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| Authors | TRAFICANTE, Alessio, Fuller, G. A., Duarte-Cabral, A., ELIA, Davide Quintino, Heyer, M. H., MOLINARI, Sergio, Peretto, N., SCHISANO, EUGENIO |
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Multiscale dynamics in star-forming regions: the interplay between gravity and turbulence

A. Traficante^{1,2*}, G. A. Fuller,² A. Duarte-Cabral³, D. Elia¹, M. H. Heyer,⁴
S. Molinari¹, N. Peretto³ and E. Schisano¹

¹IAPS - INAF, via Fosso del Cavaliere, 100, I-00133 Roma, Italy

²Jodrell Bank Centre for Astrophysics, Department of Physics and Astronomy, The University of Manchester, Oxford Road, Manchester M13 9PL, UK

³School of Physics and Astronomy, Cardiff University, Queens Buildings, The Parade, Cardiff CF24 3AA, UK

⁴Department of Astronomy, University of Massachusetts, Amherst, MA 01003, USA

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ABSTRACT

In this work, we investigate the interplay between gravity and turbulence at different spatial scales and in different density regimes. We analyse a sample of 70- μm quiet clumps that are divided into three surface density bins, and we compare the dynamics of each group with the dynamics of their respective filaments. The densest clumps form within the densest filaments, on average, and they have the highest value of the velocity dispersion. The kinetic energy is transferred from the filaments down to the clumps most likely through a turbulent cascade, but we identify a critical value of the surface density, $\Sigma \simeq 0.1 \text{ g cm}^{-2}$, above which the dynamics change from being mostly turbulent-driven to mostly gravity-driven. The scenario we obtain from our data is a continuous interplay between turbulence and gravity, where the former creates structures at all scales and the latter takes the lead when the critical surface density threshold is reached. In the densest filaments, this transition can occur at the parsec, or even larger scales, leading to a global collapse of the whole region and most likely to the formation of the massive objects.

Key words: stars: formation – stars: kinematics and dynamics – stars: massive – stars: statistics – infrared: stars – radio lines: stars.

1 INTRODUCTION

The massive star formation involves a hierarchical, multiscale process that starts in giant molecular clouds (GMCs), objects with size up to tens of parsecs (Solomon et al. 1987; Roman-Duval et al. 2010). Within GMCs, the formation of stars begins preferably across elongated features called filaments, which are ubiquitous in the Galaxy (Molinari et al. 2010; André et al. 2014; Schisano et al. 2014, 2019; Li et al. 2016). Along these filaments, parsec-scale massive and dense clumps develop an average surface density larger than 0.05–0.1 g cm^{-2} (Urquhart et al. 2014; Traficante et al. 2015b; Svoboda et al. 2016), which leads to the formation of the first protostellar cores (Zhang et al. 2015; Csengeri et al. 2017), which will finally produce a cluster of stars (e.g. Lada & Lada 2003), including the most massive ones.

These different structures are identified at the various spatial scales, but it is not yet clearly understood whether (and eventually how) they are dynamically coupled. Turbulent accretion models (e.g. McKee & Tan 2003) predict that the collapse occurs *locally*

and in a similar fashion to what happens in low-mass regions: Massive clumps are relatively isolated from the rest of the clouds. They are maintained close to the virial equilibrium, thanks to the high pressure sustained by the supersonic non-thermal motions driven by the turbulence of the local interstellar medium. Other models predict that the whole cloud is in hierarchical, *global* collapse and the supersonic, non-thermal velocity dispersions reflect infalling motions at all spatial scales (Vázquez-Semadeni et al. 2009; Ballesteros-Paredes et al. 2011; Vázquez-Semadeni et al. 2019).

Observations of massive, filamentary infrared dark clouds (IRDCs) support the global collapse mechanism (e.g. Peretto et al. 2013), with evidence that large-scale motions driven by self-gravity occur simultaneously and are coupled to the local motions, with a net increase in supersonic motions along the embedded clumps due to the longitudinal collapse of the whole cloud, like in SDC 13 (Peretto et al. 2014), or along the junctions of filaments, where dense clumps are forming, like in Monoceros OB1 (Montillaud et al. 2019). The gravitational fragmentation along a filament may also be regulated by the turbulence that acts at all spatial scales, as suggested from the dynamics observed in G26 (Liu et al. 2018). In contrast, a cloud such as G035 shows that the dynamics of the

* E-mail: alessio.traficante@inaf.it

filaments and the embedded clumps are different and they may be considered as independent structures, with non-thermal motions that are only mildly supersonic (Henshaw et al. 2014).

These few examples illustrate the need to distinguish between non-thermal motions driven by gravity or regulated by turbulence in massive star-forming regions (Krumholz, McKee & Bland-Hawthorn 2018; Motte, Bontemps & Louvet 2018), as well as the need to determine whether the dynamics of the collapse in the high-mass regime differ or not from the low-mass counterpart.

With large surveys of the Galactic plane that observed the cold dust continuum emission of thousands of star-forming regions [e.g. Hi-GAL, Molinari et al. 2010; ATLASGAL, Schuller et al. 2009; BGPS (Bolocam Galactic Plane Survey), Aguirre et al. 2011] and the relatively low density gas associated with molecular clouds and filaments, traced with the various CO isotopologues [e.g. ^{13}CO (1–0), GRS, Jackson et al. 2006; ^{12}CO (1–0) and ^{13}CO (1–0), FUGIN, Umemoto et al. 2017; ^{13}CO (2–1) and C^{18}O (2–1), SEDIGISM, Schuller et al. 2017; ^{13}CO (3–2) and C^{18}O (3–2), CHIMPS, Rigby et al. 2019], we now have a rich database to statistically analyse the dynamics of star-forming regions in different environments. These catalogues, combined with observations of high-density tracers to identify the kinematics within the embedded clumps, enable a multiscale study of the kinematics and dynamics in different mass and density regimes.

In this work, we present new data obtained with the IRAM 30-m telescope relative to 13 clumps selected from the Hi-GAL sample to be 70- μm quiet and in a range of low- to intermediate-surface densities, together with 16 high-mass 70- μm quiet clumps already studied in Traficante et al. (2017). We combine these data with the properties of their parent filaments extracted from the newly published catalogue of filaments extracted from the Hi-GAL survey (Schisano et al. 2019), for which we derived the dynamics using FUGIN and GRS data. We aim at studying the physical and kinematic properties of star-forming regions in different density regimes and at various spatial scales to investigate whether there is a critical density regime above which gravity starts to dominate over turbulence, and at which scales the collapse can be considered global. This paper is organized as follows: In Section 2, we describe the data sets we use to determine the properties of the clumps and parent filaments; in Section 3, we divide the clumps in three different density regimes, and we discuss the physical properties and the dynamics of the clumps and the parent filaments in these different regimes; in Section 4, we combine the dynamics of the clumps and the parent filaments, for a multiscale analysis of these star-forming regions; in Section 5, we discuss the scenario that emerges from our observations, and, finally, in Section 6, we present our summary.

2 OBSERVATIONS

2.1 Clump data

The clumps used in this work combine two samples of so-called 70- μm quiet clumps, due to the lack of (or extremely faint) emission at this wavelength, which is attributed to the lack of the presence of an evolved protostar (Dunham et al. 2008). We focus on this extremely young stage of the evolution to assure that any feedback from the presence of an embedded protostar is minimal. The first sample is a group of 16 massive and dense ($\Sigma \geq 0.05 \text{ g cm}^{-2}$) clumps presented in Traficante et al. (2017), and it is combined with a new sample of 23 70- μm quiet clumps extracted from the catalogue of candidate starless sources embedded in IRDCs (Traficante et al. 2015b) and selected to cover a lower range of surface densities,

with $0.005 \leq \Sigma \leq 0.05 \text{ g cm}^{-2}$. The physical properties of the clumps have been rederived from the *Herschel* Hi-GAL survey data (Molinari et al. 2010) and also including the emission at 870 μm and 1.1 mm from the ATLASGAL (Schuller et al. 2009) and the BGPS (Aguirre et al. 2011, Ginsburg et al. 2013) respectively.

The gas kinematics of each clump have been obtained with 3-mm observations centred on the N_2H^+ (1–0), HCO^+ (1–0), and HNC (1–0) lines, acquired at the IRAM 30-m telescope in 2015 December under the project 133-15. The sources were observed in the on-the-fly mode covering a 2×2 -arcmin² region centred on the peak of the clump, identified in the 250- μm Hi-GAL map. The spatial resolution is $\simeq 28$ arcsec. The system temperature of the observations was in the range $85 \leq T_{\text{sys}} \leq 125$ K, depending on the source. We used the FTS50 backend to obtain a spectral resolution of 0.2 km s^{-1} , enough to resolve the hyperfine components of the N_2H^+ (1–0) emission line. We achieved a noise per channel per pixel of $T_{\text{mb}} \leq 0.16$ K, assuming a beam efficiency at our frequencies of 0.78 (Rygl et al. 2013). We averaged the emission across each clump and considered only clumps with a signal-to-noise ratio (S/N) of the N_2H^+ (1–0) average spectra above 5σ , where the noise level in each clump is determined as the standard deviation of the emission-free channels using the GILDAS/CLASS¹ software package.

Four sources out of 23 have no detected emission in the N_2H^+ (1–0) at the sensitivity of our observations, while three more sources have an $S/N < 5$ and are therefore removed from the sample studied here. We also removed from the analysis the clumps 18.287–0.256, 33.332–0.531, and 49.57–0.192, since the N_2H^+ (1–0) spectrum showed two components along the LOS that cannot be disentangled in the dust continuum data.

We ended up with a total of 13 out of 23 sources with defined physical and kinematic properties, which, combined with the sample of 16 massive 70- μm quiet clumps of Traficante et al. (2017), provide us with a total of 29 sources covering a range of surface densities $0.005 \leq \Sigma \leq 0.25 \text{ g cm}^{-2}$. Examples of the new 70- μm quiet sources 23.076–0.209 as seen at 70 and 250 μm and of its N_2H^+ (1–0) spectrum are shown in Fig. 1.

2.2 Filament data

For each clump, we identified the parent filament from the catalogue of filaments in the Galactic plane extracted from the Hi-GAL survey (Schisano et al. 2019) through the algorithm described in Schisano et al. (2014). This algorithm extracts all the elongated structures in the Hi-GAL column density maps that have a higher contrast with respect to the surroundings and excludes all structures with the main spine smaller than $\simeq 2$ arcmin, which are automatically filtered out from the catalogue (Schisano et al. 2019). The final catalogue contains more than 30 000 filaments extracted from the whole Galactic plane.

A strong correlation between star-forming clumps and filamentary structures has been already observed, at least in nearby regions (Schneider et al. 2013; Könyves et al. 2015). In fact, only four out of our 29 clumps do not have an association with a Hi-GAL filament in the Schisano et al. (2019) catalogue. Three of these missing counterparts (20.717–0.01, 24.552+0.096, 25.254–0.166) are associated with IRDCs smaller than 2 arcmin, while IRDC 32.006–0.51 is fragmented in the Hi-GAL column density image, and this is likely the reason why it was not identified in the Schisano

¹<http://www.iram.fr/IRAMFR/GILDAS>

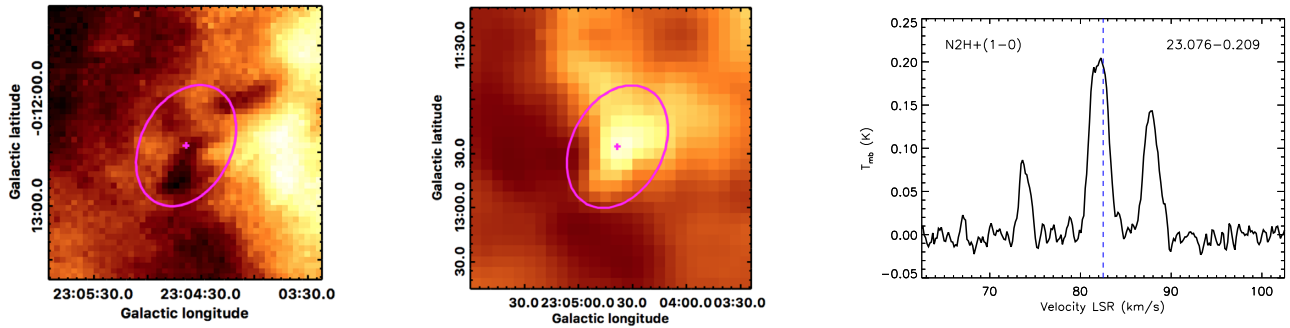


Figure 1. The 70- μm quiet clump 23.076–0.209, belonging to the Σ_{low} group, as seen in Hi-GAL at 70 μm (left-hand panel) and 250 μm (central panel). The magenta cross represents the 250- μm centroid and the magenta ellipse is the result of the 2D-Gaussian fitting performed with HYPER. In the right-hand panel, the N_2H^+ (1–0) emission averaged across the clump. The blue-dotted vertical line is in correspondence of the central velocity of the clump obtained with a hyperfine fitting of the N_2H^+ spectrum done with CLASS using the standard routines.

