘jellyfish’ galaxies located in close proximity to the ICM features associated with a merging subcluster and its shock front. Their interpretation is that the SF knots detected in the tails of stripped gas are due to the rapid increase in pressure experienced by the galaxies interacting with the shock.

Hence, the observations seem to confirm that galaxy/group accretion and cluster–cluster merger affect the cluster galaxies and that the galaxy properties can be sensitive indicators of the merger stage.

2.2 Measuring the environmental effects beyond r_{200}

The colour–density relation is observed in galaxy clusters well beyond r_{200} (e.g. Haines et al. 2009b) and the fraction of star-forming galaxies steadily increases with cluster-centric radius. Nevertheless even at $3\text{--}5 \times r_{200}$, the fraction of star-forming galaxies is below that seen in the field (e.g. Chung et al. 2011) suggesting that the environment plays a role inducing transformations well outside r_{200} .

By means of hydrodynamical cosmological simulations, Bahé et al. (2013) investigated the trend of cold and hot gas contents and SF as function of clustercentric distance in clusters and groups of galaxies. They found a large-scale trend with cold and hot gas contents and SF increasing with the clustercentric distance, but approaching the values of the field galaxy sample only at $\sim 5r_{200}$ (corresponding to 10 Mpc for a massive cluster). Moreover, the SF in low-mass cluster galaxies ($M = 10^{9\text{--}9.5} M_\odot$) still shows a significant discrepancy even at this large radius with respect to that of the field galaxies. According to the authors, three main causes can explain the observed trends: (i) ‘pre-processing’ of galaxies within infalling groups; (ii) ‘overshooting’ for those galaxies that are not falling in for the first time (also called ‘backsplash’ galaxies); (iii) ram-pressure stripping (RPS). Of course, the three mechanisms act differently depending on the galaxy and host halo mass and the clustercentric distance, but in any case the radial trend of the hot gas cannot be explained solely by overshooting and pre-processing. A direct interaction with the host group/cluster is also required and their simulations suggest that RPS can strip the hot gas from low- and high-mass galaxies out to several times r_{200} , although it cannot affect the cold gas except for the low-mass galaxies. The authors also claimed that tidal interaction does not play an important role in the gas removal, but it should be an important factor in the morphological transformation of these galaxies.

S0s differ from normal spirals by their higher bulge luminosities rather than fainter discs (Christlein & Zabludoff 2004), disfavouring simple RPS or starvation mechanisms (but see Kapferer et al. 2009). Instead, mechanisms such as merging or harassment are capable of channelling material to a central bulge, sufficient to produce the higher central mass densities seen in cluster spirals and ultimately the stellar phase densities found in S0s (Moran et al. 2007). This channeling of material will also likely fuel rapid growth of the central supermassive black hole and trigger a period of nuclear activity. Deep observations are needed to reveal tidal streams or ‘fans’, characteristic of recent merging activity (van Dokkum 2005).

In situ transformation via merging in clusters is strongly suppressed due to the high encounter velocities. Low-velocity encounters and mergers should be much more frequent in galaxy groups, leading to the suggestion that the bulge growth required to form S0s occurs primarily via pre-processing in galaxy groups, which are subsequently accreted into clusters (Balogh et al. 2004). Moreover, the S0 fractions of $z \lesssim 0.5$ groups match those seen in clusters at the same redshifts, and are much higher than found in the field (Wilman et al. 2009). Strong nuclear activity which is linked to bulge growth is also found to be effectively suppressed in cluster galaxies, with the distribution of X-ray AGN in rich clusters revealing them to be an infalling population with high relative velocities (Haines et al. 2012). This confirms that the transformation of galaxies often occurs (or begins at least) well outside the cluster cores, in groups or as infalling galaxies enter the virialized region for the first time (Moran et al. 2005).

Hints about the distribution of galaxy luminosity, colours and morphology in the filaments have been given by Pandey & Bharadwaj (2006), who found that brighter galaxies have a less filamentary distribution than the fainter ones. With respect to the morphology, the early-type galaxies are concentrated in the vicinity of the nodes, while spiral galaxies are sparsely distributed across the filaments. Fadda et al. (2008) observed that the fraction of starburst galaxies in the filaments around the cluster A1763 is twice that in other cluster regions as detected by *Spitzer* and suggested that filaments are a sort of *galaxy reservoir* for clusters.

By studying a sample of supercluster filaments in the 2dFGRS, Porter et al. (2008) observed a sudden enhancement in SF in faint dwarf galaxies outside the cluster virial radius. They interpreted their findings with a close interaction and/or harassment with other infalling galaxies along the same filaments. Another possible explanation is the very first interaction that these galaxies experienced with the ICM inducing tidal shocks and then a burst of SF. Actually, the dominant processes that quench SF in galaxies depend crucially on the galaxy mass. Dwarf galaxies, given their shallow potential wells, should be more susceptible to environmental processes such as tidal/RPS (e.g. Barazza et al. 2009). Indeed their SF histories are completely defined by their local environment, since passive dEs are only found as satellites within massive haloes (i.e. cluster, group or massive galaxy; Haines et al. 2006b, 2007). The formation of many cluster dEs is often rather recent, manifest by young stellar ages (~ 2 Gyr) and the significant populations of starburst and post-starburst dwarf galaxies in the outskirts of local clusters (Smith et al. 2009; Mahajan, Haines & Raychaudhury 2011).

2.3 Relating star formation quenching to ram-pressure stripping

RPS, as originally proposed by Gunn & Gott (1972) requires, in principle, the presence of a dense ICM. Thus, its effect would be limited to cluster cores where the gas discs of massive spirals are

rapidly truncated. Nevertheless, with a 3D hydrodynamical simulation, Marcolini, Brighenti & D’Ercole (2003) showed that RPS may extend to poorer environments for low-mass galaxies which, thanks to their lower escape velocities, are easier to strip. Roediger & Hensler (2005), by means of high resolution 2D hydrodynamical simulations, demonstrated that ram pressure effects can be observed over a wide range of ICM conditions. In particular, in high-density environments RPS severely truncates the gas disc of L^* galaxies, while in low-density environments, where moderate ram pressure is foreseen, their gas disc is clearly disturbed and bent (Roediger & Hensler 2005). The gas discs of these galaxies can be truncated to 15–20 kpc in the first 20–200 Myr of RPS.

Merluzzi et al. (2013) identified a bright ($L > L^*$) barred spiral galaxy 1 Mpc from the centre of the rich cluster A3558 in the Shapley Supercluster core, which is strongly affected by RPS. Integral-field spectroscopic observations revealed ongoing gas stripping in the form of one-sided extraplanar ionized gas along the full extent of the disc, simultaneously with a starburst triggered by gas compression along the leading edge of the galaxy. The galaxy is estimated to be being subjected to weak-moderate ram pressure, as defined by Roediger & Hensler (2005). This adds a piece of evidence to the fact that RPS is acting more efficiently on the galaxy interstellar medium (ISM) than previously foreseen and also outside of the cluster cores as also observed in the Virgo cluster by Chung et al. (2009). This new understanding of the RPS supports the view that this mechanism is the principal transformation process to quench SF in spirals although probably helped by other processes affecting the structure of the galaxies. Bösch et al. (2013) drew a similar conclusion analysing a sample of 182 disc galaxies in the cluster system Abell 901/902. The fraction of galaxies showing asymmetric gas rotation curves, and thus probably affected by RPS (Kronberger et al. 2008), turned out to be 75 per cent higher in the clusters than in the field, the majority of them being morphologically undisturbed. Although these galaxies seemed preferentially located at low clustercentric radii, they also observed a population of dusty and red spirals in the cluster outskirts, which could be also affected by ISM–ICM interactions. The relevance of RPS for quenching SF in cluster galaxies is further complicated by the fact that a non-homogeneous ICM is expected for non-virialized merging and post-merging clusters, thus the clustercentric distance is not the only parameter describing the ICM density, i.e. the ram-pressure strength.

2.4 The approach of the Shapley Supercluster Survey

Our study is based on characterizing a dynamically active environment in a multifaceted way and disentangling its effects on galaxy evolution. It has the following main objectives.

(i) To investigate the role of cluster-scale mass assembly on the evolution of galaxies mapping the effects of the environment in the cluster outskirts and along the filaments with the aim to identify the very first interactions between galaxies and their environment.

(ii) To identify and measure signs of ongoing transformation in galaxies belonging to a complex structure with the goal of improving our comprehension of what drives their SF quenching and structural modification.

(iii) To obtain detailed maps of the dark matter and baryonic matter distributions (galaxies, ICM), combining weak lensing, X-ray and dynamical analyses.

(iv) To quantify the variation in the stellar mass fractions going from cluster cores to groups, by comparing the near-infrared light distribution with the dark matter maps and dynamical masses.

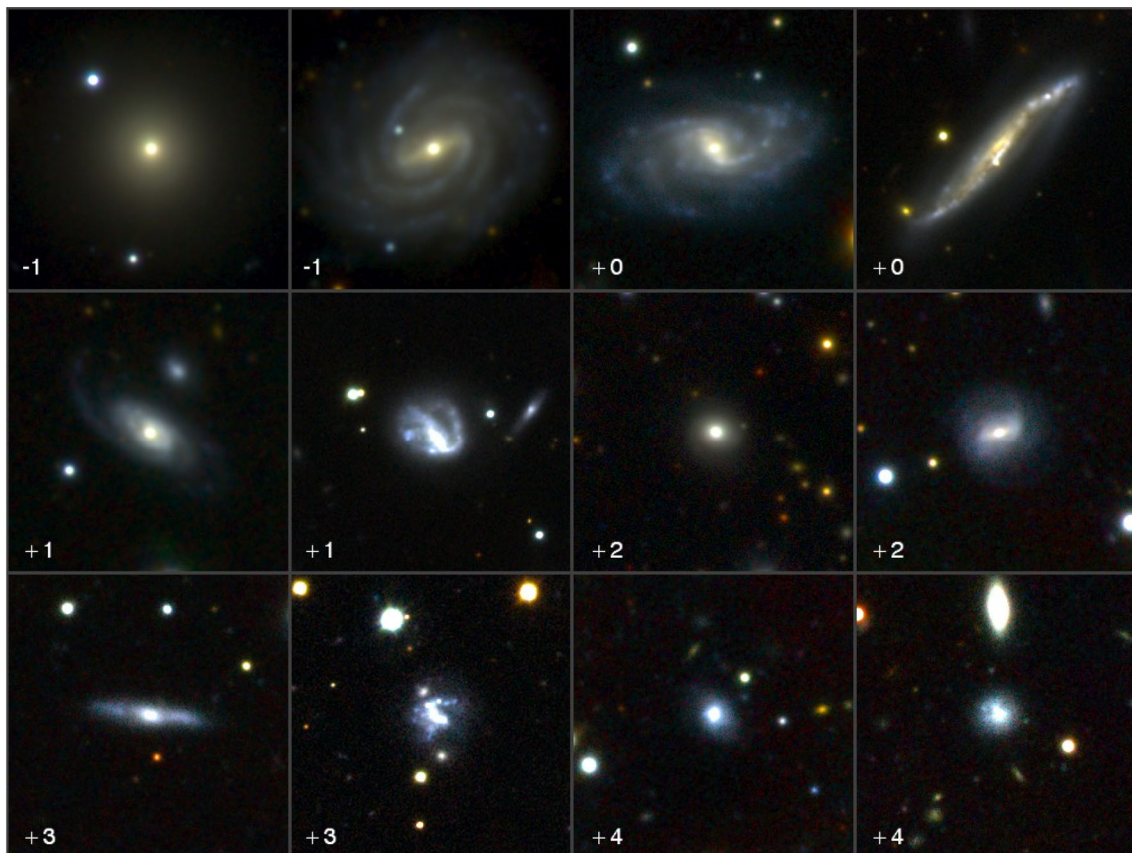


Figure 1. VST-ACCESS *gri* composite images of a sample of galaxies spectroscopically confirmed as members of the Shapley Supercluster. From left to right and top to bottom different magnitude bins from $m^* - 1$ to $m^* + 4$ as indicated. Each square frame is 1 arcmin (~ 56 kpc) wide.

(v) To build up a multiband homogeneous data set on this area of the sky made of subkiloparsec resolution imaging and magnitude-limited spectroscopy, thus providing the community with a solid background for studies of the Shapley Supercluster.

