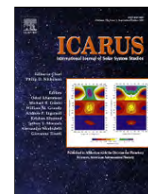




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## Global maps of Venus nightside mean infrared thermal emissions obtained by VIRTIS on Venus express

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### ABSTRACT

One of the striking features about Venus atmosphere is its temporal variability and dynamics, with a chaotic polar vortex, large-scale atmospheric waves, sheared features, and variable winds that depend on local time and possibly orographic features. The aim of this research is to combine data accumulated over several years and obtain a global mean state of the atmosphere focusing in the global structure of the clouds using the cloud opacity and upper cloud temperatures.

We have first produced global maps using the integrated radiance through the infrared atmospheric windows centred around 1.74  $\mu\text{m}$  and 2.25  $\mu\text{m}$ , that show the spatial variations of the cloud opacity in the lower clouds around 44–48 km altitude and also provide an indirect estimation of the possible particle size. We have also produced similar global maps using the brightness temperatures seen in the thermal region at 3.8  $\mu\text{m}$  and 5.0  $\mu\text{m}$ , which provide direct indication of the temperatures at the top of the clouds around 60–70 km altitude.

These maps have been generated using the complete dataset of the Visible and InfraRed Thermal Imaging Spectrometer mapping channel (VIRTIS-M) on board Venus Express, with a wide spatial and long temporal coverage in the period from May 2006 until October 2008.

Our results provide a global view of the cloud opacity, particle size and upper cloud temperatures at both hemispheres, showing the main different dynamical regions of the planet. The profiles obtained also provide the detailed dependencies with latitude, local time and longitude, diagnostic of the global circulation flow and dynamics at various altitude layers, from about 44 up to 70 km over the surface.

### 1. Introduction

#### 1.1. Venus express

Venus Express is the ESA's first mission to Earth's nearest planetary neighbour, launched on November 9th 2005 from Baikonur and arrived at its destination on April 11th 2006. The spacecraft was successfully inserted on a stable elliptical orbit of 24-hour period, with

the pericenter at about 250 km of altitude near the north pole and the apocenter at about 66,000 km of altitude over the southern hemisphere. Being designed in this way, one spacecraft orbit is completed approximately in one Earth day, and hence the total number of orbits is indicative of the days from the orbit insertion of Venus Express. Among the main science goals of the mission, the study of the complex dynamics and structure of the atmosphere was one of the key objectives, as well as the interaction with the surface and outer space (Svedhem et al., 2007, Titov et al., 2006).

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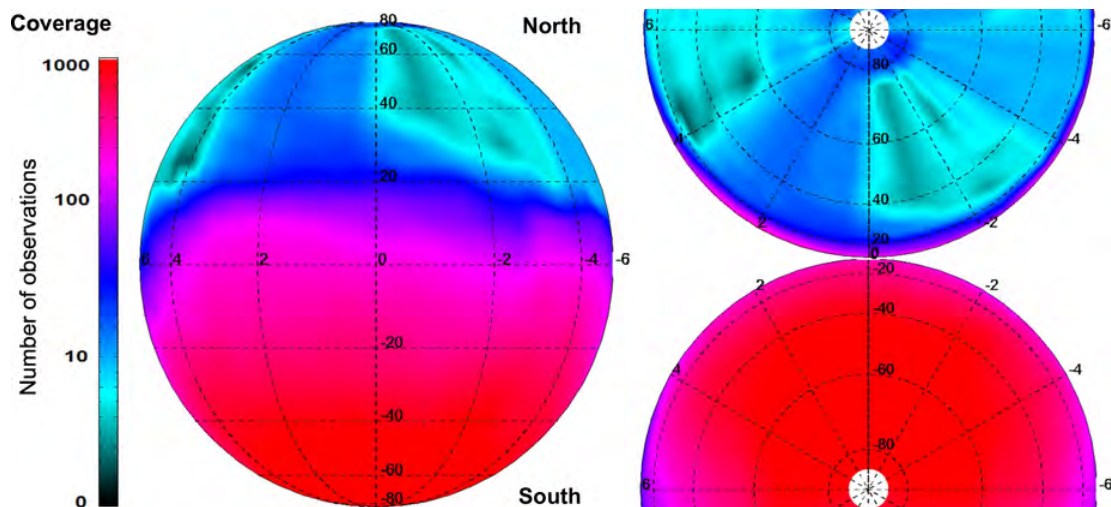


Fig. 1. Local Time coverage of VIRTIS-M-IR observations (in Logarithmic scale) of Venus nightside, as seen from the equator at midnight (left) and from both poles (right). Shown in an orthographic projection distributed along latitude (north up) and local time (increasing eastward) with a fixed grid of  $[5^\circ \times 30 \text{ min}]$ .

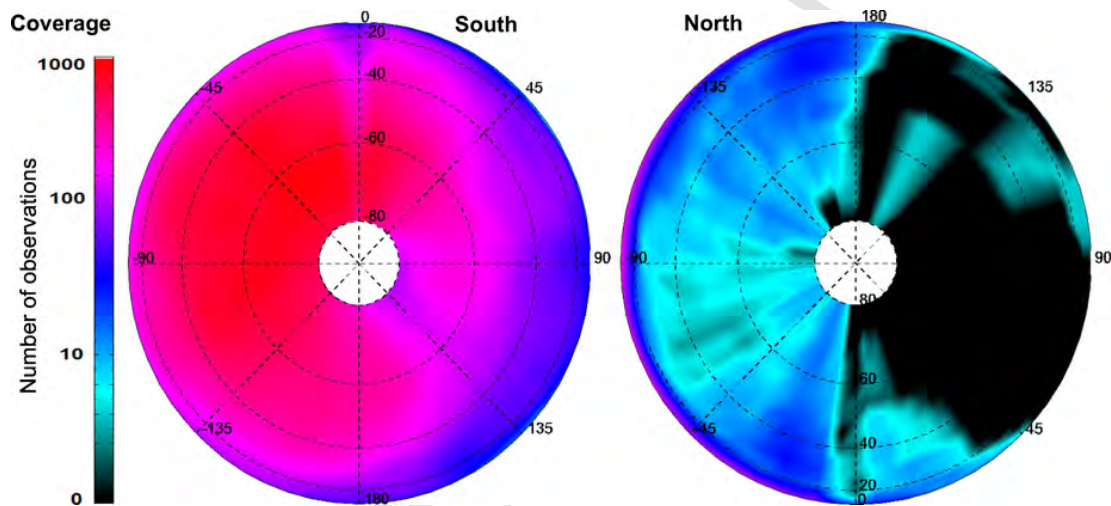


Fig. 2. Longitudinal coverage of VIRTIS-M-IR observations (in Logarithmic scale) over both Venus hemispheres, as seen from the South pole (left) and North pole (right). Shown in an orthographic projection distributed along latitude and longitude (increasing eastward) with a fixed grid of  $[5^\circ \times 5^\circ]$ .

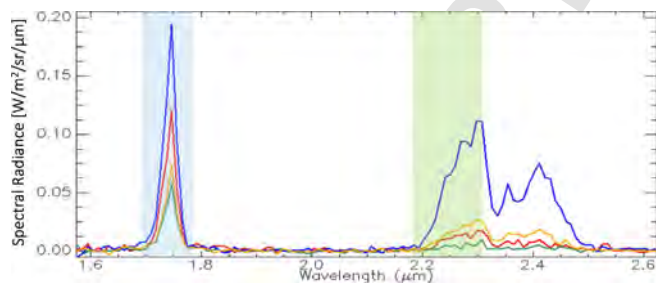


Fig. 3. Typical VIRTIS-M infrared spectrum in the nightside of Venus between 1.6 and  $2.6 \mu\text{m}$  showing the main two atmospheric windows centred at  $1.74 \mu\text{m}$  (marked in blue) and  $2.25 \mu\text{m}$  (marked in green). The different coloured plots show typical variations in the radiances seen through these windows at various locations of the Venus nightside atmosphere: Red = polar vortex; Green = cold collar, Blue = mid-latitude, Orange = low-latitude. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

## 1.2. VIRTIS instrument

VIRTIS (Visible and Infra-Red Thermal Imaging Spectrometer) was the imaging spectrometer on-board Venus Express with the

capability to map the planet from the surface up to the thermosphere. An imaging spectrometer of this type is perfectly suited to cover very different fields, from meteorology of the mesosphere to geology, through the near infrared deep windows sounding down to the surface (Drossart et al., 2007). VIRTIS was constituted of three channels in one compact instrument (Piccioni et al., 2008). One channel was devoted solely to spectroscopy at high resolution, housed in the High resolution (-H) optical subsystem ( $2\text{--}5 \mu\text{m}$ ). The other channels were housed in the Mapper (-M) optical subsystem and devoted to spectral mapping, with one detector in the visible (-M-VIS,  $0.25\text{--}1.0 \mu\text{m}$ ) with a spectral sampling of about  $2 \text{ nm}$  and another one in the infrared band (-M-IR,  $1.0\text{--}5.2 \mu\text{m}$ ) with a spectral sampling of about  $10 \text{ nm}$ . In this study we use only the VIRTIS-M-IR data, calibrated in spectral radiance units, with an uncertainty below  $10^{-3} \text{ W/m}^2/\text{sr}/\mu\text{m}$  in the spectral range between  $1.7 \mu\text{m}$  and  $5.1 \mu\text{m}$  (Piccioni et al., 2008, Cardesín-Moinelo et al., 2010).

### 1.3. Venus lower clouds in the infrared atmospheric windows

Observations at different wavelengths, including ultra-violet (UV), visible (VIS) and near-infrared (NIR) spectral ranges, and comparison of observed features are very important to investigate cloud properties. As largely demonstrated in the past and more recently shown with

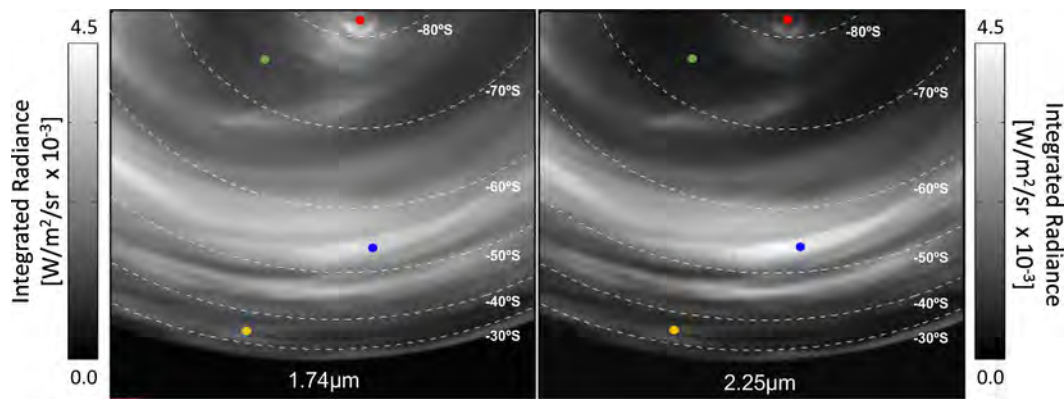


Fig. 4. Comparison of the integrated radiance for the two wavelength bands in the same observation (VI0025\_07, taken 15 May 2006 at 14:36 UTC) of the southern hemisphere of Venus. The coloured dots mark four representative latitude regions (Red = polar vortex; Green = cold collar, Blue = mid-latitude, Orange = low-latitude). (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

Table 1

List of data selection parameters used to build the mean global maps of integrated radiance in the two main infrared atmospheric windows.