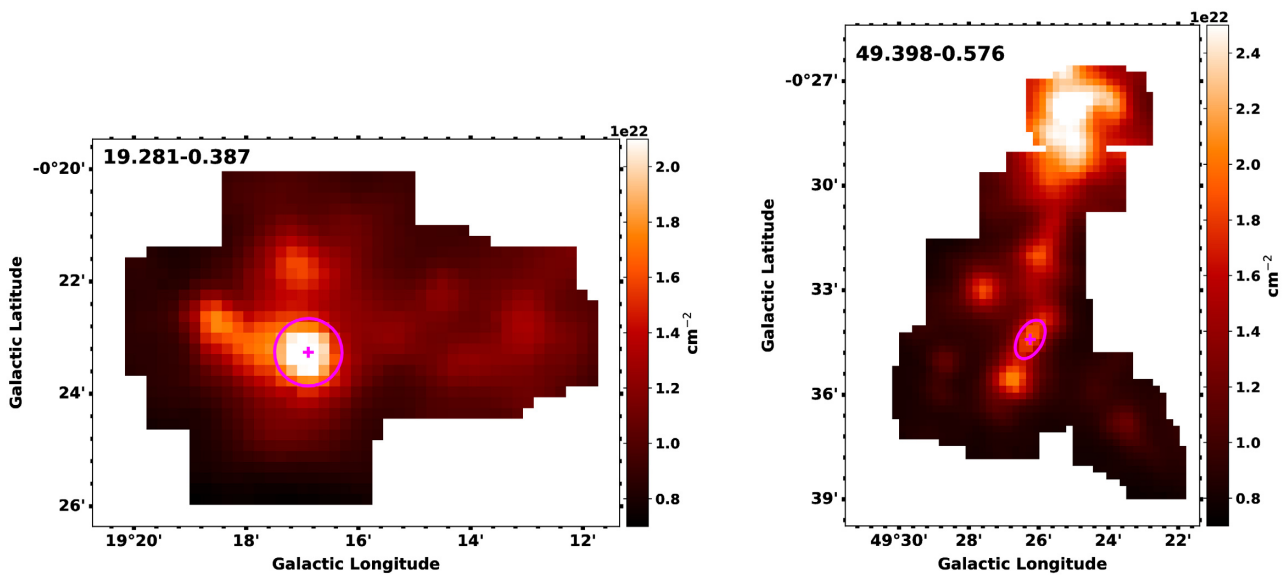


Figure 2. Two examples of filaments extracted from the Hi-GAL column density maps and presented in the Schisano et al. (2019) catalogue, associated with the sources 19.281–0.387 (left-hand image) and 49.398–0.576 (right-hand image), respectively. The magenta crosses and ellipses are in correspondence of the 70- μm quiet clumps associated with these two filaments.

et al. (2019) catalogue. This ratio is in agreement with the findings of Schisano et al. (2019), where the whole Hi-GAL catalogue of filaments was combined with the catalogue of $\simeq 15\,000$ IRDCs from Peretto & Fuller (2009), it was found that only $\simeq 28$ per cent of the IRDCs do not have a counterpart in the Hi-GAL filaments, and in most cases, it was because of the filtering of features smaller than $\simeq 2$ arcmin.

We used the filament mask to extract the ^{13}CO (1–0) data for each cloud from the GRS survey (Jackson et al. 2006). This survey observed a portion of the inner Galactic plane ($15^\circ \leq l \leq 55^\circ$, $-1^\circ \leq b \leq 1^\circ$) in the ^{13}CO (1–0) emission with a spatial resolution of 46 arcsec and a spectral resolution of 0.212 km s^{-1} . The sensitivity is $\sigma_{T_{\text{mb}}} = 0.26 \text{ K}$, considering a main beam efficiency of 0.48 (Roman-Duval et al. 2010).

From the GRS counterparts, we excluded the filament associated with the source 28.905–0.534. The central velocity of the main CO emission is shifted $\simeq 30 \text{ km s}^{-1}$ with respect to the central velocity of the clump determined from the N_2H^+ (1–0) emission, and the CO emission around the clump velocity position is below the 4σ

level used as a lower limit for the molecular emission, as discussed in Section 3.

The sample of filaments and clumps contains the physical parameters and the kinematics of 24 out of the 29 sources. Two examples of the filaments extracted from the Hi-GAL column density maps and the embedded clumps for which we have kinematic information are shown in Fig. 2.

To estimate the properties of the filaments such as optical depth and mass from the ^{13}CO data, we need to evaluate the excitation temperature T_{ex} . The excitation temperature is evaluated under the local thermal equilibrium assumption combining the ^{12}CO and ^{13}CO data taken from the FUGIN survey (Umemoto et al. 2017). This survey has been carried out with the Nobeyama 45-m telescope and observed the ^{12}CO , ^{13}CO , and C^{18}O (1–0) transitions simultaneously across 156 deg^2 of the inner Galactic plane in a 2-deg-wide longitude strip around $b = 0^\circ$. Our clumps and filaments are all in the first quadrant, where the FUGIN survey covers the longitudes $10^\circ \leq l \leq 50^\circ$. The noise is 0.24 K per channel and 0.12 K per channel for the ^{12}CO and the ^{13}CO cubes, respectively.

The FUGIN survey has a better spatial resolution than the GRS survey, $\simeq 20$ arcsec; however, it has a lower spectral resolution, 0.65 km s^{-1} .

We therefore used both the ^{12}CO and ^{13}CO FUGIN cubes to evaluate the excitation temperature, as discussed in Section 3, but we take advantage of the finer spectral resolution of the GRS data to estimate the physical parameters and the velocity dispersion of the clouds.

3 RESULTS

The goal of this work is to investigate whether the dynamics of the massive and dense clumps that can lead to the formation of high-mass stars are peculiar or not, compared to the lower density counterparts. The surface density threshold discussed in the literature above which a clump will likely form high-mass objects varies from $\Sigma_{\text{thres}} = 0.05$ (Urquhart et al. 2014) to 0.1 g cm^{-2} (Tan et al. 2014; Traficante et al. 2018b), and up to $\Sigma_{\text{thres}} = 1 \text{ g cm}^{-2}$ if we consider the radiative feedback needed to prevent the fragmentation and the formation of the massive objects in some models (Krumholz & McKee 2008; Federrath, Krumholz & Hopkins 2017; Guszejnov et al. 2018). Based on these numbers, we divide our sample into the following three different groups:

Σ_{low} : These are clumps with a (relatively) low surface density, below the minimum threshold assumed for high-mass star-formation, $\Sigma \leq 0.05 \text{ g cm}^{-2}$.

Σ_{int} : These are clumps with a higher surface density than the ones in the Σ_{l} group, but still below the 0.1 g cm^{-2} critical value. These are defined as intermediate-density clumps.

Σ_{high} : These are clumps with the surface density above the surface density threshold of $\Sigma = 0.1 \text{ g cm}^{-2}$.

We aim to investigate any possible difference between regions in these surface density regimes, in particular between clumps and filaments in the Σ_{low} and Σ_{high} groups, with a somehow intermediate behaviour for the regions in the Σ_{int} group.

3.1 Clumps analysis

The physical properties of each of the new 13 clumps have been rederived with the same approach as used for the massive $70\text{-}\mu\text{m}$ quiet clumps in Traficante et al. (2017) for consistency. Therefore, we performed a dedicated run to extract the fluxes at all Hi-GAL wavelengths with the HYPER routine (Traficante et al. 2015a), including also the emission at $870 \mu\text{m}$ from ATLASGAL and at $1100 \mu\text{m}$ from BGPS. HYPER performs a 2D-Gaussian fit at $250 \mu\text{m}$ of each source, which is used to determine the two semi-axes (the two full width at half-maximum of the fit) and the position angle of the ellipse. These parameters define the radius R of the clump and the region of integration of the source fluxes at all wavelengths.

The physical parameters of each clump are determined with a single-temperature grey-body fit of the spectral energy distribution (SED) using all wavelengths in the range $160 \leq \lambda \leq 1100 \mu\text{m}$. If the observations in the ATLASGAL and BGPS surveys are missing, we perform the SED fitting only on the Hi-GAL wavelengths ($160 \leq \lambda \leq 500 \mu\text{m}$). This is the case for six sources. The grey-body fit is done assuming a fixed spectral index $\beta = 2.0$ and with the uncertainties associated with each flux equal to 20 per cent of the integrated fluxes in the $70\text{--}500 \mu\text{m}$ wavelength range and to 40 per cent of the integrated fluxes for the 870 and $1100 \mu\text{m}$ values, to also include the uncertainties due to the combination of different surveys with the Hi-GAL fluxes (Traficante et al. 2017). The masses

of these clumps are in the range $25 \leq M \leq 531 M_{\odot}$. We obtain surface densities $\Sigma = M/(\pi R^2)$ in the range $0.006 \leq \Sigma \leq 0.07 \text{ g cm}^{-2}$. The uncertainties associated with each measurement are derived as given in Traficante et al. (2018a), which made an extensive analysis to evaluate the source of uncertainties for each parameter. A summary of the source parameters for all the 29 clumps used in this work is given in Table 1. The clumps have been divided into three groups based on the value of their surface density. We obtain 10, 8, and 11 clumps, respectively, corresponding to the Σ_{low} , Σ_{int} , and Σ_{high} groups.

In order to study the dynamics of each clump, we need to estimate their velocity dispersion. In agreement with the analysis done for the 16 clumps in Traficante et al. (2017), we first determined the N_2H^+ (1–0) spectrum averaged in all pixels within the ellipse defined by the 2D-Gaussian fit of each source, then we estimated the velocity dispersion from a hyperfine fitting of the average spectrum performed with standard CLASS routines. As shown in Table 1, all these clumps are supersonic, with $\sigma \geq 0.39 \text{ km s}^{-1}$ (the H_2 thermal velocity at temperature $T = 10 \text{ K}$ is $\simeq 0.2 \text{ km s}^{-1}$), including two clumps that are mildly supersonic ($25.254\text{--}0.166$, $\sigma = 0.29 \text{ km s}^{-1}$ and $49.398\text{--}0.576$, $\sigma = 0.34 \text{ km s}^{-1}$). There are no evident differences between the three groups, despite the difference in mass between clumps can be up to one order of magnitude.

3.1.1 The dynamics of star-forming clumps

In this section, we study the dynamics of our sample of clumps, and we divide and compare the results between the three different groups.

We start with a useful parameter to investigate the energy budget of a clump, the virial parameter α_{vir} , which describes the ratio between the gravitational and the kinetic energy of a defined region. Here, we evaluate α_{vir} following the Bertoldi & McKee (1992) definition of $\alpha_{\text{vir}} = 2aE_{\text{k}}/E_{\text{G}} = a5\sigma^2R/(GM)$, where E_{k} and E_{G} are the kinetic and gravitational energy, respectively, M and R are the mass and radius of the clump, respectively, G is the gravitational constant, and a is a constant that includes modifications due to non-spherical and inhomogeneous density distributions. We assume for simplicity $a = 1$, although accounting for a more realistic mass and kinematic distribution of the gas in the clumps and their surroundings can lead to different values of a and strongly affect the values of E_{k} and E_{G} (Federrath & Klessen 2012; Beaumont et al. 2013).

