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Planck 2018 results. II. Low Frequency Instrument data processing

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ABSTRACT

We present a final description of the data-processing pipeline for the *Planck* Low Frequency Instrument (LFI), implemented for the 2018 data release. Several improvements have been made with respect to the previous release, especially in the calibration process and in the correction of instrumental features such as the effects of nonlinearity in the response of the analogue-to-digital converters. We provide a brief pedagogical introduction to the complete pipeline, as well as a detailed description of the important changes implemented. Self-consistency of the pipeline is demonstrated using dedicated simulations and null tests. We present the final version of the LFI full sky maps at 30, 44, and 70 GHz, both in temperature and polarization, together with a refined estimate of the solar dipole and a final assessment of the main LFI instrumental parameters.

Key words. Space vehicles: instruments – Methods: data analysis – cosmic microwave background

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1. Introduction

This paper is part of the 2018 data release ('PR3') of the *Planck*[1] mission, and reports on the Low Frequency Instrument (LFI) data processing for the legacy data products and cosmological analysis. The 2018 release is based on the same data set as the previous release ('PR2') in 2015, in other words, a total of 48 months of observation (eight full-sky Surveys), more than three times the nominal mission length of 15.5 months originally planned (Planck Collaboration I 2011).

This paper describes in detail the complete data flow through the LFI scientific pipeline as it was actually implemented in the LFI data-processing centre (DPC), starting from the basic steps of handling raw telemetry (for both scientific and housekeeping data), and ending with the creation of frequency maps and validation of the released data products (similar information for the High Frequency Instrument [HFI] can be found in Planck Collaboration III 2018). Since this is the last Planck Collaboration paper on the LFI data analysis, in this introduction we provide a pedagogical description of all the data-processing steps in the pipeline. Later sections report in greater detail on those pipeline steps that have been updated, modified, or improved with respect to the previous data release. For the many steps that remain unchanged, the interested reader should consult Planck Collaboration II (2016).

Processing LFI data is divided into three main levels (see Fig. 1). In Level 1, the process starts with the ingestion of the required information from the telemetry data packets and auxiliary data received from the Mission Operation Centre; both the science and housekeeping information is then transformed into a format suitable for Level 2 processing. The goal of Level 2 is the creation of calibrated maps at all LFI frequencies in both temperature and polarization, with known systematic and instrumental effects removed. Finally, Level 3 requires the combination of both LFI and HFI data to perform astrophysical component separation (both CMB and foregrounds), extraction of CMB angular power spectra, and determination of cosmological parameters. This last level is not described in this paper: we refer readers to Planck Collaboration I (2018), Planck Collaboration IV (2018), Planck Collaboration V (2018), and Planck Collaboration VI (2018).

Level 2 includes three main blocks of the analysis pipeline: TOI processing or 'preprocessing'; calibration; and mapmaking. Preprocessing starts with flagging data that were made unusable due to lost telemetry packets and spacecraft manoeuvres. It continues with corrections for nonlinearity in the analogue-to-digital converters (ADCs) and for small spurious electronic signals at 1 Hz. The ADCs convert analogue output voltages from the detectors into digital form. Any departure from exact linearity creates a distortion in the response curve of the radiometer. The current implementation of the algorithm to correct for ADC nonlinearity includes improvements made since Planck Collaboration II (2016), which are described in Sect. 2. The 1-Hz electronic spikes result from an unwanted, low-level interaction between the electronic clock and the science data, and occur in the data-acquisition electronics after the acquisition of raw data from the radiometer diodes, and before ADC conversion (Meinhold et al. 2009; Mennella et al. 2010, 2011). They appear as a 1-Hz square wave, synchronous with the onboard time signal. The procedure for correcting the data is the same as described in Planck Collaboration II (2016), and consists of fitting and subtracting a 1-Hz square wave template from the time-domain data.

In the pseudo-correlation scheme adopted for the LFI radiometers (Bersanelli et al. 2010), each radiometer diode produces an alternating sequence of sky and reference load signals at the 4096-Hz phase-switch frequency. The $1/f$ noise of the sky and reference data streams are highly correlated. Subtracting the optimally scaled reference data stream from the sky data stream reduces the $1/f$ noise in the sky data by several orders of magnitude. We calculate this optimal gain modulation factor (GMF) using the same method as for the 2015 release (Planck Collaboration II 2016).

The last preprocessing step is diode combination. This reduces the impact of imperfect isolation between the two diodes of each LFI radiometer. The weighted combinations are unchanged since the 2015 release, and may be found in Table 3 of Planck Collaboration II (2016); typical values range between 0.4 and 0.6 (a perfect diodes isolation would yield 0.5 equal weights).

The pointing pipeline runs in parallel with the preprocessing pipeline just described. It uses the focal plane geometry, the spacecraft velocity and attitude, and 'PTCOR' a long-time-scale pointing correction (which takes account of both the distance from the Sun and thermometry from the Radiometer Electronics Box Assembly). The pointing pipeline reconstructs the pointing position and horn orientation for each sample in the data stream. PTCOR is unchanged from the 2015 release, and is described in Planck Collaboration I (2016).

The next step in the pipeline is photometric calibration. In addition to the pointing and sky-minus-load differences, auxiliary information is required to obtain accurate calibration. This includes the 4π beam response, a model of the CMB dipole together with the time-varying modulation of its amplitude induced by the motion of the spacecraft along its orbit — our primary calibration source — and a model of Galactic emission whose contribution through the beam far sidelobes is modelled and subtracted from each time line. The calibration process converts raw voltages at the output of the radiometers into thermodynamic temperatures. The basic calibration reference signal is the *Planck* dipole convolved with the full 4π beam response, properly weighted according to the bandpass of each radiometer. The *Planck* dipole used in this step is identical to the one employed in the earlier 2015 release (Planck Collaboration I 2016, see further details in the following sections). On the other hand, the calibration algorithm has been significantly improved compared to the previous release by including Galactic emis-

[1] *Planck* (http://www.esa.int/Planck) is a project of the European Space Agency (ESA) with instruments provided by two scientific consortia funded by ESA member states and let by Principal Investigators from France and Italy, telescope reflectors provided through a collaboration between ESA and a scientific consortium led and funded by Denmark, and additional contributions from NASA (USA).

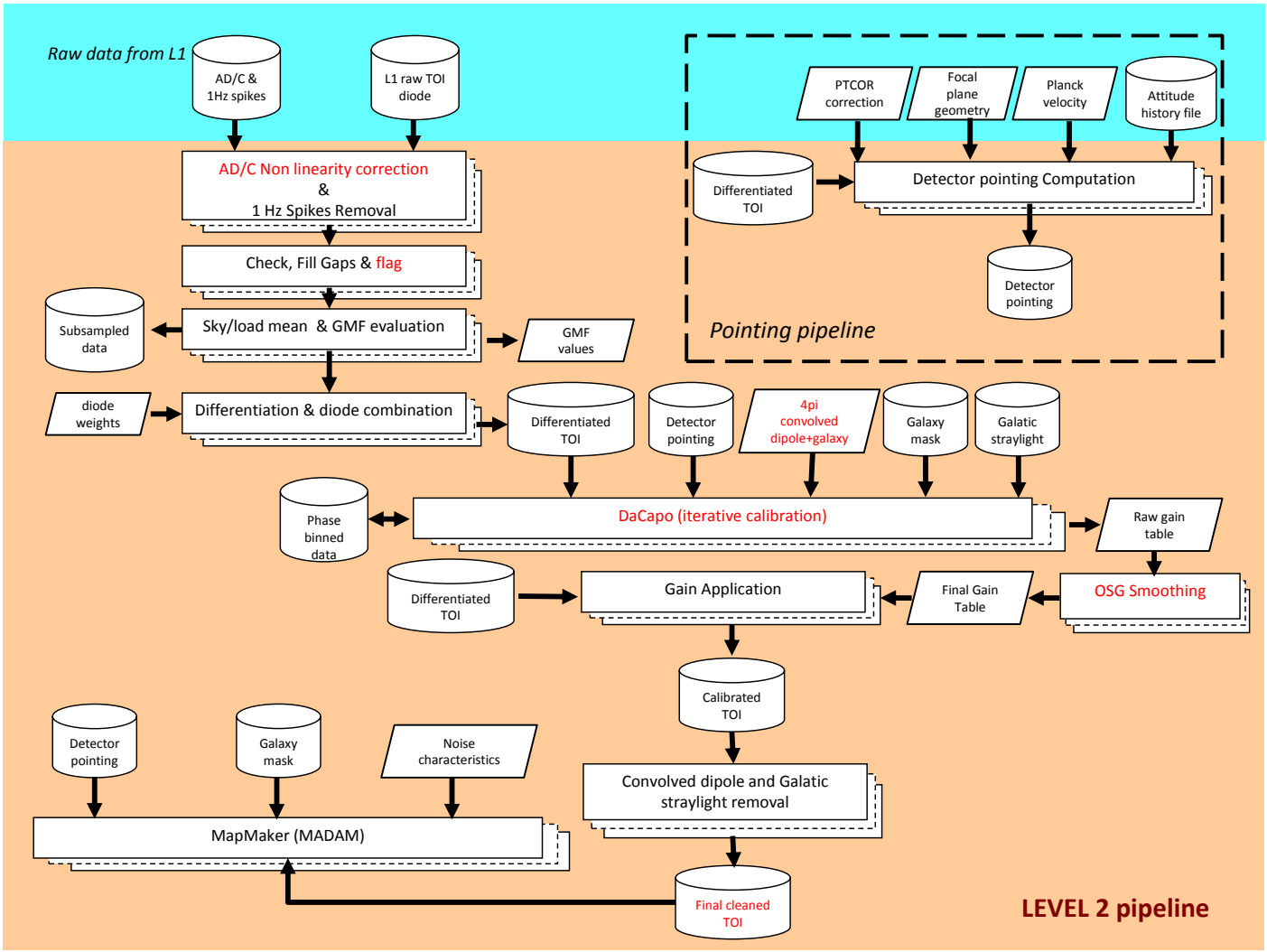


Fig. 1. Schematic representation of the LFI data-processing pipeline from raw telemetry down to frequency maps. Elements in red are those changed or improved with respect to Planck Collaboration II (2016).

sion along with the CMB dipole in the calibration model. This is particularly important during periods when the spacecraft spin axis is nearly aligned with the CMB dipole, and the variation of the dipole signal along scan circles is small ('dipole minima'). The final gain solution is obtained with an iterative destriper, `DaCapo`, which at each step determines radiometer gains, constraining the data to fit the dipole + Galaxy beam-convolved model. The output gain solutions are noisy during dipole minima (especially in Surveys 2 and 4). Therefore, as in the previous release (but with further optimization), we employ an adaptive smoothing algorithm that reduces scatter in gain solutions, but preserves real discontinuities caused by abrupt changes in the radiometer operating conditions. Finally, these smoothed gain solutions are applied to raw data streams, after subtraction of both the dipole and an estimated signal contributed by Galactic emission into the beam sidelobes.

The final step of the Level 2 pipeline is mapmaking, in other words, using the calibrated data and pointing information to create Stokes I, Q, and U maps of the sky at each frequency. The LFI mapmaking code is `Madam`, fully described in Keihänen et al. (2005) and Planck Collaboration VI (2016), which removes correlated $1/f$ noise with a destriping approach. Correlated noise is modelled as a single baseline (Maino et al. 2002). The algorithm makes use of the redundancy in the observing strategy to constrain these baselines, which are then subtracted from the time-ordered data in the creation of the sky maps. The algorithm allows a selection of baseline lengths, which is always a compromise between optimal noise removal and computational cost. As in the 2015 release, we adopt baselines of 1 s at 44 and 70 GHz, and 0.25 s at 30 GHz. The shorter baseline at 30 GHz is appropriate for the higher $1/f$ noise of the radiometers at this frequency (see Table 1), which introduces a larger correlated component in the noise.

Table 1 gives typical values for the main instrument performance parameters measured in flight; similar tabulations were given in previous releases. Beam and optical properties are derived from Jupiter transits, and are consistent with 2015 results. Major improvements in the calibration uncertainty are reflected in more stable results for noise parameters. Values reported are averages among the radiometers operating at a given frequency. At 44 GHz the FWHM is not entirely representative of the actual beamwidth, since one of the three 44-GHz horns is located on the opposite site of the focal plane from the other two (Planck Collaboration IV 2016).

Table 1. LFI performance parameters.

Parameter	30 GHz	44 GHz	70 GHz
Centre frequency [GHz]	28.4	44.1	70.4
Bandwidth [GHz]	9.89	10.72	14.90
Scanning beam FWHM[a] [arcmin]	33.10	27.94	13.08
Scanning beam ellipticity[a]	1.37	1.25	1.27
Effective beam FWHM[b] [arcmin]	32.29	26.99	13.22
White-noise level in timelines[c] [$\mu K_{CMB} s^{1/2}$]	147.9	174.0	151.9
f_{knee}[c] [mHz]	113.9	53.0	19.6
$1/f$ slope[c]	−0.92	−0.88	−1.20
Overall calibration uncertainty[d] [%]	0.17	0.12	0.20

[a] Determined by fitting Jupiter observations directly in the timelines.
[b] Calculated from the main-beam solid angle of the effective beam. These values are used in the source extraction pipeline (Planck Collaboration XXVI 2016).
[c] Typical values derived from fitting noise spectra (see Sect. 5).
[d] Difference between first and last iteration of the iterative calibration (for 30 and 44 GHz) or E2E 2015 result (for 70 GHz). In 2015, the calibration uncertainty was 0.35 % and 0.26 % at 30 and 44 GHz, respectively.

2. Time-ordered information (TOI) processing

The main changes in the Level 1 pipeline since the last release are related to data flagging and to correcting the nonlinearity in the analogue-to-digital converter (ADC).

We revised our flagging procedure to use more conservative and rigorously homogeneous criteria. The new procedure results in a slightly higher flagging rate, particularly during the first 200 operational days (ODs) of the mission; however, the fraction of flagged data remains negligible. Table 2 gives final values for the missing and unusable data for the full mission; changes from the release reported in Planck Collaboration II (2016) are a fraction of one percent. Since the fraction of flagged data is negligible, so the effect on science is also negligible. It is worth mentioning that although the LFI radiometers are quite stable, there are occasional jumps in gain that if not treated properly would impact the calibration procedure well beyond the single data point in which the jump occurs. These jumps are now properly identified and taken into account.

Table 2. Percentage of LFI observation time lost due to missing or unusable data, and to manoeuvres.

Category	30 GHz	44 GHz	70 GHz
Missing [%]	0.15425	0.15425	0.15433
Anomalies [%]	0.82402	0.50997	0.84842
Manoeuvres [%]	8.03104	8.03104	8.03104
Usable [%]	90.99069	91.30474	90.96621

Nonlinearity in the ADCs that convert analogue detector voltages into numbers distorts the radiometer response, possibly mimicking a sky signal. For the present release, we developed a new approach to the correction of this effect that produces significantly better results at 30 GHz.

The first step in the correction is calculation of the white-noise amplitude, given by the difference between the sum of the variances and twice the covariances of adjacent samples in the time-stream. Specifically, $\sigma^2_{WN} = \mathrm{Var}[X_o] + \mathrm{Var}[X_e] - 2\mathrm{Cov}[X_o, X_e]$, where X_o and X_e are data points with time-stream odd and even indices respectively. ADC nonlinearity produces a variation in the white-noise amplitude as a function of the detector voltage.

In the previous release, we fitted the white-noise amplitudes binned with respect to detector voltage with a simple spline curve, and translated the results into a correction curve as described in appendix A of Planck Collaboration III (2014). For this release, we tried a more physically motivated fitting function based on the fact that ADCs suffer from a linearity error ε on each bit. We modelled the output voltage V_0 as

$$V_0 = V_{adu} \sum_{i=0}^{nbit-1} 2^i b_i \left(1 + \varepsilon_i/2^i\right) - V_{off}\,, \quad (1)$$

where b_i is 1 if the ith bit is set and 0 otherwise, ε_i is the linearity error of the ith bit (which is between −0.5 and +0.5), V_{adu} is the voltage step for one binary level change (one analogue-to-digital unit or adu), and V_{off} allows for a possible offset (see figures 9 and 10 of Planck Collaboration III (2014)). Due to complex degeneracies in Eq. (1), we adopted an annealed optimization procedure to avoid local minima in the χ^2 fit to this model.

Even this improved model proved to be too simple, however, as it did not reproduce some of the asymmetries present in the original ADC curve, which appeared to be due to coupling between adjacent bits. We therefore add to the previous expression an extra summation for adjacent coupled bits:

$$V_0 = V_{adu} \sum_{i=0}^{nbit-1} 2^i b_i \left(1 + \varepsilon_i/2^i\right) + \sum_{i=0}^{nbit-2} b_i b_{i+1} \varepsilon_{i,i+1} - V_{off}\,, \quad (2)$$

where $\varepsilon_{i,i+1}$ is the coupled error between bits i and $i+1$.

We compared the results between this method and the previous one by means of null maps, checking the consistency of the resulting new gain solution with the new ADC correction applied. This was done by computing the rms scatter from the eight different survey maps, taking into account pixel hits and zero levels. A "goodness" parameter can then be derived from the mean level of the masked null map made between these survey scatter maps.

The null maps showed substantial improvement at 30 GHz, but little improvement at 44 and 70 GHz. Inspection of the ADC solutions revealed that the higher noise per radiometer and low ADC nonlinearity at 70 GHz did not allow for a good fit. At

44 GHz, on the other hand, the ADC effect was so large that the new model could not reproduce some of the details, and so led to some small residuals. The 30-GHz system has much lower noise and less thermal drift in the gain, meaning that more voltage levels were revisited more often, yielding a more consistent ADC model curve. It was therefore decided to keep the new solution only for the 30-GHz channels. The other two frequencies thus have the same correction for ADC nonlinearity as in the previous release.

3. Photometric calibration

The raw output from an LFI radiometer is a voltage, V, which we can write (Planck Collaboration II 2016) as

$$V(t) = G(t) \times \left[B * (D_{\rm solar} + D_{\rm orbital} + T_{\rm sky}) + T_0 \right], \qquad (3)$$

where G is the gain B encodes both convolution with a 4π instrumental beam and the observation scanning strategy, $D_{\rm solar}$ and $D_{\rm orbital}$ are the solar and orbital CMB dipoles,[2], $T_{\rm sky}$ represents the sum of the CMB and foreground fluctuations, and T_0 is the sum of the 2.7-K CMB temperature, other astrophysical monopole terms, and any internal instrumental offsets. Photometric calibration is the process of determining $G(t)$ accurately over time, which is critical for the quality of the final maps.

In the *Planck* 2013 release, based on 15 months of data, an accurate *Planck* determination of $D_{\rm orbital}$ was not possible, and $G(t)$ was estimated from $D_{\rm solar}$ alone. We used the best-fit, 9-year WMAP dipole estimate as the reference model against which to compare the measured voltages (Planck Collaboration V 2014; Bennett et al. 2013). Successive analyses (Planck Collaboration XXXI 2014; Planck Collaboration I 2016) showed that this model resulted in gain estimates that were offset by about 0.3 % (due largely to foreground contamination in the WMAP dipole), within the originally estimated error uncertainty.

In the *Planck* 2015 release, we implemented internal and self-consistent estimation of the solar dipole by using the orbital dipole for absolute calibration (see Sect. 4 for further details). The orbital dipole is much smaller than the solar dipole, but is known absolutely with exquisite accuracy from the orbital motion of *Planck* itself. This resulted in relative calibration uncertainties $\lesssim 0.3$ % (Planck Collaboration II 2016), adequate to allow high-precision cosmology based on temperature measurements. However, for polarization even a relative error of 10^{-3} is non-negligible, and a large fraction of the LFI work on data quality since the 2015 release has revolved around reducing this error further.

As discussed extensively in Planck Collaboration II (2016) and Planck Collaboration XI (2016), one of the most notable problems in the 2015 LFI processing was the failure of a specific internal null test, namely that taken between Surveys 1, 3, 5, 6, 7, and 8, and Surveys 2 and 4. In particular, Surveys 2 and 4 showed significantly larger uncertainties in their gain estimation than the other Surveys (see figure 4 in Planck Collaboration II 2016), and, critically, they also showed significant excess B-mode power on the very largest scales. Although it was well-known that Surveys 2 and 4 happen to be aligned with the *Planck* scanning strategy in such a way that the dipole modulation reaches very low minima, thus exacerbating the impact on calibration of any potential systematic effect, we could not identify the specific source of the anomaly. Nevertheless, because of the null-test failure, those two surveys were removed from the final polarization maps and likelihood analysis.

Since that time, we have performed a series of detailed end-to-end simulations designed specifically to identify the source of this null-test failure, and this work ultimately led to a minor, but important, modification of the calibration scheme outlined above and described in detail in Planck Collaboration II (2016). In short, the survey null-test failure was due to not accounting for the polarized component of the sky signal in Eq. (3). This has now been done, as described in detail below. Thus the updated calibration scheme represents the logical conclusion of Eq. (3), since we now account for all terms as far as we are able to model them.

3.1. Joint gain estimation and component separation

Before describing the updated calibration scheme, we first establish some useful intuition regarding the physical effect in question. We start with the raw gains, $G(t)$, as measured in 2015 (Planck Collaboration II 2016). Overall, this function may be crudely modelled over the course of the mission as a sum of a linear term and a 1-year sinusoidal term:

$$G_{\rm model}(t) = (a + b\,t) + c\,\sin\left(\frac{2\pi}{365\,{\rm days}}\,t + d\right), \qquad (4)$$

where the four free parameters, $\{a,b,c,d\}$, must be fitted radiometer by radiometer. From this model, we compute the "normalized fractional gain" as

$$\hat{G}(t) = 100\frac{G(t) - G_{\rm model}(t)}{a}. \qquad (5)$$

This function is simply the fractional gain excess (or deficit) relative to a smoothly varying model, expressed as a percentage.