Atmospheric window	1.74 $\mu\text{m}$	2.25 $\mu\text{m}$
Wavelength range for radiance integration	[1.68 $\mu\text{m}$ –1.78 $\mu\text{m}$ ] ( $\pm 10$ nm) bands 70–80	[2.18 $\mu\text{m}$ –2.30 $\mu\text{m}$ ] ( $\pm 10$ nm) bands 122–135
Wavelengths used as reference for continuum subtraction	[1.67 $\mu\text{m}$ , 1.79 $\mu\text{m}$ ] ( $\pm 10$ nm), bands 69,81	[2.17 $\mu\text{m}$ , 2.51 $\mu\text{m}$ ] ( $\pm 10$ nm) bands 121–157
Incidence angle filter	> 100°	> 100°
Emission angle filter	< 80° (to avoid limb observations)	< 80° (to avoid limb observations)
Limb darkening correction	$R/(0.34 + 0.66 \cdot \cos(\text{EA}))$ R: Radiance EA: Emission Angle	$R/(0.23 + 0.77 \cdot \cos(\text{EA}))$ R: Radiance EA: Emission Angle
Exposure time filter	> 0.1 s (exclude dayside observations)	> 0.1 s (exclude dayside observations)
Total observations used	3133	3133

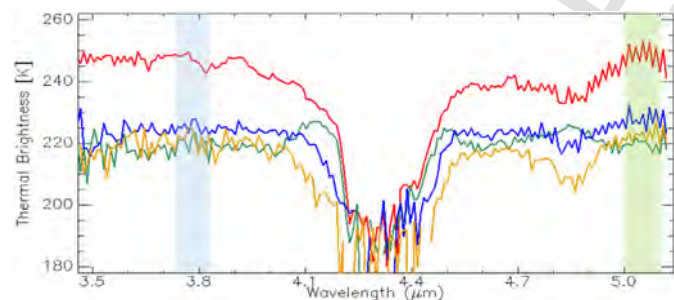


Fig. 5. Typical VIRTIS-M infrared spectrum in the nightside of Venus between 3.5 and 5.1  $\mu\text{m}$  showing the main thermal regions used in this study, around 3.8 (marked in blue) and 5.0  $\mu\text{m}$  (marked in green). The different coloured plots show typical profiles in the thermal brightness as measured at various reference points of the nightside atmosphere: Red = polar vortex; Green = cold collar, Blue = mid-latitude, Orange = low-latitude. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

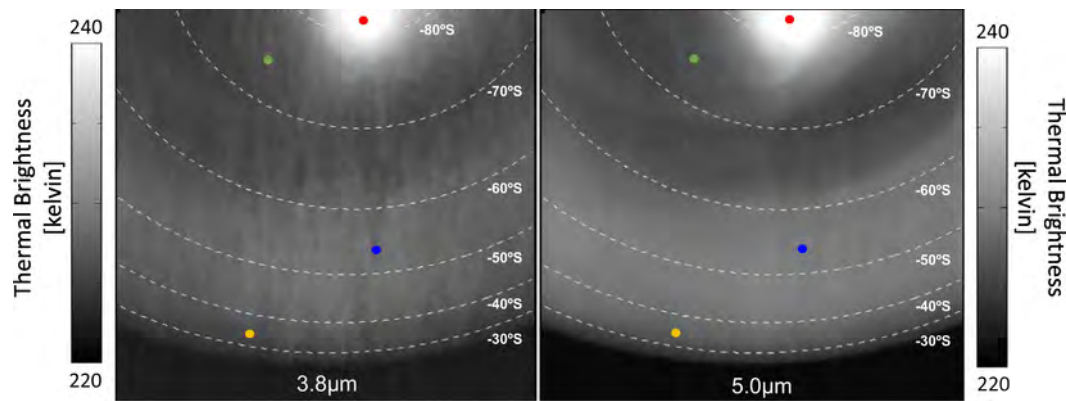
the Akatsuki data (Limaye et al., 2018a), observations in the NIR bands are able to probe clouds at lower altitudes with respect to the observations in the UV because these are less sensitive to scattering by small particles. However, the NIR  $\text{CO}_2$  transparency windows can contribute to probe deeper in the atmosphere, and hence when same structures are found in the UV, VIS and NIR these might indicate that the features extend from the cloud top to a depth of one scale height and larger cloud particles are detected. The wavelengths analysed in the present paper, centred around 1.74 and 2.25  $\mu\text{m}$ , belong to radiance coming

mainly from below the cloud layers. Radiative transfer models indicate that the radiance at these bands originates about 10–20 and 20–30 km above the surface, respectively, see (Allen and Crawford, 1984, Crisp et al., 1991, Carlson et al., 1991) and also the recent review by Titov et al. (2018). This thermal radiation from the lower layers goes up through the cloud layer and is attenuated mostly by the lower clouds with differing optical depths (Titov et al., 2018, Sánchez-Lavega et al., 2017). Therefore, the features observed at these wavelengths show mainly the spatial variations in the opacity of the lower to middle clouds, around 44–48 km altitude (Peralta et al., 2017a).

These clouds have been studied since decades by Earth telescope observations, see Tavenner et al. (2008) and references therein and by several space missions, starting with the fly-by of Venus from the Galileo NIMS instrument (Carlson et al., 1991), later by VIRTIS on Venus Express (Sánchez-Lavega et al., 2008, Hueso et al., 2012) and more recently by Akatsuki (Nakamura et al., 2016) with the IR2 camera as described by Satoh et al. (2016, 2017) and Peralta et al. (2018), which also includes a good overview of Venus cloud observations at 2.3  $\mu\text{m}$  since 1978. These clouds are known to be composed mainly of sulphuric acid covering the whole planet (Marcq et al., 2018; Titov et al., 2018), dynamically elongated by the strong zonal winds coming from the super-rotation and travelling towards the poles with the meridional winds caused by the Hadley cell (Sánchez-Lavega et al., 2008). The general structure and dynamics of the clouds have been analysed and compared extensively with General Circulation Models (Lebonnois et al., 2010, Garate-Lopez and Lebonnois, 2018, Kashimura et al., 2019, Ando et al., 2019).

Ground-based observations pointed out variability in the cloud distribution, the equator region being more cloudy than at mid-latitudes (Crisp et al., 1991), and confirmed by more recent investigation that report a minimum of cloud opacity at mid-latitudes (Tavenner et al., 2008). The same conclusions are obtained with the Akatsuki IR2 camera at 2.3  $\mu\text{m}$  through investigation of the wind speed (Peralta et al., 2018). The authors report high opacity at low latitudes as well as complex cloud features at mid-latitudes. These dark and bright structures could be explained with horizontal inhomogeneities in the clouds' opacity, as already suggested by Allen and Crawford (1984) and Crisp (1989).

The comparison of the different opacities observed at 1.74  $\mu\text{m}$  and 2.3  $\mu\text{m}$  has also been studied to provide an indication of the cloud particle size, showing a uniform spatial distribution in the low to middle latitudes, with a significant increase of the particle size towards the poles (Carlson et al., 1993, Grinspoon et al., 1993, Wilson et al., 2008, McGouldrick and Toon, 2008). The comprehensive retrieval by Haus et al. (2013) also reported a minimum of the particle size around the mid-latitudes, concluding that mid latitude clouds between 30 and 60° are covered by thinner clouds than the equatorial and polar latitudes.



**Fig. 6.** Comparison of the thermal brightness observed in the two wavelength bands in the same observation (VI0025.07, taken 15 May 2006 at 14:36 UTC) of the southern hemisphere of Venus. The coloured dots mark four representative latitude regions (Red = polar vortex; Green = cold collar, Blue = mid-latitude, Orange = low-latitude). (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

**Table 2**

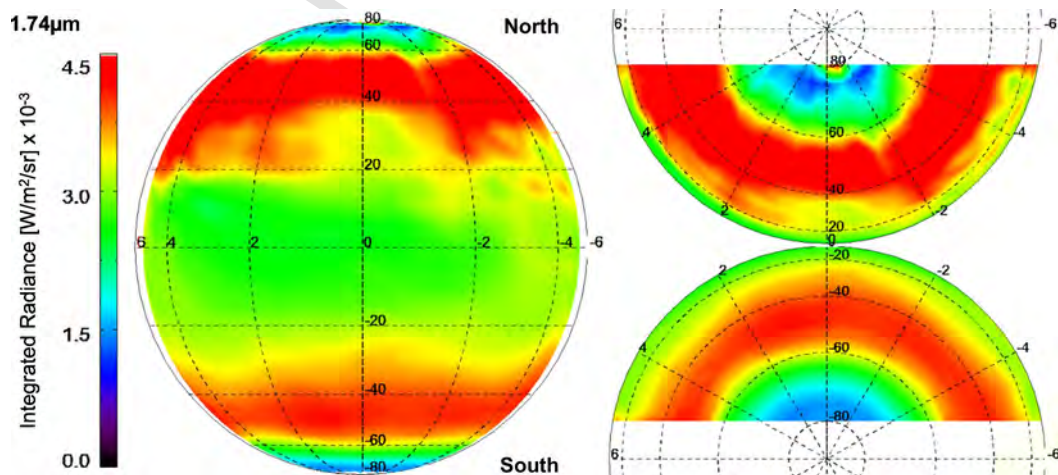
List of data selection parameters used to build the mean global maps of thermal brightness in the two main thermal emission regions.

Thermal emission region	3.8 $\mu\text{m}$	5.0 $\mu\text{m}$
Wavelength range for thermal emission	[3.73 $\mu\text{m}$ –3.82 $\mu\text{m}$ ] ( $\pm 10$ nm), bands: 287–296	[5.01 $\mu\text{m}$ –5.10 $\mu\text{m}$ ] ( $\pm 10$ nm), bands: 421–431* *Even bands are ignored to avoid “odd-even effect” (Cardesin-Moinelo et al., 2010)
Mean Thermal Brightness	All bands are converted into Thermal Brightness and then averaged together	All bands are converted into Thermal Brightness and then averaged together
Incidence angle filter	> 100°	> 100°
Emission angle filter	< 60° (to reduce limb darkening)	< 60° (to reduce limb darkening)
Limb darkening correction	$R/(0.13 + 0.87 \cdot \cos(\text{EA}))$ R: Radiance EA: Emission Angle	$R/(0.20 + 0.80 \cdot \cos(\text{EA}))$ R: Radiance EA: Emission Angle
Exposure time filter	> 0.1 s (exclude dayside observations)	0.1–2 s (to avoid data saturation)
Total observations used	3034	1524

#### 1.4. Venus cloud top temperatures seen in the thermal region

In the second part of this study, we focus on the upper cloud layer that is observed by VIRTIS-M in the thermal range from 3 to 5  $\mu\text{m}$ , around the main  $\text{CO}_2$  absorption band at 4.3  $\mu\text{m}$ , visible in Fig. 5. This wavelength region has been discussed at length in the recent review by Limaye et al. (2018b) and studied by several groups in the past years, including various specific analysis using VIRTIS data. In particular (Grassi et al., 2008, 2010, 2014) provided average temperature fields focusing mostly in the southern hemisphere and describing the local time variabilities and strong dusk/dawn asymmetries in the cloud top temperatures especially around the south cold collar region. Garate-Lopez et al. (2015) extended the analysis of this local time variability around the southern cold collar, including an extensive study of the strong variations in altitude, shape and very fine structures around the south polar vortex. Migliorini et al. (2012) provided temperature field variations using VIRTIS-H data and covering a wider latitude range, reaching the equatorial region and northern hemisphere, although limited by the low resolution at high northern latitudes. Haus et al. (2013) also performed an independent analysis of the cloud top temperatures with VIRTIS and showed a remarkable latitudinal symmetry of both hemispheres with respect to the equator.