In Fig. 3, we show the histogram of the virial parameter divided into three groups. Blue bins are used for the Σ_{low} sources, green bins for the Σ_{int} sources, and red bins for the Σ_{high} sources. The dotted lines are the median values for each group, which are $\alpha_{\text{vir}} = [1.51, 0.63, 0.54]$ for the [Σ_{low} , Σ_{int} , Σ_{high}] clumps, respectively. The low-density clumps have a wider distribution of the virial parameter, with the large majority of them with values $\alpha_{\text{vir}} > 1$ and with few objects (four out of 10) that have $\alpha_{\text{vir}} > 2$. A value of the virial parameter larger than 2 suggests that these regions are gravitationally unbound (Kauffmann, Pillai & Goldsmith 2013; Tan et al. 2014) and the gravity alone cannot overcome the strong non-thermal motions fed by the turbulence in the interstellar medium, although this value may be overestimated as a result of an underestimation of the gravitational energy budget of a given region (Ballesteros-Paredes et al. 2018). On the other hand, supercritical values of the virial parameter, i.e. $\alpha_{\text{vir}} \ll 1$ are considered as evidence that the clumps are significantly gravitationally bound and therefore prone to collapse (Urquhart et al. 2018), but also in this case the value of α_{vir} obtained from the observations of apparent supervirial clumps may be misleading due to a bias in the estimation of the local non-

Table 1. Physical and kinematic parameters of the 29 clumps used in this work, combining the 13 clumps observed with the IRAM 30-m telescope under the project 133-15 and the 16 objects presented in Traficante et al. (2017, 2018b). The sources are divided into three groups, as described in Section 3, and ordered by increasing values of their surface density within each group. Column 1: clump name; Columns 2–4: clump mass, radius, and surface density derived from the HYPER fitting at 250 μm ; Column 5: velocity dispersion measured from the hyperfine fitting of the N_2H^+ (1–0) spectrum; Column 6: virial parameter; Column 7: distance of the source obtained from the Traficante et al. (2015b) catalogue; Columns 8–9: equatorial coordinates of the 250- μm centroid obtained from the Gaussian fitting; and Column 10: group based on the clump surface density, as discussed in Section 3.

| Clump | M_{clump} (M_{\odot}) | R_{cl} (pc) | Σ_{cl} (g cm^{-2}) | σ_{clump} (km s^{-1}) | α_{vir} | Distance (kpc) | RA ($^{\circ}$) | Dec. ($^{\circ}$) | Group |
|---------------------------|---------------------------------------|-------------------------|--|---|-----------------------|-------------------|----------------------|------------------------|------------------------|
| 28.905–0.534 | 25 (7) | 0.54 (0.14) | 0.006 (0.002) | 0.45 (0.13) | 4.85 (3.15) | 4.34 | 18:46:04.3 | – 3:47:24 | Σ_{low} |
| 24.552+0.096 | 47 (13) | 0.55 (0.14) | 0.011 (0.004) | 0.64 (0.19) | 5.42 (3.52) | 3.44 | 18:35:42.2 | – 7:24:21 | Σ_{low} |
| 25.254–0.166 | 59 (16) | 0.59 (0.15) | 0.011 (0.004) | 0.29 (0.09) | 0.94 (0.61) | 3.89 | 18:37:57.6 | – 6:54:14 | Σ_{low} |
| 49.732–0.471 | 110 (30) | 0.79 (0.20) | 0.012 (0.004) | 0.43 (0.13) | 1.51 (0.98) | 5.43 | 19:24:33.1 | – 14:41:59 | Σ_{low} |
| 20.717–0.01 | 102 (30) | 0.50 (0.12) | 0.027 (0.010) | 0.97 (0.29) | 5.33 (3.46) | 4.05 | 18:28:57.6 | – 10:51:32 | Σ_{low} |
| 28.19–0.192 | 219 (63) | 0.71 (0.18) | 0.029 (0.010) | 0.59 (0.18) | 1.30 (0.85) | 4.47 | 18:43:29.0 | – 4:19:15 | Σ_{low} |
| 53.784–0.138 | 102 (35) | 0.48 (0.12) | 0.030 (0.010) | 0.54 (0.16) | 1.57 (1.02) | 4.96 | 19:31:27.1 | 18:25:40 | Σ_{low} |
| 45.531+0.042 | 212 (62) | 0.67 (0.17) | 0.031 (0.011) | 0.55 (0.16) | 1.10 (0.71) | 4.51 | 19:14:33.1 | 11:12:50 | Σ_{low} |
| 23.076–0.209 | 159 (44) | 0.55 (0.14) | 0.035 (0.012) | 0.58 (0.17) | 1.34 (0.87) | 3.72 | 18:34:04.1 | – 8:51:25 | Σ_{low} |
| 26.432–0.662 | 221 (86) | 0.64 (0.16) | 0.036 (0.013) | 0.81 (0.24) | 2.20 (1.43) | 3.92 | 18:41:40.3 | – 6:04:55 | Σ_{low} |
| 35.608+0.111 | 176 (49) | 0.50 (0.12) | 0.047 (0.016) | 0.53 (0.16) | 0.92 (0.60) | 3.16 | 18:55:55.7 | 2:26:02 | Σ_{low} |
| 30.357–0.837 ^a | 371 (110) | 0.67 (0.17) | 0.055 (0.019) | 0.57 (0.17) | 0.68 (0.44) | 4.30 | 18:49:40.6 | – 2:39:46 | Σ_{int} |
| 53.361+0.042 | 36 (9) | 0.21 (0.05) | 0.055 (0.019) | 0.39 (0.12) | 0.99 (0.65) | 2.04 | 19:29:48.2 | 18:08:02 | Σ_{int} |
| 30.131–0.644 | 379 (101) | 0.66 (0.17) | 0.058 (0.020) | 0.56 (0.17) | 0.63 (0.41) | 4.59 | 18:48:38.2 | – 2:47:06 | Σ_{int} |
| 32.006–0.51 ^a | 448 (127) | 0.70 (0.17) | 0.061 (0.021) | 0.31 (0.09) | 0.18 (0.12) | 4.24 | 18:51:34.1 | – 1:03:25 | Σ_{int} |
| 15.631–0.377 ^a | 268 (78) | 0.54 (0.14) | 0.061 (0.021) | 0.30 (0.09) | 0.21 (0.14) | 3.47 | 18:20:29.0 | – 15:31:26 | Σ_{int} |
| 49.398–0.576 | 531 (176) | 0.74 (0.18) | 0.065 (0.023) | 0.34 (0.10) | 0.18 (0.12) | 5.47 | 19:24:18.5 | 14:22:30 | Σ_{int} |
| 28.792+0.141 ^a | 449 (128) | 0.61 (0.15) | 0.080 (0.028) | 0.99 (0.30) | 1.55 (1.01) | 4.62 | 18:43:08.9 | – 3:36:17 | Σ_{int} |
| 49.433–0.22 | 385 (108) | 0.53 (0.13) | 0.091 (0.032) | 0.82 (0.25) | 1.07 (0.70) | 4.32 | 19:22:52.6 | 14:30:43 | Σ_{int} |
| 25.982–0.056 ^a | 888 (274) | 0.80 (0.20) | 0.092 (0.032) | 0.69 (0.21) | 0.50 (0.32) | 5.00 | 18:38:54.5 | – 6:12:32 | Σ_{int} |
| 33.332–0.531 | 1266 (350) | 0.94 (0.23) | 0.095 (0.033) | 0.86 (0.26) | 0.64 (0.41) | 5.37 | 18:54:03.8 | 0:07:01 | Σ_{int} |
| 19.281–0.387 ^a | 700 (206) | 0.67 (0.17) | 0.104 (0.036) | 0.47 (0.14) | 0.25 (0.16) | 3.82 | 18:27:33.8 | – 12:18:17 | Σ_{high} |
| 34.131+0.075 ^a | 480 (141) | 0.55 (0.14) | 0.106 (0.037) | 0.74 (0.22) | 0.72 (0.47) | 3.56 | 18:53:21.6 | 1:06:14 | Σ_{high} |
| 23.271–0.263 ^a | 997 (296) | 0.72 (0.18) | 0.128 (0.045) | 0.94 (0.28) | 0.75 (0.49) | 5.21 | 18:34:37.9 | – 8:40:44 | Σ_{high} |
| 31.946+0.076 ^a | 1431 (439) | 0.82 (0.20) | 0.142 (0.050) | 1.19 (0.36) | 0.94 (0.61) | 5.51 | 18:49:22.3 | 0:**:31 | Σ_{high} |
| 22.756–0.284 ^a | 655 (193) | 0.55 (0.14) | 0.144 (0.050) | 0.95 (0.28) | 0.88 (0.57) | 4.43 | 18:33:49.0 | – 9:13:04 | Σ_{high} |
| 22.53–0.192 ^a | 1579 (487) | 0.80 (0.20) | 0.164 (0.057) | 1.25 (0.37) | 0.92 (0.60) | 5.77 | 18:32:59.8 | – 9:20:02 | Σ_{high} |
| 28.537–0.277 ^a | 1152 (326) | 0.67 (0.17) | 0.171 (0.060) | 0.78 (0.23) | 0.41 (0.27) | 4.96 | 18:44:22.1 | – 4:01:40 | Σ_{high} |
| 28.178–0.091 ^a | 2092 (610) | 0.85 (0.21) | 0.193 (0.067) | 1.07 (0.32) | 0.54 (0.35) | 5.35 | 18:43:02.6 | – 4:14:52 | Σ_{high} |
| 24.013+0.488 ^a | 1957 (559) | 0.81 (0.20) | 0.198 (0.069) | 0.91 (0.27) | 0.40 (0.26) | 5.18 | 18:33:18.5 | – 7:42:25 | Σ_{high} |
| 25.609+0.228 ^a | 3098 (953) | 0.97 (0.24) | 0.219 (0.077) | 1.05 (0.31) | 0.40 (0.26) | 5.57 | 18:37:10.6 | – 6:23:31 | Σ_{high} |
| 18.787–0.286 ^a | 1915 (549) | 0.69 (0.17) | 0.267 (0.094) | 1.07 (0.32) | 0.48 (0.31) | 4.36 | 18:26:15.4 | – 12:41:34 | Σ_{high} |

^aFrom Traficante et al. (2017, 2018b).

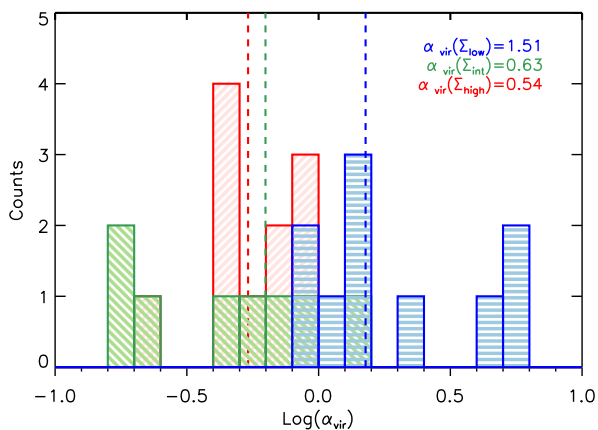


Figure 3. Histogram of the virial parameters of the 70- μm quiet clumps divided into the three groups, as discussed in Section 3.

thermal motions (Traficante et al. 2018c). These regions should indeed be naturally close to the virial value (Lee & Hennebelle 2016), and often appear virialized (Ballesteros-Paredes et al. 2011).

In general, the virial parameter gives a rough estimate of the energy balance of each region, and it is not possible to determine whether the observed non-thermal motions are dominated by local turbulence or by the gravitational collapse from solely its value. In the former scenario, we expect to observe a velocity dispersion-size relation, $\sigma \propto R^{\delta}$, which resembles the turbulent cascade of energy when $\delta = 0.5$ (Larson 1981; Heyer & Brunt 2004; McKee & Ostriker 2007; Federrath 2013). In Traficante et al. (2018a), we have already discussed how this relation breaks down at the clump scales. We draw similar conclusions from our sample of objects, as shown in Fig. 4, although our sample covers a limited range of radii. A correlation is possibly present in the high-density clumps (Pearson’s coefficient $\rho = 0.53$), but it is mild in the Σ_{int} clumps ($\rho = 0.29$) and absent in the Σ_{low} ones ($\rho = -0.22$). The combined sample of high- and intermediate-density clumps, $\Sigma_{\text{int} + \text{high}}$, which together form a sample of 18 objects, also have a mild correlation ($\rho = 0.55$).

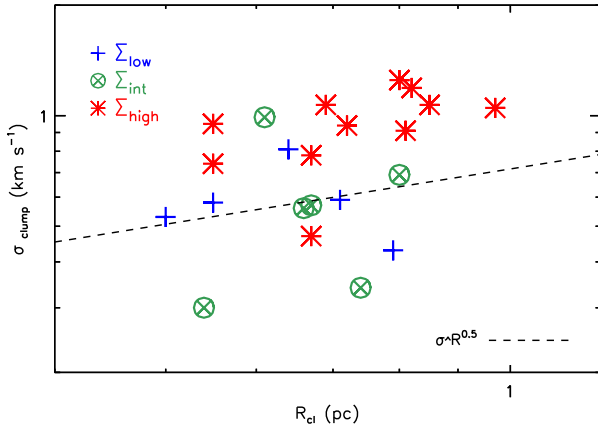


Figure 4. Larson’s relation for the clumps divided into the three groups. The black-dashed line is classical Larson’s relation $\sigma \propto R^\delta$ with $\delta = 0.5$ and normalized to 0.72 km s^{-1} at 1 pc (Heyer & Dame 2015).

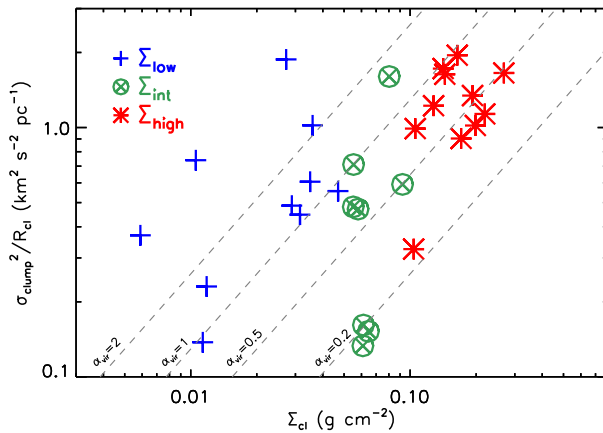


Figure 5. The quantity σ^2/R as a function of the surface density Σ for the clumps belonging to the three groups. The relation changes considering only the clumps belonging to the Σ_{high} and Σ_{int} groups, or also including the clumps in the Σ_{low} group. The grey-dotted lines are in correspondence of constant values of the virial parameter.