Each LFI horn feeds two independent polarization-sensitive radiometers with polarization angles rotated 90° with respect to each other (Planck Collaboration II 2014); these are called 'M' (main) and 'S' (side), respectively. Since two such radiometers are often susceptible to the same instrumental effects (thermal, sidelobes, etc.), it is useful to study differences between them to understand instrumental systematic effects. For example, Fig. 2 shows the difference in normalized fractional gain for two 30-GHz radiometers, namely $\hat{G}_{28\rm M} - \hat{G}_{28\rm S}$. Other 30- and 44-GHz radiometers show qualitatively similar behaviour, at the sub-percent level, whereas the 70-GHz radiometers behave differently, for reasons explained below. The following discussion therefore applies in detail only to the 30 and 44 GHz radiometers, while the 70-GHz radiometers will be treated separately.

3.2. Calibration at 30 and 44 GHz

The black curve in Fig. 2 shows the normalized gain difference for the default 2015 orbital dipole-based calibration scheme (DX11D), and exhibits a striking oscillatory pattern with a period equal to one survey. Such a pattern is very difficult to explain instrumentally, since the two radiometers reside inside the same horn, while it is consistent with the effect of polarized foregrounds. Since the polarization angles of the two radiometers are rotated by 90°, any polarized signal on the sky will at any given time be observed with opposite signs by the two radiometers.

[2] The solar dipole is the dipole anisotropy in the CMB induced by the motion of the solar System barycentre with respect to the rest frame of the CMB itself. The orbital dipole is the modulation of the solar dipole due to the orbital motion of the spacecraft around the Sun.

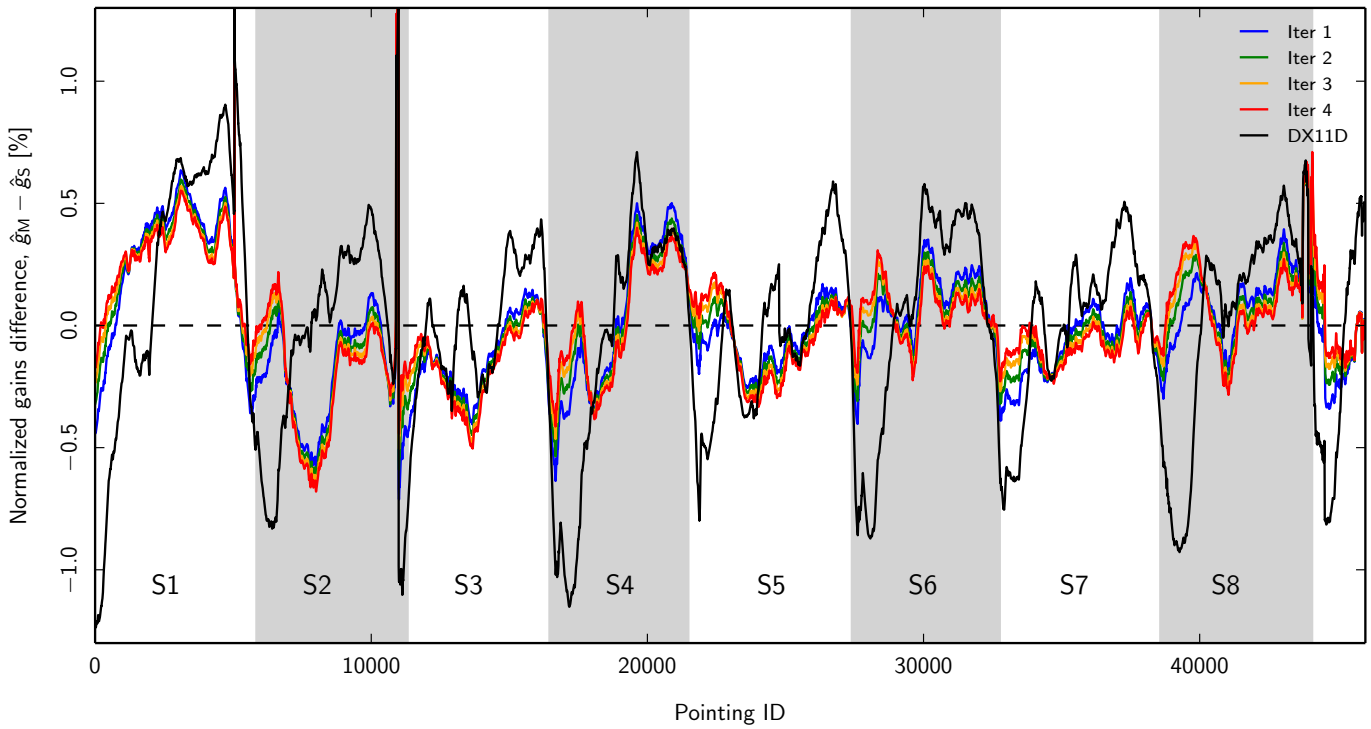


Fig. 2. Normalized gain difference (see main text for precise definition) between two radiometers inside the same horn (28M and 28S here) as a function of pointing ID for both the 2015 (DX11D) and 2018 calibration schemes. The coloured lines show this function for the various iterations in the new {gain-estimation + mapmaking + component separation} calibration scheme. Sky surveys are indicated by alternating white and grey vertical bands. Other 30- and 44-GHz radiometers show qualitatively similar behaviour, whereas the 70-GHz radiometers exhibit too much noise, and corresponding iterations for these detectors do not converge within this scheme.

Strong polarized foregrounds therefore lead to the kind of difference shown in Fig. 2, with a sign given by the relative orientation of the satellite and the Galactic magnetic field. Furthermore, this difference will be repeatable across surveys. This is confirmed by simulations — inserting a polarized foreground sky into end-to-end simulations induces precisely the same pattern as shown here.

The solution to this problem is to include the sky signal, $T_{\rm sky}$, in the calibrator, on the same footing as the orbital and solar dipoles, including both temperature and polarization fluctuations. This is non-trivial, since the purpose of the experiment is precisely to measure the polarized emission from the sky. A good approximation can be established, however, through an iterative process that alternates between gain calibration, map making, and astrophysical component separation, using the following steps.

0. Let $T_{\rm sky}$ be the full best-fit (`Commander`-based; Eriksen et al. 2008; Planck Collaboration X 2016) *Planck* 2015 astrophysical sky model, including CMB, synchrotron, free-free, thermal and spinning dust, and CO emission for temperature maps, plus CMB, synchrotron, and thermal dust in polarization.
1. Estimate G from Eq. (3), explicitly including the temperature and polarization component of $T_{\rm sky}$ in the calibration on the same footing as $D_{\rm solar}$ and $D_{\rm orbital}$.
2. Compute frequency maps with these new gains.
3. Determine a new astrophysical model from the updated frequency maps using `Commander` (at present the sky model is adjusted only for LFI frequencies.)
4. Iterate steps (1) to (3).

Since the true sky signal is stationary on the sky, while the spurious gain fluctuations are not, this process will converge, essentially corresponding to a generalized mapmaker in which the $G(t)$ is estimated jointly with the sky maps. Alternatively, this process may also be considered as a Gibbs sampler that in turn iterates through all involved conditional distributions, and thereby converges to the joint maximum likelihood point (Eriksen et al. 2008).

The process is, computationally expensive, however; each iteration takes about one week to complete. For practical purposes, the current process was therefore limited to four full iterations (not counting the 2015 model used for initialization). The normalized gain differences established in each iteration are shown as coloured curves in Fig. 2 for the same radiometer pair as discussed above (for example 28M and 28S). Here we see that most of the effect is accounted for simply by introducing a rough model, as already the first iteration is significantly flatter than the initial model (black versus blue curves). Subsequent iterations make relatively small differences, and, critically, the differences between consecutive iterations become smaller by almost a factor of 2 in each case, indicating that the algorithm indeed converges.

While the most obvious oscillatory pattern in the initial model has been eliminated by the introduction of the astrophysical sky model, it is not clear whether a smaller contribution remains. In Fig. 3 we therefore show the same functions, but now with all surveys stacked together. Explicitly, the coloured regions in Fig. 3 represent the mean and standard deviations as evaluated over the eight surveys. Since the surveys have slightly different lengths, the stacking is done such that the starting pointing periods of the surveys are aligned, and longer surveys are trun-

cated at the end. These stacked functions will tend to suppress random signals across surveys, but highlight those common to all eight surveys. As before (but now more clearly), we see that the 2015 model (grey band) exhibits a highly significant Survey-dependent pattern. This pattern is greatly suppressed simply by adding a rough model to the calibrator (blue band). And the pattern is additionally reduced by further iterations, with a convergence rate of about a factor of 2 per iteration.

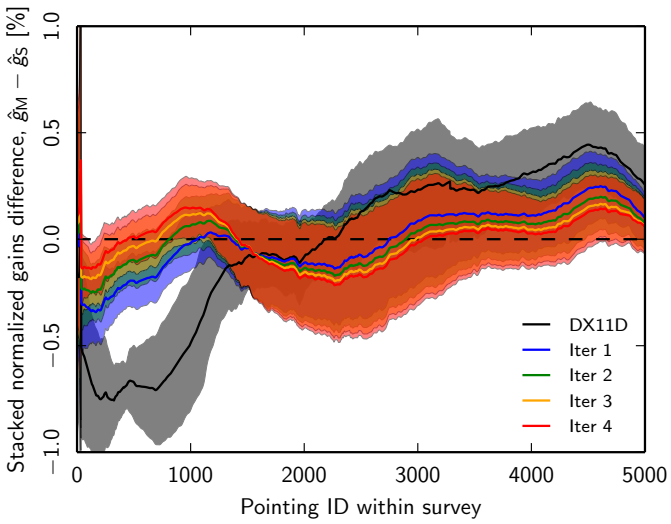


Fig. 3. Same as Fig. 2, but stacked over surveys. Each band corresponds to the mean and $1\,\sigma$ confidence region as evaluated from eight surveys.

To understand the convergence rate in more detail, we show in Fig. 4 the differences in polarization amplitude between two consecutive iterations of the full 30-GHz map. The top panel shows the difference between the second and first iterations; the middle panel shows the difference between the third and second iterations; and, finally, the bottom panel shows the difference between the fourth and third iterations. As anticipated, here we see that the magnitude of the updates decreases by a factor of 1.5–2 at high latitudes.

In addition to the decreasing amplitude with iterations, it is also important to note that the *morphology* of the three difference maps is very similar, and dominated by a few scans that align with Ecliptic meridians. In other words, most of the gain uncertainty is dominated by a few strong modes on the sky, and the iterations described above largely try to optimize the amplitude of these few modes. Furthermore, as seen in Figs. 2–4, it is clear that we have not converged to numerical precision with only four iterations. Due to the heavy computational demand of the iterative process, we could not produce further steps. As a consequence, we expect that low-level residuals are still present in the 2018 LFI maps, with a pattern similar to that of the 2015 maps, though with significantly lower amplitude. For the 2018 release, we adopt the difference between the two last iterations as a spatial template of residual gain uncertainties projected onto the sky. This template is used only at 70 GHz.

3.3. Calibration at 70 GHz

As already mentioned, the above discussion applies only to the 30- and 44-GHz radiometers, since the 70-GHz radiometers behave differently. The reason for this may be seen in Fig. 5,

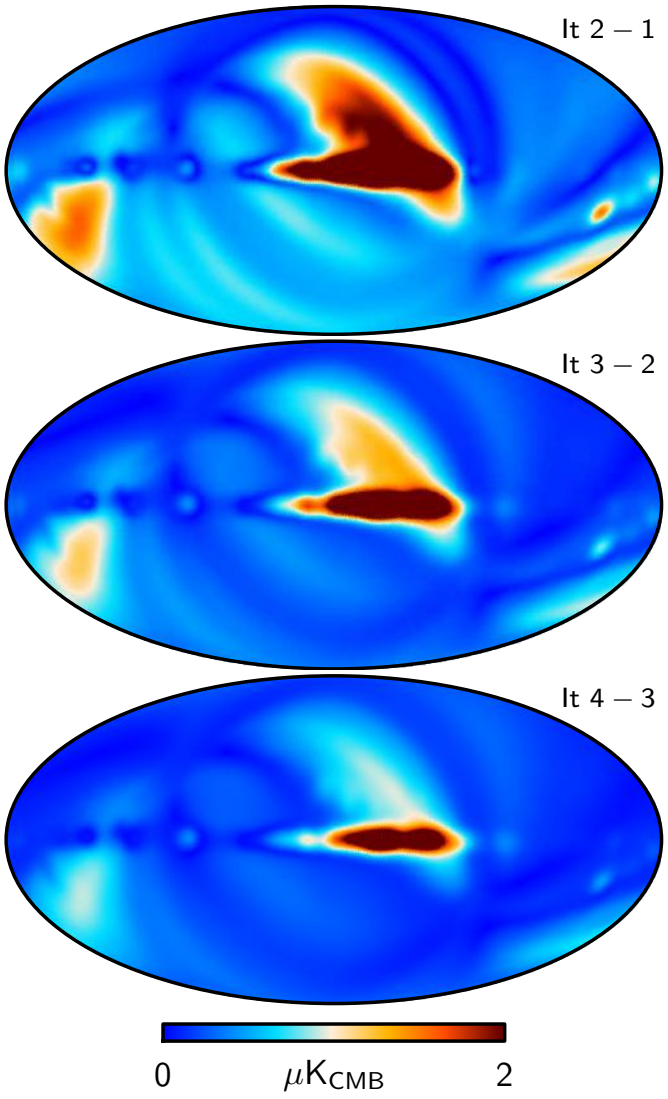


Fig. 4. Polarization amplitude difference maps between consecutive iterations of the internal foreground model evaluated at 30 GHz, as derived with `Commander`. The three panels show the differences between: the second and first iterations (top); the third and second iterations (middle); and the fourth and third iterations (bottom). All maps are smoothed to an effective angular resolution of 8° FWHM.

which simply shows the final co-added 30-, 44-, and 70-GHz frequency maps, downgraded to `HEALPix` (Górski et al. 2005) $N_{\rm side} = 16$ resolution (3°.8 × 3°.8 pixels) to enhance the effective signal-to-noise ratio per pixel. The grey regions show the Galactic calibration mask used in the analysis, within which we do not trust the foreground model sufficiently precisely to use it in gain calibration, primarily due to bandpass leakage effects (Planck Collaboration II 2016). Here we see a qualitative difference between the three frequency maps: while both the 30- and 44-GHz polarization maps are signal-dominated, the 70-GHz channel is noise-dominated. This has a detrimental effect on the iterative scheme described above, ultimately resulting in a diverging process; essentially, the algorithm attempts to calibrate on noise rather than actual signal.

For this channel, we therefore retain the same calibration scheme used in 2015, noting that the iterative algorithm fails because the foregrounds are weak at this channel. While this is a

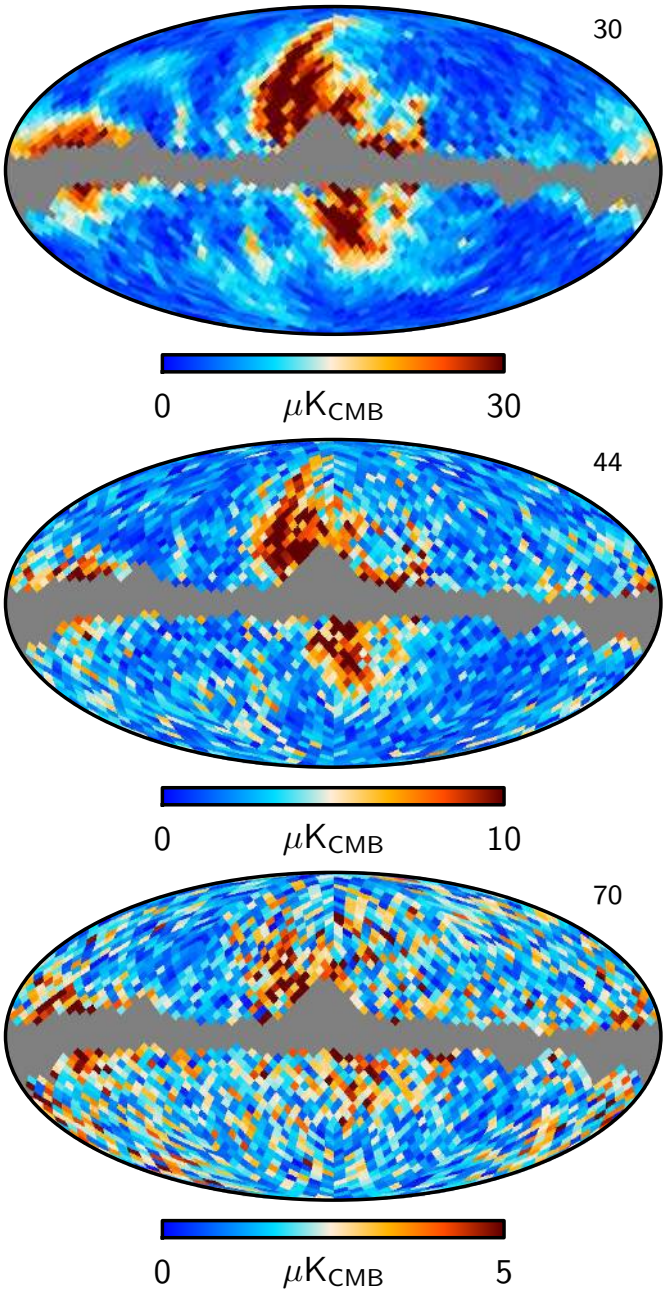


Fig. 5. Final low-resolution LFI 2017 polarization amplitude sky maps. From *top* to *bottom*, the panels show the co-added 30-, 44-, and 70-GHz frequency maps. The grey regions indicate the mask used during gain calibration. Each pixel is $3\overset{\circ}{.}8 \times 3\overset{\circ}{.}8$, corresponding to `HEALPix` resolution $N_{\rm side} = 16$.

problem when using foregrounds directly as a calibrator, it also implies that foregrounds are much less of a problem than in the other channels. In place of an iterative scheme, we adopt the corresponding internal differences described above, obtained at 30 GHz, as a tracer of gain residuals also for the 70-GHz channel (shown in Fig. 6), and marginalize over this spatial template in a standard likelihood fit in pixel space. Indeed, we provide this additive template as part of the LFI 2018 distribution, with a normalization given by a best-fit likelihood accounting for both CMB and astrophysical foregrounds. Thus, the best-fit amplitude of the provided template is unity. For all the cosmological analysis involving the 70 GHz channel and presented in this final *Planck* data release we accounted for gain residuals by subtracting this correction. Therefore we strongly recommend to do the same for any other cosmological investigation involving the 70-GHz frequency channel.

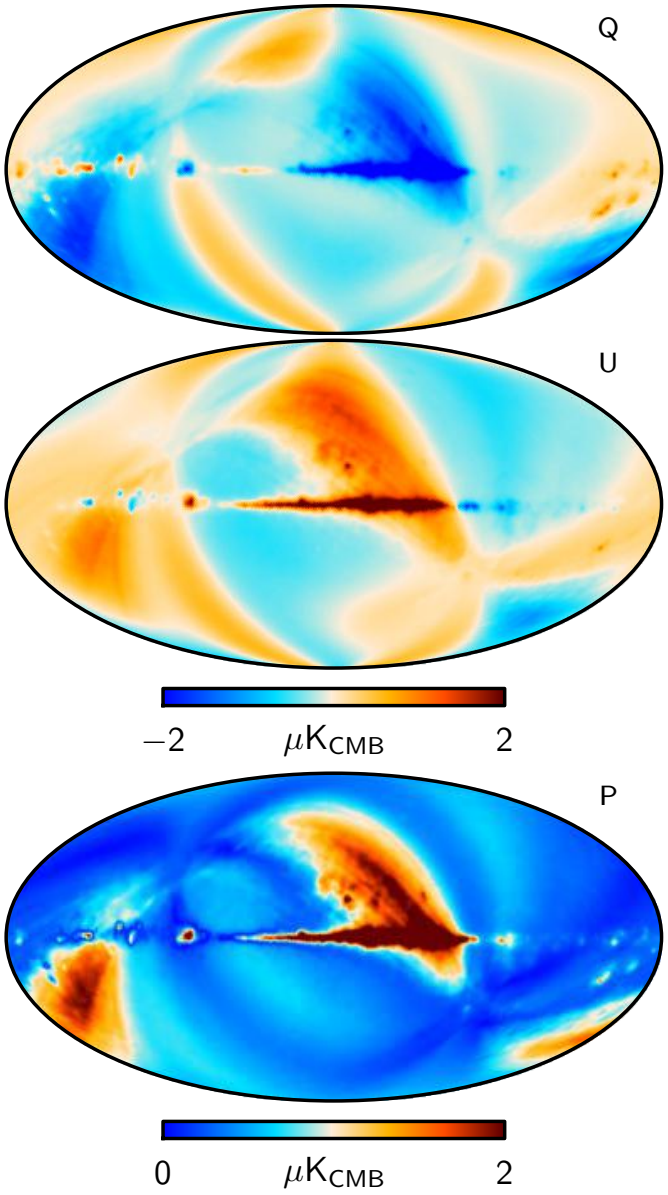


Fig. 6. Gain correction template for the 70-GHz channel in terms of Stokes Q (top) and U (middle), and the polarization amplitude, P (bottom). The template is smoothed to 2° FWHM, and its amplitude is normalized to the best-fit value derived in a joint maximum likelihood analysis of both the CMB power spectrum and template fit, as described in the text.

We started this discussion by recalling the failure reported in Planck Collaboration II (2016) of a null test between Survey sets {1,3, 5–8} and {2, 4}. In Fig. 7 we therefore compare this particular null map as derived with the old (top panel) and new (bottom panel) calibration schemes. The improvement is obvious, with most fluctuations in the latter appearing consistent with noise. The remaining excesses appear on angular scales compatible with the smoothing scale of 8° FHWM, and are consistent with the position of known variable point sources. A direct com-

parison with WMAP frequency maps suggests similar improvements; see Appendix A for details.