In this work we focus on the thermal brightness measured outside the main  $\text{CO}_2$  absorption band, at two specific wavelengths around 3.8  $\mu\text{m}$  and 5.0  $\mu\text{m}$ , shown in Fig. 6, where the atmospheric



**Fig. 7.** Map of integrated radiance in the window centred at 1.74  $\mu\text{m}$ , shown in an orthographic projection of the Venus nightside hemisphere, distributed along Latitude (north up, south down) and Local Time. The left projection is centred at the anti-solar point (midnight, local time 00h). The right projections are centred at the poles.

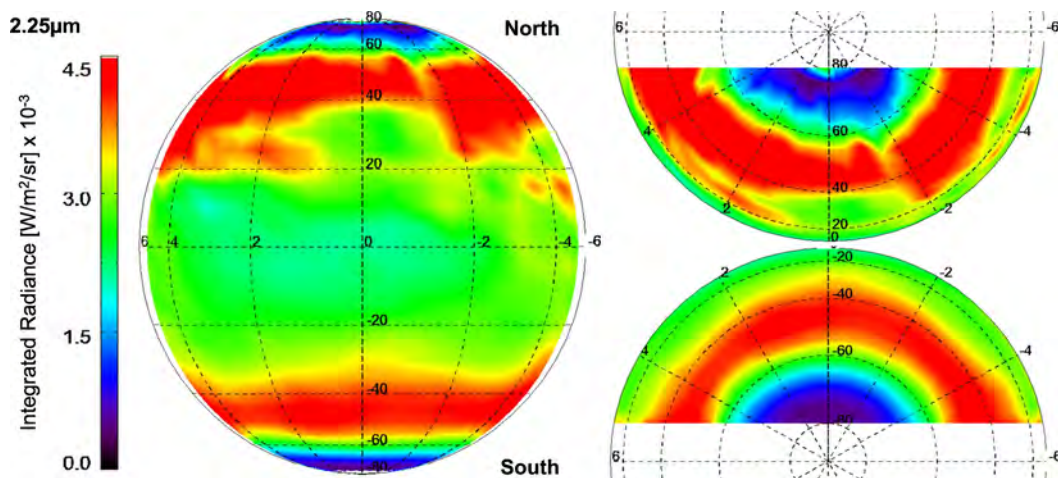


Fig. 8. Map of integrated radiance in the window centred at  $2.25\mu\text{m}$ , shown in an orthographic projection of the Venus nightside hemisphere, distributed along Latitude (north up, south down) and Local Time. The left projection is centred at the anti-solar point (midnight, Local time 00h). The right projections are centred at the poles.

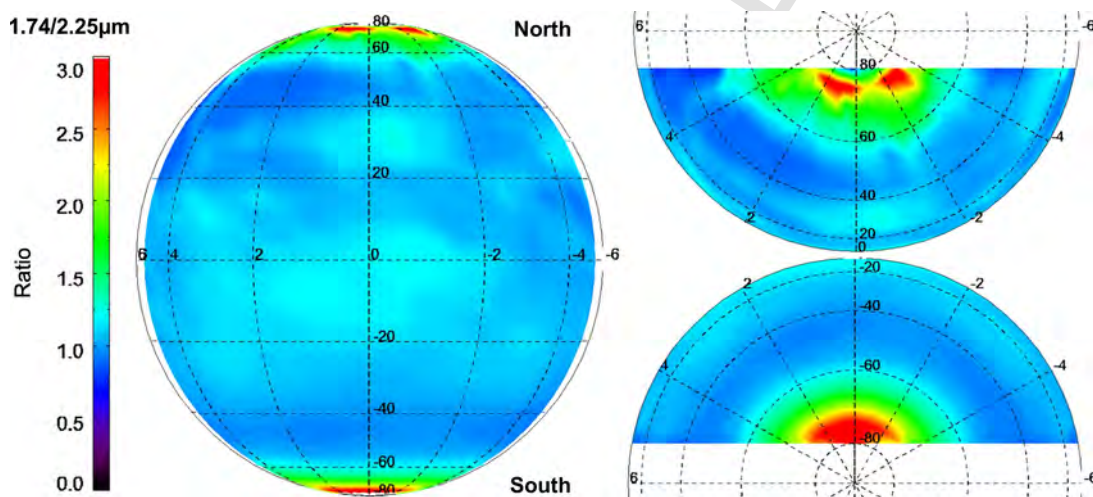


Fig. 9. Ratio between the integrated radiances measured at  $1.74\mu\text{m}$  and  $2.25\mu\text{m}$ . Shown in an orthographic projection of the Venus nightside hemisphere, distributed along Latitude (north up, south down) and Local Time. The left projection is centred at the anti-solar point (midnight, Local time 00h). The right projections are centred at the poles.

opacity from the cloud tops towards space is relatively low. Radiative transfer calculations set the main source of the radiance at these wavelengths as coming from a layer at about 60–70 km altitude, corresponding to the upper cloud layer (Peralta et al., 2017a; Limaye et al., 2018b). Moreover, these wavelengths are of particular interest because they allow the study of the polar vortex morphology and dynamics (Piccioni et al., 2007; Luz et al., 2011; Garate-Lopez et al., 2013, 2015, 2016).

The paper is organized as follows. Data are described in Section 2. Results for the lower clouds are presented in Section 3, while cloud top temperature maps are presented in Section 4. Discussion and conclusion follow in Sections 5 and 6.

## 2. Data selection

This study is based on all VIRTIS-M-IR data from the beginning of the Venus nominal science mission in May 2006, starting with the first science orbit number 23. Previous observations obtained in April 2006 during the Venus orbit insertion and commissioning phase have not been included in this study as the instrument temperatures were much higher than the nominal values, affecting the spectral and radiometric calibration which could introduce important bias in the mean values.

A total of 3133 nightside observations have been included in this study, up to October 2008 when the VIRTIS-M-IR operations

terminated due to the failure of the cryocooler, which was used to keep the cold temperature required by the infrared sensor.

The radiance from Venus at different wavelengths is very different, so there is no unique exposure time covering the full spectral range in an optimized way. The environmental temperature conditions are quite severe and a relatively high thermal background from the instrument dark current prevents the use of long exposure time without saturation (Piccioni et al., 2008; Cardesín-Moinelo et al., 2010). On the other hand, the radiance from the atmospheric windows is very weak and it requires relatively long exposure time. The VIRTIS operations used an optimized exposure time depending on the science case for each specific observation. The long exposures, typically  $>3.3\text{ s}$ , have been used roughly in  $\sim 60\%$  of the nightside observations to focus in the deep atmosphere (windows  $1.74$  and  $2.25\mu\text{m}$ ). With such an exposure time, the thermal band saturates and hence thermal investigation (including cloud top temperature and atmospheric structure) can be addressed only below  $4.2\mu\text{m}$ . The short exposures, below  $0.36\text{ s}$ , have been used in  $\sim 40\%$  of the nightside observations and allow measuring the thermal region at longer wavelengths up to  $5.1\mu\text{m}$ . The maps are constructed using all available data for each wavelength, including both long and short exposure times where possible, only excluding saturated data, in particular for the thermal region at  $5.0\mu\text{m}$  where only short exposures have been used.

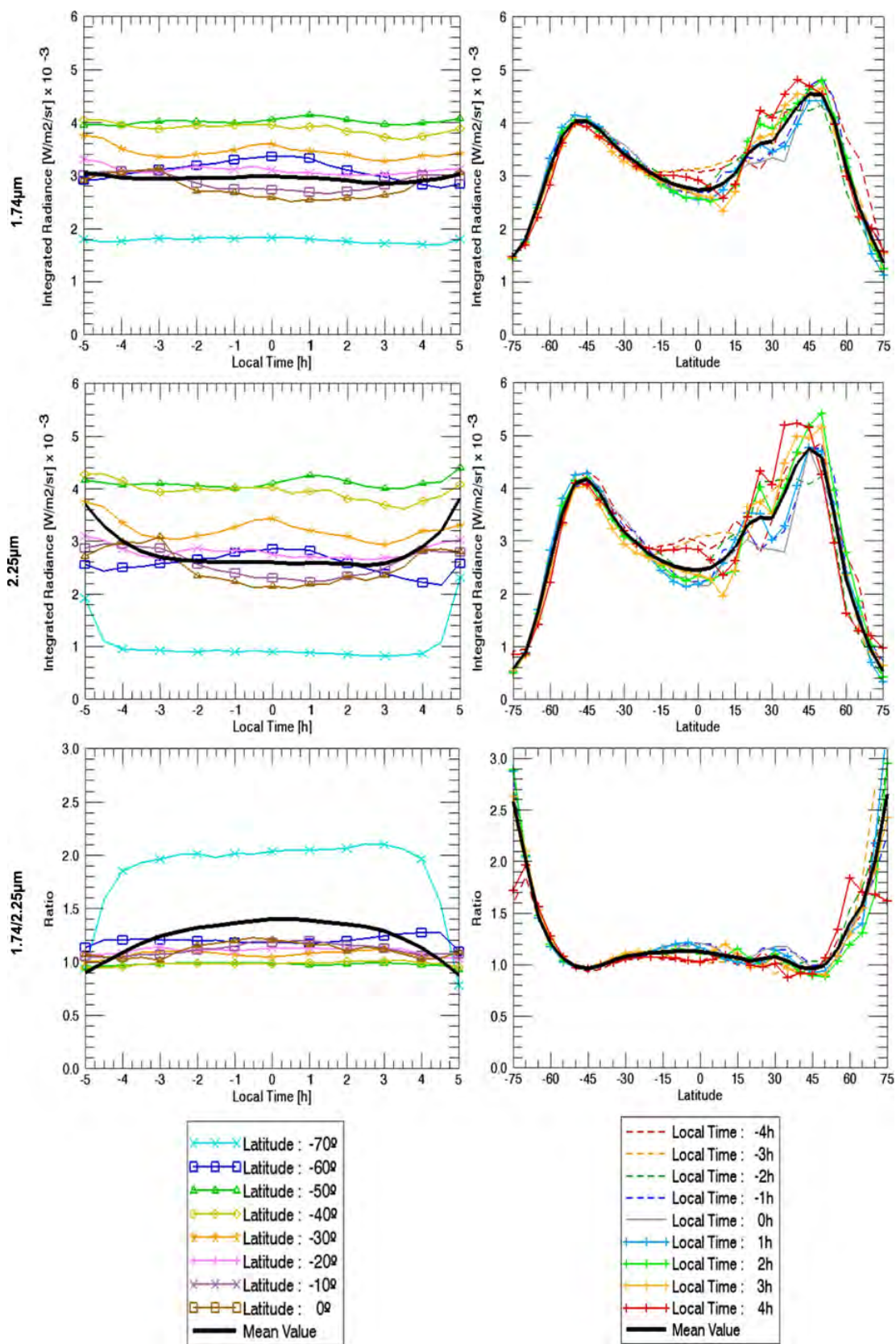


Fig. 10. Radiance profiles at 1.74 μm (top), 2.25 μm (middle) and ratio between them 1.74/2.25 μm (bottom) as a function of local time (left) and latitude (right). The regions close to the terminator and polar caps are omitted since they are affected by scattered radiation.