If it is gravity that is playing a central role in driving the observed non-thermal motions, then the clumps should follow the generalized form of Larson’s relation, $\sigma^2/r \propto \Sigma$ (Heyer et al. 2009; Ballesteros-Paredes et al. 2011; Traficante et al. 2018b). Massive clumps seem to follow this generalized form (Camacho et al. 2016), which is interpreted as evidence of a hierarchical, global collapse driven by self-gravity (Ballesteros-Paredes et al. 2011; Vázquez-Semadeni et al. 2019).

In Fig. 5, we show the $\sigma^2/r \propto \Sigma$ relation for our clumps. If we focus on the clumps belonging to the intermediate and high surface density groups, they show a good degree of correlation, with Pearson’s coefficient $\rho_{\Sigma_{\text{int+high}}} = 0.65$. The relation changes as we introduce the low-density clumps. These clumps are much sparser in the parameter space, as a consequence of the large spread in the values of their virial parameter. Pearson’s coefficient of the Σ_{low} clumps alone is also relatively low, $\rho = 0.31$, suggesting that the behaviour of these objects is different from the rest of the sample.

A similar way to explore if gravity plays a dominant role in the collapse is through the velocity dispersion–surface density diagram. This relation is similar to the generalized Larson relation, but it has a great advantage: Both Σ and σ are two distance-independent

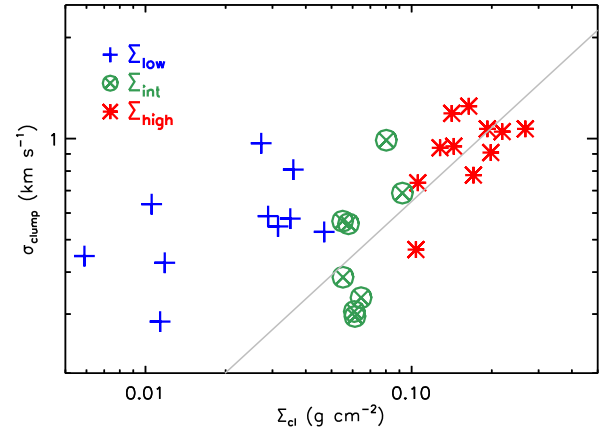


Figure 6. Velocity dispersion of the clumps as a function of the surface density. The grey line is the linear fit to the log–log space of clumps belonging to the Σ_{int} and Σ_{high} groups. The grey-dotted horizontal line is the average value of the velocity dispersion for the clumps in the Σ_{low} group.

quantities, so they are not affected by one of the strongest sources of uncertainty in their estimation. If the non-thermal motions are driven by gravity, they are expected to correlate, keeping the region in a pseudo-virial state (Ballesteros-Paredes et al. 2011).

In Fig. 6, we report the Σ versus σ diagram. The correlation is good for the clumps in the $\Sigma_{\text{int+high}}$ groups, with Pearson’s coefficient of 0.76. The solid line is the fit of the $\Sigma_{\text{int+high}}$ clumps done with the `linefit` IDL routine in the log–log space, which gives a scaling law $\sigma \propto \Sigma^{0.73}$. Analogously to what happened in the $\sigma^2/r \propto \Sigma$ plot, if we also consider the low-density clumps in the diagram, the correlation breaks down. The grey-dotted line in the plot is the average value of the non-thermal motions in the Σ_{low} clumps. The line bumps into the result of the fit for the clumps of the $\Sigma_{\text{int+high}}$ groups at a surface density $\Sigma \simeq 0.08 \text{ g cm}^{-2}$. This value is very close to the critical value of $\Sigma \simeq 0.1 \text{ g cm}^{-2}$ above which all the $70\text{-}\mu\text{m}$ quiet clumps show evidence of dynamical activities (Traficante et al. 2017). In some of these Σ_{high} regions, the HCO^+ (1–0) and HNC (1–0) spectra exhibit clear blue-asymmetries interpreted as evidence of global, parsec-scale collapse (Fuller, Williams & Sridharan 2005; Kirk et al. 2013; Rygl et al. 2013; Traficante et al. 2017), although even the red-asymmetries may be the result of the inflow of material (Smith et al. 2013).

We looked at this dynamical activity at the various surface density regimes searching for asymmetries in the HCO^+ (1–0) and HNC (1–0) (Kirk et al. 2013) spectra in all our sources. The majority of the spectra in the Σ_{low} and Σ_{int} clumps are symmetric (two examples are given in the top row of Fig. 7), with the exception of four clumps that show hints of asymmetries in only the HCO^+ (1–0) spectrum (26.432–0.662) or in both HCO^+ (1–0) and HNC (1–0) lines (24.552+0.096, 35.608+0.111, and 45.531+0.042). These spectra are shown in Fig. 7. The asymmetries are evident once we overplot the central velocity of the N_2H^+ (1–0) line, which we have assumed as the optically thin line and used to estimate the central velocity of each region. The asymmetries in the clumps 35.608+0.111 and 26.432–0.662 can be due to the dynamical activity of the regions, while the clumps 24.552+0.096 and 45.531+0.042 have more complex spectra and the interpretation is less obvious.

The previous results suggest that, moving towards denser and denser regions and approaching the critical surface density value of $\Sigma \simeq 0.1 \text{ g cm}^{-2}$, the contribution of the increasing gravitational

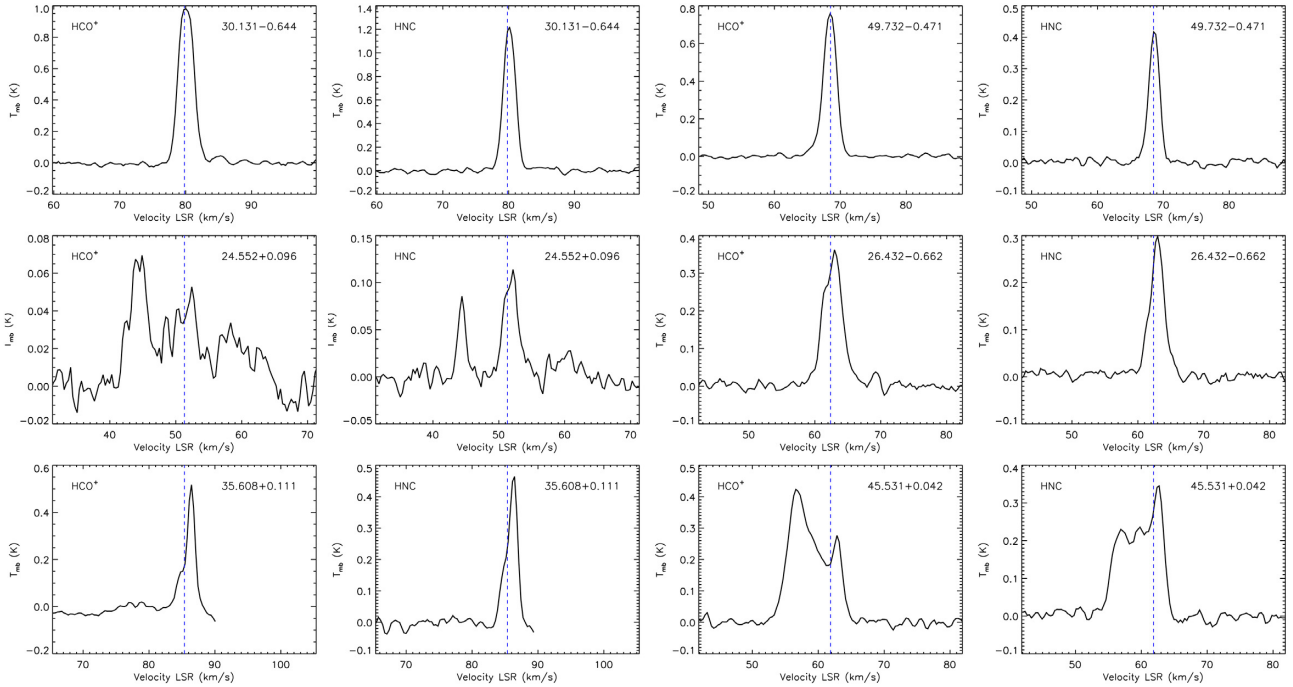


Figure 7. HCO^+ (1–0) and HNC (1–0) spectra averaged across the clump region for six representative clumps. Sources 30.131–0.644 and 49.732–0.471 have spectra symmetric around the central velocity of the clump defined by the N_2H^+ (1–0) emission (blue-vertical line). Sources 24.552+0.096, 26.432–0.662, 35.608+0.111, and 45.531+0.042 have the HCO^+ , HNC , or both asymmetric spectra with respect to the central velocity of the clump.

potential starts to dominate over the turbulence in driving the observed dynamics at the clump scales.

3.2 Filament analysis

In this section, we investigate the properties of the natal filaments and their dynamics, separating them in the same three groups identified by the properties of their embedded 70- μm quiet clumps.

The filament masks derived from the Hi-GAL column density maps (Schisano et al. 2019) delimit the ^{13}CO (1–0) regions that we extract to estimate the properties of each object. For each filament, we take the length, mass, and linear mass from the Hi-GAL filament catalogue. We obtain the length L of each filament from the angular length determined in the Schisano et al. (2019) catalogue, and assuming the same distance of the embedded clump derived in Traficante et al. (2015b). In this way we also determined the length of two filaments in the Hi-GAL catalogue with no determined distance, 49.398–0.576 and 49.732+0.471. We have also derived the filament mass from the ^{13}CO GRS data, together with the central velocity, equivalent radius, and velocity dispersion parameters. The clump 45.531+0.042 partly overlaps with two separated filaments in the Schisano et al. (2019) catalogue with the CO peak emission at very similar velocities. We have estimated the central velocity and the velocity dispersion of the cloud as the average values of the two filaments combined, but it was not possible to uniquely define the length and the linear mass of the combined cloud. For this object, we did not consider these parameters in the analysis.

In order to evaluate the filament properties from the ^{13}CO data cubes, we, first, need to determine the central velocity $v_{\text{LSR, fil}}$ and the velocity range v_{fil} that defines each filament. We initially fix the central velocity of each filament equal to the v_{LSR} of the clump determined from the N_2H^+ (1–0) fit, $v_{\text{LSR, cl}}$, and we fix v_{fil} equal to 10 km s^{-1} , defined as a $\pm 5 \text{ km s}^{-1}$ spectral window

around $v_{\text{LSR, cl}}$. This range is enough to include the emission of each filament (the velocity gradient we measure is few km s^{-1}) and to exclude the contribution from other clouds that may be present along the LOS.

Before refining the calculation of $v_{\text{LSR, fil}}$ (and redetermining the extremes of v_{fil}), we need to take into account that the ^{13}CO (1–0) emission can easily become optically thick along the LOS of the dense clumps. To minimize the uncertainties derived from the optical depth effects, we masked in the GRS cubes a box of 3×3 pixels (to consider a GRS beam) centred on each clump identified in the Elia et al. (2017) catalogue. The ^{13}CO emission in the remaining pixels is assumed to be optically thin, and these pixels are used to estimate the central velocity and v_{fil} of each region. The differences in the average values with and without including these lines of sight (LOS) are, however, relatively small, few per cent, on average. Two examples of the masked maps for the filaments 25.609+0.228 and 49.398–0.576 are shown in Fig. 8.

Starting from the first estimate of v_{fil} and the masked cubes for each filament, we have estimated $v_{\text{LSR, fil}}$ with an iterative approach, starting from the definition of $v_{\text{LSR, fil}}$ as the first moment of the ^{13}CO (1–0) cubes (e.g. Roman-Duval et al. 2010):

$$v_{\text{LSR, fil}} = \frac{\sum_{l,b,v_{\text{fil}}} T_{\text{mb}}(l, b, v) v}{\sum_{l,b,v_{\text{fil}}} T_{\text{mb}}(l, b, v)}, \quad (1)$$

where the sum is done across all (l, b) pixels contained within the filament region in the masked cubes and in a velocity range v_{fil} . We further consider only the pixels with $\text{S/N} > 4$ at the peak position, where the noise level of the map is evaluated separately for each filament in a 30 km s^{-1} spectral window along an LOS with no ^{13}CO emission.