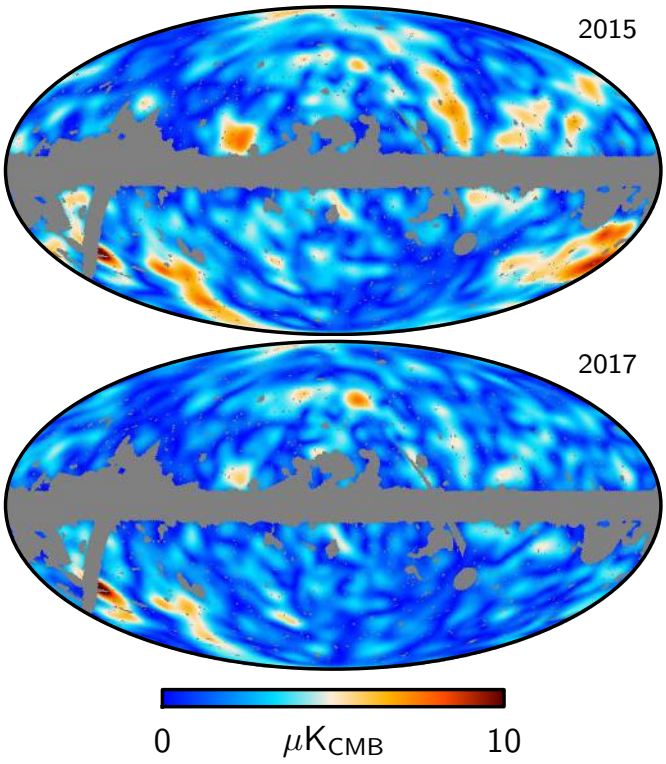


Fig. 7. 30 GHz polarization amplitude null maps evaluated between survey combinations {S1,S3, S5–8} and {S2, S4} for both the 2015 (top) and 2018 (bottom) calibration schemes. This particular survey split is maximally sensitive to residual gain uncertainties from polarized Galactic foreground contamination because of the orientation of the scanning strategy employed in Surveys 2 and 4; see text and Planck Collaboration II (2016) for further details. Both maps are smoothed to an effective angular resolution of 8° FWHM.

4. The LFI Dipole

The calibration signal for LFI is the dipole anisotropy due the motion of the solar system relative to the CMB. Precise knowledge of the amplitude and direction of the 3.3-mK solar dipole, however, requires another absolutely defined signal. This is given by the orbital dipole, the time-varying 200-μK modulation of the dipole amplitude induced by the motion of the spacecraft in its yearly orbit around the Sun (including the small velocity component due to the spacecraft orbit around L2). As the amplitude and orientation of the orbital dipole can be determined with exquisite accuracy from the satellite telemetry and orbital ephemeris, it is the best absolute calibration signal in all of microwave space astrophysics. It should be emphasized that the dipole determination is primarily a velocity measurement, and that the actual dipole amplitude is derived from the velocity assuming a value for the absolute temperature of the CMB; together with HFI (Planck Collaboration III 2018) we use the value $T_0 = 2.72548$K (Fixsen 2009).

4.1. Initial calibration to determine the amplitude and direction of the solar dipole

These two dipoles are merged into a single signal at any given time, but they can be separated over the course of the mission, since the solar dipole is fixed on the sky while the orbital dipole varies in amplitude and direction with the satellite velocity as it orbits the Sun. It is therefore possible to base the calibration entirely on the orbital dipole alone. As in the previous release (Planck Collaboration II 2016), we omit the solar dipole from the fit but retain the far-sidelobe-convolved orbital dipole and the fiducial dipole convolved again with far sidelobes, and also remove the restriction of having no dipole signal in the residual map. In this way, the solar dipole is extracted as a residual, and its amplitude and position can be determined. In this section we discuss the LFI 2018 measurements of the solar dipole and compare them to other measurements.

For accurate calibration, we have to take into account two effects that behave like the orbital dipole, in the sense that they are linked to the satellite and not to the sky: polarized foregrounds; and pick-up in the far sidelobes. While the orbital dipole calibration is robust against unpolarized foregrounds, the polarized part of the foregrounds depends on the orientation of the satellite. Similarly, the far sidelobes are also locked to the direction in which the satellite is pointing.

The corrections for polarized foregrounds are made directly in the timelines by unrolling the Q and U frequency maps from the previous internal data release, in other words, projecting them into timelines according to the scanning strategy and beam orientation, and also taking into account the gain calibration factor derived from an initial calibration run. We find that only one iteration is required to remove the polarized signal, with further iterations in this cleaning process making no difference. The amplitude of the polarized signal removed is about 40 μK at 30 GHz, 15 μK at 44, and 4 μK at 70 GHz, mainly due to the North Polar Spur and the Fan regions.

In the previous dipole analysis, the far sidelobes were removed using the `GRASP` beam model but reduced to the lowest multipoles to obtain the expected, properly convolved dipole signal for both the orbital and solar dipoles. In the calibration code now used (`DaCapo`; Planck Collaboration II 2016), we fit for the orbital dipole convolved with far sidelobes, as well as for the convolution of the solar dipole with the far sidelobes. In such a way, we force a pure dipole (without far sidelobes) into the residual map. However, we found that far-sidelobe pick-up was not completely removed, which resulted in a trend in the dipole amplitude with horn position on the focal plane, as well as small differences between the orbital and solar dipole calibration factors. By adjusting the direction of the large-scale compoment of the far sidelobes, we find a correction between 1 and 10 μK, depending on horn focal plane position, which brings both calibration factors together and, simultaneously, removes the asymmetry in the focal plane.

To calibrate on the orbital dipole, we need to mask the strong emission from the Galactic plane. The mask is generated using a 5°-smoothed 30-GHz intensity map with different threshold cut-offs, which result in different sky fractions. The orbital calibration is carried out using a sky fraction of 94 %, since this calibration is robust to the intensity of unpolarized foregrounds. For the analysis of the resultant dipole maps, we use instead an 80 % sky fraction, accounting for the presence of unpolarized foregrounds in the residual maps. An even more conservative mask with a sky fraction of 60 % gives about 15 % more scatter in the dipole position between channels, but not in any sys-

Table 3. Dipole characterization.

Radiometer	Amplitude [μK_{CMB}]	Galactic Coordinates l	b
30 GHz			
28M	3357.82 ± 1.09	264.°242 ± 0.°030	48.°174 ± 0.°009
28S	3363.98 ± 1.37	264.°252 ± 0.°198	48.°192 ± 0.°062
27M	3366.99 ± 1.46	264.°323 ± 0.°181	48.°195 ± 0.°042
27S	3368.23 ± 1.15	264.°226 ± 0.°091	48.°160 ± 0.°043
44 GHz			
26M	3363.83 ± 0.34	264.°045 ± 0.°016	48.°253 ± 0.°007
26S	3363.64 ± 0.35	264.°026 ± 0.°015	48.°255 ± 0.°006
25M	3360.22 ± 0.35	264.°021 ± 0.°014	48.°252 ± 0.°006
25S	3360.96 ± 0.36	264.°009 ± 0.°014	48.°245 ± 0.°006
24M	3363.94 ± 0.32	264.°012 ± 0.°014	48.°253 ± 0.°007
24S	3363.13 ± 0.37	264.°034 ± 0.°015	48.°257 ± 0.°006
70 GHz			
23M	3364.18 ± 0.30	264.°005 ± 0.°011	48.°267 ± 0.°005
23S	3364.69 ± 0.30	264.°003 ± 0.°012	48.°271 ± 0.°005
22M	3364.15 ± 0.28	263.°988 ± 0.°011	48.°269 ± 0.°005
22S	3364.63 ± 0.28	264.°013 ± 0.°012	48.°263 ± 0.°005
21M	3364.16 ± 0.30	263.°991 ± 0.°011	48.°267 ± 0.°005
21S	3364.20 ± 0.27	264.°007 ± 0.°012	48.°262 ± 0.°006
20M	3364.74 ± 0.27	263.°983 ± 0.°012	48.°266 ± 0.°005
20S	3364.38 ± 0.27	263.°991 ± 0.°012	48.°261 ± 0.°006
19M	3363.94 ± 0.27	263.°984 ± 0.°012	48.°265 ± 0.°005
19S	3364.58 ± 0.28	264.°019 ± 0.°013	48.°264 ± 0.°006
18M	3364.42 ± 0.28	264.°005 ± 0.°011	48.°267 ± 0.°005
18S	3364.23 ± 0.29	263.°991 ± 0.°010	48.°263 ± 0.°005
Combined[a]	3364.4 ± 3.1	263.°998 ± 0.°051	48.°265 ± 0.°015

[a] This estimate is based on the collective sum of all the Markov chain Monte Carlo (MCMC) samples for all 70-GHz channels. Final amplitude error bars include 0.07–0.11 % calibration uncertainty.

tematic way. This is consistent with a lower signal-to-noise ratio due to the poorer sky coverage. The actual fit is performed with a Markov chain Monte Carlo (MCMC) approach, where we search for dipole position and amplitude as well as the amplitudes of synchrotron, dust, and free-free templates derived from `Commander` (Planck Collaboration X 2016). We performed several tests, varying the amplitude of the mask (ranging from 60 % to 80 % sky fraction), with and without point source masks, in the derivation of the foreground templates. We found that use of a mask with 80 % sky fraction and masked foregrounds reduces the scatter in the dipole estimation at 70 GHz by about 15 %.

4.2. The solar dipole

From the MCMC samples, the 1 % and 99 % values are used to set the limits on the dipole amplitude and position. Figure 8 shows results for the three LFI frequencies for both dipole direction (upper panel) and amplitude (lower panel). There is a clear trend with frequency in dipole direction, due to foreground contamination. As expected, the 70-GHz channel has the lowest foreground signal, and it is used to derive the final LFI dipole. With respect to the previous release, the use of the small far-sidelobe correction has removed the systematic amplitude variation with focal plane position. Thus the cross-plane null pairs that were used in the previous release are not needed. This results in a smaller scatter of both dipole positions and derived amplitude, as shown in the bottom panel. For each LFI data point we report two error bars: the small (red) one is the actual error in the fit (also reported in Table 3); and the large (black) one is obtained by summing the calibration error in quadrature. The grey band represents the WMAP derived dipole amplitude, for comparison. Numerical results are summarized in Table 3, where single radiometer errors are derived from the MCMC samples. The final uncertainty in the LFI dipole derived from only the 70 GHz measurements, however, also takes also into account gain errors, estimated through the use of dedicated simulations with `DaCapo`, in the range 0.07–0.11 %. This yields a final dipole amplitude $D = 3364 \pm 3\,\mu$K and direction in Galactic coordinates $(l, b) = (263.°998 \pm 0.°051, 48.°265 \pm 0.°015)$.

5. Noise estimation

We estimated the basic noise properties of the receivers (for example, knee-frequency and white-noise variance) throughout the mission lifetime. This is a simple way to track variations and possible instrument anomalies during operations. Furthermore, a detailed knowledge of noise properties is required for other steps of the analysis pipeline, such as optimal detector combination in the mapmaking process or Monte Carlo simulations used for error evaluation at the power spectrum level.

The noise model and the approach for noise estimation is the same as described in Planck Collaboration II (2016) and

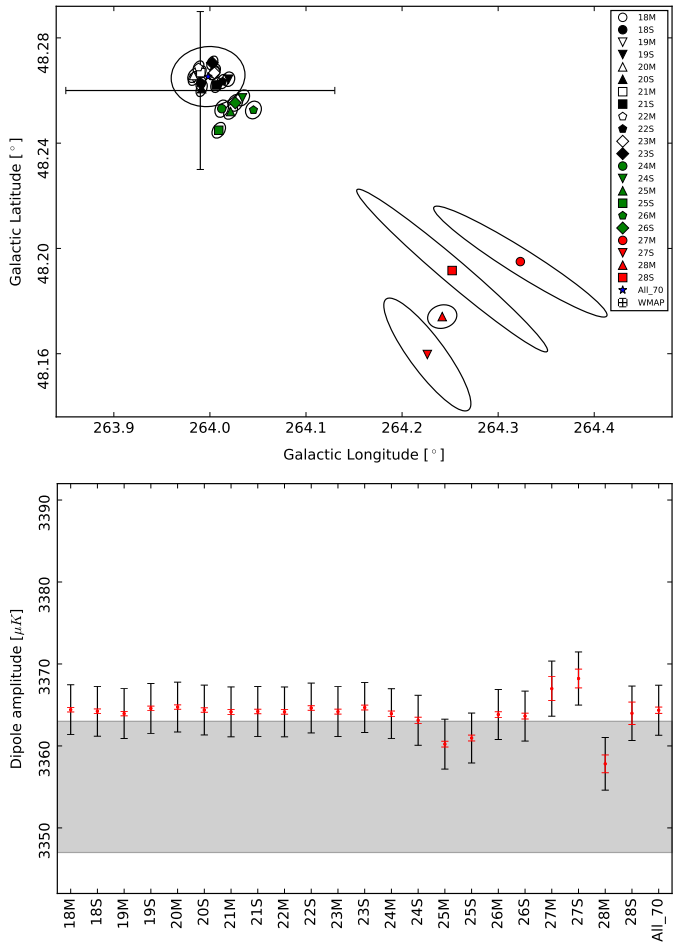


Fig. 8. Direction (top) and amplitude (bottom) of the solar dipole determined from each of the LFI detectors. Uncertainties in direction are given by 95 % ellipses around symbols, colour-coded for frequency (70 GHz black or unfilled, 44 GHz green, and 30 GHz red). The four discrepant points are at 30 GHz, where we expect foregrounds affect dipole estimation. Amplitude uncertainties are dominated by the systematic effects of gain uncertainties. The grey band in the bottom panel shows the WMAP dipole amplitude for comparison.

Planck Collaboration II (2014). We employed a noise model of the form

$$P(f) = P_0^2 \left[1 + \left(\frac{f}{f_{\rm knee}} \right)^{\beta} \right], \qquad (6)$$

where P_0^2 is the white-noise power spectrum level, and $f_{\rm knee}$ and β encode the non-white ($1/f$-like) low-frequency noise component. We estimate P_0^2 by taking the mean of the noise spectrum in the last few bins at the highest values of f (typically 10 % at 44 and 70 GHz, and 5 % at 30 GHz due to the higher knee-frequency). For the knee frequency and slope, we exploited the same MCMC engine as in the previous release. Tables 4 and 5 give white noise and low-frequency noise parameters, respectively. Comparing these results with those from the 2015 release (Planck Collaboration II 2016), we see that both the white-noise level and slope β show variations well below 0.1 %, while $f_{\rm knee}$ varies by less than 1.5 %. However, error bars (rms fitted values over the mission lifetime) are in some cases larger than before. This results from both improved TOI processing (mainly flagging) and an improved calibration pipeline that allows us to detect a sort of bimodal distribution (at the $\simeq$ 1 % level) in $f_{\rm knee}$ for three of the LFI radiometers. Figure 9 shows typical noise spectra at several times during the mission lifetime for three representative radiometers, one for each LFI frequency. Spikes at the spin-frequency (and its harmonics) are visible in the 30-GHz spectra due to residual signal left over in the noise estimation procedure (see Planck Collaboration II 2016). These are due to the combined effect of the large signal (mainly from the Galaxy) and the large value of the knee frequency, together with the limited time window (only 5 ODs) on which spectra are computed.

Table 4. White-noise levels for the LFI radiometers.

Radiometer	White-Noise Level	
	Radiometer M [$\mu K_{CMB}\ s^{1/2}$]	Radiometer S [$\mu K_{CMB}\ s^{1/2}$]
70 GHz		
LFI-18	512.5 ± 2.0	466.7 ± 2.3
LFI-19	578.9 ± 2.1	554.2 ± 2.2
LFI-20	586.9 ± 1.9	620.0 ± 2.7
LFI-21	450.4 ± 1.4	559.8 ± 1.8
LFI-22	490.1 ± 1.3	530.9 ± 2.2
LFI-23	503.9 ± 1.6	539.1 ± 1.7
44 GHz		
LFI-24	462.8 ± 1.3	400.5 ± 1.1
LFI-25	415.2 ± 1.3	395.0 ± 3.1
LFI-26	482.6 ± 1.9	422.9 ± 2.5
30 GHz		
LFI-27	281.5 ± 2.0	302.8 ± 1.8
LFI-28	317.9 ± 2.4	286.1 ± 2.1

6. Mapmaking

The methods and implementation of the LFI mapmaking pipeline are described in detail in Planck Collaboration II (2016), Planck Collaboration VI (2016), and Keihänen et al. (2010). Here we report only the changes introduced into the code with respect to the previous release.

Our pipeline still uses the `Madam` destriping code, in which the correlated noise component is modelled as a sequence of short baselines (offsets) that are determined via a maximum-likelihood approach. In the current release, the most important change is in the definition of the noise filter in connection with horn-uniform detector weighting. When combining data from several detectors, we assign each detector a weight that is proportional to

$$C_{\rm w}^{-1} = \frac{2}{\sigma_{\rm M}^2 + \sigma_{\rm S}^2}, \qquad (7)$$

where $\sigma_{\rm M}^2$ and $\sigma_{\rm S}^2$ are the white-noise variances of the two radiometers ("Main" and "Side") of the same horn. The same weight is applied to both radiometers. In the 2015 release the weighting was allowed to affect the noise filter as well: for the noise variance σ^2 in Eq. (6) we used the average value $C_{\rm w}$. For the current release we have completely separated detector weighting from noise filtering. We use the individual variances $\sigma_{\rm M}^2$ and $\sigma_{\rm S}^2$ for each radiometer when building the noise filter. In principle, this makes maximum use of the information we have about the noise of each radiometer.

We compared the previous and the current versions of `Madam`, using a single noise filter, and found excellent agree-

Table 5. Knee frequencies and slopes for each of the LFI radiometers.

	Knee Frequency $f_{\rm knee}$ [mHz]		Slope β	
	Radiometer M	Radiometer S	Radiometer M	Radiometer S
70 GHz				
LFI-18	14.8 ± 2.1	17.8 ± 1.5	-1.06 ± 0.08	-1.18 ± 0.10
LFI-19	11.7 ± 1.1	13.7 ± 1.3	-1.21 ± 0.23	-1.10 ± 0.13
LFI-20	7.9 ± 1.5	5.6 ± 1.0	-1.19 ± 0.27	-1.30 ± 0.33
LFI-21	37.9 ± 5.7	13.1 ± 1.3	-1.24 ± 0.08	-1.20 ± 0.08
LFI-22	9.5 ± 1.7	14.3 ± 6.7	-1.41 ± 0.25	-1.23 ± 0.27
LFI-23	29.6 ± 1.1	58.9 ± 1.6	-1.07 ± 0.02	-1.21 ± 0.02
44 GHz				
LFI-24	26.9 ± 1.0	89.6 ± 13.8	-0.94 ± 0.01	-0.91 ± 0.01
LFI-25	19.7 ± 1.0	46.8 ± 1.9	-0.85 ± 0.01	-0.90 ± 0.01
LFI-26	64.4 ± 1.8	70.7 ± 18.7	-0.92 ± 0.01	-0.75 ± 0.07
30 GHz				
LFI-27	173.7 ± 3.1	109.6 ± 2.5	-0.93 ± 0.01	-0.91 ± 0.01
LFI-28	128.5 ± 10.9	44.1 ± 2.2	-0.93 ± 0.01	-0.90 ± 0.02

ment. Differences were at the $0.01\,\mu$K level, and the code took the same number of iterations. We then compared the results obtained with the combined noise filter with those obtained with the separated noise filters for each radiometer. We found that using separate filters for the two radiometers of the same horn has the effect of reducing the total number of iterations required for convergence by almost a factor of 2. The net effect of using both the new version of `Madam` and the separate noise filters is that we obtain the same maps as before, but considerably faster.

In Figs. 10 to 12, we show the 30-, 44-, and 70-GHz frequency maps. The top panels are the temperature (I) maps based on the full observation period, and presented at the original native instrument resolution, `HEALPix` $N_{\rm side} = 1024$. The middle and bottom panels show the Q and U polarization components, respectively; these are smoothed to 1° angular resolution and downgraded to $N_{\rm side} = 256$. Polarization components have been corrected for bandpass leakage (see Sect. 7). Table 6 gives the main mapmaking parameters used in map production. All values are the same as for the previous data release except the monopole term; although we used the same plane-parallel model for the Galactic emission as for the 2015 data release, the derived monopole terms are slightly changed at 30 and 44 GHz for the adopted calibration procedure.

7. Polarization: Leakage maps and bandpass correction

The small amplitude of the CMB polarized signal requires careful handling because of systematic effects capable of biasing polarization results. The dominant one is the leakage of unpolarized emission into polarization; any difference in bandpass between the two arms of an LFI radiometer will result in such leakage. In the case of the CMB, this is not a problem. That is because calibration of each radiometer uses the CMB dipole, which has the same frequency spectrum as the CMB itself, and so exact gain calibration perfectly cancels out in polarization. However, unpolarized foreground-emission components with spectra different from the CMB will appear with different amplitudes in the two arms, producing a leakage into polarizaion.

In order to derive a correction for this bandpass mismatch, we exploit the *IQUSS* approach (Page et al. 2007) used in the 2015 release. The main ingredients in the bandpass mismatch recipe are the leakage maps L, the spurious maps S_k (see below), the a-factors, and the $A_{Q[U]}$ maps (see section 11 of Planck Collaboration II 2016, for definitions). With respect to the treatment of bandpass mismatch in the previous release, we introduce three main improvements in the computation of the L and A maps. Leakage maps, L, are the astrophysical leakage term encoding our knowledge of foreground amplitude and spectral index. These maps are derived from the output of the `Commander` component-separation code. In the present analysis this is done using only *Planck* data from the current data release at their full instrumental resolution. In contrast, in the earlier approach we also used WMAP 9-year data and applied a 1° smoothing prior to the component-separation process. We also exclude *Planck* channels at 100 and 217 GHz, since these could be contaminated by CO line emission. Spurious maps S_k (one for each radiometer) are computed from `Madam` mapmaking outputs (for the full frequency map creation run). Basically, spurious maps are proportional to the bandpass mismatch of each radiometer, and can be computed directly from single radiometer timelines. As described in Planck Collaboration II (2016), the output of the Main and Side arms of a radiometer can be redefined, including bandpass mismatch spurious terms, as

$$\begin{aligned}
\text{LFI 27}\begin{cases} d_{s1} = I + Q\cos(2\psi_{\rm s1}) + U\sin(2\psi_{\rm s1}) + S_1\,, \\ d_{m1} = I + Q\cos(2\psi_{\rm m1}) + U\sin(2\psi_{\rm m1}) - S_1\,, \end{cases} \\
\text{LFI 28}\begin{cases} d_{s2} = I + Q\cos(2\psi_{\rm s2}) + U\sin(2\psi_{\rm s2}) + S_2\,, \\ d_{m2} = I + Q\cos(2\psi_{\rm m2}) + U\sin(2\psi_{\rm m2}) - S_2\,, \end{cases}
\end{aligned} \tag{8}$$

or in the more compact form

$$d_i = I + Q\cos(2\psi_i) + U\sin(2\psi_i) + \alpha_1 S_1 + \alpha_2 S_2\,. \tag{9}$$

Here α_1 and α_2 can take the values -1, 0, and 1. The problem of estimating $m = [I, Q, U, S_1, S_2]$ is similar to a mapmaking problem, with two extra maps. The pixel-noise covariance matrix is therefore given by the already available `Madam`-derived covariance matrix, with two additional rows and columns, as

$\mathsf{M}_p = \sum_{i\in p} w_i \times$

Table 6. Mapmaking parameters used in the production of maps. Details are reported in Planck Collaboration VI (2016).