To address the above objectives, we will explore the global and internal properties of galaxy populations extending outside the cluster/group virial radius and aim for an accurate characterization of the environment. This will be defined through galaxy density, dark matter distribution, dynamical substructure and ICM properties. The different quantifications of the environment will allow us to disentangle the effects of local and large-scale density, cluster and group merging, dynamical state and mass of the host system on the properties of galaxies in different ranges of mass.

Recent studies investigated the effects of the local and large-scale environment on galaxy evolution, mainly considering groups of galaxies. Some of these works defined the environment through the luminosity-density only (e.g. Lietzen et al. 2012), or associate galaxy groups to a large-scale environment identified by a smoothed luminosity field (e.g. Luparello et al. 2013), or are limited to massive galaxies (e.g. Catinella et al. 2013). Therefore, although there is general agreement that the group environment is important for the evolution of the group members, a study of global and internal properties of galaxy population in groups and, in general, in different structures dynamically bound in a supercluster is still lacking. For instance, Rasmussen et al. (2012) found that both local (distance from the centre) and global (mass) group environments play a role in quenching SF, while Ziparo et al. (2014) show that global group-

scale mechanisms linked to the presence of a hot gas halo are dominant in quenching SF in group galaxies rather than purely density related processes.

In the following, we briefly outline how we will derive the quantities necessary for our study (see also Section 4).

As the main classification scheme of galaxies we will adopt the concentration, asymmetry and clumpiness parameters (Kent 1985; Conselice, Bershady & Jangren 2000; Bershady, Jangren & Conselice 2000; Conselice 2003), complemented by the M20 and Gini coefficients (Abraham, van den Bergh & Nair 2003; Lotz, Primack & Madau 2004). This parameter set (hereafter ‘CAS+MG’) has proven to robustly link the internal light distribution of galaxies to their formation and evolution (e.g. Scarlata et al. 2007; Muñoz-Mateos et al. 2009; Lotz et al. 2008, 2011; Holwerda et al. 2014; Kleiner et al. 2014, and references therein). Besides being effective in separating the different morphological types and tracing SF, the CAS+MG scheme is particularly sensitive to their recent interaction or merging activities, making it the ideal tool to obtain a census of galaxies whose structure appears disturbed by the environment, which is crucial for our project. For a selected sample of ‘normal’ galaxies (as determined through the CAS+MG parameters), it will be possible to quantify the relevance of (pseudo-)bulges and bars as a function of the environment by surface photometry fitting.

In Fig. 1, we show examples of SSC galaxies in the magnitude range from $m^* - 1$ to $m^* + 4$ ($r = 14$ – 19 mag). For galaxies in this magnitude range, we plan to perform classical (visual) morphological classification that will be used to calibrate the CAS+MG parameters against the Hubble sequence for our data. The

signal-to-noise ratio (SNR) and spatial resolution of VST images will instead allow us to obtain reliable CAS+MG parameters up to $m^* + 6$ ($r = 21$ mag), i.e. well into the dwarf galaxy regime (see Section 4.1). Finally, AAOmega and literature spectra of more luminous galaxies will provide spectral classification through line indices.

Galaxy global properties such as colours, stellar masses and SFRs will be derived. The contribution of both unobscured and obscured SFR need to be considered since dust-obscured cluster galaxies are common in merging clusters (e.g. Haines et al. 2009a). The u -band luminosity will provide a SFR indicator once calibrated using the multiwavelength data already available in the Shapley Supercluster core (see Section 4.4), by obtaining correlations between luminosity, dust extinction and metallicity. The available mid-infrared data will provide a robust and independent indicator of obscured SF (Section 4.4), and also allow us to identify AGN via their unusually red W1–W2 colours (e.g. Wright et al. 2010). AGN will also be identified from their emission-line ratios or broad emission lines from the AAOmega spectra, or their X-ray emission in the supercluster core where *XMM* data are available.

The depths of our survey are conceived to reach the necessary accuracy of all of the above quantities, as will be outlined in the next sections.

Finally, we will not only measure statistically the environmental effects on galaxy properties on such large scales, but also ‘catch in the act’ the direct interaction of supercluster galaxies with their environment into the surrounding large-scale structure. The latter is possible only using imaging with a subkiloparsec resolution and follow-up integral-field spectroscopy for a few individual cases (e.g. Merluzzi et al. 2013).

3 THE TARGET: SHAPLEY SUPERCLUSTER

The ‘remote cloud of galaxies in Centaurus’ first identified by Shapley (1930) is one of the richest superclusters in the nearby Universe, consisting of as many as 25 Abell clusters (Zucca et al. 1993) in the redshift range $0.033 < z < 0.060$ (Quintana et al. 1995, 1997). The first spectroscopic study confirming the existence of the ‘Centaurus supercluster’ was carried out by Melnick & Moles (1987), then it was rediscovered by Scaramella et al. (1989, their ‘ α -region’) as a cluster overdensity in the Abell, Corwin & Olowin (1989, ACO) cluster catalogue and identified by Raychaudhury (1989, his ‘Shapley concentration’) as an excess in galaxy number counts in the UK Schmidt Telescope Sky Survey plates.

The Shapley Supercluster (hereafter SSC) because of its peculiar richness and location, lying in the direction of the dipole anisotropy of the cosmic microwave background (CMB), was investigated as responsible of at least a fraction of the Local Group acceleration (Raychaudhury 1989; Scaramella et al. 1989; Plionis & Valdarnini 1991; Quintana et al. 1995; Kocevski, Mullis & Ebeling 2004; Feindt et al. 2013, and references therein). To assess the role of this structure in the acceleration of the Local Group, a robust estimate of its mass is required. This gave a boost to spectroscopic and X-ray observations of the supercluster devoted to (i) map the whole structure and measure the density contrast in galaxy number counts and mass/number of associated clusters; (ii) investigate its dynamical state and (iii) derive the underlying mass distribution. A detailed review of the properties of the SSC based on previous investigations will be given elsewhere (Merluzzi et al., in preparation); here, we will describe the key characteristics of the SSC – i.e. those features making the SSC, and above all its central region, *a magnifying glass*

to investigate the effect of the environment and mass assembly on galaxy evolution.

3.1 Supercluster morphology and dynamics

By analysing a spectroscopic sample across a region of 15 deg in diameter centred on the dominant cluster A3558, Quintana et al. (1995) concluded that the supercluster has a ‘cigar-shape with the eastern side being the closest to us’ (but see also Drinkwater et al. 2004) and from the flattened geometry they suggested that it is not spherical and virialized. The complex morphology of the SSC comprises a main body at $cz \sim 15\,000$ km s⁻¹ together with walls/filaments/arms of galaxies connecting the three main systems of interacting clusters (the A3558, A3528 and A3571 complexes) as well as a foreground structure connecting the SSC to the Hydra-Centaurus supercluster ($cz \sim 4000$ km s⁻¹) and hints of a background structure at $cz \sim 23\,000$ km s⁻¹ (for details see Quintana, Carrasco & Reisenegger 2000; Proust et al. 2006). The main plane of the SSC ($cz \sim 14\,800$ km s⁻¹) extends 10×20 deg² (see Drinkwater et al. 2004). Is this vast and complex structure gravitationally bound? In this case, it would be the most massive bound structure known in the Universe.

Applying a spherical collapse model, Reisenegger et al. (2000) found that the SSC is gravitationally collapsing at least in its central region within a radius of $8 h^{-1}$ Mpc centred on A3558 including 11 ACO clusters; the very inner region, associated with the massive clusters, is likely in the final stages of collapse. X-ray observations of clusters in the SSC confirmed the overdensity: a factor 10–50 cluster overdensity (Raychaudhury et al. 1991) and factor 3 baryon overdensity (Fabian 1991) over a region of 60×80 Mpc². These results indicated that the structure is gravitationally bound at least in its central region.

From the *ROSAT* All-Sky Survey in a region of 0.27 sr, De Filippis, Schindler & Erben (2005) measured the cluster number density which turned out to be more than an order of magnitude greater than the mean density of Abell clusters at similar latitudes – mainly due to an excess of low-luminosity X-ray systems in the outskirts, suggesting that the supercluster is still accreting low-mass systems. Ettori, Fabian & White (1997) analysing a mosaic of 15×20 deg² of *ROSAT* and *Einstein Observatory* X-ray observations centred on A3558, concluded that the SSC core is dynamically bound within ~ 9 Mpc, approaching the maximum expansion before collapsing, i.e. the turnaround point. On the other hand, the dynamical analysis of Bardelli et al. (2000) proposed a scenario where the SSC has already reached its turnaround radius and the final collapse will happen in ~ 1 Gyr.

Using the X-ray cluster sample of De Filippis et al. (2005) and the spherical collapse model, Muñoz & Loeb (2008) investigated the dynamics of the SSC. Their study showed that the SSC is not bound at a radius of 51 Mpc with an outer shell moving radially away from the centre, while the excess of mass becomes enough to bind a spherical region of ~ 10 Mpc radius.

Recently, Pearson & Batuski (2013) ran N -body numerical simulations of the SSC and other superclusters in order to determine what portions of the superclusters were potentially gravitationally bound. The SSC showed the most extended bound structure among the other analysed superclusters. In particular, A3554, A3556, A3558, A3560 and A3562 have a large number of close encounters in their simulations, while an additional pair, A1736 and A3559, has also some chance of being bound. In such a crowded environment, it is very unlikely that galaxies have not been affected by cluster related

processes such as RPS, starvation and/or galaxy harassment and tidal interaction.

One of the key objectives of ShaSS is to constrain the dark matter distribution and mass over this whole region via a weak lensing analysis, and in conjunction with X-ray and dynamical analyses, estimate the extent of the region that is gravitationally bound to the SSC core, and determine whether clusters such as A3559 are currently falling towards the supercluster core.

3.2 The Shapley supercluster core

We now zoom in to the very central dense region of SSC, namely the ‘core’, which consists of three Abell clusters: A3558, A3562, and A3556 and two poor clusters (SC1327–312 and SC1329–313), resolved by Breen et al. (1994) by means of X-ray observations. The evolutionary stage of the SSC core (SSCC) is still matter of debate.

The dynamical analysis of the SSCC pointed out its very complex dynamical state with interacting clusters belonging to the same structure elongated both in declination and along the line of sight (Bardelli et al. 1994, 1998b). Several subcondensations, detected also in X-ray, could be recognized. Bardelli et al. (1998a) identified 21 significant three-dimensional subclumps, including eight in the A3558 cluster alone. This ‘clumpy’ structure as well as the proximity of the clusters (e.g. the Abell radius of A3556 is overlapped to that of A3558) makes robust estimates of the cluster velocity dispersions, masses (see Bardelli et al. 1998b) and even Abell richness (see Metcalfe, Godwin & Peach 1994) difficult. Based on their spectroscopic study, Bardelli et al. (1998a) proposed two alternative evolutionary scenarios for the SSCC: (i) a cluster–cluster collision seen just after the first core–core encounter; (ii) a series of random mergings occurred among groups and clusters.

Diffuse filamentary X-ray emission has been observed across the whole SSCC (Bardelli et al. 1996; Hanami et al. 1999; Kull & Böhringer 1999) connecting the clusters. A continuous filament of hot gas connecting A3562 and A3558 was also seen in the reconstructed thermal Sunyaev–Zeldovich (SZ) Planck survey map (Planck Collaboration 2014). The distribution of the X-ray emission was found to be clearly asymmetric in A3562, A3558 and SC1329–313 and gradients of gas temperature and metallicity have been measured (Ettori et al. 2000; Akimoto et al. 2003; Finoguenov et al. 2004) suggesting that cluster–cluster mergers have occurred and/or are ongoing. By analysing *XMM–Newton* observations, Finoguenov et al. (2004) proposed a tidal interaction between SC1329–313 and A3562 to explain the observed properties of the hot gas in A3562 – the tailed shape of the X-ray emission associated with SC1329–313 as well as the sloshing of the A3562 core. Another detailed analysis of *XMM–Newton* and *Chandra* by Rossetti et al. (2005) pointed out the very complex dynamical history of A3558 having characteristics which are typical of both merging (e.g. gas temperature gradients) and relaxed (e.g. cool core) clusters. They also detected a cold front leading in the NW direction and probably due to the sloshing of the cluster core induced by the perturbation of the gravitational potential associated with a past merger.