The value of $v_{\text{LSR, fil}}$ we obtain from equation (1) is then used to determine the new $\pm 5 \text{ km s}^{-1}$ v_{fil} window across which re-evaluate $v_{\text{LSR, fil}}$. Usually, after the second iteration, the value of

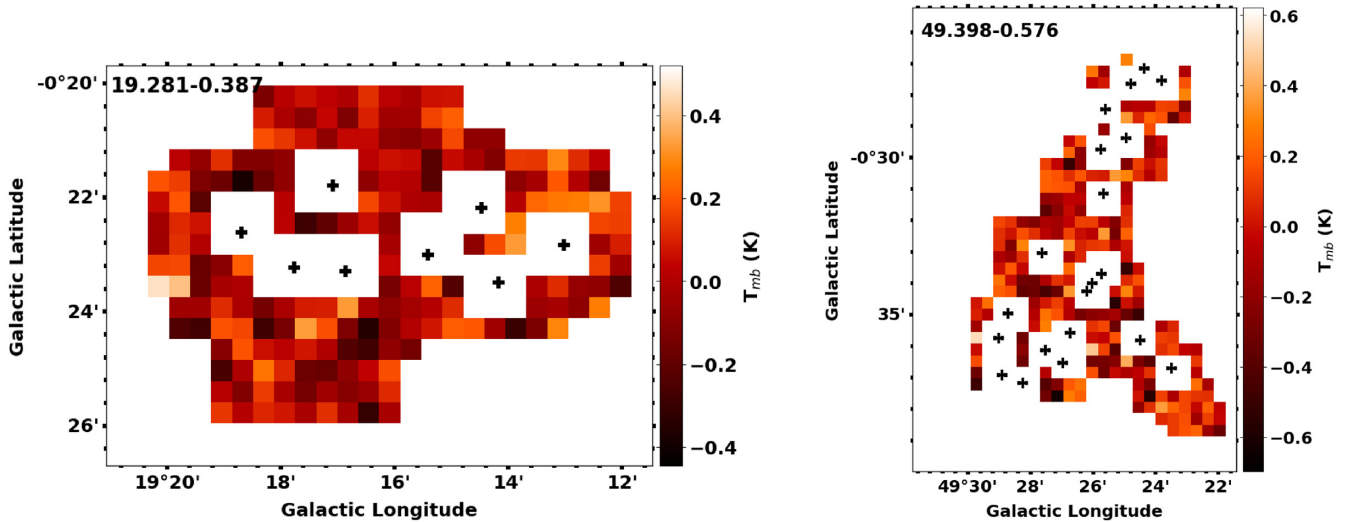


Figure 8. A spectral channel of the ^{13}CO GRS data cube for the filaments 19.281–0.387 and 49.398–0.576 (the same as Fig. 2). The black crosses are the centroids of the sources extracted from the Elia et al. (2017) catalogue. A box of 3 x 3 pixels centred on each embedded Hi-GAL source has been masked to minimize the uncertainties of the filament kinematics across LOSs where the CO emission can be optically thick. The CO emission across the unmasked pixels is assumed to be optically thin and is used to estimate the kinematics of each filament.

$v_{\text{LSR, fil}}$ reaches convergence, and we stop the run and fix the value; otherwise, we perform a third iteration and fix the value of $v_{\text{LSR, fil}}$ (and v_{fil}) at the end of this run. In this way, we can also estimate the velocity shift between the central velocity of the clumps and the average central velocity of the natal filaments. This shift is, on average, of $\simeq 1.2 \text{ km s}^{-1}$, and for only three objects the difference is $> 2 \text{ km s}^{-1}$, as shown in Table 2.

Once we have defined $v_{\text{LSR, fil}}$ (and consequently the range v_{fil}), we can derive all the other parameters of the filaments from the CO data cubes: CO mass M_{CO} , linear mass $M_{\text{line, CO}}$, equivalent radius R_{eq} , average surface density, and velocity dispersion. Only the latter is evaluated in the masked maps, since it may be affected by the optical depth effects. The other parameters are evaluated along all the good pixels, including the LOS of the Hi-GAL sources, because the mass estimation includes the effects of the optical depth, and the regions where the clumps are embedded contribute strongly to the physical parameters (mass, equivalent radius, and surface density) of the filaments. The parameters derived for each filament are listed in Table 2.

The ^{13}CO mass M_{CO} is derived following Roman-Duval et al. (2010):

$$M_{\text{CO}} = 0.27 \frac{\text{d}^2}{\text{kpc}^2} \int_{l, b, v_{\text{fil}}} \frac{T_{\text{ex}}(l, b, v) \tau_{13}(l, b, v)}{1 - e^{-\frac{5.3}{T_{\text{ex}}(l, b, v)}}} \frac{dv}{\text{km s}^{-1}} \frac{dl}{\text{arcmin}} \frac{db}{\text{arcmin}}, \quad (2)$$

where τ_{13} is the opacity and T_{ex} is the excitation temperature evaluated for each voxel (l, b, v) within the 10 km s^{-1} interval defined by v_{fil} . The integral in this case is performed along all the good pixels, including the LOS of the Hi-GAL sources. The opacity in each voxel is estimated as

$$\tau_{13}(l, b, v) = -\ln \left(1 - \frac{0.189 T_{13}(l, b, v)}{(e^{\frac{5.3}{T_{\text{ex}}(l, b, v)}} - 1) - 0.16} \right), \quad (3)$$

where $T_{13}(l, b, v)$ is the ^{13}CO (1–0) brightness temperature of each voxel.

The ^{13}CO (1–0) excitation temperature at each voxel is determined from the excitation temperature of the ^{12}CO (1–0) line, under the assumption that these two temperatures are equal along our filaments. Assuming also that the ^{12}CO line is optically thick along the filaments, the excitation temperature at each voxel $T_{\text{ex}}(l, b, v)$ can be derived from the ^{12}CO brightness temperature as

$$T_{\text{ex}}(l, b, v) = 5.53 \frac{1}{\ln \left(1 + \frac{5.53}{T_{12}(l, b, v) + 0.837} \right)}. \quad (4)$$

The ^{12}CO emission along our filaments is taken from the FUGIN data cubes (Umemoto et al. 2017). First, we created a T_{ex} cube from the ^{12}CO data at the FUGIN spatial resolution (20 arcsec), then we smoothed these maps to the GRS resolution (55 arcsec).

The velocity resolution of the FUGIN cubes (0.65 km s^{-1}) is approximately three times lower than the GRS one ($\simeq 0.2 \text{ km s}^{-1}$). We derive the voxel cubes of T_{ex} assuming at each GRS velocity channel the T_{ex} of the closest FUGIN velocity channel. In this way, the T_{ex} estimated in each FUGIN channel is associated, on average, with three GRS consecutive velocity channels. An example for a pixel of the source 23.076–0.209 is shown in Fig. 9. Since the FUGIN survey does not cover the clouds above longitude $l = 50^\circ$, for the cloud associated with the clump 53.361+0.042, we assumed a fixed excitation temperature of $T_{\text{ex}} = 10 \text{ K}$, in agreement with the average value of the rest of the sample ($T_{\text{ex, mean}} = 8.7 \pm 2.5 \text{ K}$).

The masses derived from dust and from CO are compared in Fig. 10. There is a good correlation between the two (the black-dashed line is the $y = x$ bisector). Two filaments, 28.19–0.192 and 36.608+0.111, have a significantly higher dust mass compared with the CO mass, both belonging to the Σ_{low} group. Both these filaments show multiple CO components along the LOS, and the clump identified in the N_2H^+ emission is not associated with the main CO peak. Therefore, the mass obtained from the integrated continuum infrared emission is likely to be overestimated for both filaments. Given the complexity of the CO spectrum along each LOS, which often intercepts many clouds at likely different dust temperatures, in the following, we will consider only the mass derived from the CO data.

Table 2. Physical and kinematic properties of the filaments associated with the Hi-GAL clumps. Column 1: name of the embedded clump; Columns 2–3: mass of the filament derived from the dust emission and from the CO, respectively; Column 4: length of the filament obtained from the Schisano et al. (2019) catalogue and rescaled assuming the distance of the clump; Columns 5–6: linear mass of the filament derived from the dust mass and the CO mass, respectively; Column 7: equivalent radius of the filament; Column 8: central velocity of the clump determined from the hyperfine fitting of the N_2H^+ (1–0) emission; Column 9: central velocity of the filament determined, as discussed in Section 3.2; Columns 10–11: velocity dispersion of the filament obtained without and with the contribution of the large-scale gradients, as discussed in Section 3.2; Column 12: ratio of the velocity dispersion obtained with and without the contribution of the large-scale gradients, $\sigma_{\text{ratio}} = \sigma_{\text{gl}}/\sigma_{\text{loc}}$; and Column 13: reference group.

| Clump | M_{dust} (M_{\odot}) | M_{CO} (M_{\odot}) | L (pc) | M/L ($M_{\odot} \text{ pc}^{-1}$) | M_{CO}/L ($M_{\odot} \text{ pc}^{-1}$) | R_{eq} (pc) | Σ_{fil} (g cm^{-2}) | $v_{\text{LSR, cl}}$ (km s^{-1}) | $v_{\text{LSR, fil}}$ (km s^{-1}) | σ_{loc} (km s^{-1}) | σ_{gl} (km s^{-1}) | σ_{ratio} | Group |
|--------------|--------------------------------------|------------------------------------|-------------|--|--|-------------------------|---|--|---|---|--|-------------------------|------------------------|
| 49.732–0.471 | – | 12 837 | 17.4 | – | 737.2 | 5.7 | 0.027 | 68.51 | 67.62 | 1.33 | 1.52 | 1.14 | Σ_{low} |
| 28.19–0.192 | 42 320 | 3648 | 25.9 | 1632.3 | 140.7 | 4.7 | 0.011 | 100.70 | 99.44 | 1.54 | 3.24 | 2.10 | Σ_{low} |
| 45.531+0.042 | – | – | – | – | – | – | – | 61.86 | 58.54 | 2.02 | 2.15 | 1.06 | Σ_{low} |
| 23.076–0.209 | 2516 | 4020 | 5.7 | 440.8 | 704.2 | 2.3 | 0.051 | 82.66 | 83.36 | 2.17 | 2.26 | 1.04 | Σ_{low} |
| 26.432–0.662 | 6946 | 12 405 | 15.2 | 458.3 | 818.6 | 4.1 | 0.048 | 62.38 | 63.23 | 1.64 | 1.84 | 1.12 | Σ_{low} |
| 35.608+0.111 | 36 978 | 4699 | 36.6 | 1009.1 | 128.2 | 4.4 | 0.016 | 85.37 | 81.58 | 2.28 | 2.55 | 1.12 | Σ_{low} |
| 30.357–0.837 | 4532 | 7421 | 16.8 | 270.0 | 442.1 | 4.9 | 0.021 | 78.80 | 76.37 | 1.65 | 2.10 | 1.27 | Σ_{int} |
| 53.361+0.042 | 30 788 | 10 326 | 57.7 | 533.3 | 178.9 | 4.2 | 0.038 | 22.66 | 23.48 | 0.91 | 1.04 | 1.14 | Σ_{int} |
| 30.131–0.644 | 1453 | 2049 | 5.5 | 264.9 | 373.5 | 2.6 | 0.020 | 79.81 | 81.28 | 1.44 | 1.85 | 1.28 | Σ_{int} |
| 15.631–0.377 | 1075 | 755 | 2.8 | 390.3 | 273.9 | 2.0 | 0.013 | 40.00 | 39.34 | 1.22 | 1.57 | 1.29 | Σ_{int} |
| 49.398–0.576 | – | 19 968 | 24.1 | – | 827.2 | 6.5 | 0.032 | 60.79 | 58.29 | 1.68 | 2.09 | 1.24 | Σ_{int} |
| 28.792+0.141 | 9601 | 10 988 | 17.5 | 549.0 | 628.3 | 4.8 | 0.031 | 107.20 | 107.96 | 1.90 | 2.28 | 1.20 | Σ_{int} |
| 25.982–0.056 | 2098 | 5780 | 5.5 | 378.2 | 1042.0 | 3.0 | 0.044 | 89.80 | 90.25 | 1.50 | 1.65 | 1.10 | Σ_{int} |
| 19.281–0.387 | 2603 | 1940 | 9.0 | 289.6 | 215.8 | 3.7 | 0.009 | 53.50 | 52.28 | 1.58 | 1.69 | 1.07 | Σ_{high} |
| 34.131+0.075 | 1704 | 2963 | 3.2 | 540.1 | 938.8 | 1.9 | 0.055 | 56.80 | 57.17 | 1.76 | 1.81 | 1.03 | Σ_{high} |
| 23.271–0.263 | 124 147 | 103 054 | 48.0 | 2587.8 | 2148.1 | 9.3 | 0.079 | 82.30 | 79.41 | 2.32 | 2.42 | 1.04 | Σ_{high} |
| 31.946+0.076 | 55 491 | 70 580 | 43.6 | 1272.6 | 1618.7 | 8.1 | 0.071 | 96.40 | 96.17 | 1.84 | 1.95 | 1.06 | Σ_{high} |
| 22.756–0.284 | 8159 | 16 318 | 18.8 | 434.6 | 869.2 | 4.4 | 0.057 | 105.00 | 106.18 | 1.96 | 2.21 | 1.13 | Σ_{high} |
| 22.53–0.192 | 35 907 | 65 881 | 32.2 | 1114.3 | 2044.5 | 7.7 | 0.074 | 76.20 | 77.02 | 2.25 | 2.30 | 1.02 | Σ_{high} |
| 28.537–0.277 | 42 291 | 23 505 | 13.1 | 3222.6 | 1791.1 | 4.4 | 0.081 | 88.30 | 86.95 | 1.61 | 1.78 | 1.11 | Σ_{high} |
| 28.178–0.091 | 17 266 | 33 436 | 28.0 | 616.6 | 1194.2 | 5.8 | 0.067 | 98.20 | 97.69 | 1.99 | 2.13 | 1.07 | Σ_{high} |
| 24.013+0.488 | 78 849 | 73 495 | 88.9 | 887.2 | 827.1 | 11.6 | 0.036 | 95.00 | 95.70 | 1.64 | 1.82 | 1.11 | Σ_{high} |
| 25.609+0.228 | 12 977 | 18 370 | 20.8 | 623.1 | 882.1 | 5.6 | 0.039 | 113.60 | 113.10 | 2.12 | 2.20 | 1.04 | Σ_{high} |
| 18.787–0.286 | 20 426 | 23 354 | 15.3 | 1336.2 | 1527.6 | 4.5 | 0.077 | 65.70 | 65.10 | 1.92 | 2.00 | 1.04 | Σ_{high} |