		Baseline length[b]		Resolution[c]		Monopole, B[d]
Channel	$f_{\rm samp}$ [Hz][a]	[s]	Samples	$N_{\rm side}$	[arcmin]	[$\mu K_{\rm CMB}$]
30 GHz	32.508	0.246	8	1024	3.44	$+11.9 \pm 0.7$
44 GHz	46.545	0.988	46	1024	3.44	-15.4 ± 0.7
70 GHz	78.769	1.000	79	1024/2048	3.44/1.72	-35.7 ± 0.6

[a] Sampling frequency.
[b] Baseline length in seconds and in number of samples.
[c] `HEALPix` $N_{\rm side}$ resolution parameter and averaged pixel size.
[d] Monopole removed from the maps and reported in the `FITS` header.

$$\begin{pmatrix} \dots & \dots & \dots & \alpha_1 & \alpha_2 \\ \dots & \dots & \dots & \alpha_1\cos(2\psi_i) & \alpha_2\cos(2\psi_i) \\ \dots & \dots & \dots & \alpha_1\sin(2\psi_i) & \alpha_2\sin(2\psi_i) \\ \alpha_1 & \alpha_1\cos(2\psi_i) & \alpha_1\sin(2\psi_i) & \alpha_1^2 & 0 \\ \alpha_2 & \alpha_2\cos(2\psi_i) & \alpha_2\sin(2\psi_i) & 0 & \alpha_2^2 \end{pmatrix}. \quad (10)$$

The *Planck* scanning strategy allows only a limited range of radiometer orientations. We therefore compute a joint solution with all radiometers at each frequency that helps also to reduce the noise in the final solutions. Once spurious maps are derived, we compute the a-factors from a χ^2 fit between the leakage map L and the spurious maps S_k on those pixels close to the Galactic plane, $|b| < 15^\circ$ (at higher latitudes both foregrounds and spurious signals are weak and do not add useful information).

The last improvement involves the final step in the creation of the correction maps. Recall that polarization data from a single radiometer probe only one Stokes parameter in the reference frame tied to that specific feedhorn. This reference frame is then projected onto the sky according to the actual orientation of the spacecraft, which modulates the spurious signal of each radiometer into Q and U. This modulation can be obtained by scanning the estimated spurious maps $\hat{S} = aL$, to create a timeline that is finally reprojected into a map. In the previous release, instead of creating timelines and then maps (a time- and resource-consuming operation) we built projection maps $A_{Q[U]}$ that accounted exactly for horn and radiometer orientation. The final correction maps were

$$\Delta Q[U] = L \times \sum_k a_k A_{k,Q[U]}\,, \quad (11)$$

where a_k and $A_{k,Q[U]}$ are the a-factors and the projection map for the radiometer k of a given frequency. In using this approach, however, there were two drawbacks. The first and more important one is related to a monopole term present in the leakage map L that directly impacted the correction maps. The synchrotron component, in fact, has a significant quasi-isotropic component (perhaps related to the ARCADE2-measured excess; Fixsen et al. 2011) and this contributed exactly to a monopole term in the correction map. Q and U maps, however, are produced by `Madam`, which tends to remove any possible monopole term. Therefore, with the simple approach of Eq. (11), we transfered the monopole term into the correction, and hence into the final bandpass-corrected map, resulting in an overestimation of the actual real effect. This was negligible at 70 GHz, but important in the 30-GHz map, which is used to correct foregrounds in the 70-GHz likelihood power spectrum estimation. The second drawback was that the resulting correction maps displayed sharp features, especially around the ecliptic poles. These were intrinsic to the projection maps, and caused problems with nearby point sources.

To resolve these issues, given new computing resources available, we exploit the scanning, timeline creation, and mapmaking approach. This is done using the *Planck* `LevelS` simulation package (Reinecke et al. 2006), which takes the harmonic coefficients $a_{\ell m}$ of the leakage maps, multiplied by the derived a-factors, and the actual scanning strategy, and then creates timelines accounting for proper beam convolution as well. The resulting TOD are used to create maps with the `Madam` mapmaking code. It is clear that in this way the final correction map is processed by the same mapmaking used for official map production and hence removes the presence of unwanted monopoles. In addition, accounting for beam convolution significantly alleviates the presence of sharp features in the correction maps.

Table 7 gives the estimated a-factors for the current release. They are very close to the 2015 values at 30 and 44 GHz, with larger (but within $1\,\sigma$) variations at 70 GHz. We investigated the origin of these variations by computing the a-factors with the present data, but using the old version of the leakage map. We find results in agreement with those obtained before. We also performed the same analysis with old data but with the new version of the leakage maps. In this case, we find results in line with the current estimates. These very simple tests clearly indicate that the observed variations in the a-factors are not due to the adopted calibration pipeline, but are mainly due to the changes in the leakage maps derived using *Planck* data only and excluding CO-dominated HFI channels.

Table 7. Bandpass mismatch a-factors from a fit to $S_k = a_k L$.

Horn	a-factor
70 GHz	
LFI 18	-0.0030 ± 0.0029
LFI 19	0.0197 ± 0.0030
LFI 20	0.0051 ± 0.0032
LFI 21	-0.0189 ± 0.0031
LFI 22	0.0063 ± 0.0031
LFI 23	0.0095 ± 0.0031
44 GHz	
LFI 24	0.0038 ± 0.0004
LFI 25	0.0006 ± 0.0004
LFI 26	0.0014 ± 0.0004
30 GHz	
LFI 27	0.0058 ± 0.0001
LFI 28	-0.0101 ± 0.0001

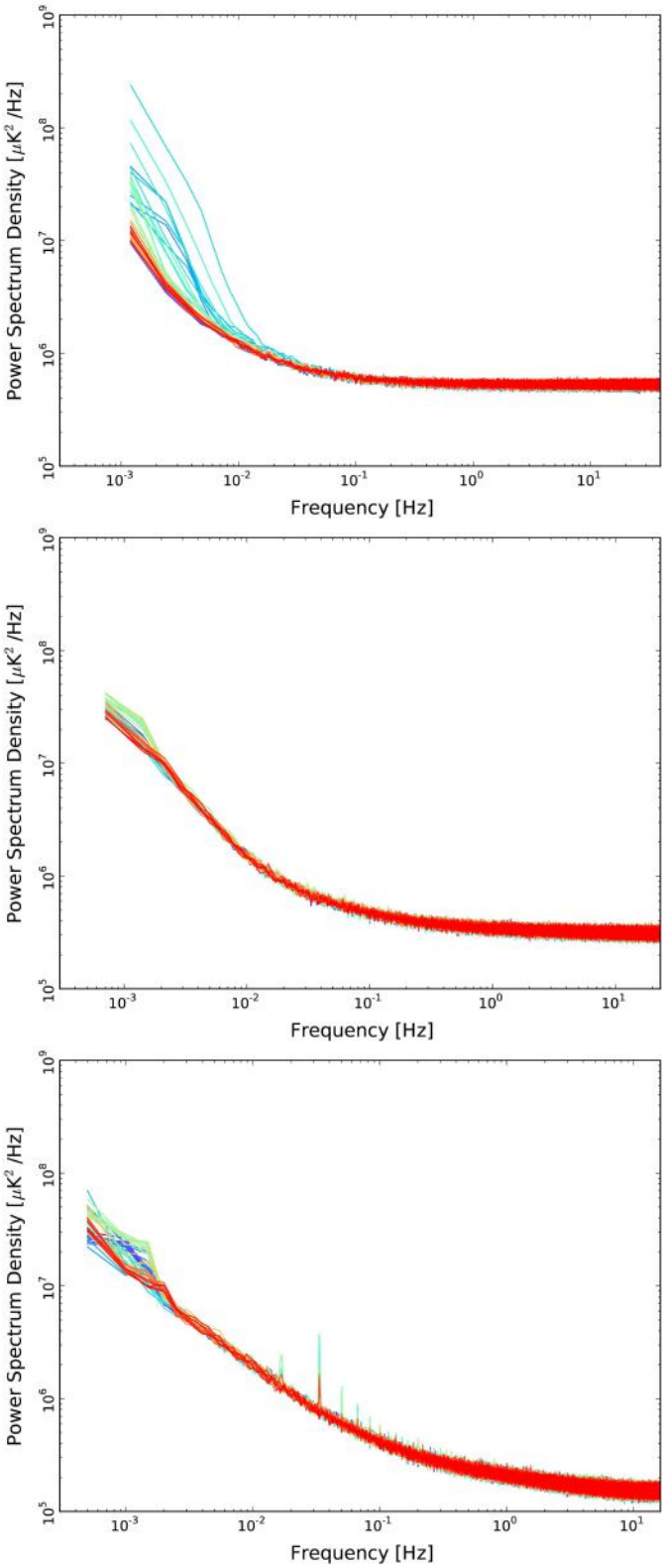


Fig. 9. Evolution of noise spectra over the mission lifetime for radiometer 18M (70 GHz top), 25S (44 GHz middle), and 27M (30 GHz bottom). Spectra are colour-coded, ranging from OD 100 (blue) to OD 1526 (red), with intervals of about 20 ODs. White-noise levels and slope stability are considerably better than in the 2015 release, being at the 0.1 % level for noise, while knee-frequencies show variations at the 1.5 % level.

8. Data validation

We verify the LFI data quality with the same suite of null tests used in previous releases and described in Planck Collaboration II (2016). As before, null tests cover different timescales (pointing periods, surveys, survey combinations, and years) and data (radiometers, horns, horn-pairs, and frequencies) for both total intensity and polarization. These allow us to highlight possible residuals of different systematic effects still present in the final data products.

8.1. Comparison between 2015 and 2018 frequency maps

Before presenting the null-test results, we compare the 2015 and 2018 maps. We expect improvements especially at 30 and 44 GHz, where the calibration procedure is significantly changed. Figure 13 shows differences between 2018 and 2015 frequency maps in I,Q, and U. Large scale differences between the two set of maps are mainly due to changes in the calibration procedure, but the exact origin of the differences is not revealed by these overall frequency maps.

A clearer indication of the origin of improvements in 2018 is given by survey differences at the frequency map level in temperature and polarization. From results for the previous release, we know that odd minus even surveys are the most problematic because of the low dipole signal in even numbered Surveys (especially in Surveys 2 and 4), which increases calibration uncertainty. This indeed was the motivation for the changes made in the calibration pipeline. In addition, since the optical coupling of the satellite with the sky is reversed every 6 months, such survey differences are the most sensitive to residual contamination from far sidelobes not properly accounted for and subtracted during the calibration process. We therefore consider the set of odd-even survey differences combining all eight sky surveys covered by LFI. These survey combinations optimize the signal-to-noise ratio, and are shown in Fig. 14 with a low-pass filter to highlight large-scale structures. The nine maps at the top show odd-even survey differences for the 2015 release, while the nine maps at the bottom show the same for the 2018 release.

The 2015 data show large residuals in I at 30 and 44 GHz that bias the difference away from zero. This effect is considerably reduced in the 2018 release, as expected from the improvements in the calibration process. The I map at 70 GHz also shows a significant improvement. In the polarization maps, there is a general reduction in the amplitude of structures close to the Galactic plane: the Galactic centre region and the bottom-right structure in Q at 30 GHz, and the rightmost region on the Galactic plane in U.

Figure 15 shows pseudo-angular power spectra from the odd-even survey differences, using the same sky mask as for the null-test spectra in Sect. 8.2, namely the union of all the single survey masks. There is great improvement in 2018 in removing large-scale structures at 30 GHz in TT, EE, and somewhat in BB, and also in TT at 44 GHz. These improvements are again expected and the reason is two-fold. First, the improved calibration now allows better tracing of the actual instrument gain, even in the low-dipole-signal period, both for 30 and 44 GHz. Second, the improved calibration enhances our ability to remove far-sidelobe contamination, resulting in significantly cleaner 30-GHz maps. At 70 GHz, even though the calibration procedure is almost unchanged from the 2015 release, we are able to reduce large-scale residuals in TT, thanks to the combined effect of data selection and the gain smoothing algorithm.

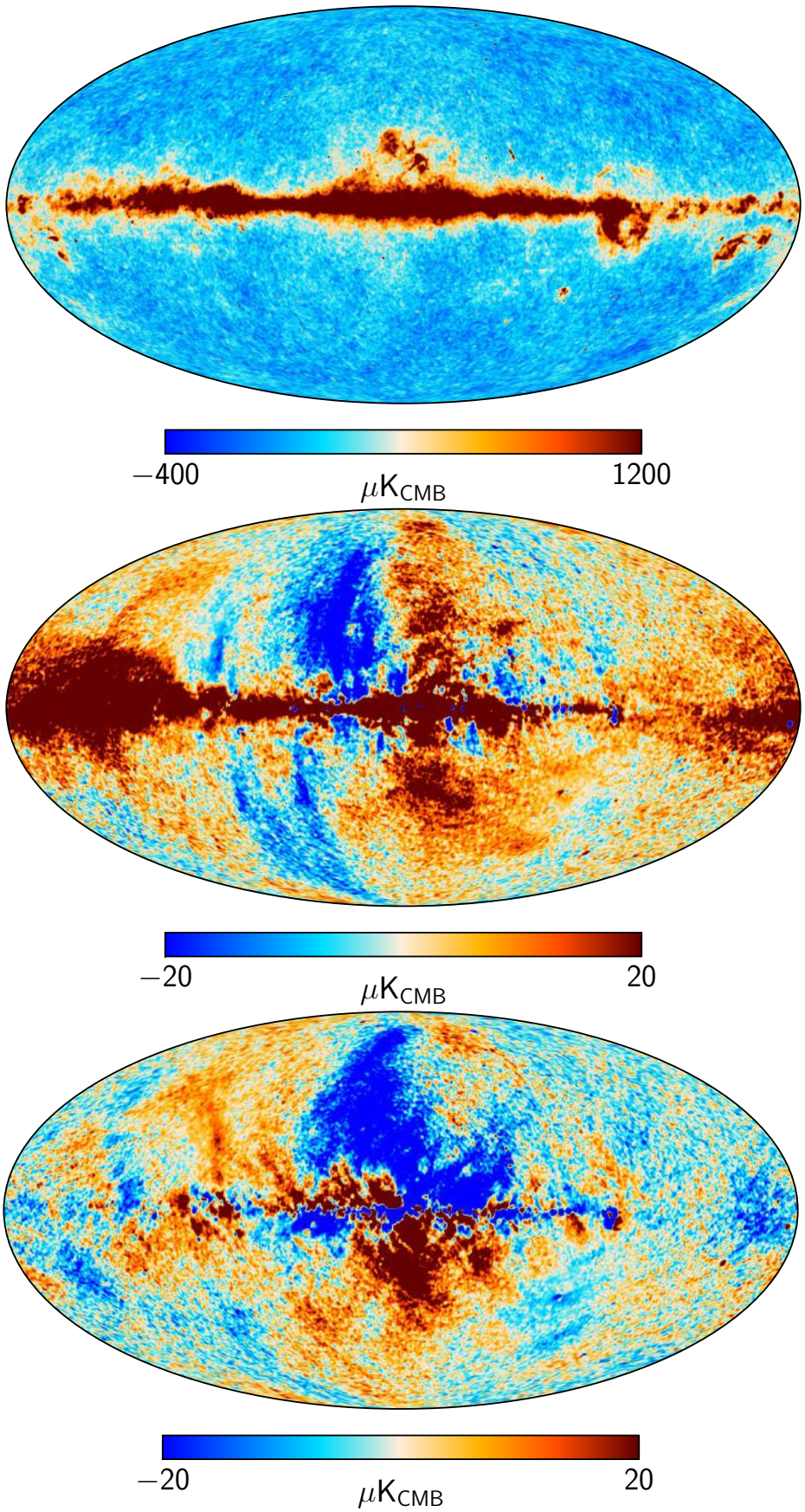


Fig. 10. LFI maps at 30 GHz: top, total intensity I; middle: Q polarization component; bottom, U polarization component. Stokes I is shown at instrument resolution and at $N_{\rm side} = 1024$, while Q and U are smoothed to 1° resolution and at $N_{\rm side} = 256$. Units are $\mu K_{\rm CMB}$. The polarization components have been corrected for bandpass leakage (Sect. 7)

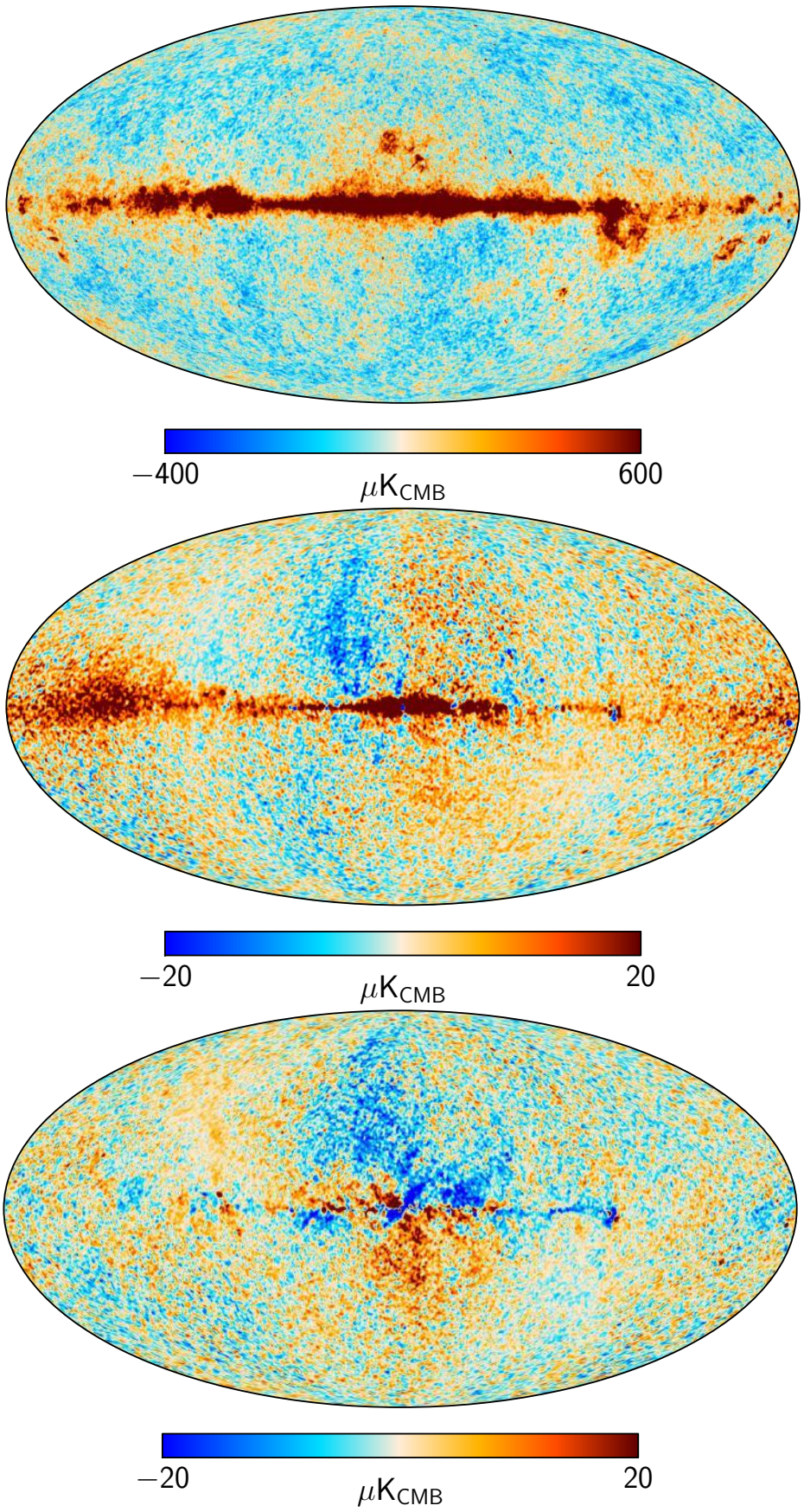


Fig. 11. Same as Fig. 10, for the 44-GHz channel.