A weak and asymmetric radio halo has been detected in A3562 (Venturi et al. 2003; Giacintucci et al. 2005). The radio halo correlates with the X-ray emission and presents a radio head-tailed galaxy embedded in it and located between the cluster core and SC1329–313, thus supporting the interaction between the two systems as proposed by Finoguenov et al. (2004).

A deficit of radio galaxies with respect to the radio-optical luminosity function (LF) of other clusters, although probably to be ascribed to A3558 alone, was observed by Venturi et al. (2000). On the other hand, Miller (2005), with a radio survey of a 7 deg² region of the SSCC, found a dramatic increase in the probability for galaxies in the vicinity of A3562 and SC1329–313 to be associated with radio emission suggesting young starbursts related to the recent merger of SC1329–313 with A3562. This observation could be explained by galaxy merging efficiently transferring gas into the galaxy centre, feeding AGN and then switching on SF (Bekki 1999).

Although the studies mentioned above are fundamental to demonstrate the complex dynamical status of the SSC and its core, none of them could systematically tackle the issue of galaxy evolution in the supercluster environment due to the lack of accurate and homogeneous multiband imaging covering such an extended structure. This prevented to collect information about the integrated (magnitudes, colours, SFR) and internal properties (morphological features, internal colour gradients) of the supercluster galaxies. Excluding observations of single clusters in SSC, the first CCD observations covering a ~ 2 deg² contiguous area of SSCC were those of the ACCESS project (PI: P. Merluzzi; Haines et al. 2006a; Mercurio et al. 2006, see Section 3.3) which analysed a galaxy sample complete down to $B = 22.5$ ($> m^* + 6$) and $R = 22.0$ ($> m^* + 7$), i.e. well into the dwarf galaxy regime.

3.3 ACCESS project: main results and open issues

The multiwavelength data set of the ACCESS¹ project (see Merluzzi et al. 2010), covering 2–3 deg² of the SSCC allowed us to obtain the complete census of stellar content and SF across the core region from A3556 to A3562. This project was dedicated to investigate the effect of the environment on galaxy evolution in the SSCC exploiting one of the first multiwavelength data sets available for such a wide area of a supercluster. The data include panoramic imaging in the UV (Galaxy Evolution Explorer, *GALEX*), optical (ESO Wide Field Imager, WFI), NIR (United Kingdom Infrared Telescope, UKIRT/WFCAM) and mid-infrared (*Spitzer*/MIPS), as well as high SNR medium-resolution optical spectroscopy (AAT/AAOmega) for 448 supercluster galaxies. In the following, we mention the main results of this project and the open issues that stimulated us to undertake a new multiwavelength study over a wider supercluster area.

By studying the optical (Mercurio et al. 2006) and near-infrared (NIR; Merluzzi et al. 2010) LFs down to $m_K^* + 6$ it was found that the slope of the LF increases from high- to low-density environments, indicating that mechanisms such as galaxy harassment and/or tidal stripping contribute to shape the LF. The stellar mass function (SMF), however, does not seem to change its trend with galaxy density in the SSCC and does not show the sharp upturn below $\mathcal{M} = 10^{9.5} M_\odot$ observed in the field galaxy population (Pozzetti et al. 2010; Baldry et al. 2012). Is this difference due to the different evolution of the supercluster and the field galaxies, or it is only an artefact due to the different range of stellar masses analysed for field and supercluster galaxies? And also, which are the different contributions to the galaxy SMF of the blue and red galaxy populations in the supercluster from the cluster cores to the filaments and field? To answer these questions, the mass range should be extended and the membership ascertained.

¹ European funded project: ACCESS: A Compete CEnsus of Star Formation in the Shapley Supercluster, PI: P. Merluzzi; www.oacn.inaf.it/ACCESS

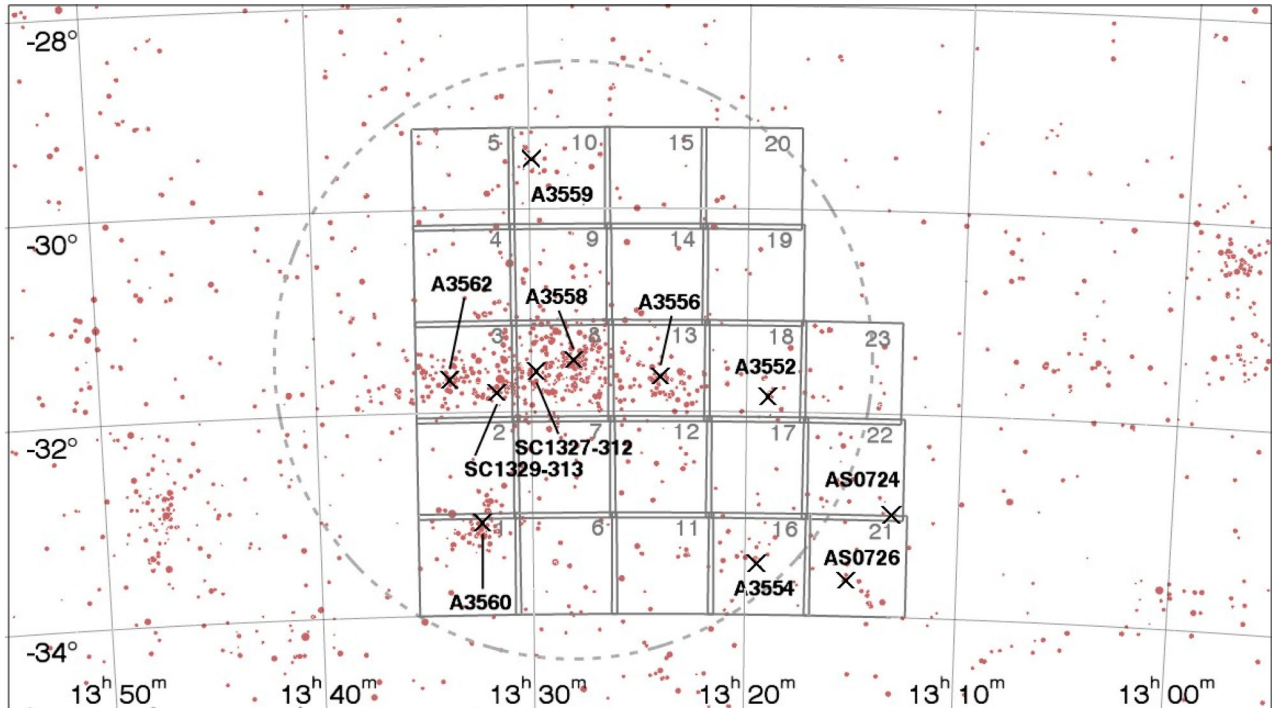


Figure 2. The 23 1 deg^2 VST fields mapping the ShaSS region. Red dots indicate the supercluster members in the range $V_h = 11\,300\text{--}17\,000 \text{ km s}^{-1}$ taken from literature. The size of the dots are proportional to the K -band flux. Black crosses show the cluster centres. The 10 Mpc radius dotted circle encloses the supercluster region believed to be dynamically bound. The SSCC corresponds to fields #3,8,13.

We also found evidence that the bulk of the star-forming galaxies have been recently accreted from the field and have yet to have their SF activity significantly affected by the cluster environment (Haines et al. 2011a) and that the vast majority of SF seen in the SSCC comes from normal infalling spirals (Haines et al. 2011b). Nevertheless, this analysis was limited to galaxies belonging to the SSCC and thus in an extremely dense environment. With the aim to reach a comprehensive picture of how and where galaxies start to quench their SF, it is important to move out from the cluster cores and analyse the properties of the infalling galaxies, i.e. to follow the cluster member from the *converging* filaments into the clusters. This will also enable us to investigate whether the formation of early-type galaxies, dominant in cluster cores, can be driven by *morphological quenching* (see Martig et al. 2009) or RPS/starvation (e.g. Crowl & Kenney 2008). It is unclear which is the most likely and effective process at work (Haines et al. 2011c). Whether the formation of passive early-type galaxies in cluster cores should involve the *prior* morphological transformation of late-type spirals into S0/Sa, mechanisms as pre-processing and tidal interaction have to be considered and quantified. In order to understand if this is the case, we need a morphological study of supercluster galaxies, both isolated and in groups, in the cluster outskirts and beyond.

4 SHASS: THE DATA

We will map a region of $\sim 260 \text{ Mpc}^2$ including the SSCC and other six galaxy clusters (AS0724, AS0726, A3552, A3554, A3559 and A3560, see Fig. 2). The supercluster region is chosen to ensure to map the structures directly connected to the SSCC. In fact, the 11 clusters in the region are all within 500 km s^{-1} of the central cluster A3558. The survey boundaries are chosen to cover all 11 clusters and the likely connecting filaments, but also to extend into the field

(fields 19 and 20 in Fig. 2). The main characteristics of the clusters and groups in the survey are listed in Table 1. We would like to point out that among the other quantities taken from the literature, the virial radii and masses should be considered only indicative for the SSC clusters studied here for which the assumptions of spherical symmetry and isotropy of the velocities are likely not applicable.

The data set includes optical ($ugri$) and NIR (K) imaging acquired with VST and VISTA, respectively, and optical spectroscopy with AAOmega. At present, the i -band imaging and AAOmega spectroscopic surveys are completed, while the other observations are ongoing. Table 2 summarizes the depths and completion of the imaging surveys which are described in Sections 4.1 and 4.3. In the table, we list both the target depths and those measured from the data.² For the characteristics of the spectroscopic survey see Section 4.2.

4.1 VST-ACCESS survey

The new optical survey (PI: P. Merluzzi), conceived in the framework of the ACCESS project and named after it VST-ACCESS, is being carried out using the Italian INAF Guaranteed Time of Observations (GTO) with OmegaCAM at the 2.6 m ESO telescope VST (Schipani et al. 2012) located at Cerro Paranal (Chile). The corrected field of view of $1^\circ \times 1^\circ$ allows the whole ShaSS area to be covered with 23 VST fields as shown in Fig. 2. Each of the contiguous VST-ACCESS fields is observed in four bands: $ugri$. Red dots in Fig. 2 denote the 1676 spectroscopic supercluster member galaxies ($11\,300 < V_h < 17\,000 \text{ km s}^{-1}$) available from literature at

² The noise inside the aperture can be estimated from the flux and its uncertainty derived by SExtractor (see equation 61 of SExtractor User Manual; www.astromatic.net/software/sextractor)

Table 1. Galaxy clusters in the ShaSS region.

Cluster	RA J2000	Dec. J2000	Redshift	σ (km s ⁻¹)	Richness ^a	R_{VIR} (Mpc)	Mass ^b [M_{\odot}]
AS0724	13 13 08.6	-32 59 38 ^c	14864 ± 157 ^d	510 ± 85 ^d	0		8.8 × 10 ¹³ ^e
AS0726	13 15 11.7	-33 38 52 ^d	14892 ± 137 ^d	578 ± 77 ^d	0	0.96 ^f	4.6 × 10 ¹³ ^f
A3552	13 19 00.7	-31 51 04 ^c	14753 ± 119 ^d	682 ± 60 ^d	1		3.6 × 10 ¹³ ^e
A3554	13 19 27.6	-33 29 49 ^c	14431 ± 94 ^d	560 ± 66 ^d	1	0.80 ^f	5.8 × 10 ¹³ ^f
A3556	13 24 00.2	-31 39 22 ^c	14357 ± 76 ^g	643 ⁺⁵³ ₋₄₃ ^g	0	0.98 ^f	1.7 × 10 ¹⁴ ^f
A3558	13 28 02.6	-31 29 35 ^c	14403 ⁺⁶⁰ ₋₅₅ ^g	996 ⁺⁴⁰ ₋₃₆ ^g	4	1.16 ^f	1.3 × 10 ¹⁵ ^g
A3559	13 29 53.1	-29 30 22 ^h	14130 ± 57 ^d	519 ± 45 ^d	3	0.31 ^f	2.0 × 10 ¹³ ^f
A3560	13 32 22.0	-33 05 24 ⁱ	14551 ± 106 ^d	793 ± 116 ^d	3	1.33 ^f	3.4 × 10 ¹⁴ ^f
A3562	13 33 47.0	-31 40 37 ^c	14455 ± 191 ^d	1197 ± 194 ^d	2	0.89 ^f	(3.9±0.4) × 10 ¹⁴ ^j
SC1327–312	13 29 45.4	-31 36 12 ^k	14844 ⁺¹⁰⁵ ₋₂₁₁ ^g	691 ⁺¹⁵⁸ ₋₁₄₆ ^g		1.30 ^f	3.0 × 10 ¹³ ^f
SC1329–313	13 31 36.0	-31 48 45 ^k	14790 ⁺¹¹⁴ ₋₆₇ ^g	377 ⁺⁹³ ₋₈₂ ^g		1.14 ^f	3.7 × 10 ¹³ ^f
			13348 ⁺⁶⁹ ₋₈₃ ^g	276 ⁺⁷⁰ ₋₆₁ ^g			

^a Abell et al. (1989).