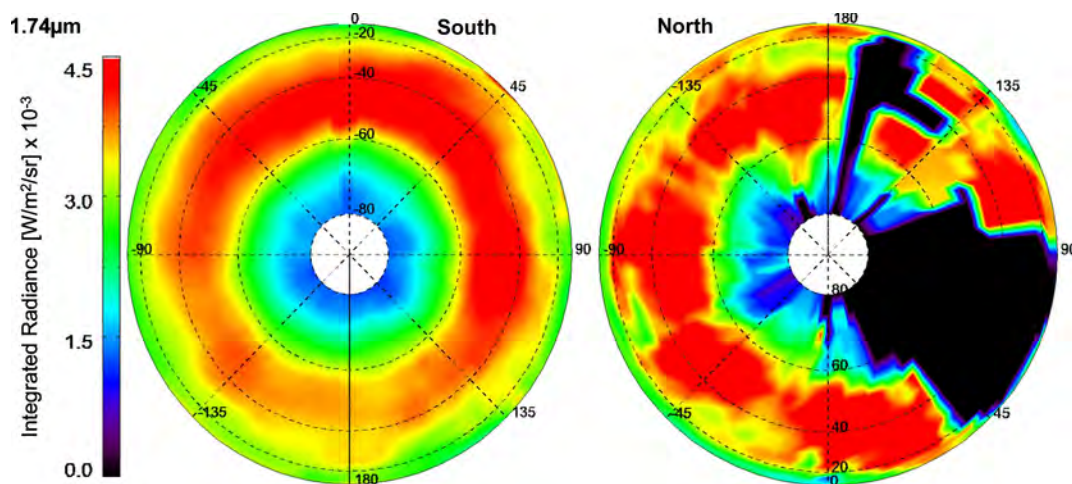


Fig. 11. Mean radiance integrated in the window at  $1.74\mu\text{m}$ , shown in an orthographic projection of the Venus southern (left) and northern (right) hemispheres, distributed along Latitude (pole in the centre) and Longitude (increasing eastward).

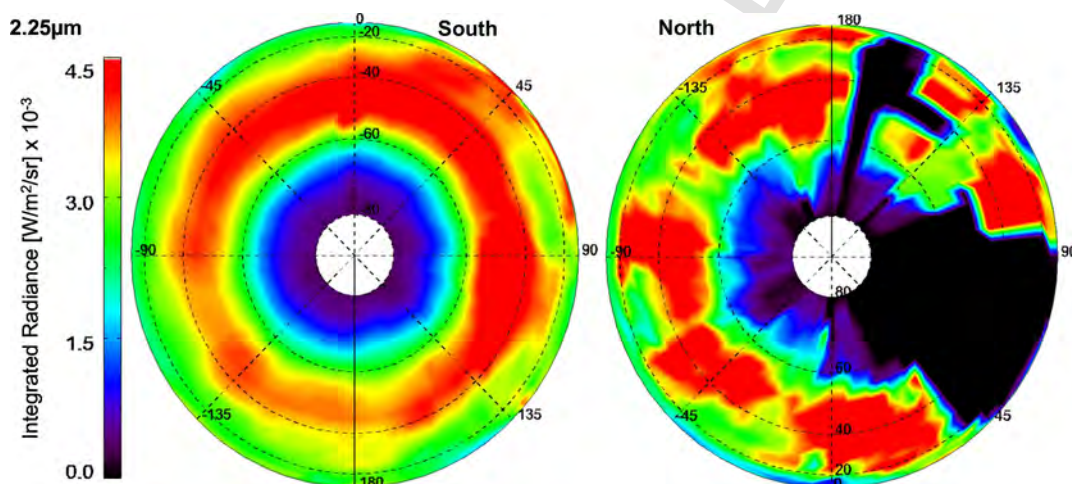


Fig. 12. Mean radiance integrated in the window centred around  $2.25\mu\text{m}$ , shown in an orthographic projection of the Venus southern (left) and northern (right) hemispheres, distributed along Latitude (pole in the centre) and Longitude (increasing eastward).

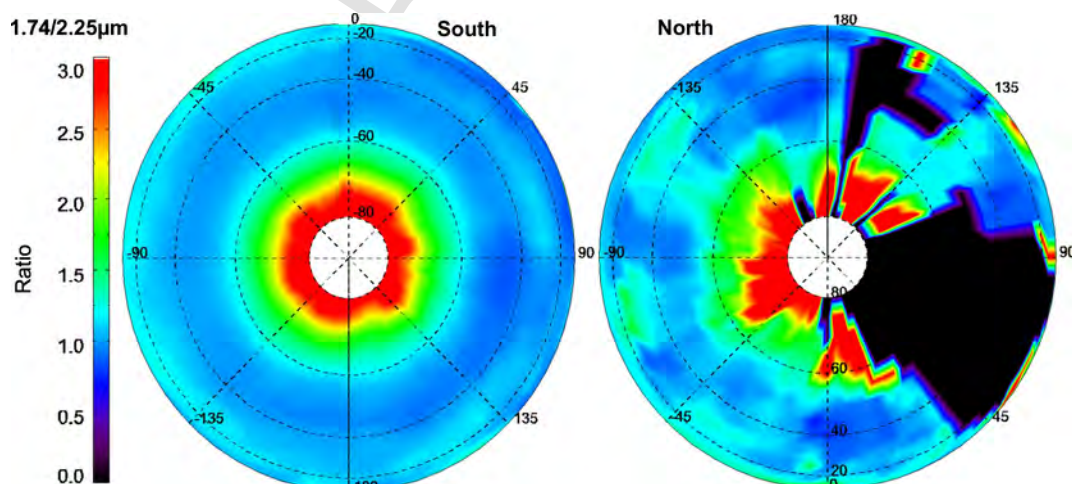


Fig. 13. Ratio between the radiances measured at  $1.74\mu\text{m}$  and  $2.25\mu\text{m}$ . Shown in an orthographic projection of the Venus southern (left) and northern (right) hemispheres, distributed along Latitude (pole in the centre) and Longitude (increasing eastward).

All input data are obtained directly from the VIRTIS calibrated data set corresponding to the nominal mission phase and first mission extension: [VEX-V-VIRTIS-2-3-V3.0] and [VEX-V-VIRTIS-2-3-EXT1-V2.0].

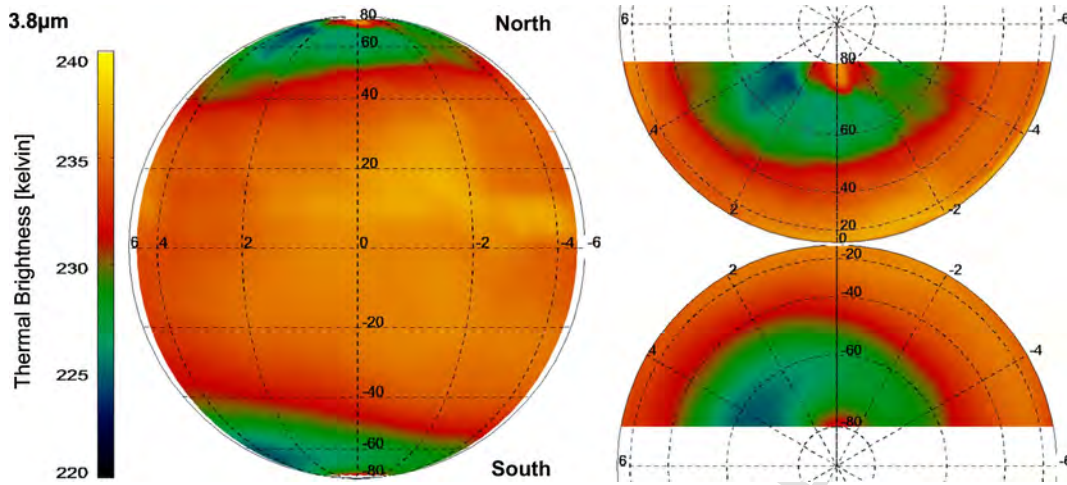


Fig. 14. Thermal Brightness mean map of Venus at  $3.8\ \mu\text{m}$ , shown in Orthographic projection, distributed along Latitude (north up, south down) and Local Time. The left projection is centred at the anti-solar point (Midnight, Local Time = 00h). The right projections are centred at the poles.

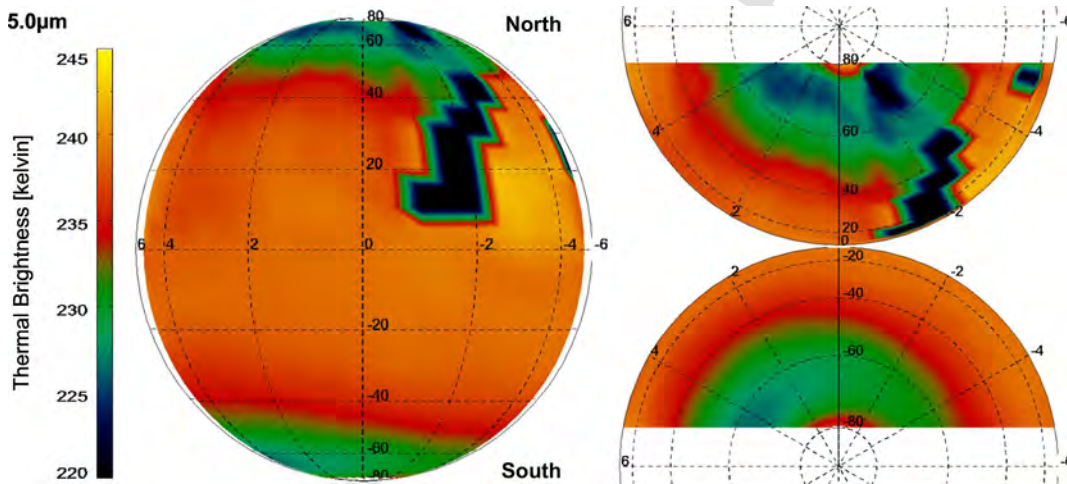


Fig. 15. Thermal Brightness mean map of Venus at  $5.0\ \mu\text{m}$ , shown in Orthographic projection, distributed along Latitude (north up, south down) and Local Time. The left projection is centred at the anti-solar point (Midnight, Local Time = 00h). The right projections are centred at the poles.