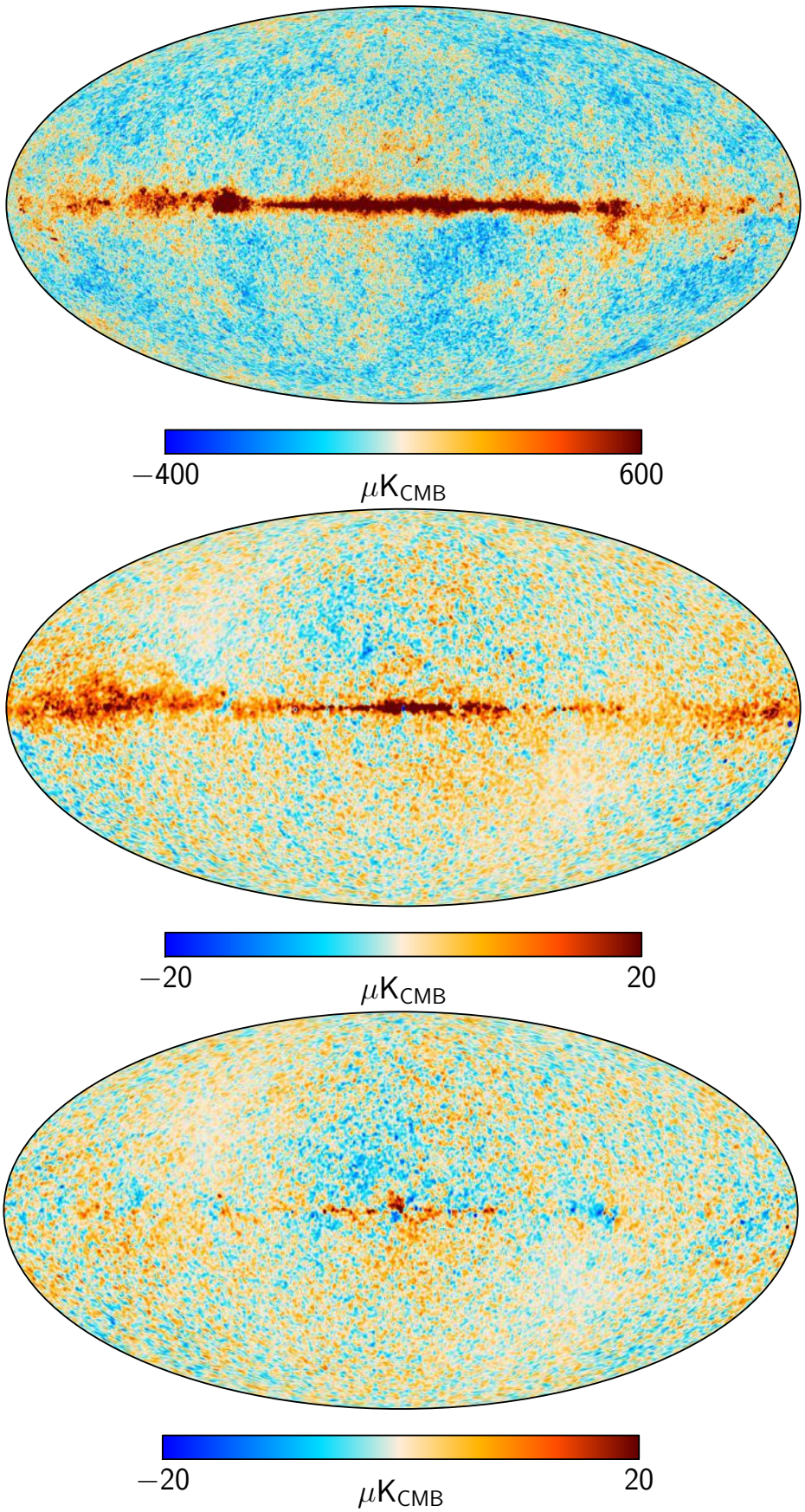


Fig. 12. Same as Fig. 10, for the 70-GHz channel.

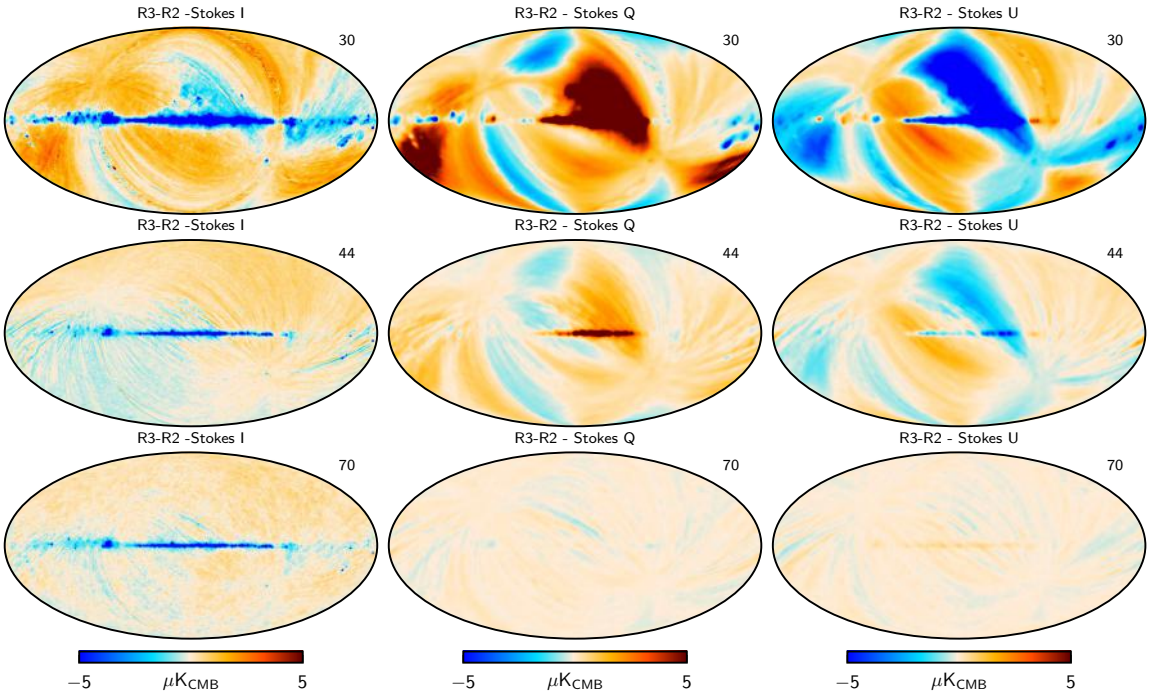


Fig. 13. Differences between 2018 (PR3) and 2015 (PR2) frequency maps in I, Q, and U. Maps are smoothed to 1° angular resolution for I and to 3° for Q and U, in order to highlight large-scale features. Differences are clearly evident at 30 and 44 GHz, and are mainly due to changes in the calibration procedure.

8.2. Null-test results

These findings are confirmed by specific null tests, taking differences of frequency maps for odd and even surveys. As for the previous release, we present differences among the first three sky surveys. Figure 16 shows the total amplitude of the polarized signal at 30 GHz (the channel with the largest expected differences), smoothed with an 8° Gaussian beam. Odd-even Survey differences reveal clear structures on large angular scales that are significantly reduced in the 2018 data set. In contrast, the Survey 1 versus Survey 3 difference map shows no large-scale features. This is expected, since for both Surveys 1 and 3 the dipole signal used for calibration is large. Moreover, the far sidelobes are oriented similarly with respect to the sky for these two surveys.

We also inspect angular power spectra of odd-even survey differences, adopting as a figure of merit the noise level derived from the 'half-ring' difference maps (made from the first and second half of each stable pointing period) weighted by the hit count. This quantity traces the instrument noise, but filters away any component fluctuating on timescales longer than the pointing period. To illustrate the general trend in null tests and the improvements in the 2018 release, Fig. 17 shows TT and EE Survey-difference power spectra for the 2015 and 2018 data sets. We compare these spectra with noise levels derived from the corresponding half-ring maps.

Results at 30 and 44 GHz are in line with expectations. In particular we see improvements at 30 GHz (survey differences are close to half-ring spectra) when considering odd-even survey differences. The better agreement results from the improved treatment of residual polarization by iterating Galactic modelling during calibration. The 44- and 70-GHz results are basically in line with the previous release findings. That is of course expected at 70 GHz, since the calibration procedure is almost the same as in the previous release, except for the gain smoothing algorithm and the foreground model adopted (now based on the `Commander` solutions using only *Planck* data).

A more quantitative way to represent null-test results, especially at low multipoles, is to compute deviations from the half-ring noise in terms of

$$\chi^2_\ell = \frac{\sqrt{2\ell+1}}{2}\left(\frac{C_\ell^{\rm SS} - C_\ell^{\rm hr}}{C_\ell^{\rm hr}}\right). \quad (12)$$

We specifically sum each single χ^2_ℓ in the range $\ell = 2$–50. Then, from the total value of χ^2 and $N_{\rm dof}$, we derive p-values of the distribution. While a proper set of noise simulations should in principle be considered, for this inspection it is adequate to use simple half-ring noise. Nonetheless we should be aware of the fact that any result derived with this approach is only indicative of possible issues and that a more detailed and refined analysis is required. Table 8 reports both χ^2 and p-values from the three survey differences, as shown in Fig. 17 for polarization spectra at the three LFI frequencies for the 2018 and 2015 data releases. A comment is in order here. On the one hand, we see that Survey 2 and Survey 4 seem to have some problems at 70 GHz, as highlighted by the poor χ^2 and p-values. However, this is expected, since we made only relatively small changes in the calibration pipeline at 70 GHz, and these surveys were known to be problematic in 2015. Nevertheless, we can anticipate that a power spectrum analysis of low-ℓ polarization at 70 GHz will find good results even including Survey 2 and Survey 4, thanks to the use of the calibration template described in Sect. 3. We note, however, that such a template does not help to improve the χ^2 and p-values, since these are derived from survey differences; any global template applied to data from both surveys would cancel out and leave χ^2 results unaffected. While at 44 GHz the picture is practically unchanged with respect to 2015, results at 30 GHz show in general a good trend of improvement for even survey differences, as indicated by χ^2 values, and underlining again the

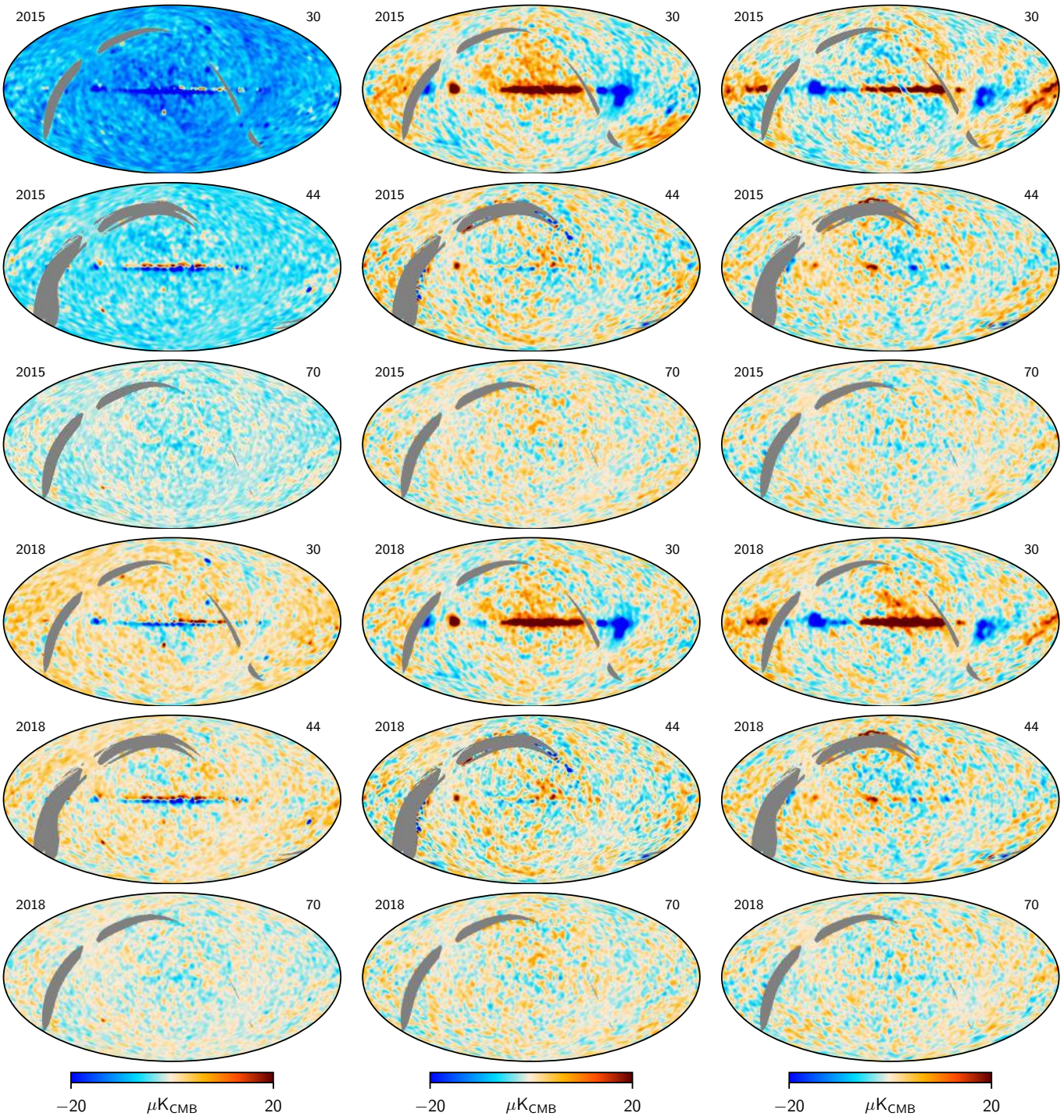


Fig. 14. Differences between odd (i.e., Surveys 1, 3, 5, and 7) and even (Surveys 2, 4, 6, and 8) surveys in I, Q, and U (from left to right) for the 2015 (upper nine maps) and 2018 (lower nine maps) data releases. These maps are smoothed to 3° to reveal large-scale structures.

benefit of the new calibration scheme. However, such values are far from being optimal and may indicate the presence of residuals showing up in the difference maps. Moreover we stress that this kind of analysis is only indicative and is used internally as an additional validation test.

8.3. Half-ring test

The actual noise in the LFI data is given directly by the half-ring difference maps. Detailed noise characterization is of paramount importance for the creation of adequate noise-covariance matrices (NCVMs), as well as for the noise MC realizations that are required in subsequent steps of the data analysis. Any noise model has to be validated against such half-ring difference maps. In the current release, we follow the same processing steps as in the previous releases. Specifically, we compute `anafast` auto-

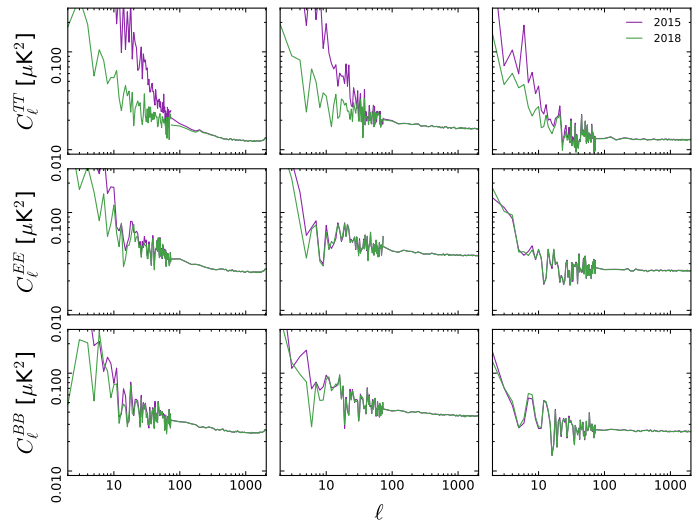


Fig. 15. Angular pseudo-power spectra of the odd-even survey difference maps for 30 (left column), 44 (middle column), and 70 GHz (right column), with the 2015 data in purple and 2018 in green.

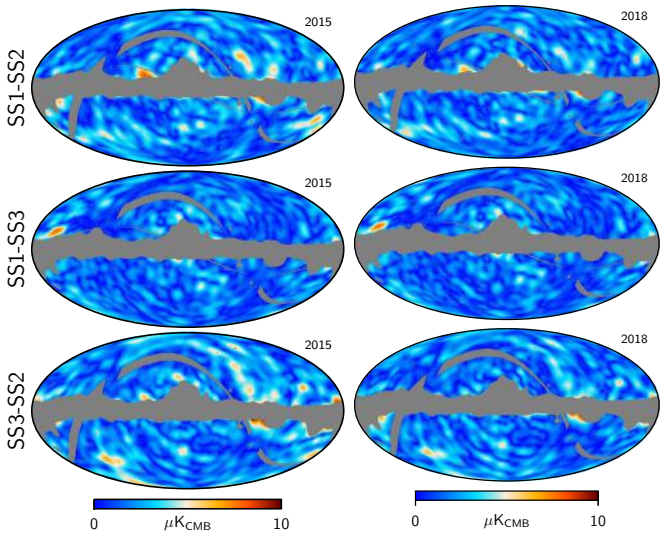


Fig. 16. Survey difference maps of polarization amplitude at 30 GHz for the current 2018 (right) and 2015 (left) releases. The improvement is evident, especially in odd-even difference maps, showing lower residuals due to the new calibration approach. Maps are smoothed with an 8° Gaussian beam to show large-scale structures.

spectra in temperature and polarization of the half-ring difference maps for the period covering the full mission. This is also done on MC noise simulations produced using noise estimation at the TOI level taken from FFP10.[3] Half-ring noise power spectra are compared with the distribution of noise spectra derived from the noise simulations and with the white-noise level computed from the white-noise covariance matrices (WNCVM) produced during the mapmaking process.

Figure 18 shows such a comparison for TT, EE, and BB spectra. The grey bands represent the 16 % and 84 % quantiles of the noise MC, while the black solid line is the median (50 % quantile) of these distributions. The half-ring spectra are depicted in red, and for $\ell \geq 75$ are binned over a range of $\Delta\ell = 25$. Even by eye the agreement is extremely good, and makes us confident about proper noise characterization in LFI data.

[3] This is the latest version of the Full Focal Plane *Planck* simulations similar to the FFP8 version used for the 2015 releases (see Planck Collaboration X 2016 for futher details)

Table 8. Odd-even surveys χ^2 and p-values ($2 \leq \ell \leq 50$).

	2015		2018	
Survey Differences	χ^2	p-value	χ^2	p-value
30 GHz				
S1 − S2	89.91	1.6×10^{-4}	63.65	0.0539
S1 − S3	40.00	0.755	41.36	0.705
S1 − S4	246.3	$< 1\times10^{-10}$	146.1	4×10^{-12}
44 GHz				
S1 − S2	44.69	0.568	44.28	0.585
S1 − S3	100.4	9×10^{-6}	108.3	9×10^{-7}
S1 − S4	46.48	0.494	59.5	0.105
70 GHz				
S1 − S2	82.01	0.0012	86.48	3.97×10^{-4}
S1 − S3	63.11	0.0582	64.39	0.0467
S1 − S4	74.06	0.0071	74.28	0.0068

We futher investigate the noise properties in the high-ℓ regime, taking the average of C_ℓ in the range $1150 \leq \ell \leq 1800$ for both temperature and polarization, and then comparing with the WNCVM. Figure 19 displays the result. As already shown in previous releases, there is still an excess of $1/f$ noise in this high-ℓ regime, meaning that both the real data and the noise MCs predict slightly larger noise than the WNCVM. It is important to note that such noise excess is reduced considerably with resepct to the 2015 release, thanks mainly to the new and more accurate calibration procedure adopted. Residuals are $\lesssim 1.4$ % at 30 GHz, $\lesssim 1$ % at 44 GHz, and $\lesssim 0.6$ % at 70 GHz, for both temperature and polarization. In addition, agreement between actual noise data and MC simulations is extremely good, with deviations of only fractions of a percent.

8.4. Intra-frequency consistency check

Data consistency can also be checked by means of power spectra, as done in previous releases (Planck Collaboration II 2016, 2014). We consider frequency maps at 30, 44, and 70 GHz, and take the cross-spectra between half-ring maps at each frequency for the full mission time span. Taking cross-spectra has the advantage that we do not need to consider noise bias at the power spectrum level. We make use of the `cROMAster` code, a pseudo-C_ℓ cross-spectrum estimator (Hivon et al. 2002; Polenta et al. 2005). Results obtained are sub-optimal with respect to a maximum likelihood approach, but are less computationally demanding and accurate enough for our purposes.

The actual spectra are computed using a Galactic mask obtained with the combination of the *Planck* G040, G060, and G070 masks at 30, 44, and 70 GHz, respectively, and accounting for the proper frequency-dependent masks for resolved point sources. Figure 20 shows cross-spectra from the half-ring maps. They agree well among the three frequencies, which is remarkable since foregrounds have not been removed from the maps, except in the masked regions. The red-dashed lines are the 2015 *Planck* best-fit TT spectrum, augmented by the contribution from unresolved point sources. The data are in good agreement with this model at all three frequencies. More quantita-

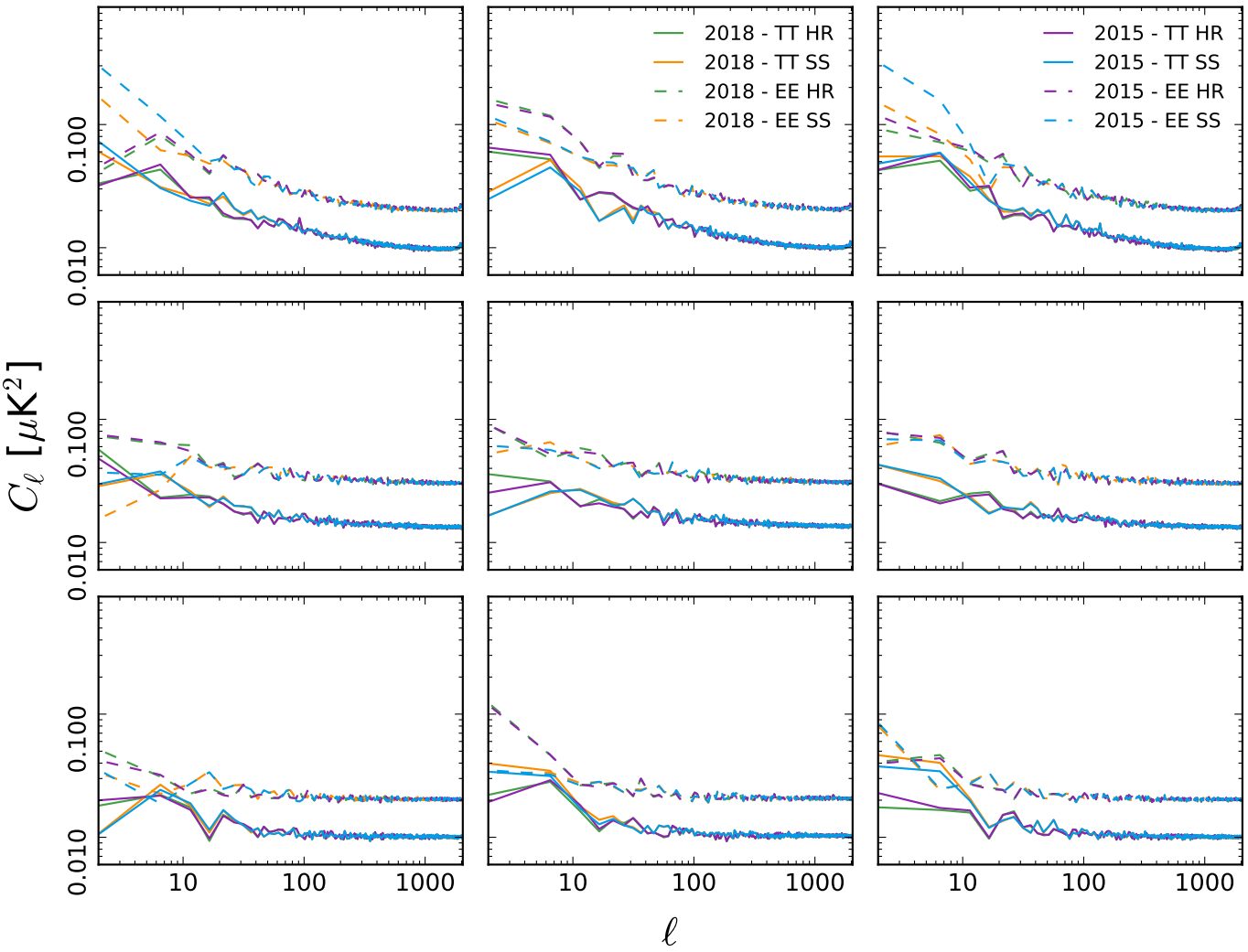


Fig. 17. Null-test results comparing power spectra from survey differences to those from the half-ring maps. Differences are: left, Survey 1 − Survey 2; middle, Survey 1 − Survey 3; and right, Survey 1 − Survey 4. These are for 30 GHz (top), 44 GHz (middle), and 70 GHz (bottom), for both TT and EE power spectra. There is a significant improvement in Survey 1 − Survey 2 and Survey 1 − Survey 4 at 30 GHz, especially in EE.

tively, we build a simple Gaussian likelihood (without any beam or foregrounds modelling) and consider multipole bins up to where the beam function blows up. We obtain the following $p-$values: 0.196 at 70 GHz ($50 \lesssim \ell \lesssim 1300$); 0.262 at 44 GHz ($50 \lesssim \ell \lesssim 500$); and 0.017 at 30 GHz ($50 \lesssim \ell \lesssim 500$). This shows that even with this simplified approach 70 and 44 GHz are consitent with the model. The 30 GHz channel is marginally consistent with the null hypothesis and clearly requires a more detailed treatment of foregrounds.