^b All the masses are dynamically derived except for that of A3562 which is derived from X-ray observations.

^c De Filippis et al. (2005).

^d Prout et al. (2006), the redshift is corrected with respect to the CMB.

^e Reisenegger et al. (2000), cluster mass within a radius enclosing an average density 500 times the critical density.

^f Ragone et al. (2006), estimates of virial radius and virial mass.

^g Bardelli et al. (1998b), the mass value is transformed into the adopted cosmology.

^h David, Forman & C. (1999).

ⁱ Bardelli et al. (2002).

^j Ettori et al. (2000), the mass value is transformed into the adopted cosmology.

^k Breen et al. (1994).

For SC1329–313 mass and virial radius are estimated for the whole system, while redshifts and velocity dispersions are given for the two different clumps by Bardelli et al. (1998b).
Uncertainties are quoted when available.

Table 2. ShaSS imaging: depths and current coverage.

Band	Exp. Time (s)	Depth [5 σ] ^a		Seeing ^b (arcsec FWHM)	Complete (per cent)
		target	measured		
<i>u</i>	2955	24.5	24.3	0.8–1.1	48
<i>g</i>	1400	24.2	24.8	0.6–1.0	43
<i>r</i>	2664	24.2	24.3	0.6–0.8	61
<i>i</i>	1000	22.4	23.2	0.5–1.0	100
<i>K</i>	1620	20.4	20.3	0.6–1.0	60

^a Within a 3 arcsec diameter aperture.

^b Range of seeing estimated from the already observed fields.

Table 3. Assumed m^* values.

Band	m^*	Reference
B^a	15.35	Mercurio et al. (2006)
R^a	14.52	Mercurio et al. (2006)
K^a	11.70	Merluzzi et al. (2010)
r_{AB}^b	15.00	Mercurio et al. (2006)

^a Vega photometric system.

^b Derived from the R -band value once converted in r_{AB} with Blanton & Roweis (2007) and Fukugita, Shimasaku & Ichikawa (1995).

the time of the survey planning. The X-ray centres are indicated by crosses for all the known clusters except AS0726, whose centre is derived by a dynamical analysis.

Survey depths. We will achieve our scientific goals studying the global and internal physical properties of SSC galaxies down to $m^* + 6$. In particular, we need to (i) derive accurate morphology

as well as structural parameters ($\delta \log r_e \sim 0.04$ and $\delta n_{\text{Ser}} \sim 1$) and detect some of the observational signatures related to the different processes experienced by supercluster galaxies (e.g. extraplanar material); (ii) estimate accurate colours, photo- z ($\delta z < 0.03$, see Christodoulou et al. 2012) and stellar masses; (iii) evaluate the SFRs and resolve the star-forming regions at least for the subsample of brighter galaxies. The required SNR depends on which galaxy property is measured: it is higher for the morphological analysis and resolving internal properties/structures (SNR \sim 100 in our 3 arcsec aperture, see Conselice et al. 2000; Häussler et al. 2007), but it can be significantly lower for accurate measurements of aperture photometry and colours (SNR \sim 20). In the latter case, however, the minimum required SNR should be achieved in all bands.

We chose to use mainly the r -band imaging for the morphological analysis, so the r band defines the survey depths in all bands. In r band $m_r^* \sim 15$ (AB magnitude, see Table 3) and at $m_r^* + 6 = 21$, we require SNR = 100 for the morphological/structural analysis. The completeness magnitude of the catalogue in r band is instead defined by the star/galaxy separation (see Section 6) which we estimate to be robust down to $r = 23.5$ mag (SNR \sim 10 within a 3 arcsec aperture), corresponding to a limiting magnitude $r = 24.2$ mag (SNR = 5). We are collecting the r -band imaging under the best observing conditions having a median seeing full width at half-maximum (FWHM) ~ 0.8 arcsec corresponding to 0.75 kpc at $z \sim 0.05$. Additionally, the r imaging is fundamental to our weak lensing analysis, to ensure a sufficient density of lensed background galaxies with shape measurements.

The typical colours at $z \sim 0.048$ for red sequence galaxies in the AB photometric system are: $u - g \sim 1.4$, $g - r \sim 0.8$, $r - i \sim 0.35$ according to stellar population models (Bruzual & Charlot 2003, $\tau = 3.0$ Gyr, $Z = Z_{\odot}$). This approximation allowed us to estimate the

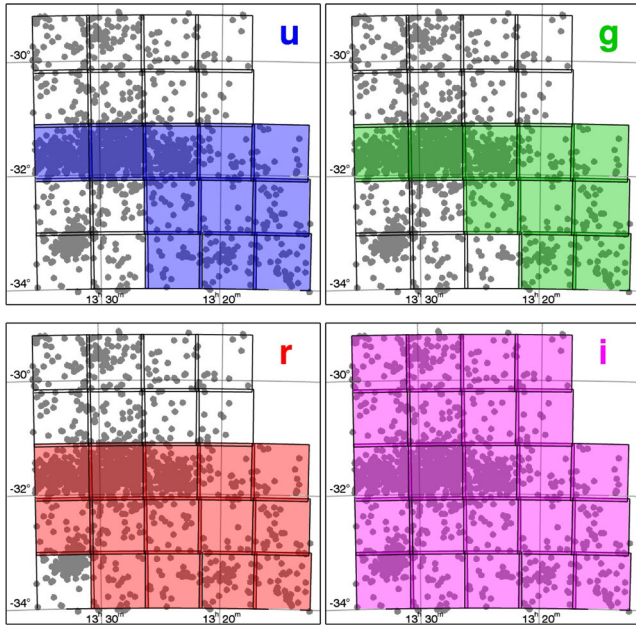


Figure 3. The coloured squares show the VST-ACCESS field coverage at the end of ESO-P91 (September 2013). Grey dots are the supercluster members from literature.

required depths for the other three bands which all should provide complete galaxy samples down to $m^* + 6$ with $\text{SNR} \geq 20$.

Observation strategy. The 1 square degree unvignetted field of view is sampled at $0.21 \text{ arcsec px}^{-1}$ by OmegaCAM with a $16 \text{ k} \times 16 \text{ k}$ detector mosaic of 32 CCDs which constitute the science array. The detector mosaic presents gaps up to 25 and 85 arcsec wide in X and Y direction, respectively. To bridge the gaps, we chose the dither offsetting mode with a diagonal pattern and five exposures for i , g and u bands. In order to reach the required depth, cover the gaps and avoid saturation in the centre of bright galaxies, the r -band images are instead obtained with nine exposures and smaller offsets. The contiguous pointings of VST-ACCESS are overlapped by 3 arcmin. This strategy allows, for each band, to use a few pointings observed under photometric conditions in each run to calibrate all the other fields as well as to check the photometric accuracy.

Survey strategy. To carry out the optical survey with VST about 100 h of telescope time are foreseen. The observations started in 2012 February and are in progress. In the first four ESO semesters, a total of 48 h have been allocated of which 87 per cent carried out. We show the coverage at the end of the first two years of the survey (2013 September) in Fig. 3 and in Table 4 the distribution of the observations across this period. The following strategy is chosen to set the priority of the fields to be observed.

(1) To map the whole area in the i band in order to have a magnitude-limited galaxy catalogue with high astrometric and photometric accuracy. This catalogue was mandatory to carry out the spectroscopic survey (see below).

(2) To collect multiband optical imaging in the region of the SSCC where we have already available the multiband data set (see Section 4.4) which allows a cross-check of some of the quantities derived from the optical data, e.g. SF indicators.

(3) To map in optical bands the southern 15 deg^2 first, including the two Abell clusters A3554 and A3560 and probable filaments

Table 4. VST-ACCESS observations P88–P91.

Field	Bands	ESO Periods ^a
F01	<i>i</i>	P89
F02	<i>ir</i>	P89 P91
F03	<i>irgu</i>	P90 P91 P88 P88
F04	<i>i</i>	P90
F05	<i>i</i>	P90
F06	<i>ir</i>	P89 P91
F07	<i>ir</i>	P89 P91
F08	<i>irgu</i>	P88 P89 P88 P90
F09	<i>i</i>	P91
F10	<i>i</i>	P90
F11	<i>iru</i>	P89 P91 P91
F12	<i>irgu</i>	P89 P90 P91 P91
F13	<i>irgu</i>	P90 P90 P89 P88
F14	<i>i</i>	P90
F15	<i>i</i>	P90
F16	<i>irgu</i>	P89 P90 P91 P91
F17	<i>irgu</i>	P89 P90 P91 P91
F18	<i>irgu</i>	P88 P90 P90 P89
F19	<i>i</i>	P91
F20	<i>i</i>	P91
F21	<i>irgu</i>	P88 P90 P91
F22	<i>irgu</i>	P88 P90 P90
F23	<i>irgu</i>	P88 P90 P90

^aFor each band in column 2 the period of observations is indicated.

connecting these clusters to the SSCC and which is already covered by the VISTA K -band survey (see Section 4.3).

(4) To complete then the wavelength coverage of the northern 8 deg^2 starting from the eastern side (A3559).

4.2 AAOmega spectroscopic survey

A spectroscopic survey of 21 deg^2 of the ShaSS area³ was carried out with the AAOmega spectrograph at the 3.9 m Anglo Australian Telescope in May 2013 (PI: P. Merluzzi). The main aim of this survey is to trace the structure of the SSC, beyond the known galaxy clusters, including filaments and groups, in order to describe the local environment in as much detail as possible. We also aimed to reach high completeness in stellar mass and SFR, and for such reason the priorities given to the targets were based on WISE W1 $3.4 \mu\text{m}$ (as a proxy for stellar mass) and W3 $12 \mu\text{m}$ (as a proxy for SFR) magnitudes. The detailed analysis of the spectroscopic survey will be presented in a forthcoming article. Here, we note brief details of this process and the pertinent numbers of galaxies recovered.

AAOmega is a dual-beam fibre-fed spectrograph, allowing the acquisition of up to 392 simultaneous spectra of objects in a 2° diameter field on the sky. Our observations were carried out with gratings 580V and 385R on the blue and red arm, respectively, covering the wavelength range $370\text{--}880 \text{ nm}$ at a resolution of $\lambda/\Delta\lambda \sim 1300$. The survey consisted of 14 AAOmega pointings of 1.5 h integration time (including several repeats to constrain the pair-wise blunder

³ At the time of the AAOmega observations two of the VST fields had not yet been observed.

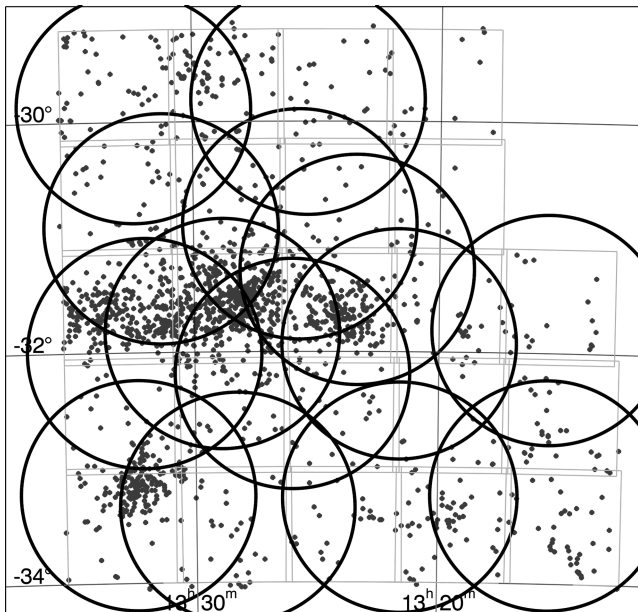


Figure 4. Overlaid to the VST-ACCESS coverage (grey boxes), the 14 2deg-diameter AAOmega pointings (circles). Grey dots as in Fig. 3. Two VST-ACCESS pointings (top-right) are not covered by the spectroscopic survey.

rate⁴), whose centres were determined using a simulated annealing algorithm in the same manner as Drinkwater et al. (2010) which allows us to optimize the number of targets. Fig. 4 shows the configuration of the 14 AAOmega pointings superimposed to the VST fields.