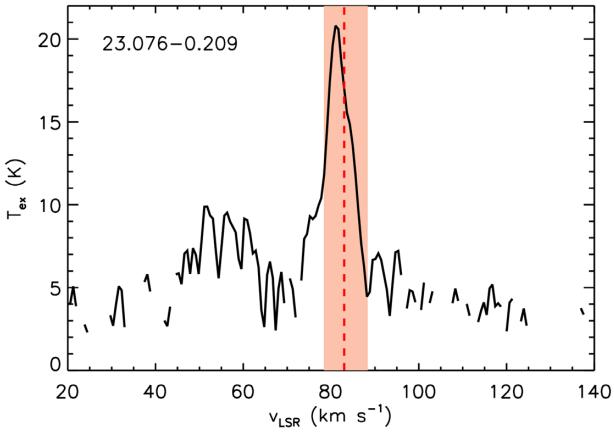


Figure 9. Excitation temperature T_{ex} along a pixel of the filament 23.076–0.209, estimated from the FUGIN data cubes and smoothed at the GRS resolution. The missing points are in correspondence of low-S/N pixels in the ^{12}CO FUGIN data cubes. The red-dotted vertical line is the central velocity of the filament $v_{\text{LSR, fil}}$ and the red-shaded area is the corresponding v_{fil} estimated as discussed in the main text.

The linear mass $M_{\text{fine, CO}}$ is obtained as M_{CO}/L . For each filament, we have also derived an equivalent radius R_{eq} , defined as $R_{\text{eq}} = \sqrt{A/\pi}$, with A equal to the area of the ^{13}CO pixels with well-defined emission (including pixels along the LOSs of the Hi-GAL

clumps). From the M_{CO} and R_{eq} , we determined the average surface density of each filament $\Sigma_{\text{fil}} = M_{\text{CO}} / (\pi R_{\text{eq}}^2)$.

To study the filament kinematics, we have derived the velocity dispersion from the CO data cubes in two different ways. First, we evaluated the second moment of the ^{13}CO (1–0) maps, i.e. the intensity-weighted velocity dispersion (Roman-Duval et al. 2010) in the v_{fil} velocity range defined as

$$\sigma_{\text{gl}}^2 = \frac{\sum_{l,b,v_{\text{fil}}} T_{\text{mb}}(l, b, v)(v - v_{\text{LSR, fil}})^2}{\sum_{l,b,v_{\text{fil}}} T_{\text{mb}}(l, b, v)}. \quad (5)$$

This is the average velocity dispersion of the whole filament, including the contribution of the large-scale velocity gradients that broaden the width of the average spectrum and have several different origins: They can originate from the dynamics of the gas, like, e.g. large-scale infall motions (Kirk et al. 2013) or the gravitational collapse along the filament (Peretto et al. 2014). They may be due to large-scale motions not related to the internal dynamics of the gas, such as the rotation (Kirk et al. 2013) or compression from an external source (Williams et al. 2018), and, for the most elongated objects, they may also originate from Galactic scale motions such as Galactic rotation, shear, or compression (e.g. Duarte-Cabral & Dobbs 2016). Alternatively, they may have a non-physical origin and be the resultant from the blending of the emission from unresolved sub-structures that may broaden the observed width (Beuther et al. 2013) and in some cases mimic a large-scale gradient (Henshaw et al. 2014).

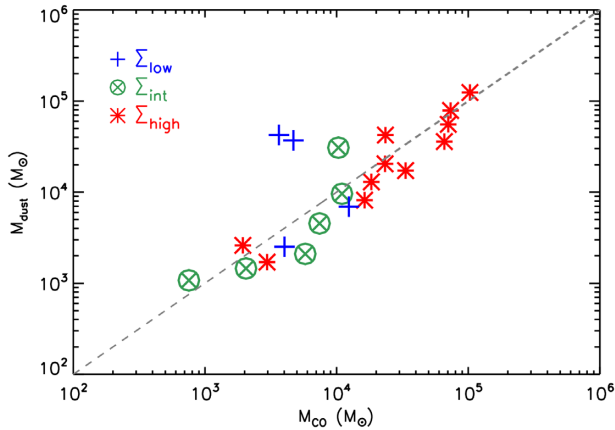


Figure 10. Mass of the filaments derived from dust, M_{dust} , versus mass derived from CO, M_{CO} . The black-dashed line is the $y = x$ bisector. The values for the three groups discussed in Section 3 are represented with different symbols/colours. Objects associated with clumps in the Σ_{low} , Σ_{int} , and Σ_{red} groups are represented as blue crosses, green circles with X, and red asterisks, respectively.

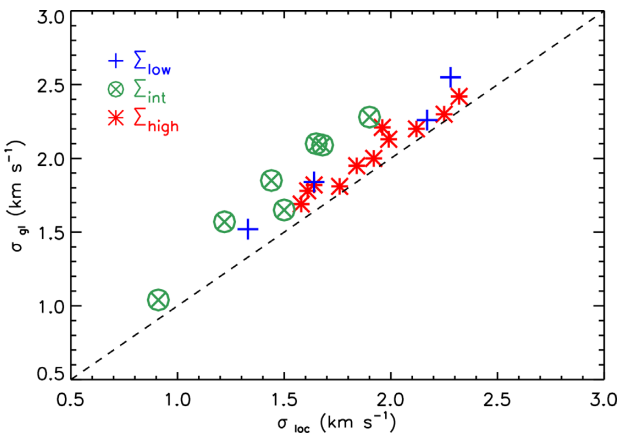


Figure 11. Velocity dispersion of the filaments derived, including the large scale gradients, σ_{gl} , versus velocity dispersion derived after the removal of these gradients, σ_{loc} . The results for the clump 28.19–0.192 are not included since the σ_{gl} value is significantly larger than σ_{loc} because of a second CO component along the LOS, as explained in Section 3.2, and it goes outside the plotted ranges.

In order to consider the motions in the filaments without these large-scale gradients, we have also derived the velocity dispersion of each filament without including their contribution. To do this, first, we derived the map of the second moment per pixel, defined as

$$\sigma_{\text{loc}}^2(l, b) = \frac{\sum_{v_{\text{fil,loc}}} T_{\text{mb}}(l, b, v)(v(l, b) - \langle v_{\text{loc}}(l, b) \rangle)^2}{\sum_{v_{\text{fil,loc}}} T_{\text{mb}}(l, b, v)}, \quad (6)$$

where $\langle v_{\text{loc}}(l, b) \rangle$ is equivalent to the central velocity $v_{\text{LSR, fil}}$ defined in equation (1) but now evaluated separately for each pixel. Since $v_{\text{fil,loc}}(l, b)$ is estimated separately for each pixel, effectively treating all of them as independent pixels, we can get an estimate of the average non-thermal motions per filament without a large-scale contribution, (σ_{loc}), by evaluating the mean value of $\sigma_{\text{loc}}(l, b)$.

The values of σ_{gl} and σ_{loc} and their ratio are given in Table 2, and in Fig. 11, we show σ_{loc} against σ_{gl} for each filament. The values of σ_{loc} are systematically lower than σ_{gl} , and the difference between them also provides an indication of the contribution of

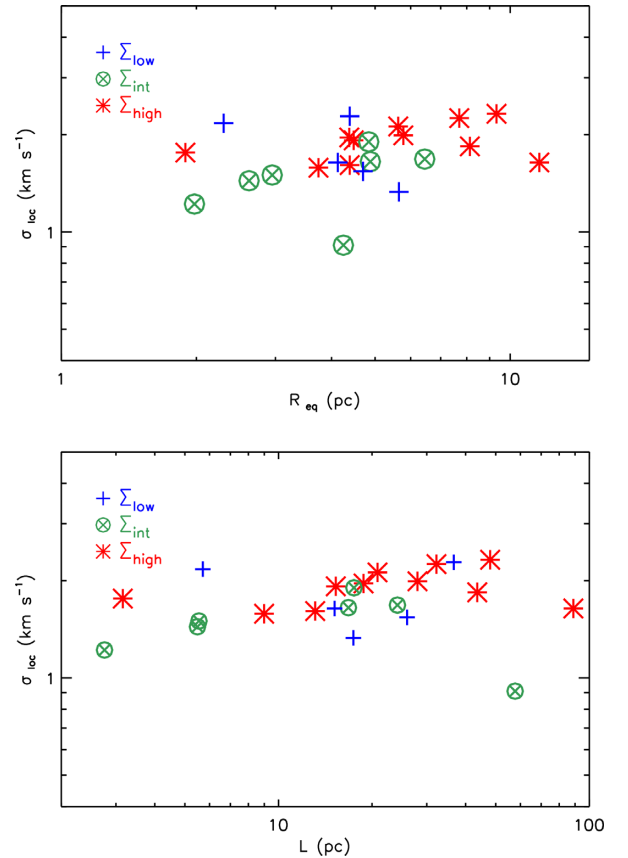


Figure 12. Upper panel: Larson relation evaluated for the filaments. The filaments size is considered as their equivalent radius R_{eq} . Lower panel: same as before, but the size of the filament is equal to its length.

the large-scale kinematics to the measured non-thermal motions in each cloud. These differences are more pronounced in the Σ_{low} and Σ_{int} clumps rather than in Σ_{high} clumps, and they are of the order of $\simeq 30$ per cent at most. The only exception is for the filament 28.19–0.192, where the difference is larger than 100 per cent. This is due to the contribution of a second CO component along the LOS, which contributes to the estimate of σ_{gl} , and it is likely the same component that also contributed to the overestimate of the mass of this filament from the dust emission maps shown in Fig. 10.

These relatively small differences between σ_{loc} and σ_{gl} imply that the observed motions are dominated by the internal motions of the cloud, and not by the mechanism(s) responsible for the observed large-scale gradients.

3.2.1 The dynamics in filaments

Similar to what we have done for the clumps, we analyse here some relations useful for investigating the dynamics of the filaments.

In order to evaluate Larson’s size–linewidth relation for our filaments, we have considered two possible definitions of the filament size. In Fig. 12, upper panel, we show the relation where the size of the filament is defined from its equivalent radius R_{eq} . However, these objects are particularly elongated, and an equivalent radius, although a useful parameter to define the size of a structure, may not be representative of the larger scales that may contribute to drive the non-thermal motions. The largest scale at which the kinetic energy can be injected in a filament is its whole length, L ,

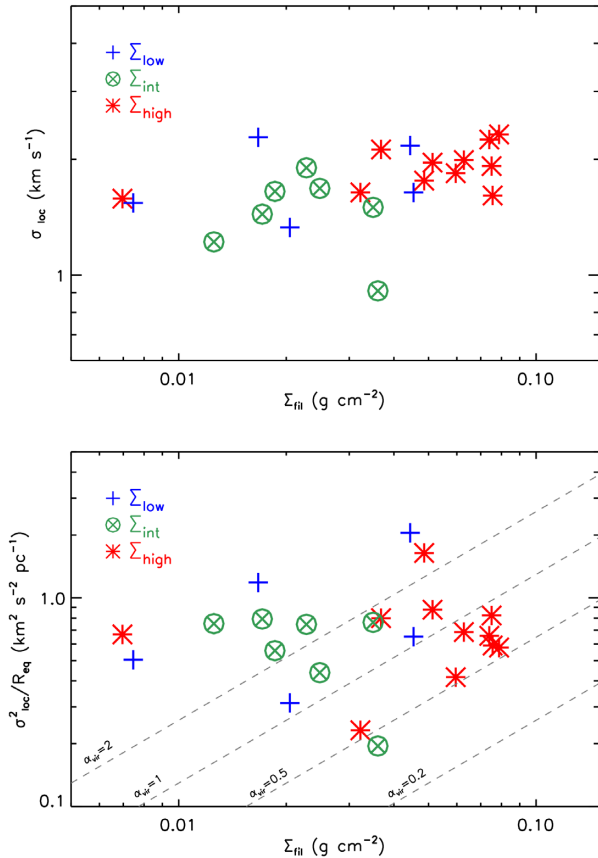


Figure 13. Upper panel: velocity dispersion of the filament (estimated after the removal of the large-gradient contribution) as a function of the surface density of the filaments. Lower panel: generalized Larson relation for the filaments. The grey-dashed lines are in correspondence of constant values of the virial parameter.

and this is the other quantity that we have considered to derive the Larson relation (Fig. 12, lower panel). In both these diagrams, there is no evident correlation between size and velocity dispersion, and this is true for all the filaments.