Input calibrated data are given in spectral radiance units ( $\text{W}/\text{m}^2/\text{sr}/\mu\text{m}$ ) as described by Cardesín-Moinelo et al. (2010).

### 2.1. VIRTIS M-IR coverage

The full coverage of all VIRTIS M-IR observations between May 2006 and October 2008, corresponding to orbits from 23 to 921, is shown in Figs. 1 and 2, distributed along local time and longitude respectively. The coverage distribution provides the visual statistical meaning of the data population for the different location of Venus.

As one can see, the southern hemisphere has very good coverage with several hundred observations per sample bin, whereas the coverage over the northern hemisphere is more patchy, with typically  $<10$  observations per bin and many gaps especially in the longitudinal coverage. The dense coverage of the southern hemisphere is caused by the apoapsis of the Venus Express orbit, which reaches its highest altitude ( $\sim 66,000\text{ km}$ ) near the south pole. In nominal observation modes, the instrument Field Of View (FOV) is typically a square of  $64 \times 64\text{ mrad}$ , which means that one single cube of observation can typically cover a projected square of  $\sim 4000 \times 4000\text{ km}$  as seen from the apoapsis. Indeed, at this distance the whole planet can be covered simply with nine observations, usually known as a “mosaic” and very common in the operations planning (Titov et al., 2006).

The more limited coverage in the northern hemisphere shows some of the gaps with regions that have never been observed during the mission. The pericenter situated over the North pole implies a spacecraft altitude very low over the northern hemisphere, spanning from  $5000\text{ km}$  at equator down to  $250\text{ km}$  or less near the north pole. At these low distances, the  $64\text{ mrad}$  wide FOV covers only  $320\text{ km}$  for a nadir observation at the equator and only  $16\text{ km}$  near the pericenter, with relatively high spatial resolution. In the northern hemisphere, the high relative speed of the spacecraft with respect to Venus, prevents the use of the scanning mirror to build an image, due to the low persistence of the target. For this reason, the observations near the pericenter are usually performed in push-broom mode, with the scanning mirror fixed at the boresight central position, pointing towards nadir or nearby, with increasing speed as it approaches the pericenter. This makes the footprints at the northern hemisphere appear like narrow stripes very extended in latitude.

Note also that there is a bias in the east-west longitudes which appear both in the northern and southern hemispheres, especially with big coverage gaps in the eastern-northern longitudes [ $0^\circ$ – $180^\circ\text{E}$ ]. These gaps are even larger if we consider only the low exposure observations, which are roughly 40% of the total nightside observations and will impact the maps produced at  $5.0\ \mu\text{m}$  as we will see later. Coverage gaps are constrained by the planetary and spacecraft ephemeris, all

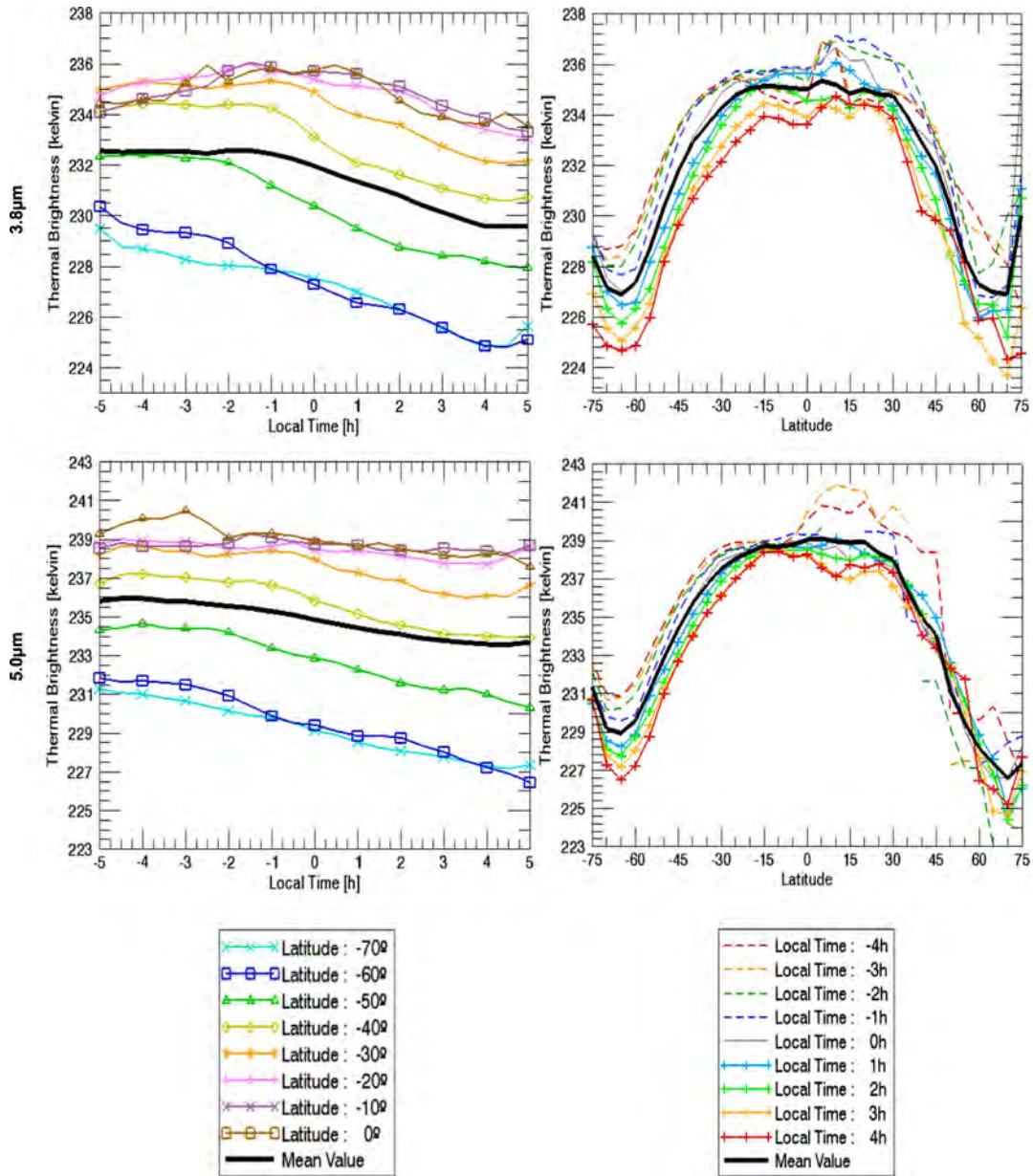


Fig. 16. Profiles of Thermal Brightness at 3.8 μm (top), 5.0 μm (bottom) plotted with respect to local time (left) and latitude (right). The regions close to the terminator and the polar caps are omitted since they are affected by scattered radiation.

together impacting the observations during certain periods. Unfortunately, these longitude coverage gap could not be fulfilled with the Infrared channel observations.

2.2. Data selection for infrared atmospheric windows

The radiance measured at 1.74 μm is a direct measurement of cloud opacity (Grinspoon et al., 1993; Titov et al., 2008) and so the whole band can be used for this analysis. However, the window around 2.3 μm contains important secondary components from the CO distribution and other chemicals (H<sub>2</sub>O, OCS, SO<sub>2</sub>) (Marcq et al., 2008). During our analysis we have performed various tests using different band widths to consider the effect of various species and we have seen that the impact of these trace gases is negligible for our study. In any case, to avoid any potential ambiguity, we focus our results only in the region 2.18–2.30 μm that is dominated by CO<sub>2</sub> as proposed by Tsang et al. (2008, 2009). Note also the advantage that the radiance

integrated over the two proposed windows has comparable values and so they can be shown using the same scale as in Fig. 4 (Fig. 3).

In Fig. 4 we can distinguish similar cloud features appearing at both wavelengths, although with slightly different contrast. Note in particular the darker polar region from the cold collar to the pole (latitudes from -70°S to -80°S) seen around 2.25 μm (right) with respect to the 1.74 μm (left). These images are shown as obtained directly from the calibrated data file in spectral radiance, without any projection or limb darkening correction. The input spectral radiance obtained from the calibrated data is integrated over the length of the windows and multiplied by the spectral width of the band (10<sup>-2</sup> μm) to convert to integrated radiance units (W/m<sup>2</sup>/sr).

The main data selection parameters used for the integrated radiance of both atmospheric windows around 1.74 μm and 2.25 μm are listed in Table 1.

The selection of wavelengths is done using a different range of bands for each case. The variability in the wavelength range (±10 nm) needs to be considered in each individual observation to correctly

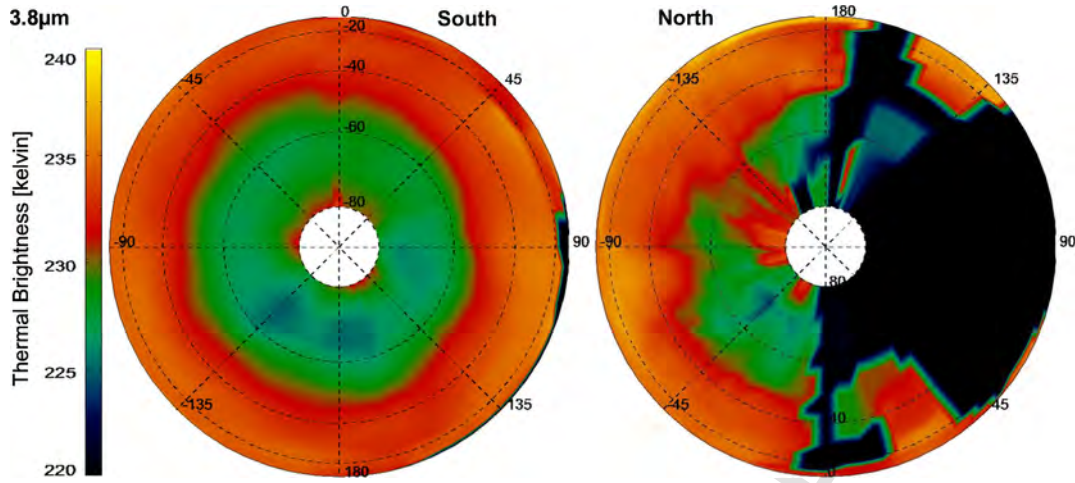


Fig. 17. Map of mean thermal brightness at  $3.8\mu\text{m}$ , shown in an orthographic projection of the Venus southern (left) and northern (right) hemispheres, distributed along Latitude (south pole in the centre) and Longitude (increasing eastward).