The data were reduced in a standard manner using 2DFDR (Lewis et al. 2002) which included a Laplacian edge detection to remove incident cosmic rays. Redshifts were obtained using the RUNZ code common to 2dFGRS, 2SLAQ, GAMA, WiggleZ (Colless et al. 2001; Cannon et al. 2006; Driver et al. 2011; Drinkwater et al. 2010) which involves Fourier space correlation of each spectrum to a wide batch of template spectra.

We obtained 4037 new redshift measurements in the whole area, which, combined with pre-existing measurements, give a total of 6130 redshifts. In the redshift range of the SSC assumed here (see Fig. 7), the available redshifts are now 2281, of which 915 have been obtained with our AAOmega survey (see Section 7.1). Beyond the SSC, 3014 AAOmega redshifts encompass the large-scale background structure, unveiling remarkable peaks in the range $z = 0.07$ – 0.25 which were mostly unknown before.

The AAOmega survey is 80 per cent complete down to $i = 17.6$ mag, $W1 = 14.7$ mag and $W3 = 10.7$ mag. These two last translate into a completeness of 80 per cent in stellar mass and SFR of, respectively, $\mathcal{M}_* \sim 8.7 \times 10^9 M_\odot$ and $\text{SFR} \sim 0.7 M_\odot \text{ yr}^{-1}$ at the supercluster redshift. The value of the stellar mass is obtained by means of the $W1$ -stellar mass calibration as determined by matching SDSS DR7 galaxies at the same distance as the SSC, with stellar mass estimates from Salim et al. (2007) and Blanton & Roweis (2007), to the WISE all sky-catalogue, and taking the best-fitting linear relation between $W1$ flux and stellar mass. The SFR is

⁴ This is the incidence rate where one measures and re-measures (independently) the redshift of a target galaxy. If these two values are different by $>600 \text{ km s}^{-1}$, it is a pair-wise blunder (Colless et al. 2001).

evaluated from the $W3$ magnitude using equation 1 of Donoso et al. (2012).

4.3 Shapley-VISTA K -band survey

The NIR survey (PI: C. P. Haines) is being carried out with VIRCAM at the 4 m ESO telescope VISTA located at Cerro Paranal (Chile). These K -band observations are accomplished using the Chilean GTO.

VIRCAM covers 0.59 deg^2 per single pointing (paw-print) with $16\,048 \text{ pixel} \times 2048 \text{ pixel}$ detectors. The gaps between the arrays amount to 90 and 42.5 per cent of the detector size along the X and Y axis, respectively. So, to obtain a contiguous coverage of the $1.5^\circ \times 1^\circ$ field of view six offsetting paw-prints (a tile) are needed. The mean pixel size is 0.339 arcsec. Almost all the ShaSS area can be covered with 15 VIRCAM tiles. At present, the nine tiles shown in Fig. 5 have been observed.

Combining the K -band data with the VST optical imaging ($ugri$), we will derive accurate stellar masses by means of stellar population models constrained by the observed optical and infrared colours. This will allow us to robustly measure the SMF and distinguish the contributions of star-forming and passive galaxies to the SMF in different mass ranges down into the dwarf regime.

VISTA survey depths. To achieve our scientific goals, we need to estimate stellar masses down to $\mathcal{M} = 10^7 M_\odot$ (see Section 3.3) corresponding to $K \sim 19.6$ ($\sim K^* + 8$, see Table 3) which is therefore the requested completeness limit for accurate aperture photometry ($\text{SNR} > 10$ for a point source in 3 arcsec aperture). In order to estimate the NIR-optical colour gradients with 20–30 per cent accuracy using VST-ACCESS and Shapley-VISTA data, an $\text{SNR} \sim 40$ – 50 is required and according to the survey depth this can be achieved at magnitudes brighter than $K \sim 17.7$ corresponding to $K^* + 6$ ($\geq 10^{8.75} M_\odot$; Merluzzi et al. 2010).

VISTA observations and survey strategy. The 13.5 deg^2 southern ShaSS regions was observed in the K_s band by VISTA in 2014 April–May (093.A-0465: 18 hours allocated), covered by a 3×3 mosaic of VIRCAM tiles as shown in Fig. 5. Each stacked paw-print image consists of $9 \times 9 \text{ s}$ exposures, repeated using five-point jitter pattern with a maximal offset from the central position of 30 arcsec. Each tile of six paw-prints was observed a second time, such that each point within the tile region was covered by four stacked

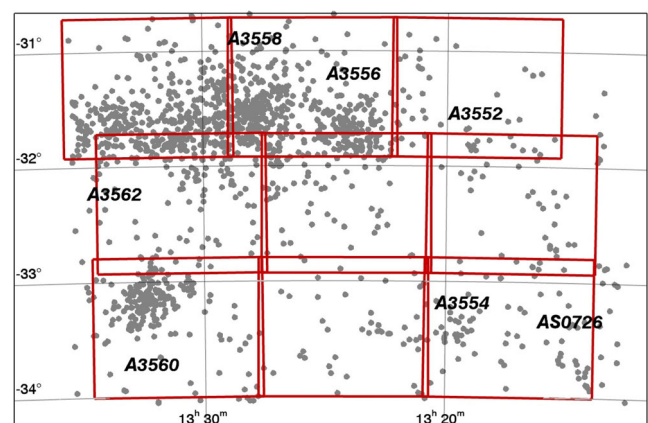


Figure 5. The nine VISTA tiles (red rectangles) mapping the southern SShaSS area. Grey dots as in Fig. 3.

paw-prints,⁵ giving an exposure time of 1620 s pxl^{-1} . As for the VST-ACCESS survey, the tiles are slightly overlapped (~ 8 arcmin in Y and ~ 3 arcmin in X) to confirm the consistency of the photometry from one tile to another. Initial data reduction steps were performed at the Cambridge Astronomical Survey Unit (CASU) using a software pipeline developed specifically for the reduction of VIRCAM data, as part of the VISTA Data Flow System⁶ (VDFS; Irwin et al. 2004). VISTA K_s magnitudes are calibrated on to the Vega magnitude photometric system using unsaturated 2MASS stars in the image, based on their magnitudes and colours in the 2MASS point source catalogue. The resultant stacked K -band images have FWHMs in the range 0.6–1.0 arcsec. The K -band magnitude detection limit at 5σ within a 3 arcsec aperture turns out to be 20.3 mag.

We aim to complete observations of the remaining northern 8 deg^2 region with VISTA in 2015.

4.4 Complementary data

The Wide-field Infrared Survey Explorer (WISE; Wright et al. 2010) is a NASA Explorer mission that observed the entire sky in 2010 in four near/mid-infrared bandpasses: $3.4 \mu\text{m}$ (W1), $4.6 \mu\text{m}$ (W2), $12 \mu\text{m}$ (W3) and $22 \mu\text{m}$ (W4). All the data have been reduced, calibrated and released to the public.⁷ The WISE satellite made twice as many passes of the region covered by the ShaSS survey as it did on average for most areas of the sky (the ecliptic poles excluded), and so the limiting magnitudes are slightly deeper than the typical WISE depths reaching $W1 = 16.96$ (Vega magnitude) at $\text{SNR} = 10$. This corresponds to $\sim m^* + 5$ for galaxies in the SSC. The W2 filter ($4.6 \mu\text{m}$) is a magnitude shallower. We typically reach $W2 = 15.26$ at an SNR of 10. The W3 filter ($12 \mu\text{m}$) reaches a depth of $W3 = 11.12$ mag (1.0 mJy) at an SNR of 10. This corresponds to an SFR of $0.46 \text{ M}_\odot \text{ yr}^{-1}$ for galaxies in the SSC (equation 1 of Donoso et al. 2012). The W4 filter ($22 \mu\text{m}$) reaches a depth of $W4 = 7.58$ (7.7 mJy) at SNR of 10. This corresponds to an SFR of $\sim 2.1 \text{ M}_\odot \text{ yr}^{-1}$ for galaxies in the SSC. The resolution of the WISE bands are 6.1, 6.4 and 6.5 arcsec FWHM in bands W1, W2 and W3. W4 instead has an FWHM of 12.0 arcsec.

To cross-correlate WISE and optical catalogues, we use the software STILTS⁸ searching for the closest match within a 3 arcsec radius between the VST-ACCESS and WISE detections. Due to higher spatial resolution, the astrometry of the optical images is more accurate and source deblending is checked to avoid multiple detections of a single extended source. In the WISE catalogue, excluding the very extended and bright sources, i.e. few resolved nearby galaxies, there are very few multiple detections of extended sources. This allows us to associate each optical (*gri*) detection only with one IR detection. The few cases of multiple detection of extended sources in the IR images are then fixed in the final cross-correlated optical-IR catalogue. This approach has been adopted for the *i*-WISE catalogue used for the spectroscopic survey. We note also that nearby IR resolved galaxies were not spectroscopic targets, since their redshifts were already available from the literature.

The ShaSS data are complemented in the central $2\text{--}3 \text{ deg}^2$ by *Spitzer*/MIPS $24/70 \mu\text{m}$ photometry and *GALEX* near-ultraviolet and FUV imaging which allowed us to produce a complete census of SF (both obscured and unobscured), extending down to SFRs

$\sim 0.02\text{--}0.05 \text{ M}_\odot \text{ yr}^{-1}$, i.e. down to levels comparable to the Small Magellanic Cloud. Although these data do not cover the whole ShaSS region they are fundamental for our survey. The W3 data will enable us to measure the SFR down to $0.2 \text{ M}_\odot \text{ yr}^{-1}$ at 5σ , while the W1–W3 colour allows us to reproduce and map the bimodal galaxy distributions seen in the f_{24}/f_K galaxy colours in the supercluster core (Haines et al. 2011a), but over the entire supercluster region, splitting the supercluster galaxies into star-forming, transitional and passive populations. This is possible because we verified that a strong linear correlation exists between W1–W3 colour and f_{24}/f_K .

NIR K -band imaging from the UKIRT with the Wide Field Infrared Camera are also available for 3 deg^2 in the SSC providing a complete galaxy sample down to $K = 18$ (Merluzzi et al. 2010). *XMM-Newton* archive data are available for the SSC region and for A3560 and will enable us to map the ICM gas and to identify possible shock fronts due to cluster mergers and then to investigate the effects of such events on galaxy properties.

For a subsample of supercluster galaxies, we have been obtaining integral-field spectroscopy with WiFeS (Dopita et al. 2007) on the Australian National University 2.3 m telescope at Siding Spring in Australia. We are targeting a few bright ($m < m^* + 1.5$) supercluster galaxies showing evidence of undergoing transformation. All these galaxies are selected by either disturbed morphology, such as asymmetry and tails, or evidence of SF knots (Merluzzi et al. 2013). For these objects, we also have obtained $\text{H}\alpha$ imaging with Maryland–Magellan Tunable Filter (MMTF; Veilleux et al. 2010) on the Magellan–Baade 6.5 m telescope at the Las Campanas Observatory in Chile.

5 VST DATA REDUCTION

The VST images have been processed using the VST-Tube imaging pipeline (Grado et al. 2012), developed for the data produced by the VST.

After applying the overscan correction and bias subtraction, we divide by the master flat – a normalized combination of the dome and twilight flats, in which the twilight flat is passed through a low-pass filter first.