In Fig. 13, upper panel, we show the relation between the velocity dispersion σ_{loc} and the surface density of the filaments. The correlation is very poor in all the groups, and this lack of correlation is also present in the generalized Larson relation, in the lower panel of Fig. 13.

The lack of correlation in the previous diagrams could indicate that the average kinetic energy of each filament is determined by the energetics of the ambient environment, which can differ from region to region. Instead, in the next section, we will show how, at the filament-to-clump scales, the turbulent cascade is the most likely explanation for the observed non-thermal motions within each filament.

The filamentary structures themselves can accrete from the ambient gas, and the inflow of new material contributes to generating a cascade of energy down to the star-forming sites, which resembles the turbulent cascade, the so-called accretion-driven mechanism (Klessen & Hennebelle 2010). In this case, filaments must evolve with time with the accretion of their mass per unit length M_{line} and, correspondingly, their internal velocity dispersion. This scenario is observed in filaments detected in the *Herschel* Gould Belt Survey (Arzoumanian et al. 2011), with a power-law relation between M_{line} and σ of the form $M_{\text{line}} \propto \sigma^{0.31}$ (Arzoumanian et al. 2013), although

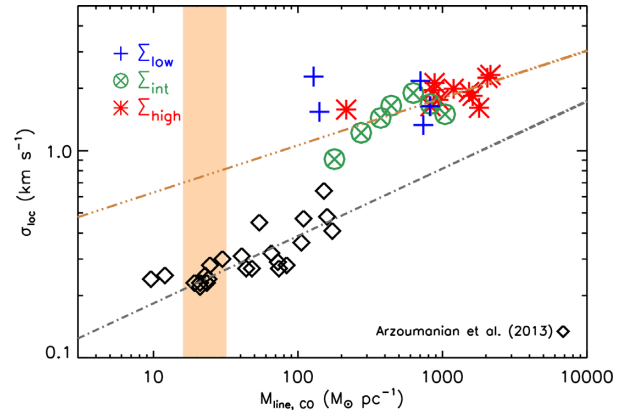


Figure 14. Velocity dispersion of the filaments estimated without the contribution of the large-scale gradients as a function of the linear mass derived from the CO data. The orange dash-dotted line is the fit to the filaments with clumps in the $\Sigma_{\text{int}} + \Sigma_{\text{high}}$ groups. The black open diamonds are the nearby filaments studied in Arzoumanian et al. (2013). The orange rectangle delimits the region above which the filaments are considered thermally supercritical, around the critical value of $M/L \simeq 16 M_{\odot} \text{pc}^{-1}$.

they are different objects compared with the ones analysed in this work. Arzoumanian et al. (2013) observed a collection of nearby, relatively low-mass filaments with sizes ≤ 1 pc, while the objects in the Schisano et al. (2019) catalogue are distributed across the whole Galaxy and they can reach lengths up to tens of parsecs.

To investigate if a similar mechanism can be responsible for the dynamics observed in our parsec-scale filaments, in Fig. 14 we plot the linear mass $M_{\text{line,CO}}$ as a function of the velocity dispersion σ_{loc} for our objects, superimposed with the results of nearby regions taken from Arzoumanian et al. (2013).

First, note that there is an offset between the velocity dispersion measured in nearby regions and in our filaments. This is likely to be attributed to the different gas tracers used to measure the motions. In the work of Arzoumanian et al. (2013), the gas tracers are C^{18}O (2–1) and N_2H^+ (1–0), both with a higher critical density than the ^{13}CO (1–0) used in this work and more suited to trace only the inner, denser regions of the filament.

The correlation of the filaments in the Σ_{high} group alone is mild (Pearson’s coefficient is 0.56), and it increases if we combine the filaments in the $\Sigma_{\text{int}} + \Sigma_{\text{high}}$ groups ($\rho = 0.71$). The fit to the data of these two groups together gives a relation $M_{\text{line}} \propto \sigma^{0.23}$, a shallower slope but still comparable with the results obtained in Arzoumanian et al. (2013). The few points of the regions belonging to the Σ_{low} group are uncorrelated in this diagram.

The same mechanism that drives the motion in nearby filaments could also drive the non-thermal motions observed at the scales of our filaments, but more data are needed to investigate this relation in a statistically significant way.

4 THE CORRELATION BETWEEN CLUMPS AND FILAMENTS

In this section, we combine the results obtained separately for clumps and filaments to attempt a multiscale analysis of our data.

The relation between the surface density of the clumps, Σ_{cl} , and the average surface density of the parent filament, Σ_{fil} , is shown in Fig. 15. On average, the densest clumps in our sample are embedded in the filaments with the highest surface density (with the evident exception of 19.281–0.387, where the massive

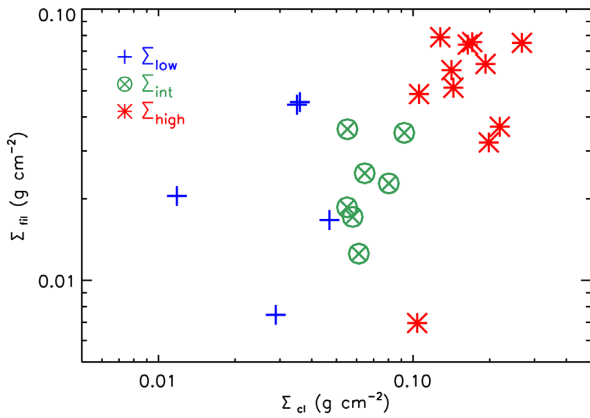


Figure 15. Average surface density of the filament Σ_{fil} versus the surface density of each clump Σ_{cl} for the three groups, as discussed in Section 3.

clump is embedded in a particularly low density filament). The large scatter observed for the values of the low-density clumps is likely to be due to an observational bias. To show that, we compared the properties of our clumps with the properties of other clumps embedded in each filament, derived from the Elia et al. (2017) Hi-GAL clump catalogue (the parameters of each clump

has been rescaled consistently at the distance of each filament). All our filaments embed more than one clump, the large majority of them being low- and intermediate-density clumps. The bias arises because 5/7 and 6/11 of the clumps identified in the Σ_{int} and Σ_{high} groups are the most massive ones embedded in their corresponding filaments, and they are the main responsible for the observed correlations. In contrast, all the filaments belonging to the Σ_{low} group embed clumps more massive than the 70- μm quiet one analysed here. We will come back to this point later in this section.

This result is a further indication that the environment plays a significant role in the formation of massive stars: Densest and more massive clumps seems to form preferentially along massive, dense filaments.

Next, we investigated the correlation between the kinematics of the clumps and the parent filaments. This can be done without any bias since we are tracing the kinematics of the clumps and filaments separately with different gas tracers. The ^{13}CO (1–0) line, with a critical density of $\simeq 3 \times 10^3 \text{ cm}^{-3}$, is a good tracer of the relatively low column density material that delimits the filament, as opposed to the N_2H^+ (1–0) line, which, with a critical density of $\simeq 6 \times 10^4 \text{ cm}^{-3}$ (and an effective critical density of $\simeq 1 \times 10^4 \text{ cm}^{-3}$, Shirley 2015), traces the denser parts of the filament, including the embedded clumps.

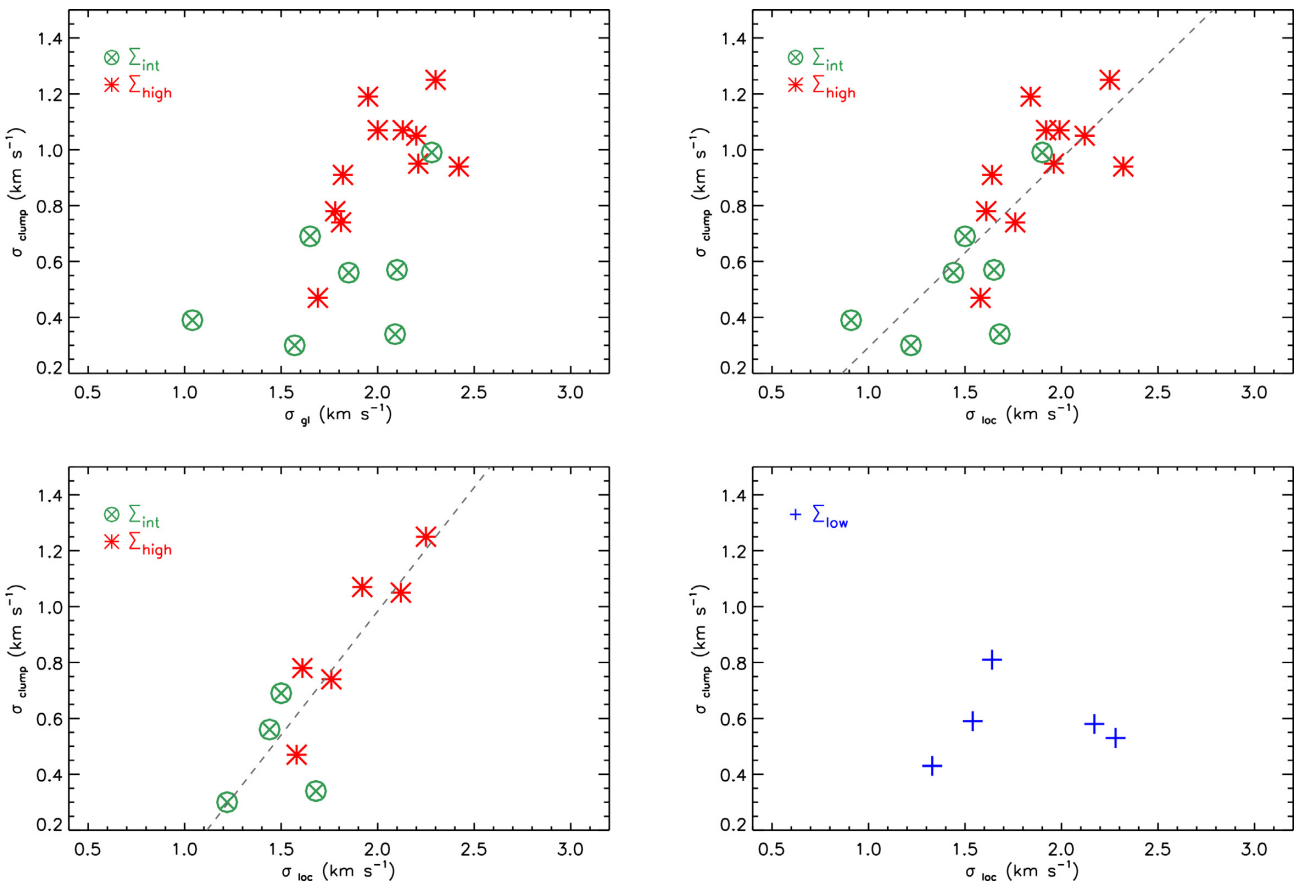


Figure 16. Velocity dispersion of the clumps as a function of the velocity dispersion of the parent filaments. Upper left-hand panel: clumps and filaments belonging to the Σ_{int} and Σ_{high} groups. The velocity dispersion of the filaments σ_{fil} is estimated, including the large-scale gradients. Upper-right panel: same as the previous plot, but using the velocity dispersion of the filaments estimated after the subtraction of the large-scale gradients, σ_{loc} . The grey-dotted line is the fit of all the sources in log–log space. Lower left-hand panel: same σ_{clump} versus σ_{loc} plot but including only the 10 regions where the observed clumps are the most massive of each filament. Lower right-hand panel: same relation but for sources belonging to the Σ_{low} group. The velocity dispersion of the filament is estimated after the removal of the large-scale gradients, σ_{loc} .