As a more quantitative test of data consistency, we build the usual scatter plots of angular power spectra for the three frequency pairs. To ensure a proper comparison, we remove the frequency-dependent contribution from unresolved sources, and perform a linear fit, accounting for errors on both axes. Figure 21 shows results in the multipole range around the first peak, where the effect of different angular resolutions at the three frequencies is still manageable. The agreement is extremely good: spectra are consistent with deviations between 0.9 % and 0.1 %, a result in line with findings from the previous release (Planck Collaboration II 2016). This agreement is also an indication of the calibration accuracy, which is at the sub-percent level. This is a sign of the good health of our data, given that fact that in these tests we have not accounted for any errors in the window function, calibration, or foreground removal.

8.5. Internal consistency check

We also check the internal consistency of the 70-GHz data used for the cosmological analysis. As done for the 2015 release, we create cross-spectra by taking different data splits. Specifically, we consider the half-ring maps and the year combination maps Year 1 − Year 3 and Year 2 − Year 4 maps. Figure 22 shows the residuals of the combination, compared with the expectations derived from FFP10 Monte Carlo simulations subjected to the same procedure. Residuals are clearly compatible with the null hypothesis. To quantify the agreement we apply the Hausman test (Polenta et al. 2005) to 70 GHz cross-spectra from half-ring and year maps. As reported in Planck Collaboration II (2016) we define the statistic

$$H_\ell = \left(\hat{C}_\ell - \bar{C}_\ell\right) \Big/ \sqrt{\mathrm{Var}\left[\hat{C}_\ell - \bar{C}_\ell\right]}, \qquad (13)$$

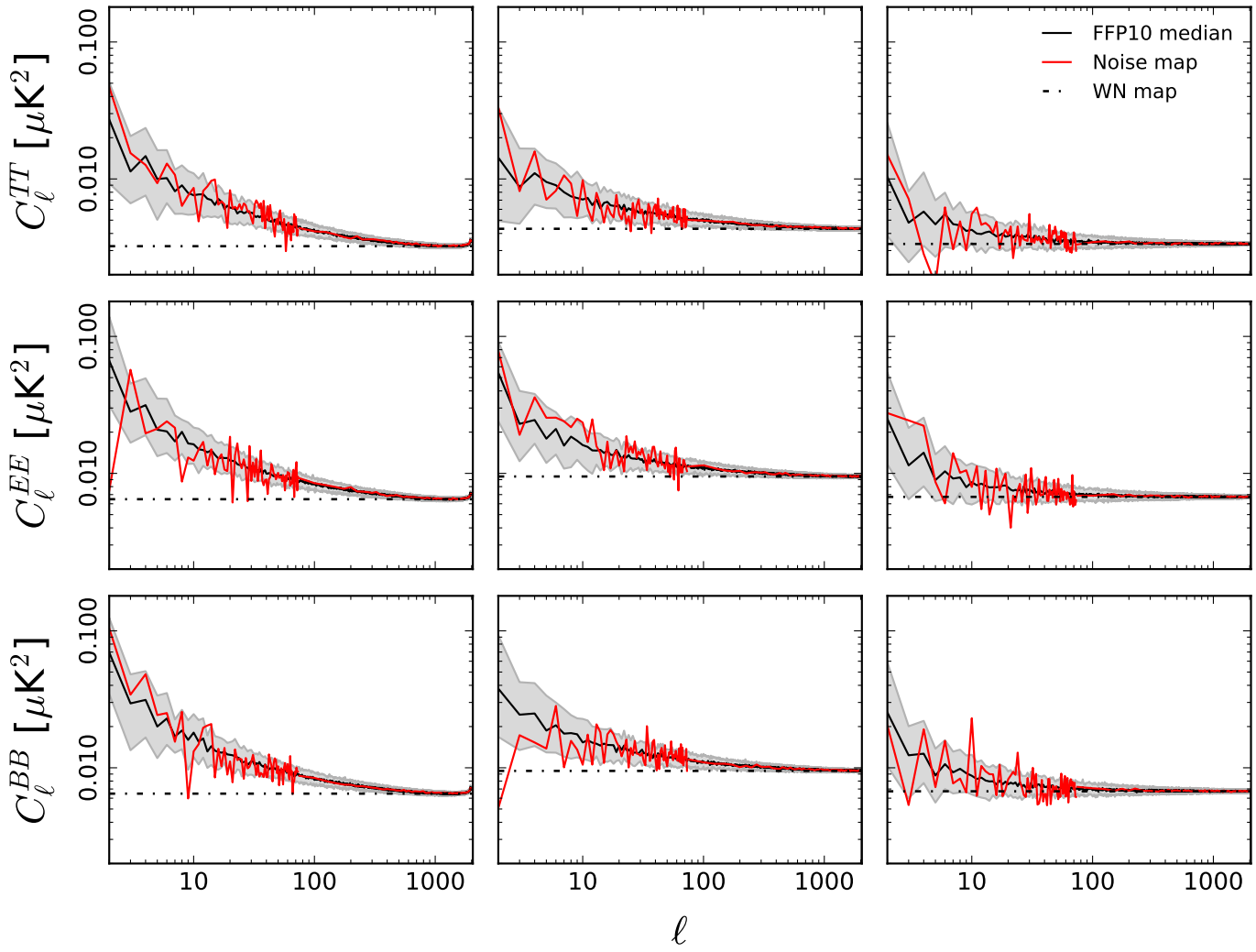


Fig. 18. Consistency check at the power spectrum level between half-ring difference maps (red), white-noise covariance matrices (black dash-dotted lines), and 100 full-noise MC simulations (grey band showing range for 16 % and 84 % quantiles of noise simulations, and the black solid lines giving the median, i.e., 50 % quantile, of the distributions). From top to bottom we show TT, EE, and BB power spectra for 30 GHz (left), 44 GHz (centre), and 70 GHz (right). Half-ring spectra are binned with $\Delta\ell = 25$ for $\ell \geq 75$.

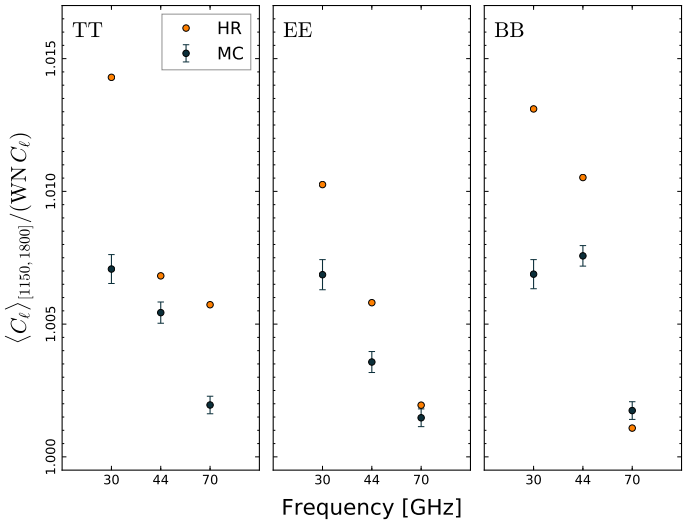


Fig. 19. Ratio of mean noise at high multipoles ($1150 \leq \ell \leq 1800$) to white noise, derived from white-noise covariance matrices produced during the mapmaking process

where $\hat{C}_\ell$ and $\bar{C}_\ell$ represent the two different cross-spectra. We further compress the multipole information with

$$B_L(r) = \frac{1}{\sqrt{L}} \sum_{\ell=2}^{[Lr]} H_\ell, \quad r \in [0,1], \tag{14}$$

where the operator [.] returns the integer part. It can be shown that the distribution of $B_L(r)$ converges to Brownian motion, which can be simply studied by means of three statistics: $s_1 = \sup B_L(r)$; $s_2 = \sup |B_L(r)|$; and $s_3 = \int_0^1 B_L^2(r)dr$. Results are reported in Fig. 23, where the three statistics are compared to their expected distributions derived from FFP10 simulations. The corresponding p-values for s_1, s_2, and s_3 are, respectively, 0.11, 0.19 and 0.13, showing again that results are perfectly compatible with the null hypothesis, and confirming the high level of internal consistency in the 70 GHz data.

8.6. Validation Summary

At the end of this detailed validation process, the improved quality of the 2018 data release, especially at 30 and 44 GHz, is clear.

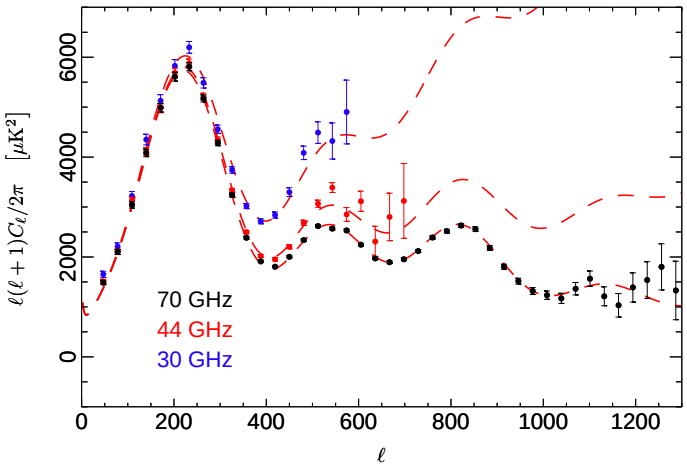


Fig. 20. Binned TT cross-spectra from half-ring maps at 30, 44, and 70 GHz. Foregrounds are accounted for only through the use of the proper frequency-dependent Galactic masks. Red-dashed lines are the *Planck* best-fit TT spectrum, to which a contribution from unresolved point sources (specific to each frequency) has been added.

This improvement is mainly due to the new calibration procedure implemented. It is also evident that residuals at very low multipoles (large angular scales) are still present at 30 and at 44 GHz. One of the reasons could be the fact that the initial foreground model used for the iterative calibration, mapmaking and component-separation process was entirely based on LFI data. We took this approach to avoid dependence on WMAP and to avoid ingesting any HFI systematic effects. However, we realize that this results in a foreground model with less power than expected. This choice should in principle be investigated further, but the actual level of residuals is small and does not prevent the use of the 30-GHz data as a synchrotron template in cosmological analyses involving the 70-GHz LFI channel (see Planck Collaboration V 2018, for a detailed analysis of the low-ℓ likelihood)

9. Updated systematic effects assessment

9.1. General approach

Analysis of LFI systematic effects from the start of the mission to the 2015 release (Planck Collaboration III 2014; Planck Collaboration II 2016; Planck Collaboration III 2016) identified uncertainty in calibration as the dominant source of systematic error. This source of error, and all other known systematic effects, are at least four orders of magnitude below the measured CMB power spectrum in total intensity for all relevant multipoles (see figures 24–26 of Planck Collaboration III 2016). However, this is not the case for polarization. Imperfect gain reconstruction, obtained independently for the Main and Side radiometers of each horn, translates directly into leakage of total intensity to polarization. In particular, at the large angular scales relevant for probing the reionization bump, the systematic error from calibration is comparable with the expected signal, for reasonable values of the optical depth parameter τ.

Our assessment of the overall systematic error budget remains essentially unchanged from the 2015 release, as summarized in table 1 of Planck Collaboration III (2016) and references therein. For the present release, we have concentrated on developing a detailed simulation programme to model all known instrumental and astrophysical effects that produce uncertainty in the gain for polarization data. We first identify those parameters in the whole calibration process that are affected by uncertainties, and then set up an ad hoc Monte Carlo simulation strategy to judge the effect of varying these parameters. This process has to be both accurate and realistic. There are two possible sources of error that are quite different in nature: *statistical* uncertainties, related to instrumental noise in conjunction with variations in the dipole amplitude on the sky; and *systematic* uncertainties due to effects and assumptions in the calibration process that are not completely known. The simulations require a trade-off between a full, physically representative set of simulations and a realistic, feasible number of simulations, given the computation resources available. Ideally, we would assume that all the effects are indeed correlated with each other, and therefore the total number of simulations would be the product of the number of simulations required for each single effect. However, to reduce computer resources needed, it is more convenient to assume that all the effects are un-correlated, and the final error is obtained by summing in quadrature the error derived from each single effect. This uncorrelated assumption is clearly not completely true: it is impossible to separate the effect of different parameter values of the gain smoothing algorithm from the actual noise realization. Therefore, we follow a hybrid approach in the systematic effect simulations.

9.2. Monte Carlo of systematic effects

The FFP10 (see Planck Collaboration X 2016) simulation pipeline is the basis of our simulations pipeline, and the sky signal we include makes use of the Planck Sky Model (Delabrouille et al. 2013). A detailed description of the sky model and components is given in Appendix B. In the current implementation, the FFP10 pipeline is partially executed in full timelines (data as a function of time at the original detector sampling frequency), and partially in the ring-set domain (data binned into a partial map for every pointing period covering full circles on the sky, as identified by `HEALPix` pixel indices).

We start by creating separate ring-sets for each signal of interest, using the actual pointing information. We have ring-sets for CMB, Galactic foregrounds, extragalactic diffuse signals, point sources convolved with the measured LFI beams, and finally the sum of the solar and orbital dipoles, also convolved with the measured LFI beams. In the second stage of the pipeline, we perform the iterative calibration. First we divide ring-sets by a fiducial set of gains (those estimated from actual data), and then we process ring-sets to estimate the gain value for each pointing period using the same pipeline as described (although using a different implementation) in Planck Collaboration V (2016). The reconstructed gains, however, are noisy and biased when the orientation of the line of sight of the telescope scans a region where the dipole is close to its minimum. Therefore we apply the same smoothing algorithm developed by the DPC for the real data. It is important to note here that the smoothing algorithm is optimized through a detailed study of real data, and we therefore do not expect that it will perform optimally on each single simulation. Observed differences in the final maps could be used to estimated directly the impact of the smoothing process on the reconstructed gains.

The final stage of simulations produces calibrated maps with a pipeline based on `TOAST`.[4] We combine timelines of the desired

[4] Time Ordered Astrophysics Scalable Tools, `http://hpc4cmb.github.io/toast/`, developed by the Computational Cosmology

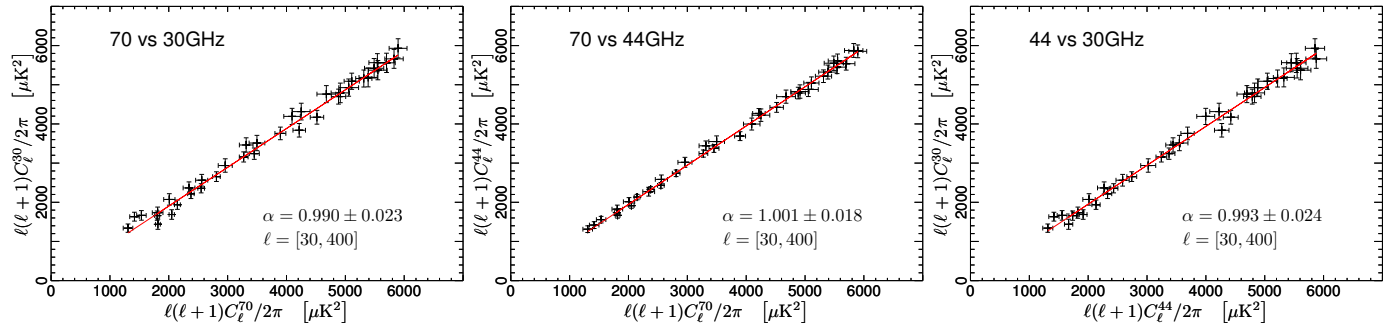


Fig. 21. Scatter plots of LFI angular cross-spectra around the first acoustic peak for frequency pairs: left, 70 GHz and 30 GHz; middle, 70 GHz and 44 GHz; right, 44 GHz and 30 GHz. The solid line is the linear fit obtained accounting for error bars on both axes. Slopes, within expected uncertainty, show good consistency among data sets

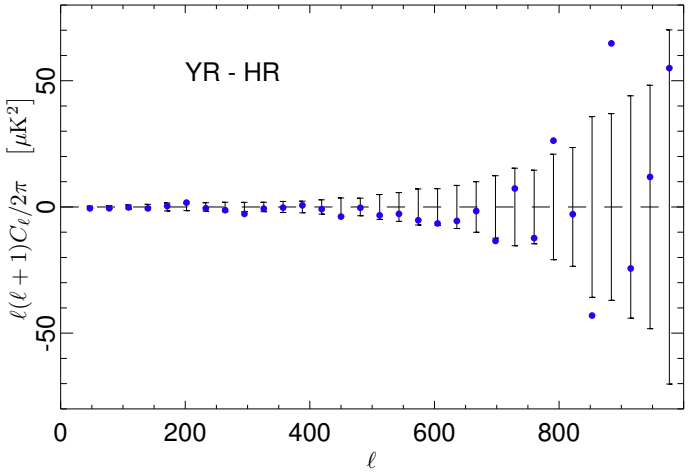


Fig. 22. Residuals between cross-spectra at 70 GHz from half-ring (HR) and year maps (YR). Error bars come from FFP10 Monte Carlo simulations.

input signals and noise (exactly the same as already used to create ring-sets for the calibration process), and divide by the fiducial gains, converted into kelvin with the previously computed set of reconstructed gains. From these calibrated timelines we subtract the simulated dipole signal, and then create maps with `Madam`. Any discrepancy between the fiducial and reconstructed gains shows up as a residual dipole and a mis-calibration of the Galactic signal and the noise. This gives an estimate of the calibration error that we believe also affects the data.

We have also examined in more depth the impact of known systematic effects on the calibration process, using targeted simulations. These simulations include noiseless maps for a few different realizations of each individual systematic effect: different solar dipole amplitudes and directions based on the expected error on the *Planck* dipole; different masks used inside the calibration pipeline during the gain fitting process; different beams used to convolve foreground signal and dipoles, based on the expected error on bandpasses and its impact on beams; ADC nonlinearities, based on the model fitted to real data; and finally different sky signals for input and calibrator, in order to simulate a discrepancy between input data and our sky model. We create simulations of each single effect plus a set of 20 simulations that include many different effects together by randomly sampling all the available options from dipole, beam, and ADC systematics, plus realistic noise simulations. It should be clear from this description that we are following a hybrid approach where, although simulated independently, systematic effects can be combined together to produce a new set of gains and maps.

In the following paragraph we describe results from these simulations of systematic effects. There are two outcomes. First, the results show how well we are able to simulate instrumental and data-processing effects in a direct comparison (mainly at power spectra level) with null tests on the data. Second, they provide a quantitative estimation of the amplitude of systematic effects. We will also summarize the overall systematic error budget, accounting for effects not directly simulated in 2018 by using simulations performed for the 2015 data release. In the evaluation of the various effects, we also used a perfect calibration simulation, where the pipeline is exactly the same except that we use the same set of fiducial gains in both the de-calibration and calibration processes.

9.2.1. Gain smoothing error

One of the most delicate steps in the calibration pipeline is the smoothing of the raw gains obtained from the dipole/sky fitting procedure. The smoothing algorithm is described in detail in Planck Collaboration V (2016), and has been tailored to the LFI data. Its performance is strictly linked to the actual noise in the data, the level of dipole signal with respect to the the sky signal, and the contribution of the far-sidelobe pickup. An algorithm tuned to the specifics of the real data might not perform equally well on any single simulation, and one might be tempted to optimize the smoothing procedure for each one. We choose not to do that, both because it would be a lengthly process, and also because it would introduce the gain smoothing algorithm itself as a new variable that would vary among simulations, making a direct comparison much more cumbersome. Instead using exactly the same smoothing procedure applied to real data allows us to evaluate its impact on data when compared to a perfect calibration simulation.