Due to differences between the electronic amplifiers, the CCDs do not have the same gain levels. In order to have the same photometric zero-point (ZP) for all the mosaic chips, a gain harmonization procedure has been applied. The procedure finds the relative CCD gain coefficients which minimizes the background level differences in adjacent CCDs. A further correction is needed for the light scattered by the telescope and instrumental baffling. The resulting uncontrolled redistribution of light adds a component to the background and the flat-field will not be an accurate estimate of the spatial detector response. Indeed, after flat-fielding, the image background will appear flat but the photometric response will be position-dependent. This error in the flat-fielding can be mitigated through the determination and application of the illumination correction (IC) map. The IC map is determined by comparing the magnitudes of photometric standard fields with the corresponding SDSS DR8 point spread function (PSF) magnitudes. The differences between the VST and SDSS magnitudes are plotted versus X and Y axis in Fig. 6 before (top panels) and after (bottom panels) the IC. Usually, the same IC map can be used for observations carried out on the time-scale of a month.

In the case of the *i* band, it is required to correct for the fringe pattern due to thin-film interference effects in the detector of sky emission lines. Also this is an additive component that must be

⁵ Each point within a tile is covered by two of the six paw-prints.

⁶ <http://casu.ast.cam.ac.uk/surveys-projects/vista/vdfs>

⁷ <http://irsa.ipac.caltech.edu/Missions/wise.html>

⁸ <http://www.star.bris.ac.uk/mbt/stilts>

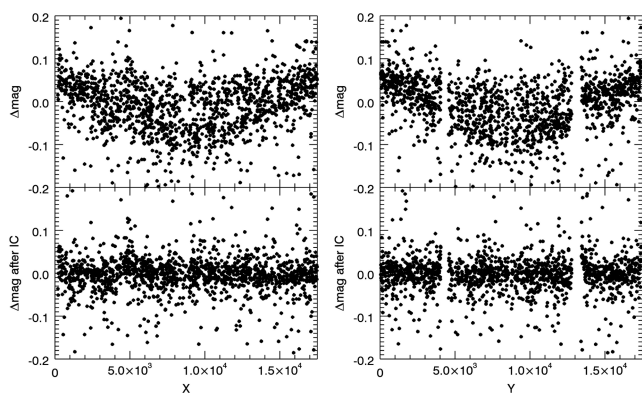


Figure 6. Top: differences between the r -band magnitudes of stars in a standard field observed with VST and SDSS magnitudes versus X and Y image axis. Bottom: the same Δmag after the IC correction has been applied.

subtracted. The fringing pattern is estimated using the $\frac{\text{SuperFlat}}{\text{TFlat}}$ ratio where *SuperFlat* is obtained by overscan and bias correcting a sigma-clipped combination of science images and *TFlat* is the sky flat. To this aim, the dithering amplitudes of the science frames used must be larger than the extended object sizes in the same images. This is mandatory to allow the sigma-clip procedure to efficiently remove such objects. The difference between the above ratio across the image and its model obtained with a surface polynomial fit is the fringing pattern. The fringe pattern is subtracted from the image applying a scale factor which minimizes the absolute difference between the *peak* and *valley* values (maximum and minimum in the image background) in the fringe corrected image.

The relative photometric calibration among the exposures contributing to the final mosaic image is obtained comparing the magnitudes of bright unsaturated stars in the different exposures, i.e. minimizing the quadratic sum of differences in magnitude using SCAMP (Bertin 2006). The absolute photometric calibration is computed on the photometric nights comparing the observed magnitude of stars in photometric standard fields with SDSS photometry. For those fields that are not observed in photometric but clear nights, we take advantage of the sample of bright unsaturated stars in the overlapping regions between clear and photometric pointings and, by using SCAMP, each exposure of the clear fields is calibrated on to the contiguous photometrically calibrated field. In Table 5, we give examples of the fitted values for the ZP and colour term obtained using the PHOTCAL tool (Radovich et al. 2004). In several cases, the photometric standard star fields were observed with insufficient span in airmass to do a suitable fit. The extinction coefficient was then taken from the extinction curve provided by ESO. The errors listed in Table 5 are those of the fit. The actual errors for the ZP is given by the rms of the detection among the different exposures of each pointing and turns out to be less than 0.03 mag in all bands.

The relative astrometry among the pointings is derived using SCAMP, while for the absolute astrometry we refer to the astrometric catalogue 2MASS obtaining $\lesssim 0.3$ arcsec astrometric accuracy in all bands. The image resampling, application of astrometric solution and co-addition is performed using the tool SWARP (Bertin et al. 2002) which produces the final stacked image with a weighted average. At the stage of the co-addition the scale factors for relative and absolute photometric calibration are applied.

6 ANALYSIS OF THE OPTICAL IMAGES

Once the mosaic images were obtained, we proceeded with the catalogue extraction and measured the photometric accuracy and completeness. Details of ShaSS catalogue extraction and release will be given in a companion article (Mercurio et al. in preparation).

We used the software package SExtractor (Bertin & Arnouts 1996) estimating the background locally and using a Gaussian filter for source detection. For each source, we measured magnitudes in different apertures as well as Kron and PSF magnitudes.

To distinguish between stars and galaxies, we adopted a progressive approach (Annunziatella et al. 2013) using the following parameters provided by SExtractor: (i) the stellarity index to select point-like sources; (ii) the half-flux radius as a measure of source concentration; (iii) the new SExtractor parameter which takes into account the difference between the model of the source and the model of the local PSF; (iv) the peak surface brightness above background; (v) a final visual inspection for objects with ambiguous values of the stellarity index.

We estimated the completeness magnitudes using both trends of the galaxy number counts and the method by Garilli, Maccagni & Andreon (1999). The catalogues turned out to be 100 per cent complete at the total magnitudes of 23.9–24.1, 23.8–24.0, 23.3–23.5, 21.8–22.0 mag in *ugri* bands, respectively. The ranges of magnitudes are due to small differences in seeing among the VST fields.

We measured the SNR inside a 3 arcsec aperture as function of magnitudes achieving $\text{SNR} = 20$ at *ugri* at 22.8, 23.3, 22.8, 21.7 mag, respectively. $\text{SNR} \sim 100$ is reached at 21.1 mag in r band ($m^* + 6$) as required. These depths enable us to study the galaxy population down to $m^* + 6$.

7 RESULTS

In this section, we derive the first quantitative characterization of the environment in the survey area in terms of stellar mass density. We also derive the dark matter distribution in the very centre of the SSC including the cluster A3558 and the poor cluster SC1327–312.

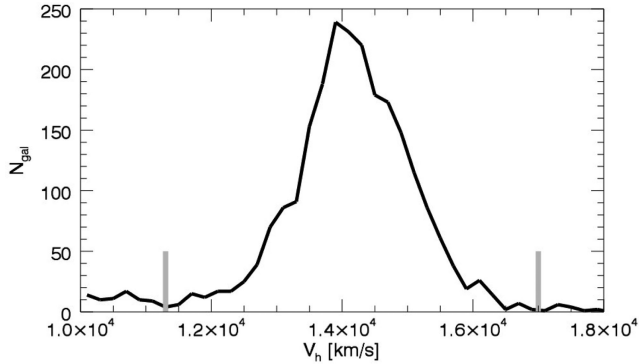
7.1 Galaxy density

To map the structure of the supercluster, and determine its extent in redshift space and across the plane of the sky, we take advantage of our redshift survey which allows us to demarcate the supercluster in redshift space as lying within the recession velocities 11 300–17 000 km s^{-1} (see Fig. 7). These cuts select a supercluster sample of 2281 galaxies across the ShaSS area.

Each galaxy was weighted (w_i) according to the inverse probability of it having been observed spectroscopically. First, each galaxy which could have been targeted for spectroscopy ($W1 < 15.0$, $i < 18.0$) or was a spectroscopic member of the supercluster, was given an initial equal weight of 1.0. For each of these galaxies lacking a redshift, its weight was transferred equally to its 10 nearest neighbours with known redshift that also had the same priority level in our AAOmega spectroscopic survey. This results in galaxies without redshifts having zero weight, while galaxies in regions where the spectroscopic survey is locally 50 per cent complete having weights of 2.0. The transferring of weight only within priority levels ensures that we can account statistically for the systematic differences in spectroscopic completeness from one level to another, as well as mapping the local spatial variations in completeness.

Table 5. Typical values for the absolute photometric calibration.

Night	Band	ZP	Colour term	Extinction
2013-02-13	<i>u</i>	23.261 ± 0.028	$0.026 \pm 0.019 (u - g)$	0.538
2013-03-16	<i>g</i>	24.843 ± 0.006	$0.024 \pm 0.006 (g - i)$	0.18
2012-04-29	<i>r</i>	24.608 ± 0.007	$0.045 \pm 0.019 (r - i)$	0.1
2012-02-27	<i>i</i>	24.089 ± 0.01	$-0.003 \pm 0.008 (g - i)$	0.043

**Figure 7.** Redshift distribution of galaxies in the ShaSS around the SSC. Grey vertical lines indicate the redshift range adopted here for the SSC: 11 300–17 000 km s⁻¹.

Each galaxy j belonging to the supercluster is represented by a Gaussian kernel whose transverse width is iteratively set to $\sigma_0[\rho_j(\mathbf{x}, z)/\bar{\rho}]^{-1/2}$, where $\bar{\rho}$ is the geometric mean of the ρ_j , and a fixed radial width of 700 km s^{-1} . Each galaxy is normalized by the weight parameter w_j to account for spectroscopic incompleteness. The transverse kernel width for each galaxy is initially set to $\sigma_0 = 6$ arcmin, and then iteratively adjusted to account for variations in the spatial density of galaxies, such that it typically encloses the 5–10 nearest neighbours of the galaxy, irrespective of its location within the supercluster.

Fig. 8 shows the resulting density map in which each galaxy is further weighted by its W1 flux as a proxy for its stellar mass (e.g. Jarrett et al. 2013; McGaugh & Schombert 2014; Meidt et al. 2014). The density map draws our attention to several features of the whole structure.

(a) Across the SSCC, the density of supercluster galaxies is always higher than 315 mJy deg^{-2} . This corresponds to $\sim 7 \times 10^{12} M_{\odot} \text{ deg}^{-2}$ based on the W1-stellar mass calibration obtained for the SSC galaxies (see end of Section 4.2). Assuming a depth of 5000 km s^{-1} or 71 Mpc the stellar mass density is $9.3 \times 10^9 M_{\odot} \text{ Mpc}^{-3}$. Both the stellar mass estimates include a 12 percent correction to account for the low-mass galaxies with $M < 10^{9.8} M_{\odot}$ or $W1 > 15$ not targeted in our spectroscopic survey, based on the SMF of Baldry et al. (2012). This is a mean overdensity of $\sim 40\times$ with respect to the cosmic total stellar mass density for galaxies in the local Universe ($z < 0.06$) estimated by Baldry et al. (2012) from the GAMA survey (Driver et al. 2011). We obtain a similar overdensity of 35–40 \times if we simply compare the observed W1 flux density of SSCC galaxies with the average W1 flux density of $W1 < 15$ galaxies within the same redshift range selected from the SDSS DR7 (8032 deg²). Although the SDSS is r -band selected, it should be complete to $W1 = 14.5$ at $z = 0.048$, and only marginally incomplete to $W1 = 15.0$. The higher density peaks in the SSCC, corresponding to the cluster cores but also to groups and cluster substructures, are interconnected forming a single system. This implies that the galaxy distribution follows the

hot gas distribution observed across the whole SSCC (e.g. Kull & Böhringer 1999).

(b) There is clear evidence of a filament (~ 7 Mpc in projection) heading north from the SSCC, connecting it with A3559, with W1 flux densities of $\gtrsim 150 \text{ mJy deg}^{-2}$ ($\gtrsim 20\times$ overdensities). Quintana et al. (2000) qualitatively indicated a ‘broad arm running north’ from the SSCC, but in their description this feature should point to A3557a which is located NW of the ShaSS region. In this direction we do not detect any clear overdensity.

(c) Several overdensities are detected across the density map. Some of them can be associated to cluster dynamical substructures already identified by Bardelli et al. (1998a, see Fig. 8), but other density peaks are also detected, i.e. that W of A3556. The positions of the system centres as defined by X-ray peak and galaxy overdensity appear slightly different, as expected in dynamically active systems. The highest discrepancy between these determinations is found for AS0724, where the X-ray centre is located 35 arcmin from the galaxy overdensity centre. Notice that AS0726 was not detected by *ROSAT* and its centre was derived by a previous spectroscopic survey (Proust et al. 2006), likely affected by incompleteness. Our newly detected overdensity is actually located 16 arcmin from the previous centre.