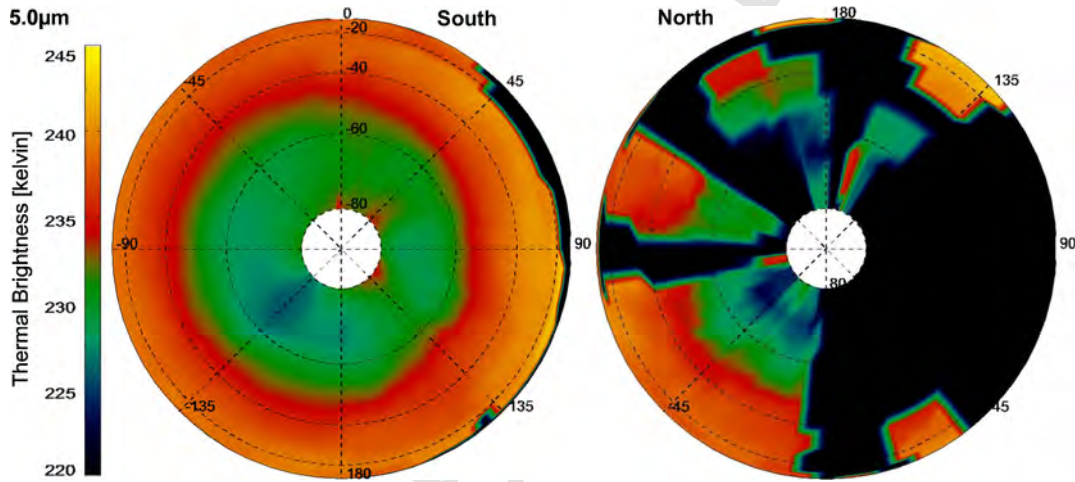


Fig. 18. Map of mean thermal brightness at  $5.0\mu\text{m}$ , shown in an orthographic projection of the Venus southern (left) and northern (right) hemispheres, distributed along Latitude (south pole in the centre) and Longitude (increasing eastward).

accommodate the temperature differences that may affect the spectral registration in different measurements (Cardesín-Moinelo et al., 2010). In any case the selection of bands has been done carefully with enough margins to accommodate for temperature variations between or within observations. The extreme bands outside each window have been used to subtract the continuum and remove the effect of background thermal contribution, avoiding any additional potential problems related to temperature.

In our study, we have used all the available data but we have used a filter of incidence angle to make sure we only consider data from the nightside. The filter value used ( $>100^\circ$ ) includes a margin of  $10^\circ$  over the terminator ( $90^\circ$ ) in order to exclude not only the data from the dayside but also any possible radiance scattered close to the terminator or the poles. For this reason we focus our analysis within  $\pm 5\text{h}$  in local time and  $\pm 75^\circ$  in latitude.

The limb darkening correction has been performed to the data following the functions and average coefficients given by Longobardo et al. (2012) and applied to the data at all latitudes over the planet. In order to avoid problems with very high emission angles, an extra filter is applied to take into account only the data below  $80^\circ$ , that is ignoring all data close to the limb (emission angle  $90^\circ$ ).

### 2.3. Data selection for thermal brightness analysis

The two thermal regions of interest,  $3.8\mu\text{m}$  and  $5.0\mu\text{m}$ , are shown in Fig. 5. VIRTIS observations at these wavelengths, shown in Fig. 6 usually show a strong correlation as the radiance at these bands depends almost purely on the thermal emission from the cloud tops at about the same level.

We can see similar features from the  $3.8\mu\text{m}$  (Fig. 6 left) and the  $5.0\mu\text{m}$  (Fig. 6 right), although with slightly different contrast. The images are shown as obtained directly from the calibrated data file without any projection or correction, only averaged and converted to thermal brightness in kelvin.

The main data selection parameters used for the thermal brightness both at  $3.8\mu\text{m}$  and  $5.0\mu\text{m}$  are reported in Table 2.

The selection of wavelengths is done using a different range of bands for each case. Again the variability in the wavelength range ( $\pm 10\text{nm}$ ) is mentioned to accommodate for possible variations in the spectral registration due to temperature differences.

The input data are obtained directly from the VIRTIS calibrated cube radiance, converted into Thermal Brightness [K], supposing an emissivity equal to 1.

Again we have used all the available data and applied the same filter to the incidence angle ( $>100^\circ$ ) to consider only nightside data

and avoid any possible scattered radiance from the terminator or poles, focusing our analysis within  $\pm 5$  h in local time and  $\pm 75^\circ$  in latitude.

The limb darkening correction for the thermal emission coming from the upper clouds is more complex than in the case of the lower clouds and has several limitations. The available methods proposed in Longobardo et al. (2012), Grinspoon et al. (1993), Roos et al. (1993) are valid only in a restricted range of emission angles and the proposed coefficients are highly dependent on latitude. In order to address this problem we have first reduced the filtering threshold for emission angles below  $60^\circ$ , adding extra margin to avoid deviations at higher angles. Then we have performed tests with many correction coefficients and we have selected the values that, when applied to the whole planet, produced the optimal results based on consistency in comparison with recent thermal studies by Grassi et al. (2008, 2010, 2014), Migliorini et al. (2012), Haus et al. (2013) and the review by Limaye et al. (2018b). In the case of  $3.8\ \mu\text{m}$  we have selected the coefficients (0.13/0.87), which were already proposed both by Longobardo et al. (2012) and Grinspoon et al. (1993) for the  $3.72\ \mu\text{m}$  wavelength. For  $5.0\ \mu\text{m}$  we have selected the coefficients (0.20/0.80), originally proposed by Longobardo et al. (2012) for the  $4.0\ \mu\text{m}$  wavelength at low latitudes.

As previously mentioned, the thermal radiance around  $5\ \mu\text{m}$  is heavily affected by the instrument thermal noise that saturates at high exposure times, so only the short exposure observations ( $< 0.36$  s) can be used for this wavelength, with some important implications. First the number of observations at  $5\ \mu\text{m}$  is much lower than for other wavelengths, so the coverage is reduced significantly, especially in the northern hemisphere where the coverage is already limited by the observing geometry. Secondly, the low exposure observations are heavily affected by an instrumental anomaly called “odd-even effect”, caused by small differences in the responsivity of the infrared sensor (Cardesín-Moinelo et al., 2010). This affects especially the wavelength region around  $5\ \mu\text{m}$ , producing a “sawtooth” spectral irregularity, well visible in Fig. 5. In order to limit this effect, we have used only odd bands, ignoring all even bands in the  $5\ \mu\text{m}$  region. Note this instrumental effect does not affect other wavelength regions, where the overall effect can be ignored because the exposure times can be longer as well.

### 3. Results: integrated radiance in the atmospheric windows

#### 3.1. Global maps of the infrared atmospheric windows

The mean global maps distributed along Latitude and Local Time are shown in Fig. 7, for the integrated radiance in the window centred at  $1.74\ \mu\text{m}$  and Fig. 8 for the window centred at  $2.25\ \mu\text{m}$ . These maps have been produced using a grid of  $5^\circ$  in latitude and 30 min in local time. To facilitate the comparison between these two bands we have also produced a map with the ratio between  $1.74/2.25\ \mu\text{m}$ . This ratio is meant to facilitate the comparison between the two bands and it also has the advantage that it can provide an indirect estimation of the particle size as described by previous studies (Carlson et al., 1993; McGouldrick and Toon, 2008; Wilson et al., 2008; Haus et al., 2013). Note however that a detailed analysis of the particle size is out of the scope of this study and we only provide here a basic estimation of the relative size spatial distribution assuming that the averaged region has a uniform cloud cover. This is a very simplistic approach with many limitations but it can still be useful to identify the main global trends in the atmosphere that we will discuss in the following sections compared to previous studies.

Maps for both infrared windows are well correlated and show basically the same global characteristics, in particular the increase of radiance at mid-latitudes, with a relatively good symmetry in both hemispheres. The more patchy distribution in the north is likely caused by the more limited coverage as previously mentioned. The increase of

radiance translates into a more transparent clouds deck in the same region, very likely associated to downwelling due to the meridional cell. On the contrary, the low radiance is where the clouds thickness is maximum, like in the equatorial region and even more pronounced in both polar regions.

Note that nightside data has been filtered with an incidence angle threshold  $> 100^\circ$ , including a margin to reduce any possible straylight or scattering from the dayside terminator. Note in particular that the data around both poles ( $> 80^\circ$  and  $< -80^\circ$ ) and around the terminators is outside this threshold and therefore it is not shown to avoid misinterpretations.

The general features and trends can be similarly seen in both spectral regions but with small and important differences in the details. For example, there is a different contrast especially at high latitudes, due to the cold collar around the polar region (Piccioni et al., 2007; Carlson et al., 1993) and the increase of abundance of large particles within the polar vortex, which enhances a lot the contrast in the  $2.25\ \mu\text{m}$  with respect to the  $1.74\ \mu\text{m}$  (Wilson et al., 2008).

#### 3.2. Detailed infrared atmospheric window profiles

We now show a more detailed trend distribution of the integrated radiances, with mean profiles with respect to latitude and local time, as in Fig. 10.

The uncertainty of the spectral radiance measurements by VIR-TIS-M-IR is below  $10^{-3}\ \text{W}/\text{m}^2/\text{sr}/\mu\text{m}$  (Piccioni et al., 2008; Cardesín-Moinelo et al., 2010) with a bandwidth of  $\sim 10$  nm. This translates into an integrated radiance below  $10^{-5}\ \text{W}/\text{m}^2/\text{sr}$ , therefore random measurement noise is negligible for our study and no error bars are shown in the profiles. Other possible sources of non-random errors, in particular systematic offsets or potential errors in the data calibration or the data processing may still be present in the results. In any case, the main source of error in our profiles is driven by the large fluctuations of the atmosphere. This is well visible in the latitude profiles for the northern hemisphere with very large variations caused by the lower coverage in the north, which make the atmospheric fluctuations more evident with respect to the smoothed, averaged values of the southern hemisphere. Also the scattered radiation coming from the strong illuminated day side affects all the averaged measurements near the terminator and the poles and therefore the data close to the terminator is not used in the analysis.

These profiles show the general latitudinal trend with low radiance (dense lower clouds) at the equator and the poles, and a relatively higher radiance (thinner lower clouds) at mid latitudes, around  $\pm 45^\circ$ . Interestingly, the variability from the relative maximum peaks at mid latitudes and the minimum at the equator is about the same for both wavelengths. Indeed the latitudinal variation of the ratio is very flat in mid-to-low latitudes, indicative of a very similar opacity and particle size in these regions. This is very different at high latitudes around the minimum radiance values at  $\pm 75^\circ$ , where  $2.25\ \mu\text{m}$  integrated radiance gets 2.5–3 times dimmer than at  $1.74\ \mu\text{m}$ , indicating a large particle size, very symmetrical at both hemispheres despite the coverage limitations. Another interesting feature is the slight asymmetry of the peaks at mid latitudes, with higher values in the north ( $+75^\circ$ ) than in the south ( $-75^\circ$ ).