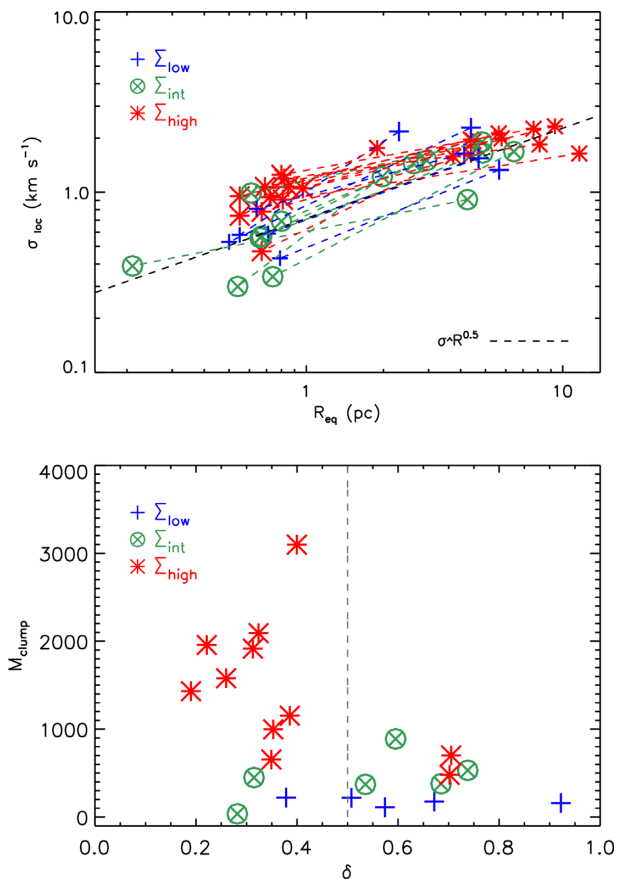


Figure 17. Upper panel: Larson relation for the clumps and the parent filaments belonging to the three groups. The dotted lines connect each clump (left-hand points) with the parent filament (right-hand points). The black-dotted line is the $\sigma \propto R^{0.5}$ relation in the case of a supersonic turbulent cascade of energy from large down to small scales. Lower panel: $M_{\text{clump}}-\delta$ relation for our sample. The dashed vertical line is in correspondence of $\delta = 0.5$, the fiducial value in the case of supersonic turbulent cascade.

As shown in Fig. 16, top left-hand panel, the velocity dispersion of the clumps in Σ_{int} and Σ_{high} correlates quite well with the velocity dispersion that includes the large-scale gradients, σ_{gl} , with Pearson’s coefficient of $\rho_{\text{int}} = 0.52$ and $\rho_{\text{high}} = 0.65$, respectively. However, the two distributions occupy slightly different regions of the parameter space. While the velocity dispersion in the Σ_{int} clumps is lower than in the Σ_{high} clumps, the global non-thermal motions of the parent filaments are very similar, in the range $1.0 \leq \sigma_{\text{gl}} \leq 2.4 \text{ km s}^{-1}$.

In Fig. 16, top right-hand panel, we show the same plot but we consider the non-thermal motions without including the large-scale gradients in the filaments, σ_{loc} . The diagram changes in particular for the Σ_{int} regions, because the contribution of the large-scale term with respect to σ_{loc} is, on average, four times higher in this group with respect to the Σ_{high} clumps (see Table 2). As a consequence of the large-scale gradient removal, the correlation for the Σ_{int} clumps increases. The two samples together have a significant Pearson’s coefficient of $\rho_{\text{int} + \text{high}} = 0.79$, and the correlation has a power-law form $\sigma_{\text{clump}} \propto \sigma_{\text{loc}}^{0.68}$.

This tight correlation between the kinematics of the filaments and the embedded clumps is confirmed when we plot the same relation but limited to the sample of 10 filaments for which the clumps we observed are the most massive ones (Fig. 16, bottom left-hand

panel). The relation is stronger, with a very high Pearson’s coefficient of 0.86. The correlation has a power-law form $\sigma_{\text{clump}} \propto \sigma_{\text{loc}}^{0.89}$.

On the other hand, although the statistics are rather poor (only five points), the non-thermal motions of the filament appear disentangled from the non-thermal motions of the low-density clumps (Fig. 16, bottom right-hand panel) if we consider either the global velocity dispersion value for the filament (with Pearson’s coefficient of $\rho_{\text{low}} = 0.02$) or σ_{loc} (Pearson’s coefficient of $\rho_{\text{low}} = 0.05$), but this is possibly due to the same observational bias discussed before.

To further explore the origin of this correlation, in Fig. 17, upper panel, we again plot the velocity dispersion–size relation, but this time with dashed lines that connect the values of each clump with the corresponding values of the parent filament (in these plots, we consider the equivalent radius R_{eq} for both clumps and filaments).

Each region now shows an increase in the velocity dispersion moving from the clump to its corresponding filament. If we average the slope of each clump–filament pair to estimate a $\sigma \propto R^{\delta}$ relation for the three groups, we obtain an average value of $\bar{\delta} = 0.50 \pm 0.23$ (the uncertainty is the standard deviation of the distribution), in agreement with the scaling relation expected from a turbulent cascade (McKee & Ostriker 2007). The average value of δ for the three groups separately is $\bar{\delta}_{\text{low}} = 0.61 \pm 0.20$, $\bar{\delta}_{\text{int}} = 0.60 \pm 0.27$, and $\bar{\delta}_{\text{high}} = 0.38 \pm 0.17$, which indicates that the densest clumps have larger internal motions at similar size-scales than the least dense clumps of our sample. To emphasize this trend, in Fig. 17, we report the mass of clump as a function of the exponent δ .

We conclude that the dynamics of the filaments and of the embedded clumps are highly correlated. The cascade of turbulence transfers the kinetic energy from the scales of the filaments down to parsec scales within each region. This cascade in each filament produces clumps of various densities, sizes, and velocity dispersions imprinted with the kinematics of the parent filament. However, the densest clumps have the largest velocity dispersion, the scaling with size indicating that they have excess motions compared with that expected from supersonic turbulence, where $\sigma \propto R^{0.5}$.

5 DISCUSSION: THE INTERPLAY BETWEEN GRAVITY AND TURBULENCE AT DIFFERENT SPATIAL SCALES

The results discussed in the previous sections suggest that energy is transferred from large down to small scales initially with a continuous cascade of energy, but there is a gradual variation in the dynamics moving from the low-density regime to the high-density one. Fig. 18 shows a cartoon of this scenario. Turbulence contributes to building up the material in density perturbations, and once these have reached a critical surface density of $\Sigma \simeq 0.1 \text{ g cm}^{-2}$, the effect of gravity becomes apparent in their motions. In the most extreme and dense regions, dynamic motions driven by accretion flows are observed at the parsec and larger scales (Peretto et al. 2013; Traficante et al. 2017), suggesting the 0.1 g cm^{-2} surface density is an effective threshold independent of the size scale. If the clouds are embedded in an environment that allows the gas configuration to reach the critical density threshold already at the clump/filament scales, there will be enough material to allow gravity to drive the motions and build up massive star-forming sites (in agreement with the global collapse model, Vázquez-Semadeni et al. 2019). In this scenario, the ability to form the most massive stars is already imprinted in the clouds from the very beginning. At the same time, turbulent eddies are likely to be responsible for the formation of the low- and intermediate-density clumps in all

filaments, leaving at the gravity the role to form objects only at the sub-clump size scales. The cascade of turbulence breaks down in any case at the transition regime between supersonic and transonic motions, reached at scales around $\simeq 0.1$ pc (Federrath 2016).

The lack of relatively isolated high-mass starless cores (Ginsburg et al. 2012; Motte et al. 2018) can be a corollary of this scenario: For these objects to be formed, the environment should be very dense. In this case, however, gravity has already overtaken turbulence at parsec scales, and the whole clump/filament is already in a state of global collapse, which triggers the hierarchical fragmentation, as described in the Vázquez-Semadeni et al. (2019) model.

6 SUMMARY

In this work, we have analysed the dynamics of a sample of 70- μm quiet clumps combined with the properties of their parent filaments. The clump analysis has been done on a new sample of 70- μm quiet clumps selected to be in a low- to intermediate-density regime, with $\Sigma \leq 0.05 \text{ g cm}^{-2}$, combined with the sample of massive 70- μm quiet clumps (with $\Sigma \geq 0.05 \text{ g cm}^{-2}$) presented in Traficante et al. (2017). The parent filaments have been extracted from the newly published catalogue of Hi-GAL filaments (Schisano et al. 2019) and their dynamics analysed using ancillary CO data.

By dividing the clumps and the parent filaments into three groups, chosen to represent three different surface density regimes for the clumps, we have obtained the following results:

(i) The kinematics of intermediate and high surface density clumps ($\Sigma \geq 0.05 \text{ g cm}^{-2}$) correlate with their surface density, with an increased level of non-thermal motions for higher Σ . This correlation is not present in the low surface density clumps ($\Sigma < 0.05 \text{ g cm}^{-2}$). Combining the results of the three regimes, we identify a critical surface density value, $\Sigma \simeq 0.1 \text{ g cm}^{-2}$, above which the velocity dispersion starts to correlate with the surface density. This is a similar threshold above which the clumps are all dynamically active and show evidence of parsec-scale infall motions (Traficante et al. 2018b).

(ii) The densest star-forming regions are found predominantly in the densest filaments. Similarly, the clumps with the highest velocity dispersion, σ_{clump} , are observed in filaments with the highest velocity dispersion, σ_{loc} (after correction for the large-scale velocity gradients). These two quantities show a power-law form $\sigma_{\text{clump}} \propto \sigma_{\text{loc}}^{0.89}$. In general, there is a strong correlation between the kinematics of the clumps and that of the parent filaments in regions that embed intermediate- and high-density objects.

(iii) The velocity dispersion–size scaling relation, when evaluated connecting each filament with its embedded clump, shows an average correlation of the form $\sigma \propto R^\delta$ with $\delta = 0.50 \pm 0.23$, in agreement with the prediction for a turbulent cascade. This multiscale dynamics, with energy transferred from large down to small scales, are observed in all regions. In the densest filaments, this cascade may be driven by the accretion itself at all scales.

(iv) On similar size scales, the velocity dispersion is, on average, higher in the densest clumps. If the velocity dispersion of the low-density clumps traces the effect of the pure turbulent cascade from larger scales, then the densest clumps have an additional source of internal motions. This would be consistent with gravity playing an enhanced role in these most extreme regions, and would support the global collapse scenario towards extremely massive clumps, like, e.g. SDC 335 (Peretto et al. 2013).

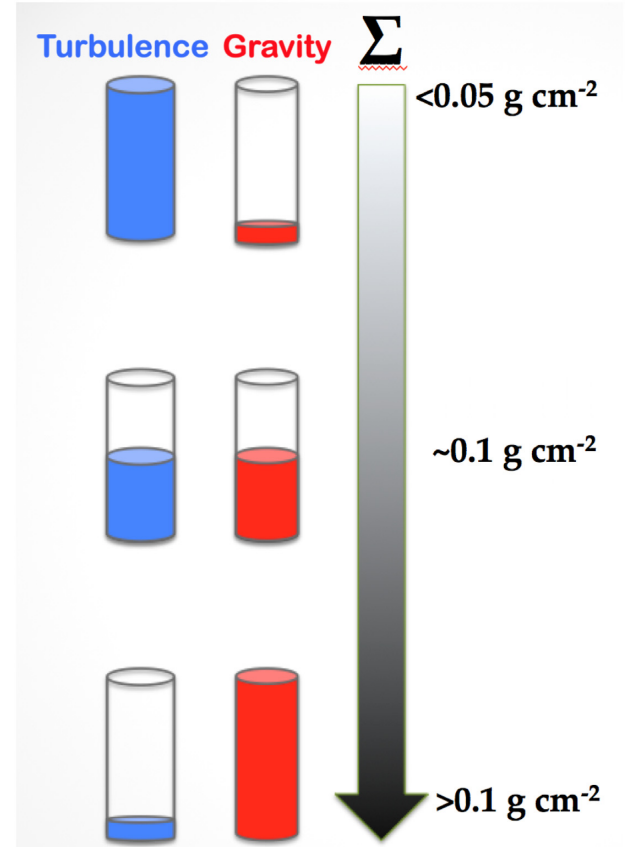


Figure 18. A cartoon representation of the interplay between gravity and turbulence in star-forming regions determined from our results. A cloud begins to be compressed due to the effect of the strong interstellar medium turbulence, with a minimal contribution from gravity. As the material starts to accumulate in different eddies due to turbulence, these regions increase their surface density and with it the role of gravity in driving the non-thermal motions. As the density reaches the critical value of $\Sigma \simeq 0.1 \text{ g cm}^{-2}$, the motions induced by gravity dominate over the turbulence, which, from this point, becomes sub-dominant. This critical value of the surface density can be reached at any spatial scales, depending on the initial conditions of the turbulent cloud. Gravity can therefore start to drive the motions at the core, the clump, or even the filament scales in star-forming regions embedded in significantly dense clouds.

In conclusion, we observe multiscale dynamics within filaments and clumps, which we interpret as a continuous interplay between turbulence and gravity. The former drives the non-thermal motions from filaments down to small scales, where gravity begins to dominate the dynamics once the regions reach a surface density above a critical value of $\Sigma \simeq 0.1 \text{ g cm}^{-2}$. This value is reached at the clump (or larger) scales within the densest filaments observed in our sample.

Finally, it should be noted that this analysis does not consider the role of the magnetic fields, which could be significant at any spatial scales and requires further investigation to fully understand the nature of the observed non-thermal motions as a function of size and density.

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