(d) In A3560, two main density clumps are sharply defined together with a western clump. The northern clump can be related to the X-ray emission, elongated towards A3558, detected in the very centre of the cluster which suggested a minor merger scenario between the main cluster and a group located ~ 8 arcmin N of the cluster centre (Venturi et al. 2013). Also the galaxy distribution shows an asymmetric shape elongated in the N direction, but extending a factor 2 in projection with respect to the X-ray emission (cf. fig. 5 of Venturi et al. 2013), taking into account only the two denser clumps in our map. Thus, the group involved in the claimed merger could be significantly more massive than that previously identified in the APM catalogue.

In a companion paper, we will assign galaxies to supercluster structures/overdensities. This will be achieved through the dynamical analysis enabling us to detect and measure the amount of cluster dynamical substructures, and to identify possible pre-merging clumps or merger remnants. In particular, we will (i) select cluster members (e.g. Fadda et al. 1996; Girardi et al. 1996); (ii) derive the centre of the most significant peaks of each identified system applying the 2D adaptive kernel technique to galaxy positions; (iii) detect possible subclumps and assign objects to groups (Ashman, Bird & Zepf 1994); (iv) identify the 3D substructures combining velocity and position information (Dressler & Shectman 1988).

7.2 Weak lensing mass distribution of A3558

We conducted a weak lensing analysis of the VST-ACCESS field 8 (see Fig. 2) following Kaiser, Squires & Broadhurst (1995) as modified by Okabe et al. (2013, 2014). We measure image ellipticity of objects detected in the r -band data, e_{α} , from the

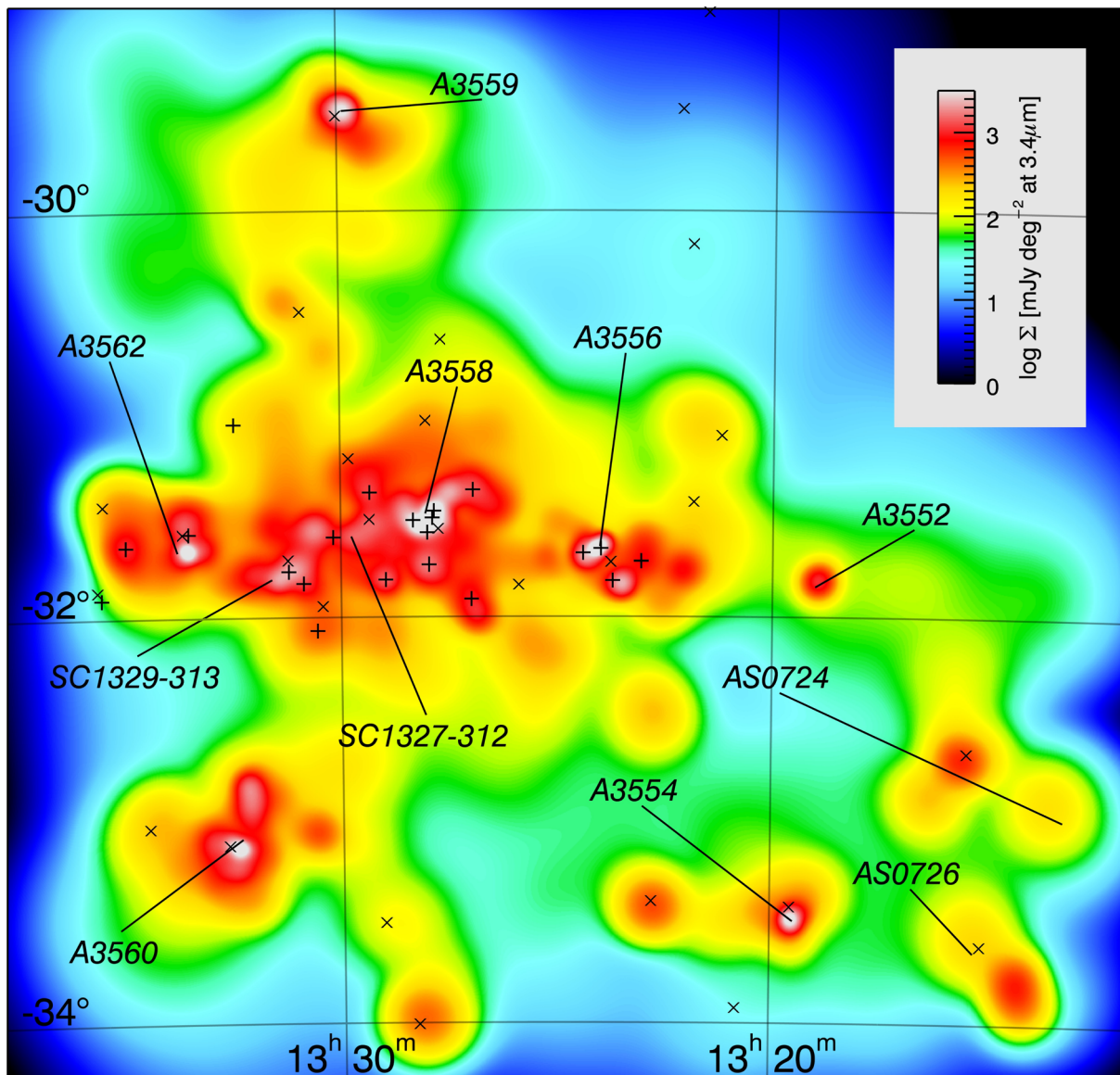


Figure 8. ShaSS density map in unit of mJy deg^{-2} at $3.4 \mu\text{m}$. Abell clusters and groups are labelled, the black straight lines pointing on the X-ray centre for all the systems except AS0726. Cluster substructures and groups identified by Bardelli et al. (1998a, +) and by Ragone et al. (2006, x) are shown in the map. The upper-right corner is not covered by ShaSS.

weighted quadrupole moments of the surface brightness. We then correct a PSF anisotropy of galaxy ellipticities by the function of second-order bi-polynomials of the stellar anisotropy kernel. A cross-correlation function of residual stellar ellipticities and the corrected galaxy ellipticities does not show overcorrection or insufficient correction as shown in Fig. 9. Then, a reduced distortion signal, g_α , is estimated with a correction of isotropic smearing effect. We select background galaxies following Okabe et al. (2013) and adopt a new method for minimizing the contamination of member galaxies extending to the colour–colour plane. We measure two colours, $g - i$ and $g - r$, and fit the red sequence with a linear function. A colour offset is defined for each galaxy by $\Delta C \equiv (\Delta(g - i)^2 + \Delta(g - r)^2)^{1/2}$, where $\Delta(g - i) \equiv (g - i) - (g - i)_{\text{ES0}}$ and $\Delta(g - r) \equiv (g - r) - (g - r)_{\text{ES0}}$ and ‘ES0’ denotes the red sequence galaxies. We select background galaxies by $\Delta C > 0.53$ which is the lower limit to allow 1 per cent contamination level. The

number density of background source galaxies is 2.2 arcmin^{-2} , and their mean redshift is $\langle z_s \rangle \simeq 0.49$.

The mass map is reconstructed as described in Okabe & Umetsu (2008). The reduced shear is pixelized into a regular grid with a Gaussian smoothing, $G(\theta) \propto \exp[-\theta^2/\theta_g^2]$. The FWHM resolution of the map is 11.7 arcmin . The smoothed shear pattern is estimated with a Gaussian kernel, $G(\theta)$, and a statistical weight $w_{g,i} = (\sigma_{g,i}^2 + \alpha^2)^{-1}$ for the i th galaxy, where σ_g is the rms error of the shear estimate and the softening constant variance. We choose $\alpha = 0.4$ as a typical value of the mean rms of σ_g . We next invert the smoothed shear field with the kernel (Kaiser & Squires 1993) in Fourier space to obtain the projected mass distribution. The resulting mass map, κ , for field 08 of VST-ACCESS is shown in Fig. 10. The mass map is elongated along the east–west direction which is parallel to the large-scale filamentary structure. A main peak is associated with distributions of member galaxies and diffuse X-ray

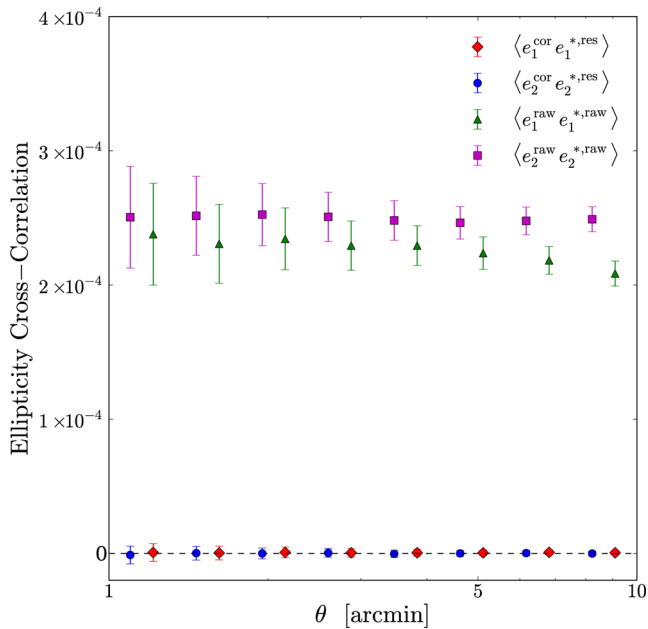


Figure 9. Cross-correlations for galaxy and stellar ellipticities, $\langle e_\alpha e_\alpha^* \rangle$. Red diamonds ($\alpha = 1$) and blue circles ($\alpha = 2$) denote the correlation function for corrected galaxy ellipticities, e_α^{cor} and residual stellar ellipticities, $e_\alpha^{*,\text{res}}$ after correcting the PSF anisotropy. Green triangles ($\alpha = 1$) and magenta squares ($\alpha = 2$) are the correlation function for raw galaxy ellipticities, e_α^{raw} and raw stellar ellipticities, $e_\alpha^{*,\text{raw}}$ before the correction. The separation radii are offset by ± 5 per cent.

emission of A3558. Another possible clump, located E of A3558, is likely associated with diffuse X-ray emission of the galaxy group SC1327–3136.

We compute the tangential distortion component, g_+ , with respect to the projected cluster-centric radius from the BCG of A3558, shown in Fig. 11. The lensing signal g_+ is decreasing as the radius, r_{cl} increases, and becomes flat or increases at $r_{\text{cl}} > 20$ arcmin. To understand this feature, we fit the two-dimensional shear pattern with two mass components: A3558 and the east clump (e.g. Okabe et al. 2011). We use the Navarro Frenk & White (NFW) profile (Navarro, Frenk & White 1996, 1997) as the mass model represented by two parameters: M_Δ which is the enclosed mass within a sphere of radius r_Δ , and the halo concentration $c_\Delta = r_\Delta/r_s$. Here, r_Δ is the radius inside of which the mean density is Δ times the critical mass density, $\rho_{\text{cr}}(z)$, at the redshift z . Since it is difficult to constrain the concentration for the east clump, we assume a mass–concentration relation (Bhattacharya et al. 2013). The centre of A3558 is fixed at the BCG position, while the position of the secondary mass component is treated as a free parameter. The model is described by five parameters in total. The best-fitting M_{500} for A3558 is $7.63^{+3.88}_{-3.40} \times 10^{14} M_\odot$. The total lensing signal of the two components (solid line in Fig. 11) well describes the observed lensing signals. Since A3558 and the east clump are embedded in the SSCC, surrounding data are essential to further constrain the mass.