The mean radiance distribution in local time can also be seen in Fig. 10. The steep increase of radiance on the left and right sides close to the terminator is explained again by scattering contamination from the day side, which affects the weak signal from the night side. The central part, from  $-4$  h to  $+4$  h, is relatively flat with the exception of sub-tropical ( $\pm 30^\circ$ ) and sub-polar ( $\pm 60^\circ$ ) latitude regions, which show an interesting peak around midnight. The higher contrast in radiance at  $2.25\ \mu\text{m}$  with respect to  $1.74\ \mu\text{m}$  between the polar and equatorial regions is obviously present also here.

### 3.3. Longitudinal views of infrared atmospheric window

We will now present the nightside integrated radiance distributed in a longitudinal view for the southern and northern hemispheres. The maps are shown in Figs. 11 and 12 for the bands 1.74 and 2.25  $\mu\text{m}$  respectively. Both bands are projected spatially with respect to Latitude and Longitude using a grid of  $5 \times 5^\circ$  and centred at each pole. This perspective visually enhances the global extension of the high radiance at mid-latitudes, as well as the difference in relative radiance near the poles between the two atmospheric windows.

A preliminary version of these polar views was previously presented for the southern hemisphere by Cardesin-Moinelo et al. (2008) and briefly discussed by Titov et al. (2008). These new maps have been reprocessed covering both hemispheres and include various updates, in particular we now use better emission angle corrections based on Longobardo et al. (2012), new grid ( $5 \times 5^\circ$  in latitude/longitude) and some other adjustments and filtering of incorrect and meaningless data.

All these improvements provide better estimations of the mean radiance values and allow for a better comparison of both wavelengths. We can remark again the global latitudinal trends that are common to both wavelengths, and the much larger depletion region around the pole for the 2.25  $\mu\text{m}$  radiance with respect to the 1.74  $\mu\text{m}$  one.

We note here that the longitudinal maps show some irregularities and asymmetries with respect to longitude, but any assumption regarding surface-cloud interactions is out of the scope of this study and in any case not obvious from our results. Several other investigations have already discussed interesting dependencies between the surface and the clouds, wind variations localized around certain regions as identified with Venus Express data by Bertaux et al. (2016) and Peralta et al. (2017b), and the recent studies based on Akatsuki bow-shape features by Fukuhara et al. (2017), and Kouyama et al. (2017), complemented also with modelling analysis by Kashimura et al. (2019).

## 4. Results: brightness temperature of infrared thermal emission

### 4.1. Global maps of infrared thermal emission

We will now present the global maps that have been done for the thermal brightness. These maps are shown in Fig. 14 for 3.8  $\mu\text{m}$ , and Fig. 15 for 5.0  $\mu\text{m}$ , with orthographic projections centred at the anti-solar point and using a grid of  $5^\circ$  in latitude and 30 min in local time.

Note again that nightside data has been filtered with an incidence angle threshold of  $>100^\circ$ , including a margin to reduce any possible straylight or scattering from the dayside terminator. Data around both poles ( $>80^\circ$  and  $<-80^\circ$ ) and around the terminators is outside this threshold and not shown in the maps.

The 5.0  $\mu\text{m}$  map has large gaps especially in the northern hemisphere, as this wavelength region was saturated for long exposure times and therefore it includes a smaller number of observations with respect to the 3.8  $\mu\text{m}$  map. The two maps have a very similar appearance, also quantitatively, and both show consistently a nice distribution of the cloud tops temperature of the full Venus night side.

Very interesting features are evident in the maps. First the uniform temperatures at low latitudes, with no evident variability in latitude or local time. Then at mid-to-high latitudes we start seeing the temperatures decreasing smoothly from the evening to the morning side in both hemispheres, especially the collar region centred at  $\pm 60^\circ$  around both poles. This clear symmetry in latitude and local time for both hemispheres is consistent with the behaviour seen by Grassi et al. (2010, 2014), and Migliorini et al. (2012) and the modelling simulations by Haus et al. (2013), and Garate-Lopez and Lebonnois (2018).

### 4.2. Detailed thermal brightness profiles

The mean thermal brightness profiles with respect to latitude and local time are reported in Fig. 16. In particular, we show the mean profiles at 3.8  $\mu\text{m}$  and 5.0  $\mu\text{m}$  at different latitudes versus local time and different local times versus latitude. The profiles of both bands are qualitatively and also quantitatively consistent with each other with very small differences. As previously mentioned, the uncertainty of the input spectral radiance values measured by VIRTIS-M-IR is negligible compared to the large fluctuations of the atmosphere, therefore error bars are not shown in these profiles. We remark that both bands have the minima in the cold collar at about  $\pm 60\text{--}70^\circ$  of latitude and the relative maximum around the equator. Regarding the differences between both bands, we note a lower thermal brightness of the northern cold collar seen at 5.0  $\mu\text{m}$  and many other small differences in the northern hemisphere, which can all be explained by the lower coverage.

The local time profiles shown in Fig. 16 are also consistent for the two wavelengths. The more evident behaviour of these profiles are the colder temperatures seen in the morning side (right) with respect to the evening side (left) with a very smooth and regular trend from one side to the other, with a difference of about 3 to 5 K in average. It is also remarkable also the fact that the largest dependency with local time corresponds to high latitudes where the Sun's elevation is very low and the average rotation period of Venus atmosphere is smaller with respect to lower latitudes, therefore the differences between day and night at high latitudes would not be expected to be so strong.

### 4.3. Longitudinal views of thermal brightness

We will now present the nightside thermal brightness distributed in a longitudinal view for the southern and northern hemispheres. These maps are shown in Fig. 17 for 3.8  $\mu\text{m}$  and Fig. 18 for 5.0  $\mu\text{m}$ . All maps are projected spatially with respect to latitude and longitude using a grid of  $5 \times 5^\circ$  and centred at each pole.

These polar views illustrate the uniform temperature brightness at low-to-mid latitudes and the sharp decrease in the cold collar around both poles. We can note an interesting asymmetry around the south pole with the cold collar not being uniform along the longitudes, visible in both bands. We cannot take any conclusions regarding this potential shift as this could be a long term evolution of the meridional displacement of the cold collar, not enough covered by the Venus Express data range with time, and the asymmetry of the data coverage in the eastern/western hemispheres as visible in Fig. 2. As previously mentioned, any further assumption regarding surface-cloud interactions is out of the scope of this work and one should refer to other relevant studies mentioned in Section 3.3.

## 5. Discussion

All the maps indicate generally four very well-defined regions: the equatorial belt, the mid-latitude regions, the cold collars and the polar regions.

### 5.1. Equatorial region (latitudes $0\text{--}30^\circ$ )

The maps show an equatorial belt with relatively high cloud opacity. The equatorial belt of Venus is dominated by generally thick clouds (Titov et al., 2018). The low integrated radiance observed all around the equator is explained by the thermal radiation coming from the layers below the lower clouds, getting absorbed by the thick equatorial clouds, causing lower values in average at the equator. Both atmospheric windows show a very similar profile and indeed the band ratio shows no significant variability in this region, implying a uniform

spatial distribution of the particle sizes with local time or latitude around the equator, consistent with Wilson et al. (2008).

The thermal emission maps show a uniform brightness temperature at the cloud tops, with very high values and low variability around the equator consistent with the intersection of the expected thermal profiles (Limaye et al., 2018b). This is well in agreement with the smooth symmetric temperature profiles reported by Migliorini et al. (2012) and Haus et al. (2013) at these latitudes, and it is also consistent with the altimetry of the clouds retrieved in Ignatiev et al. (2009). Regarding the local time variations in this low latitude region, the profiles are also very flat and there are very small variations between evening and morning hours, except those attributed to the limited coverage at these latitudes.

### 5.2. Mid-latitude region (latitudes 30–60°)

The mid latitude region, centred between  $-30^\circ$  and  $-60^\circ$  of latitude, shows values of integrated radiance higher than the surroundings in average. The lower opacity of these thinner clouds are due to a more regular laminar flow motion described in Sánchez-Lavega et al. (2017) and the deep radiance penetrates the clouds towards the outer space relatively easier in this region (Titov et al., 2018). This is consistent with previous observations from Earth, see Tavenner et al. (2008), and results from other missions as summarized by Peralta et al. (2018). The mid latitude is then the region where the highest dissipation of the surface or close to surface heat can take place, an important point for the differential (meridional) and absolute thermal balance of the planet (Limaye et al., 2018b).

This region is very evident in the integrated radiance maps (Figs. 7 and 8) and profiles (Fig. 10) getting to the maximum values around  $\pm 45^\circ$  latitude, with values exceeding  $4 \times 10^{-3} \text{ W/m}^2/\text{sr}$  for both the  $1.74 \mu\text{m}$  and  $2.25 \mu\text{m}$  bands, with no significant variation in the ratio between them. This is a transition zone where the atmospheric flow is more laminar and where there is a “clearing” effect on the clouds from the boundary of the cold collar and polar vortex dynamics. In this respect, Kashimura et al. (2019) demonstrates the morphological similarity between the night-side bright streaks and down-welling region in high-resolution numerical simulation. Our maps of integrated radiance maps show a symmetric behaviour in both hemispheres. The maximum values are slightly asymmetric, higher in the north in both wavelengths ( $1.74 \mu\text{m}$  and  $2.25 \mu\text{m}$ ) and is consistent with recent Akatsuki IR2 images at  $2.3 \mu\text{m}$  that showed differences between both hemispheres at low latitudes (Peralta et al., 2018, 2019), although these reports were not systematic analyses of long term observations like the results presented in this work.

We may also discuss here the interesting bulge in radiance observed around midnight in the sub-tropical ( $\pm 30^\circ$ ) and sub-polar ( $\pm 60^\circ$ ) latitudes, shown in Fig. 10. This higher values at the boundaries of the mid-latitude region appear consistently in both bands,  $1.74 \mu\text{m}$  and  $2.25 \mu\text{m}$ , with no clear variation in the ratio between them, and it is very reliable since the coverage in the southern hemisphere is very good with nearly a thousand observations. On the other hand this is not visible in the northern hemisphere, although this could be a consequence of the lower coverage and higher variability. This interesting increase in the radiance around midnight can be explained by an enlargement in the meridional direction of the mid-latitude thin cloud belt towards the equator. We remark that this potential expansion has not been reported in previous studies and shall be studied in future simulations with General Circulation Models.

It is worth mentioning that despite the great variation in the integrated radiance from low to mid-latitudes, the variations of band ratio stay almost flat and symmetric between  $60^\circ\text{S}$  and  $60^\circ\text{N}$ . In general this is consistent with Wilson et al. (2008) and Barstow et al. (2012) that reported that latitudinal variations of cloud thickness are not directly correlated with the particle size in this region. When looking at our

detailed profiles, we can identify a small decrease of the ratio by 10–15% around  $\pm 60^\circ$  with respect to the equator, and again symmetric in both hemispheres. This decrease in the ratio is only a potential indication of smaller particle size, but is in good agreement with the thinner mid latitude clouds reported by Haus et al. (2013).