Figure 24 compares angular power spectra of odd-even-year null tests (in other words, the difference between the sum of Year 1 and Year 3 and the sum of Year 2 and Year 4) with the distribution (median plus 16 % and 84 % quantiles) derived from simulations. The simulations include the impact of the gain smoothing algorithm on the gain set that has been derived using the full set of systematic simulations, but applied to the noise-only timestream. We see that there is overall good agreement between data null tests and these pure-noise simulations. However, there are clearly some multipoles for which data null values are

Center C3 at Lawrence Berkeley National Laboratory to handle MPI-based parallel data processing.

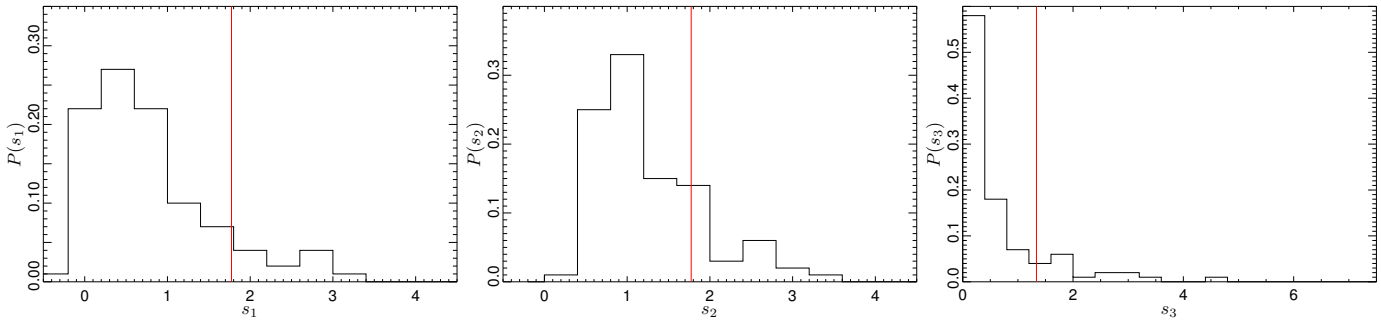


Fig. 23. Empirical distributions derived from FFP10 simulations for the s_1, s_2, and s_3 statistics of the Hausman test on the 70-GHz data. Vertical lines show values derived from half-ring and year maps at 70 GHz.

outside the range of simulations. This tells us that there are effects in the data other than the gain smoothing error. Indeed, we made the same comparison for perfect calibration simulations, and found a very similar plot with very small differences in the distribution of simulated spectra. In order to evaluate the impact of the gain smoothing error, we take the amplitude (for both perfect calibration and gain smoothing error) of the $1\,\sigma$ bands for some specific range of multipoles. Considering $\ell = 4$–5, the overall effect of the gain smoothing algorithm for the three frequency channels is to increase noise by $\lesssim 3 \times 10^{-3}\,\mu\mathrm{K}^2$ in total intensity and $\lesssim 6 \times 10^{-3}\,\mu\mathrm{K}^2$ in polarization, almost the same for all three frequencies. Considering higher multipole ranges, $\ell \simeq 100$ and $\ell \simeq 500$, errors are reduced by at least one order of magnitude with respect to the $\ell = 4$–5 range, ranging between 10^{-4}–$10^{-5}\,\mu\mathrm{K}^2$, both in temperature and polarization.

9.2.2. ADC nonlinearities

The ADCs convert analogue detector outputs into numbers, and any nonlinearities in their response could mimic a sky signal thereby affecting data calibration. As mentioned in Sect. 2, for the current release we implemented a new method to track and correct ADC nonlinearities that improves the quality of the data at 30 GHz.

As well as the analysis done in the definition of the new correction method, we also performed specific simulations of the effects of nonlinearities in the ADCs. We created 10 noiseless simulations with the ADC effect, based on a model fitted with real data. In these simulations, we randomized the errors in the voltage steps associated with each binary bit of the ADC in a way that is consistent with what we found from the data. Results are in line with the findings in 2015, with improvements at 30 GHz where residuals decrease by almost an order of magnitude. Again considering the range $\ell = 4$–5, the ADC effects now contribute an increase in the noise of $\lesssim 10^{-4}\,\mu\mathrm{K}^2$ at 30 GHz and around $10^{-5}\,\mu\mathrm{K}^2$ at 44 and 70 GHz. At higher multipoles the effect drops below $10^{-6}\,\mu\mathrm{K}^2$. In polarization the effect is at least one order of magnitude smaller than in total intensity.

9.2.3. Full systematic simulations

The final set of simulations we consider includes all the effects that we expect to directly impact calibration accuracy. We now present the results obtained, together with a final comprehensive table of the estimated impact of systematic effects. In addition, we create a summary plot like the ones in Planck Collaboration III (2016), but with updates for those effects being simulated for the current data release, including plots of EE and BB.

We begin by comparing the full systematic error budgets between the current and the 2015 data release. Figure 25 shows TT and EE angular power spectra at the three LFI frequencies. Systematic simulations made for the two data releases are in very good agreement; this justifies our claim that the overall systematic error budget remains essentially unchanged between 2015 and 2018. Moreover, the null spectra from odd-even-year differences are very close to the the systematic error expectations, both in temperature and polarization. There are a very few multipoles where null spectra lie outside the $1\,\sigma$ band of the simulations. Specifically, for TT they are $\ell = 2$ and 3 at 30 GHz, and $\ell = 2$ for both 44 and 70 GHz. The extra variance in the null tests is $\lesssim 0.04\,\mu\mathrm{K}^2$ at 30 and 70 GHz, and around $0.065\,\mu\mathrm{K}^2$ at 44 GHz. At higher multipoles, the agreement with simulations, dominated by instrumental noise, is extremely good.

For the EE power spectra, we can see again very good agreement between systematic simulations performed for the 2015 data release and those done for the current release. We see that the overall account of systematics estimated in 2015 falls very close to the median of the present simulations and, in any case, well within the $1\,\sigma$ limits. In addition, the odd-even-year null difference is extremely well represented by our simulations, except at a few multipoles. This again tells us that we understand well the instrumental and systematic effects actually present in the data. Evaluating the extra variance in the data not accounted for by simulations, we see that at 30 GHz the agreement is extremely good, at 44 GHz, there is an excess at $\ell = 4$–6 of $\lesssim 0.05\,\mu\mathrm{K}^2$, while at 70 GHz, the very low multipoles are in line with expectations (although for $\ell = 10$ there is an excess of about $0.10\,\mu\mathrm{K}^2$).

Table 9 summarizes systematic effects at the power spectrum level for three multipole ranges: $\ell = 4$–5; around the CMB first peak $\ell \simeq 100$; and the almost noise-dominated regime $\ell \simeq 500$.

A comment is in order here: for multipoles $\ell \gtrsim 100$, it is clear that the extra noise induced by systematic effects is well described by simulations, and accounted for in the overall estimation of the noise budget. At large angular scales, the improvement with the new calibration scheme is evident from the comparison of null-test spectra on different data combinations between the 2015 and 2018 data releases. However, there are a few multipoles that deserve particular attention, since specific null tests still show a noise excess that is not completely traced by the simulations of systematics. Nevertheless, results presented in this section are intended to be a useful indication of the overall amplitude of systematic effects. More scientifically-oriented analysis could in principle make use of the simulations to build

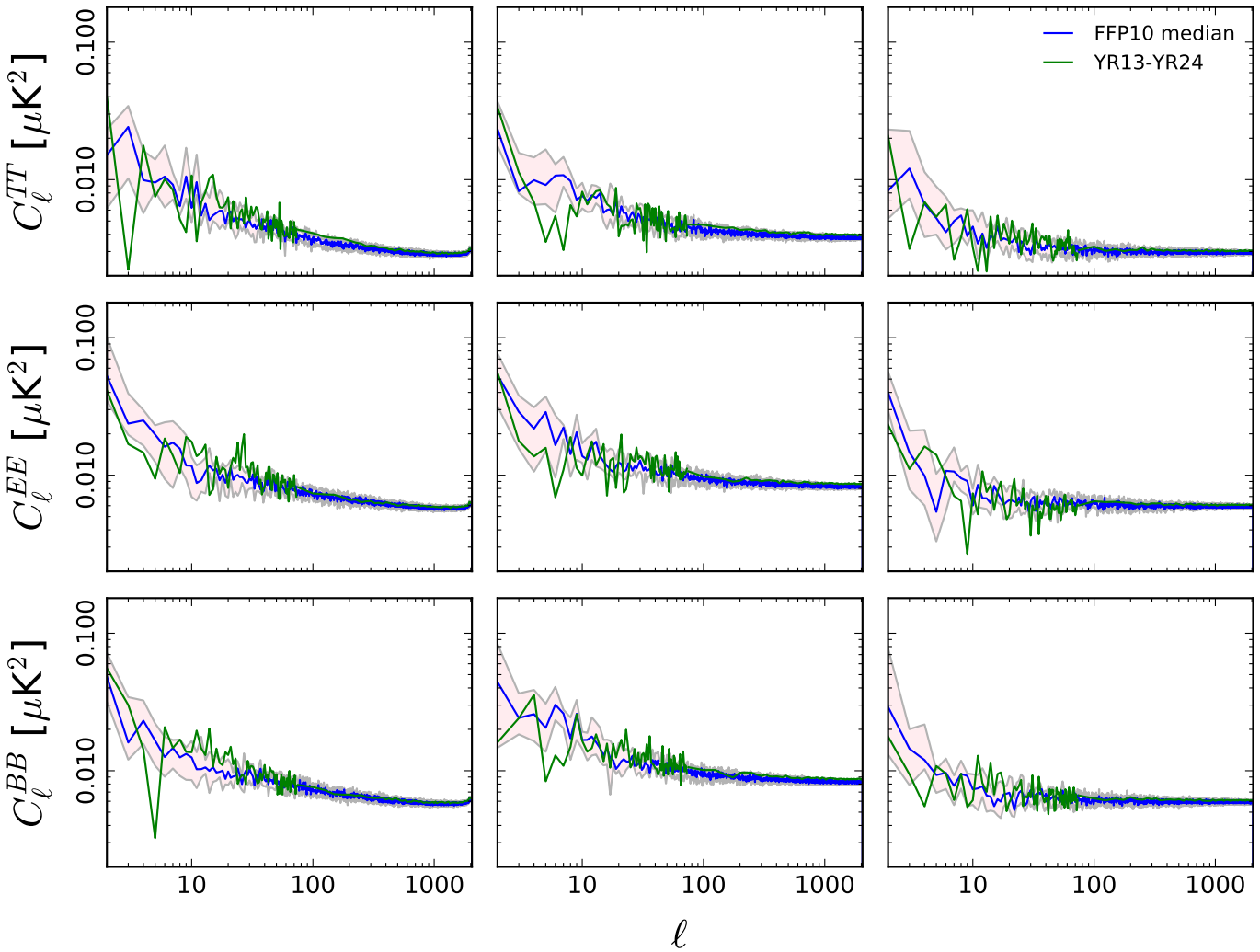


Fig. 24. Odd-even-year null spectra compared with simulations including the gain smoothing algorithm at 30 (left), 44 (middle), and 70 GHz (right). The pink band shows the 16 %–84 % quantile range of the simulations, with the median traced by the blue line. On large angular scales the data show larger variations with multipole than the simulations.

a template that could be fitted for in a maximum likelihood approach (along the same lines as foreground emission templates), and where errors in the simulations could be properly propagated into the final low-ℓ angular power spectrum and cosmological parameters.

10. LFI data products available through the *Planck* Legacy Archive

Before concluding, we provide a list and short description of the *Planck* LFI data products available through the *Planck* Legacy Archive,[5] based on the data covering the operational lifetime of the instrument from 12 August 2009 to 23 October 2013 (for further details on the data format refer to the Explanatory Supplement, Planck Collaboration ES 2018).

- Pointing timelines: identical to the 2015 release. One FITS file for each OD for each frequency. Each FITS file contains the OBT (on-board time) and the three angles, θ, ϕ, and ψ that identify each sample on the sky.
- Time timelines: identical to the 2015 release. One FITS file for each operational day (OD) for each frequency. Each FITS file contains the OBT and its corresponding TAI (International Atomic Time) value (without leap second) in modified Julian-day format. The user can thus cross-correlate OBT with UTC.
- Housekeeping timelines: identical to the 2015 release. All housekeeping parameters with their raw and calibrated values, separated by the housekeeping sources, for each OD.
- Timelines in volts: raw scientific data in volts for each detector at 30, 44, and 70 GHz, and each OD, before any calibration procedure and with no instrumental systematic effects removed.
- Cleaned and calibrated timelines: provided in K_{CMB} units, for each detector at 30, 44, and 70 GHz, and each OD, after scientific calibration and with convolved dipoles and convolved Galactic straylight removed.
- Scanning beam: 4π beam used in the calibration pipeline.
- Effective beam: sky beam representation as a projection of scanning beam on the maps.
- Full-sky maps at each frequency: maps of the observed sky at 30, 44, and 70 GHz in temperature and polarization at $N_{side} = 1024$, and also at $N_{nside} = 2048$ for 70 GHz. Maps are delivered for different data periods. We note that the *Planck* adopted polarization convention is not the one proposed by "IAU" but the one used more generally in

[5] http://archive.esac.esa.int/pla2

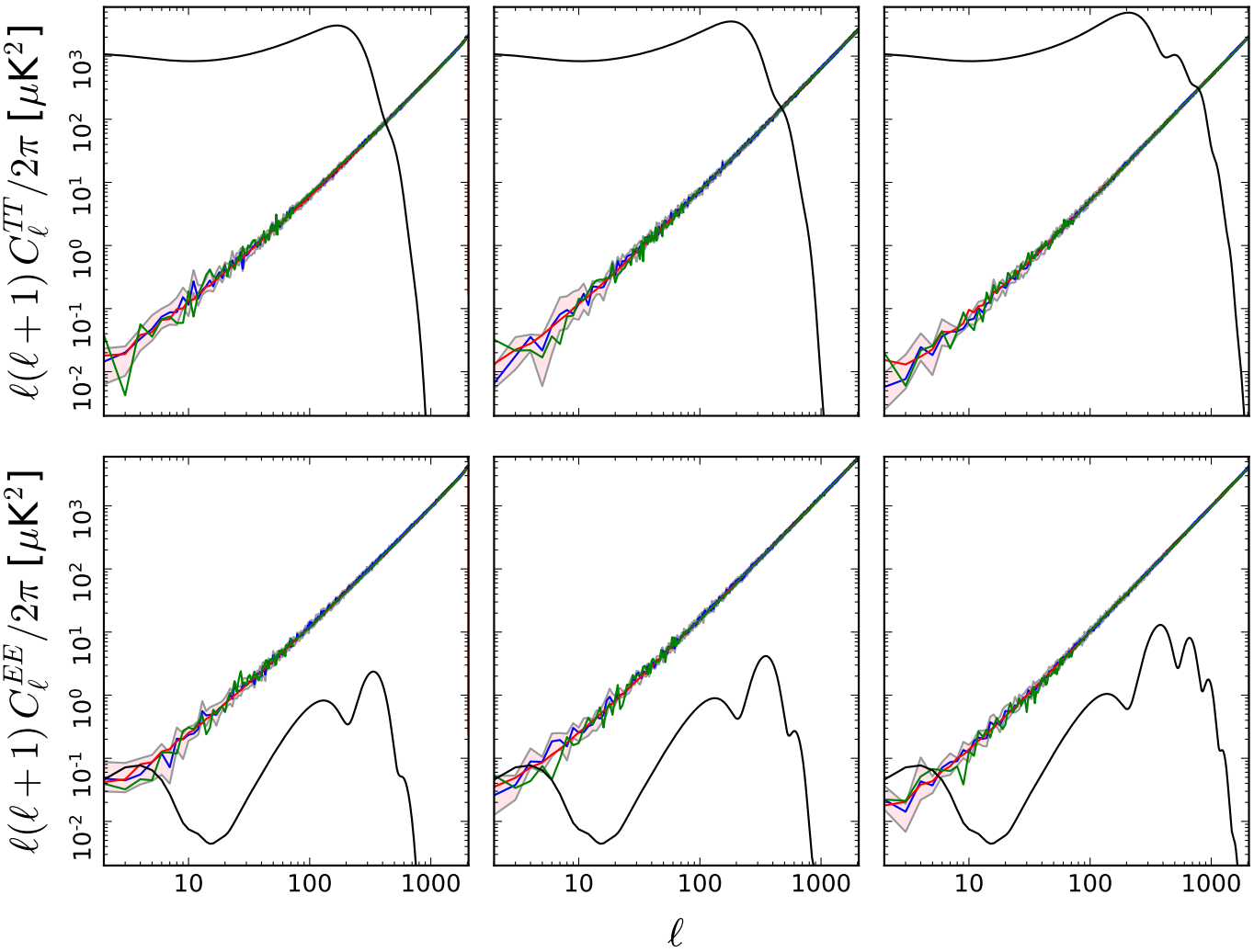


Fig. 25. Angular TT (top) and EE (bottom) power spectra at 30 (left), 44 (middle), and 70 GHz (right) of a complete set of systematic effects for the 2015 data release (red) and the 2018 simulations (blue, 50 % percentile plus 1 σ band from simulations). This shows the good agreement between the two releases. In green we show the null spectra from odd-even year difference. The CMB angular power spectra (black lines) are convolved with the frequency dependent window functions.

CMB full-sky maps and referred to as "COSMO" (see Planck Collaboration ES 2018).
- Bandpass correction maps at each frequency and maps with the bandpass correction applied to delivered frequency sky maps (one specific example of each different data period) computed according to the prescription detailed in Sect. 7.
- Gain Correction Template: template map to be subtracted from the delivered full sky map at 70 GHz in order to account for calibration uncertainties. This template has to be removed prior of any cosmological exploitation of the 70 GHz map. For completeness we also deliver the 70 GHz map with both bandpass and gain correction template applied.
- RIMO (Reduced Instrument MOdel): model which includes parameters for the main instrument properties, including noise, bandpass, and beam function.

11. Discussion and conclusions

We have presented a comprehensive description of the LFI data analysis pipeline that produces final frequency maps to be used for scientific exploitation. The major improvements with respect to the 2015 data release are in the new iterative calibration procedure at 30 and 44 GHz, which uses sky estimation and component separation to create a sky model to be fed into the calibration algorithm. Other minor improvements are in the redefinition of the data flags that allow better selection of data.

The validation and improvements of the current data release are performed with the usual battery of null tests on data with different observing periods (half-rings, odd-even survey differences, odd-even year differences). In addition, we performed an exhaustive comparison of the results on such null tests obtained with the 2015 and 2018 data sets. These tests, more than any others, clearly show the improvements in the data quality, especially at 30 and 44 GHz, thanks to the new calibration scheme. The better data selection in 2018 is also able to marginally improve the quality of the 70 GHz.

For the analysis of systematic effects, we chose here not to consider the whole set of effects simulated in 2015, but select only those expected to contribute to the final calibration accuracy. As a result, the overall systematic error budget remains unchanged with respect to 2015, and we have verified that this assumption is indeed true by comparing present and past simulations with current null test spectra. Specifically, we consider the

Table 9. Additional noise in μK^2 from systematic effects for three different multipole ranges (ℓ = 4–5, $\ell \simeq 100$, and $\ell \simeq 500$).

Effect	Procedure	30 GHz	44 GHz	70 GHz
ℓ = 4–5				
Gain smoothing error	Simulations + odd-even years	$\lesssim 3\times10^{-3}$ (TT)	$\lesssim 3\times10^{-3}$ (TT)	$\lesssim 3\times10^{-3}$ (TT)
		$\lesssim 6\times10^{-3}$ (EE)	$\lesssim 6\times10^{-3}$ (EE)	$\lesssim 6\times10^{-3}$ (EE)
ADC	Simulations + half-rings	$\lesssim 1\times10^{-4}$ (TT)	$\lesssim 1\times10^{-5}$ (TT)	$\lesssim 1\times10^{-5}$ (EE)
		$\lesssim 1\times10^{-5}$ (EE)	$\lesssim 1\times10^{-6}$ (EE)	$\lesssim 1\times10^{-6}$ (EE)
Full (4π beam + dipole params)	Simulations + odd-even years	$\lesssim 7\times10^{-3}$ (TT)	$\lesssim 4\times10^{-3}$ (TT)	$\lesssim 1\times10^{-3}$ (TT)
		$\lesssim 1\times10^{-2}$ (EE)	$\lesssim 1\times10^{-2}$ (EE)	$\lesssim 6\times10^{-3}$ (EE)
$\ell \simeq 100$				
Gain smoothing error	Simulations + odd-even years	$\lesssim 1\times10^{-4}$ (TT)	$\lesssim 1\times10^{-4}$ (TT)	$\lesssim 1\times10^{-5}$ (TT)
		$\lesssim 1\times10^{-4}$ (EE)	$\lesssim 1\times10^{-4}$ (EE)	$\lesssim 1\times10^{-4}$ (EE)
ADC	Simulations + half-rings	$\lesssim 1\times10^{-6}$ (TT)	$\lesssim 1\times10^{-6}$ (TT)	$\lesssim 1\times10^{-6}$ (TT)
		$\lesssim 1\times10^{-7}$ (EE)	$\lesssim 1\times10^{-7}$ (EE)	$\lesssim 1\times10^{-7}$ (EE)
Full (4π beam + dipole params)	Simulations + odd-even years	$\lesssim 1\times10^{-4}$ (TT)	$\lesssim 1\times10^{-3}$ (TT)	$\lesssim 3\times10^{-5}$ (TT)
		$\lesssim 1\times10^{-3}$ (EE)	$\lesssim 6\times10^{-4}$ (EE)	$\lesssim 4\times10^{-4}$ (EE)
$\ell \simeq 500$				
Gain smoothing error	Simulations + odd-even years	$\lesssim 1\times10^{-5}$ (TT)	$\lesssim 1\times10^{-5}$ (TT)	$\lesssim 1\times10^{-5}$ (TT)
		$\lesssim 1\times10^{-5}$ (EE)	$\lesssim 1\times10^{-5}$ (EE)	$\lesssim 1\times10^{-5}$ (EE)
ADC	Simulations + half-rings	$\lesssim 1\times10^{-6}$ (TT)	$\lesssim 1\times10^{-6}$ (TT)	$\lesssim 1\times10^{-6}$ (TT)
		$\lesssim 1\times10^{-7}$ (EE)	$\lesssim 1\times10^{-7}$ (EE)	$\lesssim 1\times10^{-7}$ (EE)
Full (4π beam + dipole params)	Simulations + odd-even years	$\sim 1\times10^{-5}$ (TT)	$\lesssim 6\times10^{-5}$ (TT)	$\lesssim 7\times10^{-6}$ (TT)
		$\lesssim 3\times10^{-4}$ (EE)	$\lesssim 4\times10^{-4}$ (EE)	$\lesssim 2\times10^{-4}$ (EE)

ADC nonlinearity effect, the impact of the gain smoothing algorithm, and the impact of parameters in both the 4π beams and the direction and amplitude of the solar dipole. The end result is that for multipoles $\ell \gtrsim 100$, the impact of systematic effects is well described by simulations, and is well-accounted for in the overall error budget. At large angular scales there are still a very few multipoles that show a noise excess with respect to simulations. It is important to note that our power spectrum analysis is performed on a masked sky with a pure pseudo-C_ℓ approach, which is known to be sub-optimal at very low multipoles. Therefore our analysis should be regarded as an overall indication of the amplitude of systematic effects. The 2018 release includes our end-to-end simulations, which allow those interested in systematic effects to create templates of the various systematics, similar to the gain-correction template described in Sect. 3.2.