Although this is a preliminary result, the derived dark matter mass can be compared with that evaluated from the *XMM-Newton* data set. To this aim, we computed the Y_X mass proxy, defined as the product of gas mass $M_{\text{gas}, 500}$ and average temperature kT (Kravtsov, Vikhlinin & Nagai 2006). As described in Bourdin & Mazzotta (2008), we inverted a gas mass profile from the

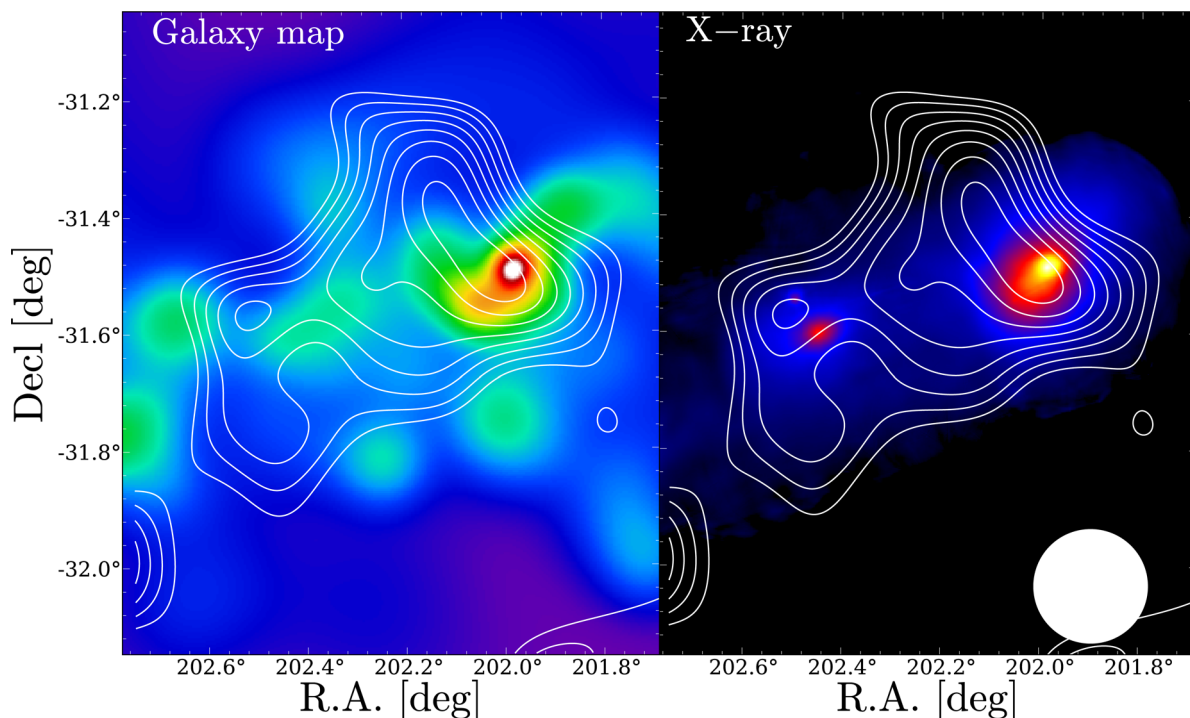


Figure 10. Weak-lensing mass map of the field 8 of VST-ACCESS overlaid on the stellar mass density map (left) and the X-ray surface brightness (right). The contours of the lensing κ -field are in units of 0.5σ reconstruction error $\delta\kappa \simeq 0.010$, above 1σ . The resolution of the WL mass map is $\text{FWHM}=11.7$ arcmin, as shown in the white circle of the bottom-right. The X-ray surface brightness is derived from a curvelet analysis of *XMM-Newton* images extracted in the 0.5–2.5 keV energy band, that have been corrected for spatially variable effective exposure and background components. The projected mass distribution is elongated along the large-scale X-ray filamentary structure. A main peak of the map is associated with the overdensity region of member galaxies and diffuse X-ray emission of A3558. A possible secondary peak is associated with diffuse X-ray emission of SC1327–312 in the east.

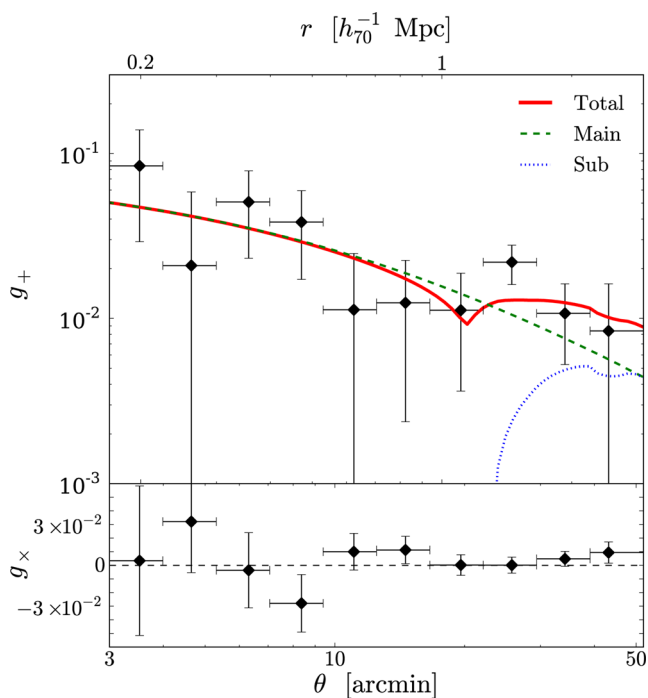


Figure 11. Top panel: the tangential distortion component, g_+ , as a function of the projected cluster-centric radius from the BCG, is estimated by azimuthally averaging the measured galaxy ellipticities. A bump in lensing signals is found around $r > 20$. The profile is well described by two NFW components of A3558 and the east clump. The red solid, green dashed and blue dotted lines are the best-fitting NFW profile for the total mass, A3558 and the east clump, respectively. Bottom panel: the 45 degree rotated component, g_x , is consistent with a null signal.

radially average surface brightness of A3558, then iterated about the $Y_X - M_{500}$ scaling relation calibrated from hydrostatic mass estimates in a nearby sample of clusters observed with *XMM-Newton* (Arnaud et al. 2010). Note that the surface brightness of A3558 has been evaluated within an angular sector excluding the eastern filament connecting A3558 to SC1327–312. This yielded estimates of gas mass $M_{\text{gas},500} = 0.62 \pm 0.01 \times 10^{14} M_{\odot}$, average temperature $kT = 4.91 \pm 0.14$ keV, and total mass $M_{500} = (4.62 \pm 0.24) \times 10^{14} M_{\odot}$ within $r_{500} = 1160 \pm 20$ kpc, which is consistent with the WL estimate. Previous X-ray mass estimates based on *ROSAT* data gave for M_{500} : $8.7 \times 10^{14} M_{\odot}$ (Ettori et al. 1997); $6.09 \times 10^{14} M_{\odot}$ (Reisenegger et al. 2000); $8.4 \times 10^{14} M_{\odot}$ (30 per cent uncertainty; Muñoz & Loeb 2008). The dynamical mass derived by Bardelli et al. (1998b) of A3558 and listed in Table 1 turns out to be a factor 1.7 higher than the WL mass, but also a factor 3 higher than X-ray mass determination previously available as the authors stated in their work. Ragone et al. (2006) computed the dynamical mass inside the virial radius obtaining $M_{\text{vir}} = 6.7 \times 10^{14} M_{\odot}$.

8 SUMMARY AND CONCLUSIONS

The Shapley Supercluster Survey aims to assess the role of cluster-scale mass assembly on galaxy evolution searching for possible connections between the properties of the cosmological structures (density, dynamical status, hot-gas content, dark and luminous matter distribution) and those of the associated galaxies (morphology, internal structure, SF, nuclear activity). This requires that we extend the investigation from the cluster cores to their outskirts, to

the infalling galaxies and groups along the filaments in a dynamically bound network. The centre (~ 3 deg radius) of the SSC at $z \sim 0.05$ is the optimal target to undertake such a study enclosing a massive and dynamically active structure showing signs of cluster-cluster mergers, enhancing the probability to observe evidence of environmental effects on galaxy evolution, but also providing an extraordinary variety of environments concentrated in a small survey volume.

ShaSS includes nine Abell clusters (A3552, A3554, A3556, A3558, A3559, A3560, A3562, AS0724, AS0726) and two poor clusters (SC1327–312, SC1329–313) covering a region of $\sim 260 \text{ Mpc}^2$. The survey includes the following data sets.

(i) Optical (*ugri*) imaging acquired with the VST (PI P. Merluzzi) provides a galaxy catalogue complete down to $r = 23.3$ (SNR ~ 10 , limit for star/galaxy separation as measured from the collected data) corresponding to $\sim m^* + 8.3$ at the supercluster redshift. The multiband *gri* catalogue is complete to $\sim m^* + 7.1$ and the *u*-band catalogue to $\sim m^* + 6.7$. The achieved SNRs allow to study the galaxy population down to $\sim m^* + 6$ and derive morphological parameters (CAS+MG, see Section 6) in *r* band to the same depth. The VST-ACCESS survey is ongoing.

(ii) NIR (*K*) imaging acquired with VISTA (PI C. P. Haines) reaching the depth of $K \sim 19.6$ ($\sim m^* + 7.9$ at the supercluster redshift) allows us to study the galaxy population down to the magnitude limit of the optical catalogue. The survey started in 2014 April.

(iii) The spectroscopic survey with AAOmega at the Anglo-Australian Telescope (PI P. Merluzzi) collected 4037 new redshifts across 21 deg^2 of ShaSS. Together with the already available redshifts, the spectroscopic sample is now 80 per cent complete down to $r \sim 18$ ($\sim m^* + 3$ at the supercluster redshift).

(iv) The above dedicated surveys are complemented by near/mid-infrared data from WISE in four bands (W1–W4) 3.4, 4.6, 12 and $22 \mu\text{m}$. Over the ShaSS area, we reach W1 = 16.96 (Vega magnitude) and W3 = 11.12 mag (1.0 mJy) at SNR = 10 corresponding to $\sim m^* + 5$ and SFR of $0.46 M_{\odot} \text{ yr}^{-1}$, respectively, for galaxies in the supercluster.

For the central $2\text{--}3 \text{ deg}^2$, *XMM-Newton*, *Spitzer*/MIPS 24 and $70 \mu\text{m}$, *GALEX* near-ultraviolet and FUV data are also available, as well as targeted observations of single galaxies providing $H\alpha$ imaging with the MMTF on the Magellan–Baade 6.5 m telescope at the Las Campanas Observatory in Chile and integral-field spectroscopy with WiFeS on the Australian National University 2.3 m telescope at Siding Spring in Australia.

In this first article, we derived the stellar mass density distribution based on supercluster members weighted by the W1 flux. This first quantitative characterization of the environment for the whole region covered by ShaSS shows a clumpy structure both in the SSCC and the surrounding clusters with several substructures, most of them already identified in previous works. All the clusters in the ShaSS area are embedded in a common network. This was suggested, but only for the SSCC where supercluster galaxies continuously populate the core following the gas filaments connecting the Abell clusters as mapped by X-ray observations (Kull & Böhringer 1999). We estimate the mean overdensity across the SSCC being $\sim 40\times$ with respect to the cosmic total stellar mass density for galaxies in the local Universe ($z < 0.06$).

Some new substructures with respect to previous works have been identified in the ShaSS density map such as those associated with A3560 N from the cluster centre in the direction of A3558, towards which also the X-ray emission is elongated (Venturi et al. 2013).

The most important new feature is however the filament connecting the SSC and the cluster A3559 as well as the less pronounced overdensity extending from the SSC towards A3560.

The other environment indicator analysed here is the dark matter distribution derived from the weak-lensing analysis of VST imaging. Using this approach, we studied the central 1 deg^2 field including A3558 and SC1327–312. The derived WL map shows that the dark matter is concentrated in two peaks which correspond to the two clusters, although the centres seem slightly offset with respect to the X-ray emission and the galaxy density. This can be due to the lower resolution of the WL map and/or ascribed to the complex dynamical state of the SSC. The estimated mass of A3558 is $M_{500} = 7.63^{+3.88}_{-3.40} \times 10^{14} M_{\odot}$, consistent with the X-ray estimate of $M_{500} = (4.16 \pm 0.19) \times 10^{14} M_{\odot}$. We notice that the WL mass determination will be improved extending the analysis to a larger region of the supercluster as planned in the ShaSS project.

We conclude pointing out that the VST imaging is the first CCD photometry covering homogeneously and continuously such a large portion of the SSC, with this depth and resolution (corresponding to 0.75 kpc at the supercluster redshift). Taking advantage of the *i* band and WISE W1 photometry, the AAOmega survey has been designed and carried out to obtain a magnitude-limited redshift sample, which was never achieved before (e.g. Quintana et al. 2000; Drinkwater et al. 2004; Proust et al. 2006). With these characteristics, ShaSS will build up the first multiband homogeneous data set of a vast region of the SSC and provide a fundamental local counterpart to the supercluster surveys at higher redshifts.

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