Regarding the local time dependencies, the integrated radiances do not present significant variations and all the mid latitude range has similar values. On the contrary, the thermal brightness profiles (Fig. 16) show a clear trend of the cloud top temperatures decreasing with local time at these latitudes (see for example the local time profile at  $-50^\circ$  latitude in both bands), decreasing by about 2K from the evening terminator towards midnight and then decreasing furthermore up to 5K towards the morning terminator. This is in line with the local time variations described for the southern hemisphere by Grassi et al. (2010, 2014) and our results show a good symmetry with the northern hemisphere, especially at  $3.8 \mu\text{m}$  (Fig. 14). The causes of this behaviour and the correlation with thermal tides have recently been investigated with modelling simulations (see (Ando et al., 2016) and references therein). Also recent observations with Akatsuki/LIR data, reported by Kouyama et al. (2019), show the global structure of the thermal tides at the cloud top level, and are consistent with the local variation of temperature obtained from our VIRTIS thermal emission data.

The distribution of radiance and thermal brightness in the longitudinal maps (Figs. 11, 12, 17 and 18) offer a good illustration of the global extension of the thin clouds at mid-latitudes in both hemispheres and the sharp decrease, both in radiance and in cloud top temperatures, when and entering the cold collar around  $\sim 60^\circ$  latitude.

### 5.3. Cold collar region (latitudes 60–75°)

The convection-driven Hadley cell extends up to about  $60^\circ$  latitude in both hemispheres moving with a global trend from the equator to higher latitudes, feeding a mid-latitude jet at its poleward extreme, inside which there is a circumpolar belt characterized by remarkably low temperatures and dense high clouds. This is the so-called “cold collar” region (Taylor et al., 1980; Piccioni et al., 2007; Titov et al., 2018; Sánchez-Lavega et al., 2017), located between  $60^\circ$  and  $75^\circ$  latitude, which is clearly visible in all global maps for its lower radiance and brightness temperature.

The cold collar represents the important dynamical boundary between the polar region and the mid to low latitudes of Venus. The cold collar is suggested to be caused by a zonal jet produced by the mean meridional circulation, as simulated through a General Circulation Model (Lee et al., 2005). More recently, the models by Ando et al. (2016, 2019) and Garate-Lopez and Lebonnois (2018) described the adiabatic expansion due to the Hadley cell and the polar vortex, with also vertical motion of air and the superposition of the zonally averaged basic field, thermal tides and transient waves.

The profiles we have obtained at  $60^\circ\text{S}$  and  $70^\circ\text{S}$  for both thermal emissions (Fig. 16) suggest a local time variability of the temperature in the cold collar region with an evident decreasing trend from the evening towards morning hours. The brightness temperature reach the lowest values in these latitudes, decreasing from  $\sim 232\text{K}$  at  $-5\text{h}$  down to  $\sim 226\text{K}$  at  $+5\text{h}$  for the  $5.0 \mu\text{m}$  band, and from  $\sim 230\text{K}$  to  $\sim 225\text{K}$  for the  $3.8 \mu\text{m}$  band. This cooling trend was first reported in the northern hemisphere by Taylor et al. (1980) and then studied with detailed temperature fields for the southern hemisphere by Grassi et al. (2010, 2014). Our results provide a very good latitudinal symmetry between northern and southern cold collars, especially evident at  $3.8 \mu\text{m}$  (Fig. 14). A latitudinal symmetry was already described by Migliorini et al. (2012) and Haus et al. (2013) considering only averaged temperature fields without local time variation, and our results now confirm also that the local time profiles of both south and north cold collars have equivalent cooling trends from the evening towards the morning side. The causes of this local time dependence and the correlation with

thermal tides are again described with modelling simulations (see (Ando et al., 2016, 2019), (Garate-Lopez and Lebonnois, 2018) and references therein).

The cold collar region is also the region where the opacity of the clouds starts to increase with consequent lower radiance seen through the atmospheric windows. The integrated radiances for both atmospheric windows decrease from their maximum values at 60°S latitude (both around  $\sim 4 \times 10^{-3} \text{ W/m}^2/\text{sr}$ ), to their minimum values at 75°S,  $\sim 1.5 \times 10^{-3} \text{ W/m}^2/\text{sr}$  in the 1.74  $\mu\text{m}$  band, and even lower to  $\sim 0.5 \times 10^{-3} \text{ W/m}^2/\text{sr}$  in the 2.25  $\mu\text{m}$  band. These difference in the trend is clearly visible in the ratio latitudinal profile shown in Fig. 9. The ratio is almost uniform and flat between 60°S and 60°N but increases sharply when entering the cold-collar region. This is caused by the different spatial distribution of size particles as explained by Carlson et al. (1993), with bigger particles depleting more the radiance at 2.25  $\mu\text{m}$  due to the higher imaginary part of the refractive index. Our maps of band ratio (Figs. 9 and 13) illustrate this particle size distribution and confirm that the concentration of bigger particles are symmetric in the southern and northern polar regions as reported by Wilson et al. (2008) and Haus et al. (2013).

In general we do not observe a dependence with local time inside the cold collar for the integrated radiance windows. In any case we can remark a 10% increase in the local time profile of the ratio at 70°S (Fig. 10, going from 1.9 at LT - 4h to 2.1 at LT + 4h). At this high latitude we cannot exclude a possible contamination by the scattered light from the terminator, but this could be an indication of a possible increase in the size of the particles from the evening towards the morning following the general atmospheric circulation and downwelling inside the cold collar (Titov et al., 2008; Sánchez-Lavega et al., 2017; Titov et al., 2018; Garate-Lopez and Lebonnois, 2018; Lebonnois et al., 2018).

#### 5.4. Polar region (latitudes 75–90°)

The study of this latitude region in our global maps is highly affected by the scattered radiation coming from the illuminated day side, so we avoid any interpretations close to the terminators and especially above 75° in latitude. What is evident is the big particle size and the higher thermal brightness seen in the polar regions up to at least 75° latitude in both hemispheres, consistent with the subsidence with clouds depression reported by Titov et al. (2008), Ignatiev et al. (2009), and Cottini et al. (2012, 2015). Most of the studies by Venus Express have focused in the Southern polar region but here we can remark the very good symmetry between the South and North polar regions, with no significant differences observed neither in the cloud top temperatures nor in the cloud opacity, except for the higher variability in the northern hemisphere due to lower coverage. This good symmetry is visible in most of the polar maps, and also in the latitudinal profiles in Figs. 10 and 16.

The polar regions are characterized by clouds thicker than in the cold collar, as reported by Lee et al. (2012) and possibly explained with mean meridional circulation responsible of cloud material transport (Imamura and Hashimoto, 1998), yielding to higher temperature and adiabatic compression of the vortex dynamics, as discussed by Garate-Lopez et al. (2015). Recent observations with IR2 on Akatsuki at latitudes  $> 60^\circ$  showed features darker than what observed at 2.02  $\mu\text{m}$ , close to the cold collar borders, which are indicative of lower cloud top altitudes, due to the  $\text{CO}_2$  absorption of reflected light (Limaye et al., 2018a).

The polar vortex is only partially visible in our maps, and its study is out of the scope of this work. However, based on the good symmetry we observe between both polar regions, we can assume that the many attempts to understand the nature of the south polar vortex with VIRTIS data, starting by Piccioni et al. (2007), are also applicable to the North. We note in particular the study of the vortex motions discussed in great detail by Luz et al. (2011) and Garate-Lopez et al. (2013),

its thermal structure investigated by Grassi et al. (2008) and Garate-Lopez et al. (2015) and its potential vorticity described by Garate-Lopez et al. (2016).

## 6. Conclusions

The large dataset of the VIRTIS instrument, operated from May 2006 to October 2008 on board the Venus Express mission, provides an extensive and unprecedented view of the Venus nightside, in particular in the near infrared atmospheric windows and thermal emission radiance. Despite the known large variability of the Venus atmosphere and clouds morphology, our mean global maps show a clear distribution of the different dynamical regions of the planet seen throughout their modification of the appearance of the clouds opacity and cloud tops thermal brightness.

A very good and comprehensive coverage is obtained for the southern hemisphere, while a more limited coverage is achieved in the North due to orbital constraints. However the general overall features observed appear symmetric on both hemispheres.

The maps show an equatorial belt with relatively high cloud opacity, spatially homogeneous particle size and a uniform brightness temperature at the cloud tops, in agreement with the smooth symmetric temperature profiles reported by Migliorini et al. (2012) and Haus et al. (2013) and also consistent with the altimetry of the clouds retrieved in Ignatiev et al. (2009).

The mid latitude region shows values of integrated radiance higher than the surroundings in average due to the lower opacity of the thinner clouds, more regular and with laminar flow motion as described in Sánchez-Lavega et al. (2017) and Titov et al. (2018). This region also shows the smallest particle size compared to any other latitude range, confirming the results by Haus et al. (2013). Here the integrated radiances do not present significant variations, but we can remark a potential enlargement of the southern mid-latitude belt towards the equator, not reported in previous studies. The thermal brightness shows a clear trend of the cloud top temperatures decreasing with local time in both hemispheres.

In the cold collar region, the thermal emissions of the cloud tops show an evident decreasing trend from the evening towards morning hours visible at both southern and northern hemispheres with very good symmetry, especially noticeable at 3.8  $\mu\text{m}$ . This confirms both the average latitudinal symmetry reported by Migliorini et al. (2012) and Haus et al. (2013) and the local time variations that had only been studied separately for the northern (Taylor et al., 1980) and southern cold collar (Grassi et al., 2010, 2014). Also in the cold collar region the opacity of the clouds starts to increase reducing the radiance seen through the atmospheric windows, consistent with the mean meridional circulation responsible of cloud material transport (Imamura and Hashimoto, 1998; Sánchez-Lavega et al., 2017), yielding to higher temperature and adiabatic compression of the vortex dynamics (Garate-Lopez et al., 2015). We also observe a major increase in the relative abundance of big particles at both hemispheres at high latitudes, confirming previous results by Carlson et al. (1993), Wilson et al. (2008) and Haus et al. (2013) with a potential increase in particle size from evening to morning hours, possibly explained by the dynamical forcing of the polar vortex and the downwelling inside the cold collar (Sánchez-Lavega et al., 2017; Garate-Lopez and Lebonnois, 2018; Lebonnois et al., 2018).

The detailed analysis of the polar areas is out of the scope of this study but our results confirm a higher cloud opacity, higher accumulation of big particle sizes and higher thermal brightness in the polar region up to at least 75° latitude in both hemispheres, consistent with the subsidence clouds depression reported in other works (Titov et al., 2008; Ignatiev et al., 2009; Cottini et al., 2012, 2015) and the global atmospheric structure and dynamics summarized in the recent reviews by Sánchez-Lavega et al. (2017), Titov et al. (2018) and Limaye et al. (2018b).

The detailed physical interpretation of these maps can be done unambiguously only with the support of a full general circulation model of Venus but at the same time this work can support in refining the main parameters used for the model itself, especially for persistent long term mean features.

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## Data availability

All the infrared calibrated data from VIRTIS Venus Express used in this study is publicly available through the Planetary Science Archive hosted by the European Space Agency (<https://archives.esac.esa.int/psa>) in the datasets corresponding to the nominal mission phase and first mission extension: [VEX-V-VIRTIS-2-3-V3.0] and [VEX-V-VIRTIS-2-3-EXT1-V2.0].

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