In conclusion, we have demonstrated substantial improvement in the calibration of the LFI data over previous releases, achieving an overall calibration accuracy of 0.1–0.2 %. We have provided a comprehensive description of the uncertainties, including systematic effects. Additional improvements are still possible, and can be anticipated in the future.

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Appendix A: Comparison of LFI 30 GHz with WMAP K and Ka bands

In order to provide a further check on systematic effects, we present here a comparison in polarization at large scales between the LFI 30-GHz channel and the K and Ka bands from WMAP (Bennett et al. 2013). As a first step we subtract from the *Planck* 2015 and 2018 maps the corresponding bandpass corrections. We then filter *Planck* and WMAP maps with a 10°-FWHM Gaussian beam in order to suppress high-frequency noise and highlight the large-scale structures. We rescale the K and Ka bands to the LFI 30-GHz effective frequency, assuming synchrotron emission with a spectral index of −3. Finally, we calculate Q and U differences between the *Planck* 2015 and 2018 maps and the WMAP maps, as shown in Figs. A.1 and A.2.

The amplitude of the large-scale structure in the difference maps decreases by roughly a factor of two between the 2015 and the 2018 maps, reflecting the improvements in the new calibration scheme. This improved calibration has reduced or eliminated some of the features noted by Weiland et al. (2018) in their comparison of WMAP and *Planck* 2015 Q and U maps. However, as pointed out in Sect. 3.2, the calibration could not be run to convergence for practical reasons, and we expect residuals in the 2018 maps with a pattern similar to that of the 2015 maps. The fact that the difference maps shown in Figs. A.1 and A.2 have some common features with the two *Planck* releases is therefore not a surprise. There is also some agreement between features in the difference maps and the 70-GHz gain correction template (Fig. 6), which is based in part on the 30-GHz maps. Despite the substantial improvement in the 2018 30-GHz

polarization maps, residual systematics are still present, and included in our estimated error budget. Additional improvements are planned in a future paper.

Appendix B: Simulations of systematic effects

B.1. Input Sky Model

The input FFP10 sky includes the following components:

- Galactic thermal dust, spinning dust, synchrotron, free-free, and CO line emission;
- the cosmic infrared background;
- Galactic and extragalactic point sources (radio and infrared);
- thermal and kinematic Sunyaev Zeldovich effects from Galaxy clusters; and
- the CMB.

B.1.1. Thermal dust

Galactic thermal dust emission is modelled by scaling across frequencies a polarized template of emission at 353 GHz. The intensity map has been obtained using the Generalized Needlet ILC (GNILC) method of Remazeilles et al. (2011), applied to the 2015 (PR2) release *Planck* HFI maps, as described in Planck Collaboration Int. XLVIII (2016). The 353-GHz *Planck* dust map obtained in this way is colour-corrected to obtain a template at 353 GHz. An overall offset of 0.13 MJy sr^{-1} (corresponding to the residual CIB monopole) is subtracted from the GNILC dust map.

Polarization templates are generated using the prescription of Vansyngel et al. (2017) to generate random dust polarization maps. The method relates the dust polarization sky to the structure of the Galactic magnetic field that is responsible for aligning elongated emitting dust grains, building on the analysis of dust polarization properties described in Planck Collaboration Int. XLIV (2016). The Galactic magnetic field is modelled as a three-dimensional superposition of a mean uniform field and a Gaussian random (turbulent) component with a power-law power spectrum of exponent −2.5. Polarization maps are obtained from the superposition of emission from four emitting layers that all share the same intensity template.

In the vicinity of the Galactic plane, the simulated polarization data are replaced by real *Planck* 353-GHz data. The transition between real data and simulations is made using a Galactic mask with a 5° apodization (which leaves 68 % of the sky unmasked), taken from the set of *Planck* common Galactic masks available in the *Planck* explanatory supplement. [6] The first harmonic multipoles, corresponding to $\ell < 10$, come from the 353-GHz polarized sky map.

The scaling in frequency of the dust templates uses the prescription used for FFP8. A different modified black-body emission law is used in each of the $N_{\rm side} = 2048$ pixels. In addition, the dust spectral index used for scaling in frequency is different at frequencies above and below 353 GHz: for frequency above 353 GHz the modified black-body parameters obtained in Planck Collaboration Int. XLVIII (2016) are used; below 353 GHz the temperature map is unchanged, but a map of emissivities computed as described in Planck Collaboration Int. XVII (2014) is used instead.

[6] http://wiki.cosmos.esa.int/planckpla2015/index.php/Frequency_Maps#Masks

B.1.2. Other Galactic emission

Synchrotron intensity is modelled by scaling in frequency the 408-MHz template map from Haslam et al. (1982), as reprocessed by Remazeilles et al. (2015), using one single power law per pixel. The pixel-dependent spectral index is derived in the analysis of the WMAP data by Miville-Deschênes et al. (2008). The generation of synchrotron polarization follows the prescription of the original PSM paper (Delabrouille et al. 2013).

Free-free, spinning dust models, and Galactic CO emission are essentially the same as used for FFP8 simulations (Planck Collaboration XII 2016), but the actual synchrotron and free-free maps used for FFP10 are obtained with a different realization of small-scale fluctuations of the intensity. CO maps do not include small-scale fluctuations, and are generated from the spectroscopic survey of Dame et al. (2001). None of these three components is polarized in the FFP10 simulations.

B.1.3. Cosmological parameters

The generation of all extragalactic components in the FFP10 sky depends on the assumed cosmological scenario, and in particular values for all cosmological parameters that impact the CMB power spectra (see Planck Collaboration XIII 2016) and/or the statistical distribution of extragalactic objects. We use the parameter values listed in Table B.1, which are used as inputs to the Cosmic Linear Anisotropy Solving System (CLASS) code (Blas et al. 2011; Dio et al. 2013) for generating CMB TT, EE, BB and lensing Φ power spectra, as well as density-contrast shells in auto- and cross-spectra at a set of redshifts (used to model CIB emission). We assume 2.0328 massless neutrinos, and one massive neutrino with a mass of 60 meV, compatible with a standard neutrino hierarchy. The pivot scale for the scalar perturbations is 0.002, and the tensor spectral index and the running of the scalar spectral index are zero.

Table B.1. Cosmological parameters adopted in the PSM.

$T_{\rm CMB}$[K]	2.7255
h	0.6701904
$\Omega_{\rm m}$	0.26782
$\Omega_{\rm b}$	0.0493498
$\Omega_{\rm k}$	0
σ_8	0.8162
$n_{\rm s}$	0.9636
r	0.01
τ	0.060
$Y_{\rm P}$	0.2453
w	−1
$10^9 A_{\rm s}$	2.119

B.1.4. CMB

The CMB is a stationary Gaussian random field on the sphere, generated from CLASS output power spectra using the HEALPix package, and lensed using the Ilens software of Basak et al. (2009). The joint generation of density contrast shells used for the CIB simulation and of the lensing potential map provides correlation between the CMB lensing and the CIB maps. This correlation has been used to generate 25 independent CMB and CIB realizations.

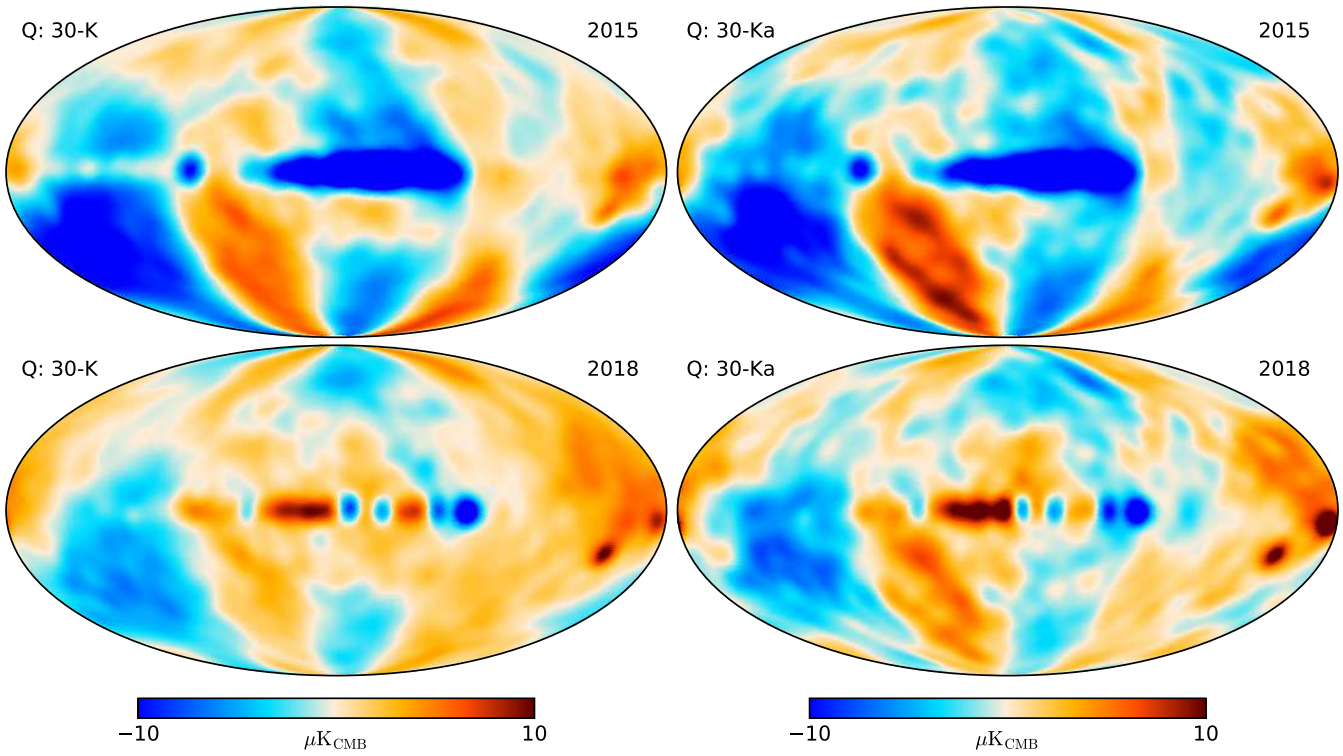


Fig. A.1. Stokes Q difference maps between *Planck* 30 GHz 2015 (top) and 2018 (bottom) and WMAP K-band (left) and Ka-band (right). The WMAP K and Ka band maps are rescaled to match the *Planck* 30-GHz effective frequency assuming synchrotron emission with a spectral index of −3. All maps used in this comparison have been smoothed with a 10° Gaussian beam.

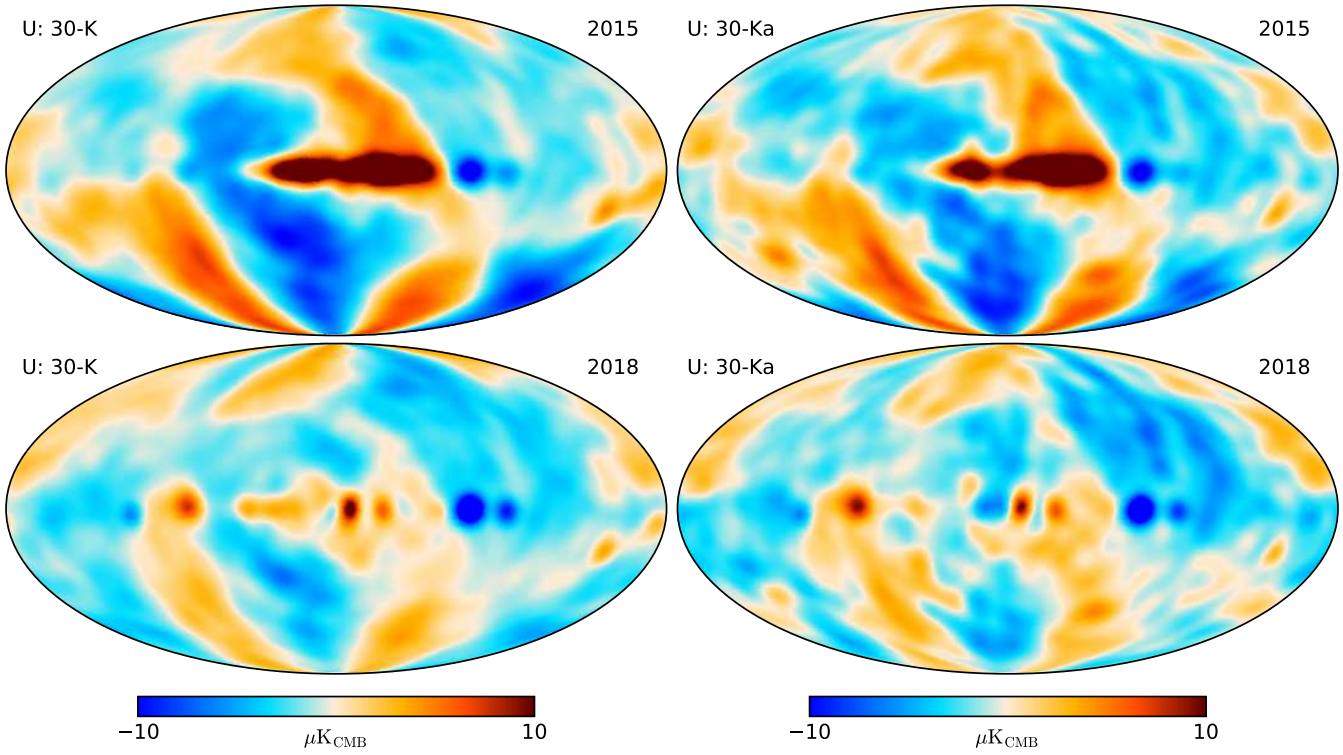


Fig. A.2. Same as in Fig. A.1, for Stokes U.

B.1.5. Unresolved sources and the cosmic infrared background

Catalogues of individual radio and low-redshift infrared sources are generated in the same way as for the FFP8 simulations (Planck Collaboration XII 2016), but use a different seed for random number generation. The generation of the cosmic infrared background (CIB), due to the integrated emission of tens of billions of distant dusty galaxies, was substantially revised to al-

low for the simulation of correlations between the lensing potential maps and the CIB emission. Number counts for three types of galaxies — early-type proto-spheroids, along with spiral and starburst galaxies — are based on the model of Cai et al. (2013). The entire Hubble volume up to $z = 6$ is cut into 64 spherical shells. For each shell, we generate a map of density contrast integrated along the line of sight between $z_{\rm min}$ and $z_{\rm max}$, such that the statistics of these density contrast maps (power spectrum of linear density fluctuations, and cross-spectra between adjacent shells and with the CMB lensing potential) agree with those computed by `CLASS`. For each type of galaxy (spiral, starburst, proto-spheroid), a catalogue of randomly-generated galaxies is generated for each shell, following the appropriate number counts. These galaxies are then distributed in the shell to generate a single intensity map at a given reference frequency, which is scaled across frequencies using the prototype galaxy spectral energy distribution at the appropriate redshift.

B.1.6. Galaxy clusters

A full-sky catalogue of galaxy clusters is generated based on number counts, following the method of Delabrouille et al. (2002). The mass function of Tinker et al. (2008) is used to predict number counts. Clusters are distributed in redshift shells proprotionally to the denstity contrast in each pixel with a bias $b(z, M)$ in agreement with the linear bias model of Mo & White (1996). For each cluster, we assign a universal profile based on XMM observations, as described in Arnaud et al. (2010). Relativistic corrections are included to first order, following the expansion of Wang & Steinhardt (1998). We use a mass bias $M_{\rm X-Ray}/M_{\rm mass-fn} = 0.63$ to match actual cluster number counts observed by *Planck* for the best-fit cosmological model based on CMB observations.

The kinematic SZ effect is computed assigning to each cluster a radial velocity that is randomly drawn from a centred Gaussian distribution, with a redshift dependent standard deviation computed from the power spectrum of density fluctuations. This neglects correlations between cluster motions such as bulk flows or pairwise velocities of nearby clusters.

B.2. Bandpass integration

The model of sky emission is finally integrated in frequency according to *Planck* bandpasses from the HFI and LFI Reduced Instrument Models (RIMO, version R2.00 for HFI and R2.50 for LFI).

B.3. From TOI to gains and maps

The LFI systematic effect simulations are done partially at timeline and partially at ring-set level, with the goal of being as modular as possible, in order to create a reusable set of simulations. From the input sky model and according to the pointing information, we create single-channel ring-sets of the pure sky convolved with a suitable instrumental beam. To these we add pure noise (white and $1/f$) ring-sets generated from the noise power spectrum distributions measured from real data one day at a time. The overall scheme is given in Fig. B.1.

In the same manner, we create ring-sets for each of the specific systematic effects we would like to measure. We add together signal, noise, and systematic ring-sets, and, given models for straylight (based on the `GRASP` beams) and the orbital dipole, we create "perfectly-calibrated" ring-sets (calibration constant

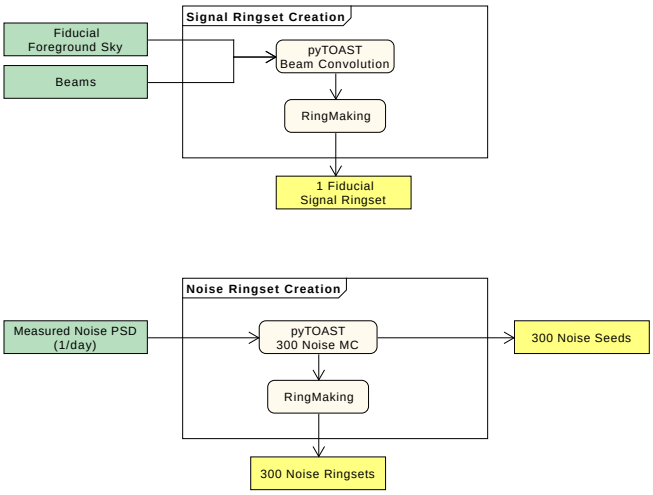


Fig. B.1. Signal and noise ring-set creation pipelines.

= 1). We use the gains estimate from the 2018 data release to "de-calibrate" these timelines, in other words, to convert them from kelvins to volts. At this point the calibration pipeline starts, and produces the reconstructed gains that will be different from the ones used in the de-calibration process due to the presence of simulated systematic effects. The calibration pipeline is algorithmically exactly the same as that used at the DPC for product creation, but with a different implementation (based principally on `python`). The gain-smoothing algorithm is the same as used for the data, and has been tuned to the actual data. This means that there will be cases where reconstructed gains from simulations differ significantly from the input ones. We have verified that this indeed happens, but only for very few pointing periods, and we therefore decided not to consider them in the following analysis. The overall process for estimating gains is given in Fig. B.2.

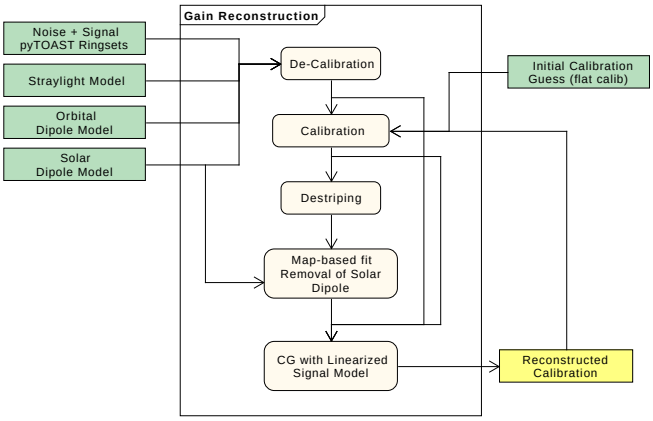


Fig. B.2. Gain reconstruction pipeline.

At this point we are able to generate maps for full mission, half-ring, and odd-even-year splits) that include the effects of systematic errors on calibration. In the final step, we produce timelines (which are never stored) starting from the same fiducial sky map, using the same model for straylight and the orbital dipole as in the previous steps, and from generated noise-only timelines created with the same seeds and noise model used be-

fore. We then apply the official gains to "de-calibrate" the timelines, which are immediately calibrated with the reconstructed gains in the previous step. The nominal destriping mapmaking algorithm is then used to create final maps. The complete data flow is given in Fig. B.3.

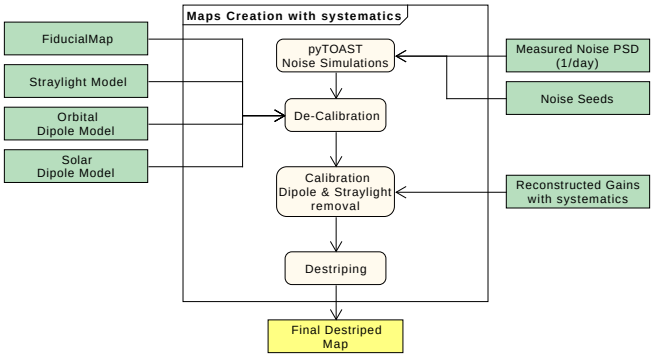


Fig. B.3. Final simulation step, where calibrated maps are created from reconstructed gains, including the impact of systematic effects.

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